

A brief overview of Efimov physics and related mathematical problems

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Review

Efimov physics: a review

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Abstract

This article reviews theoretical and experimental advances in Efimov physics, an array of quantum few-body and many-body phenomena arising for particles interacting via short-range resonant interactions, that is based on the appearance of a scale-invariant three-body attraction theoretically discovered by Vitaly Efimov in 1970. This three-body effect was originally proposed to explain the binding of nuclei such as the triton and the Hoyle state of carbon-12, and later considered as a simple explanation for the existence of some halo nuclei. It was subsequently evidenced in trapped ultra-cold atomic clouds and in diffracted molecular beams of gaseous helium. These experiments revealed that the previously undetermined three-body parameter introduced in the Efimov theory to stabilise the three-body attraction typically scales with the range of atomic interactions. The few- and many-body consequences of the Efimov attraction have been since investigated theoretically, and are expected to be observed in a broader spectrum of physical systems.

quant-ph] 7 Apr 2017

Plan

1. The Efimov effect for three particles
 1. A physical description
 2. Physical examples (experiments)
 3. A mathematical description
2. More than three particles
 1. Universal tetramers (4 identical bosons)
 2. The four-body Efimov effect (3 fermions + 1)

Physical description

What is the Efimov effect?



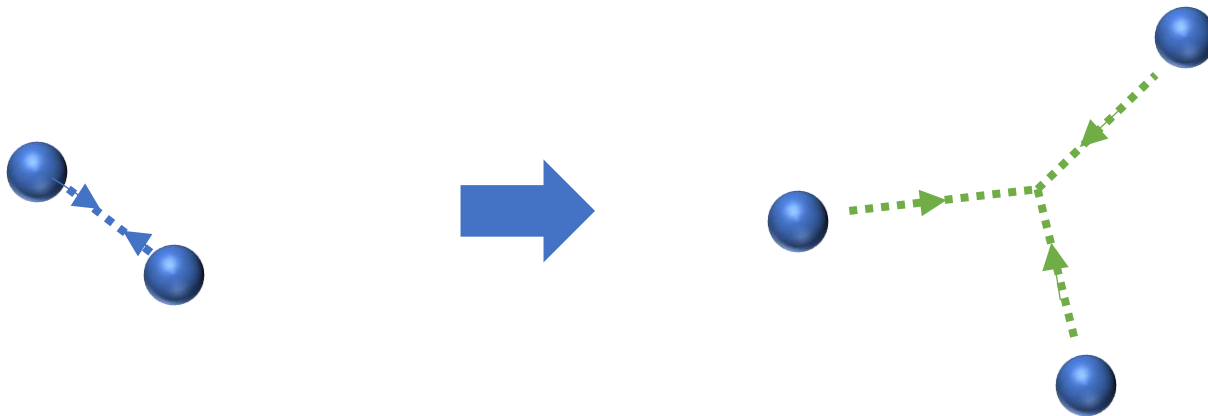
Vitaly Efimov in 1977

V. Efimov, “Weakly-bound states of three resonantly-interacting particles,” **Yad. Fiz.**, **12**, 1080–1091, November 1970, [**Sov. J. Nucl. Phys.** **12**, 589-595 (1971)].

V. Efimov, “Energy levels arising from resonant two-body forces in a three-body system.” **Physics Letters B**, **33**, 563 – 564, 1970.

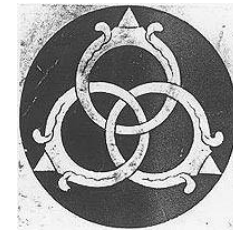
What is the Efimov effect?

The appearance of an effective **long-range three-body attractive force** between three particles interacting via **short-range two-body attractive forces**

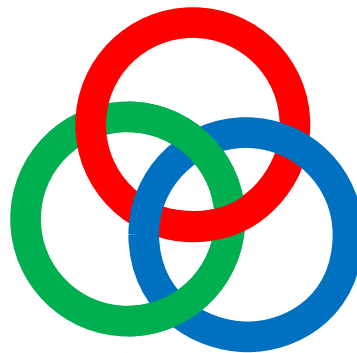


This leads to remarkable properties:

Existence of **an infinite number of Borromean states**,
i.e. three-body bound states in the absence of two-
body bound state.

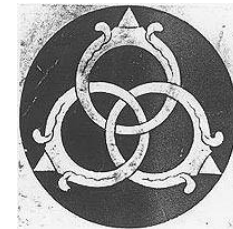


Borromeo family emblem

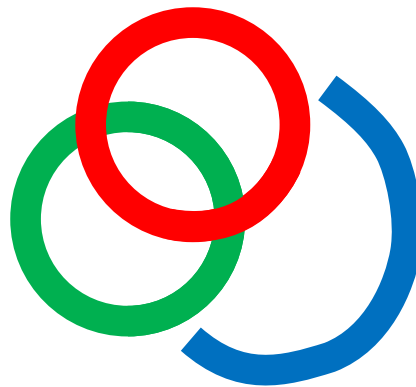


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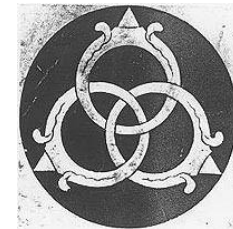


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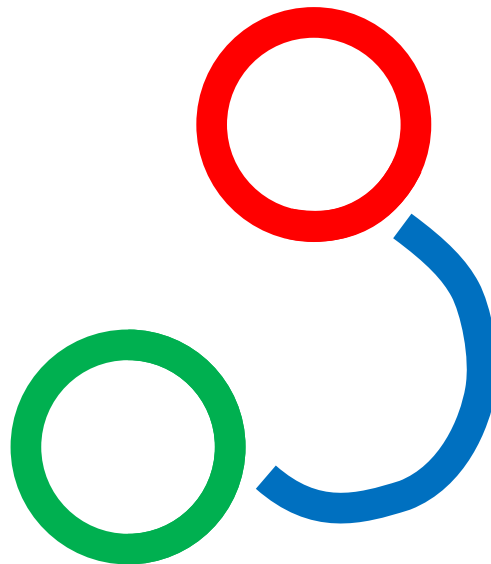


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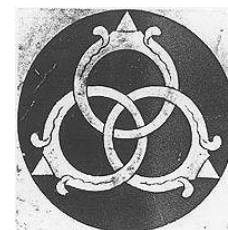


Borromeo family emblem



This leads to remarkable properties:

Existence of **an infinite number of Borromean states**, i.e. three-body bound states in the absence of two-body bound state.



Borromeo family emblem

Discrete scale invariance of the spectrum: each three-body state can be obtained from another by a scale transformation



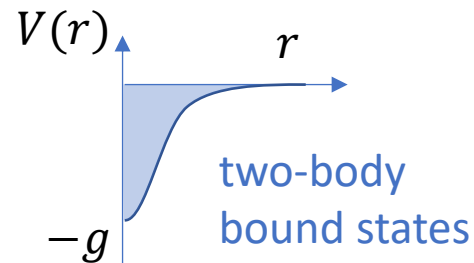
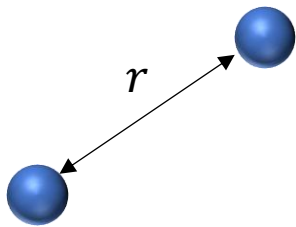
Matryoshka (Russian nesting dolls)

Universality of the physical properties: they do not depend on the details of the two-body force, but **only on two parameters**.

What is the origin of the Efimov effect?

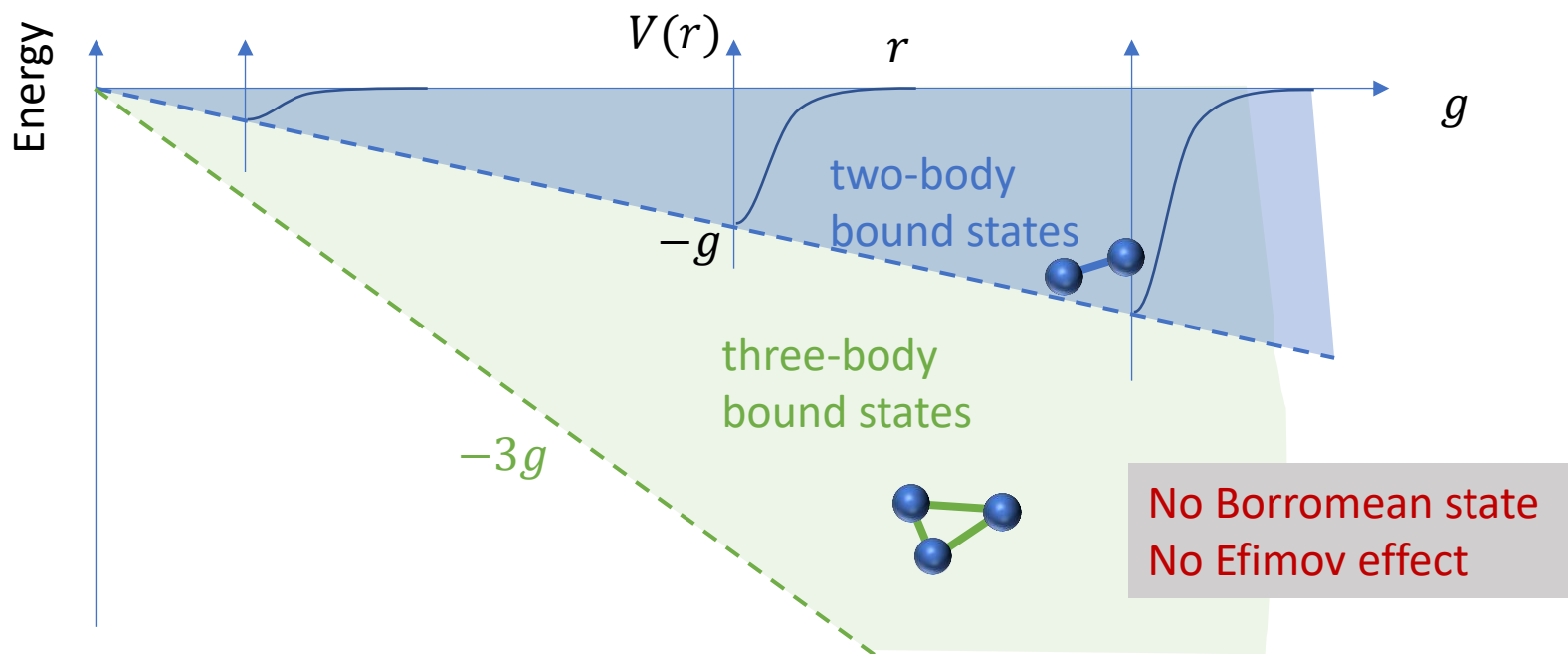
- Is there an Efimov effect in classical physics?

Short-range attractive two-body interaction potential



What is the origin of the Efimov effect?

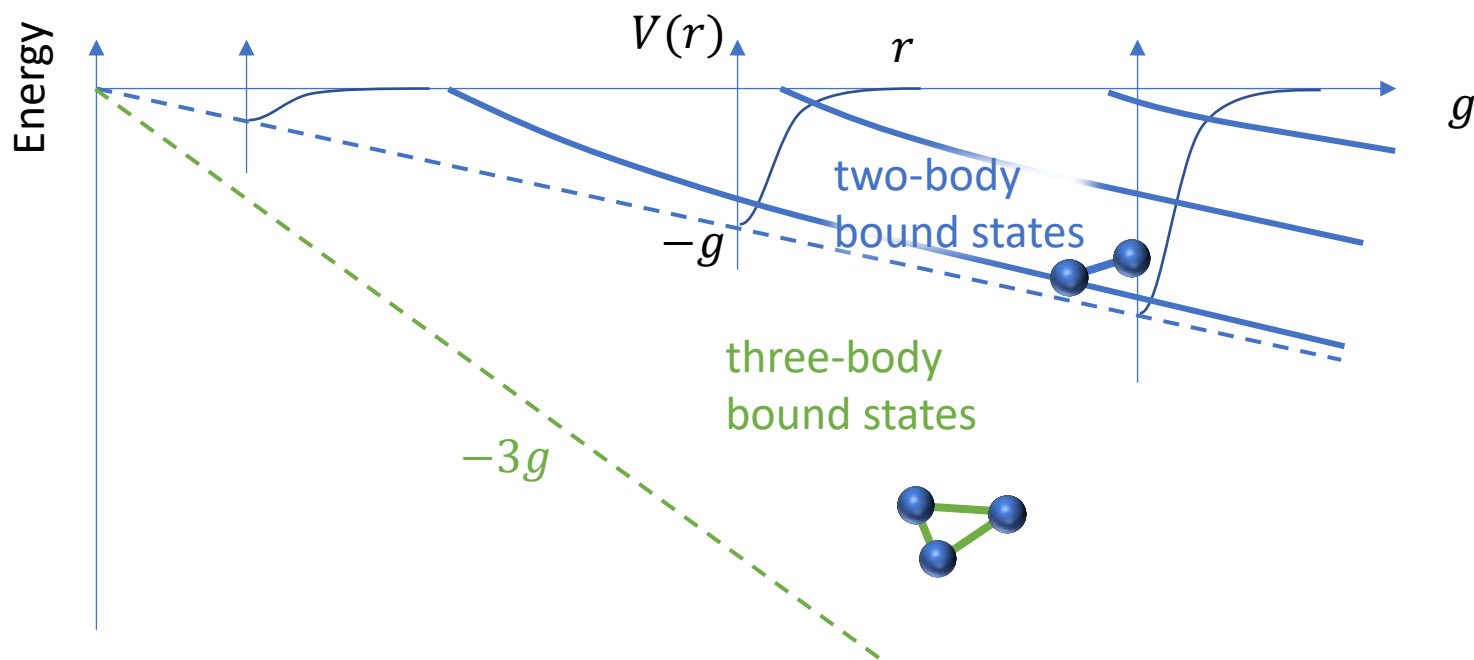
- Is there an Efimov effect in classical physics?



What is the origin of the Efimov effect?

- The quantum origin of the Efimov effect

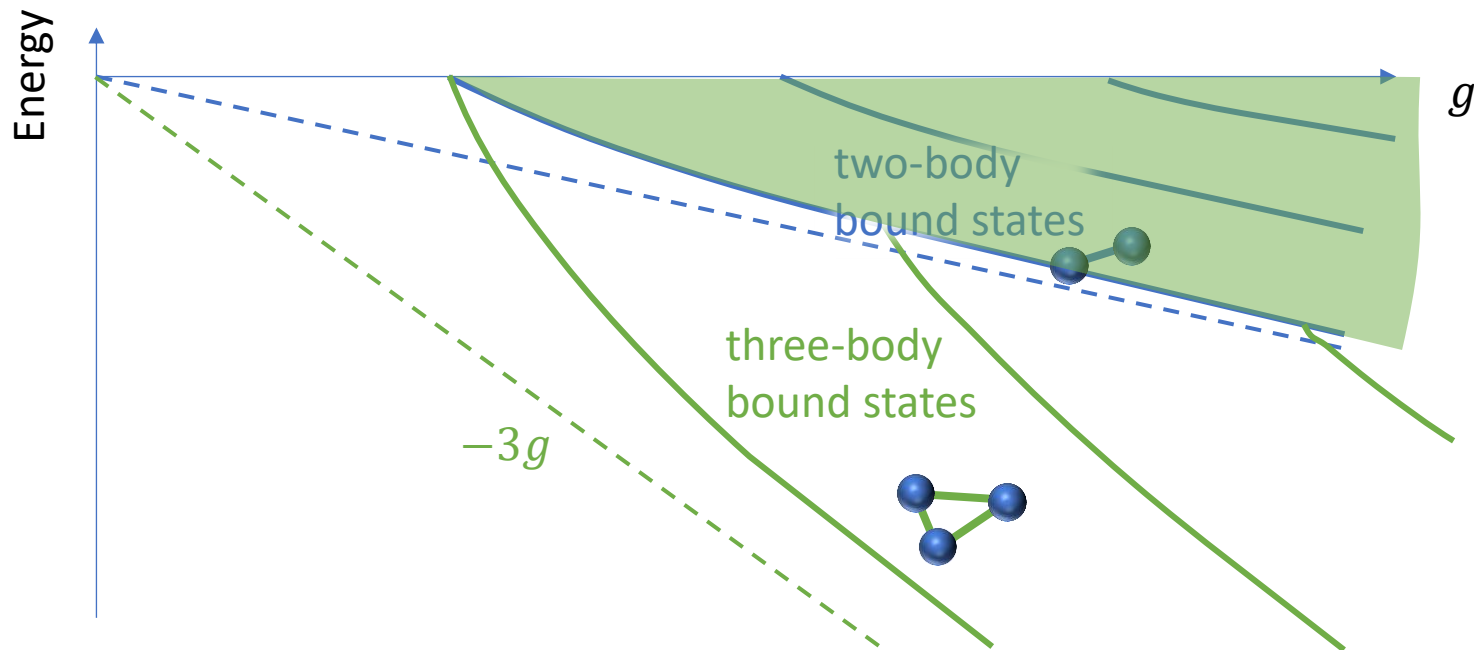
Quantum fluctuations + interactions \Rightarrow critical strength g (zero-point) and quantisation



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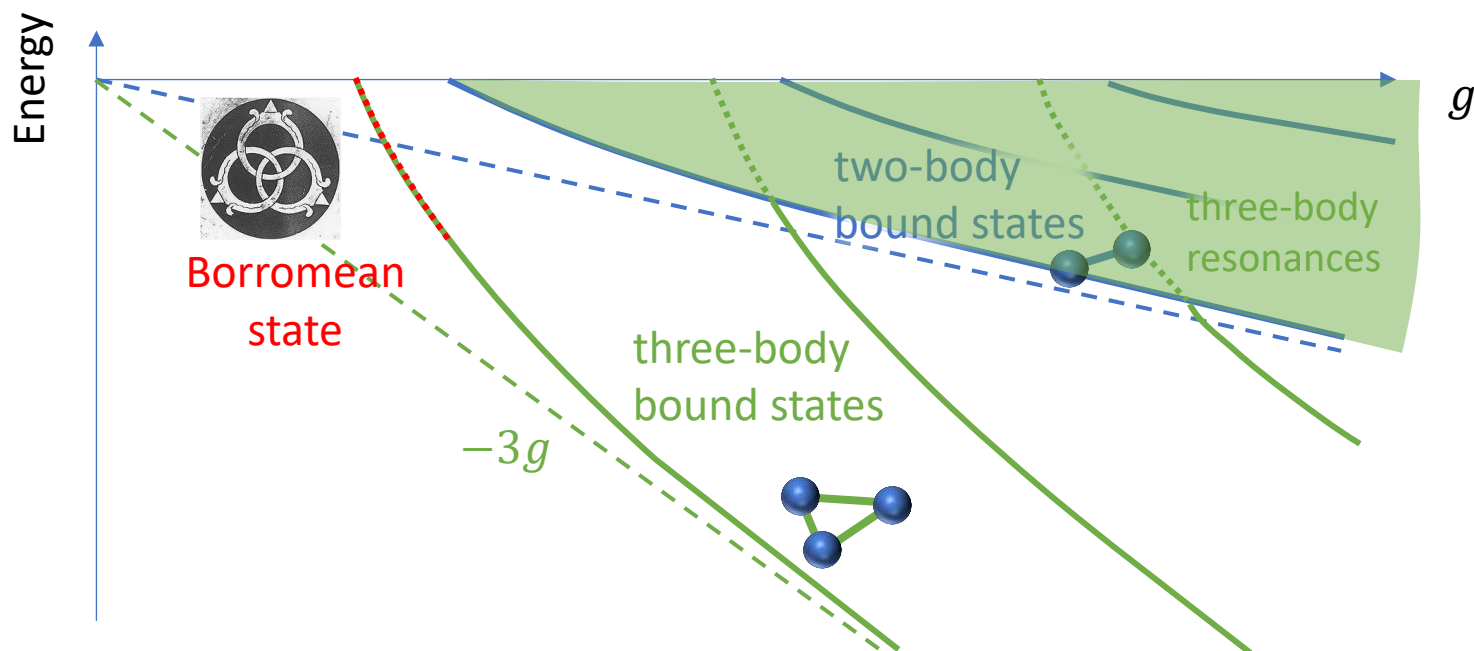
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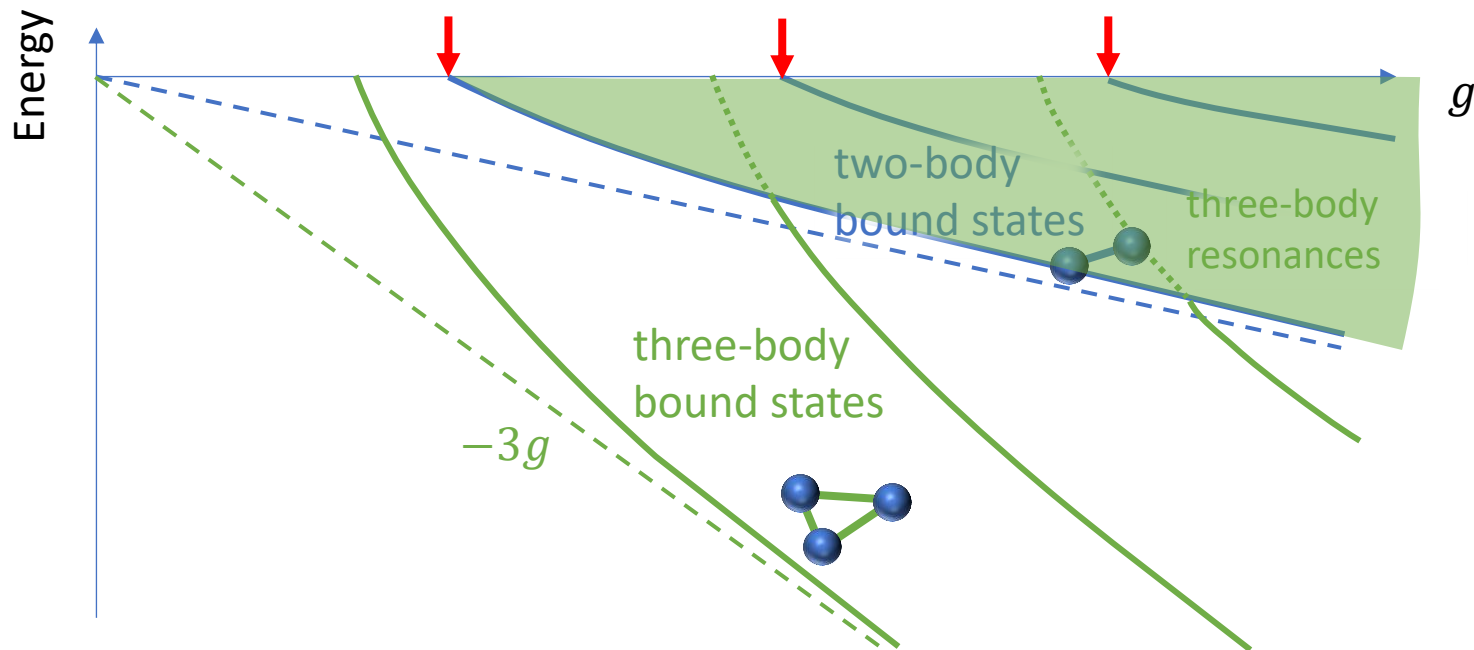
Quantum fluctuations + interactions \Rightarrow critical strength g (zero-point) and quantisation



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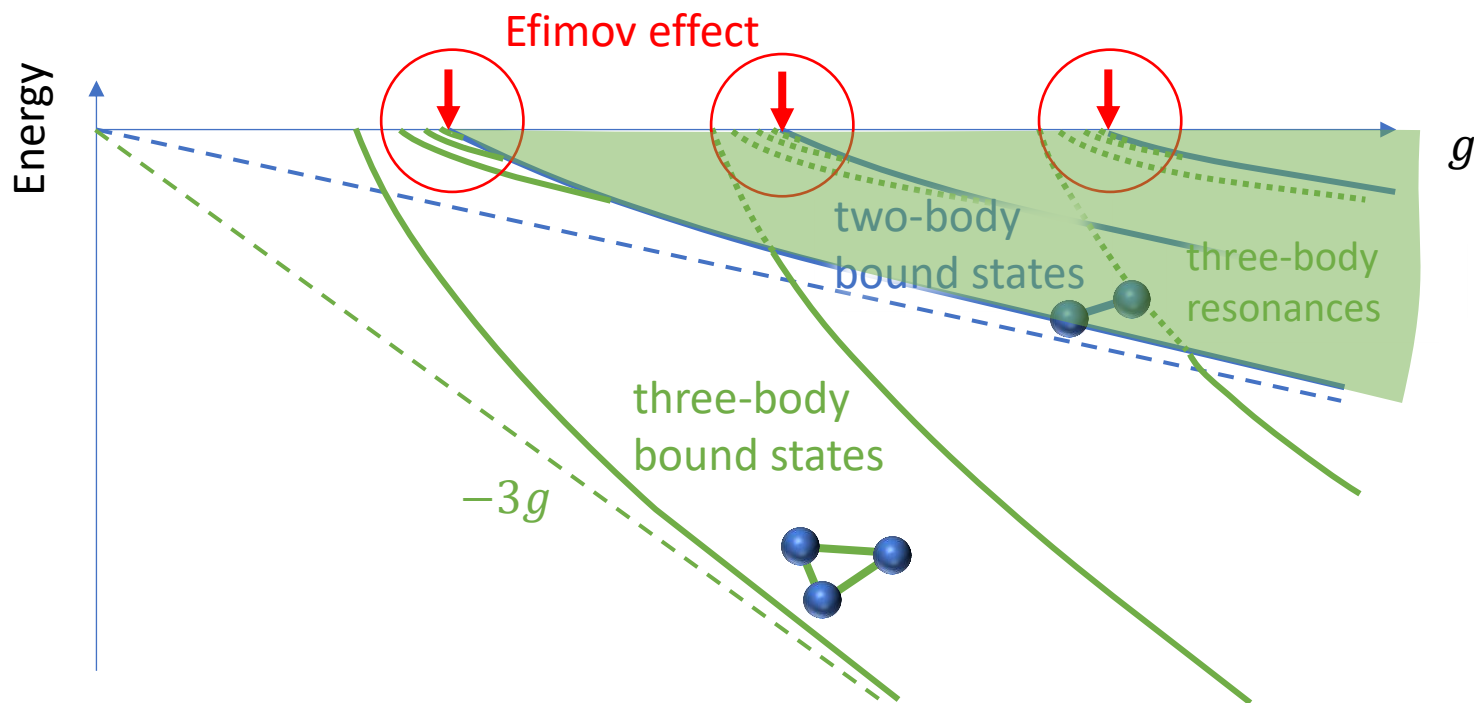
Critical strengths g for binding two particles \Rightarrow “unitarity”



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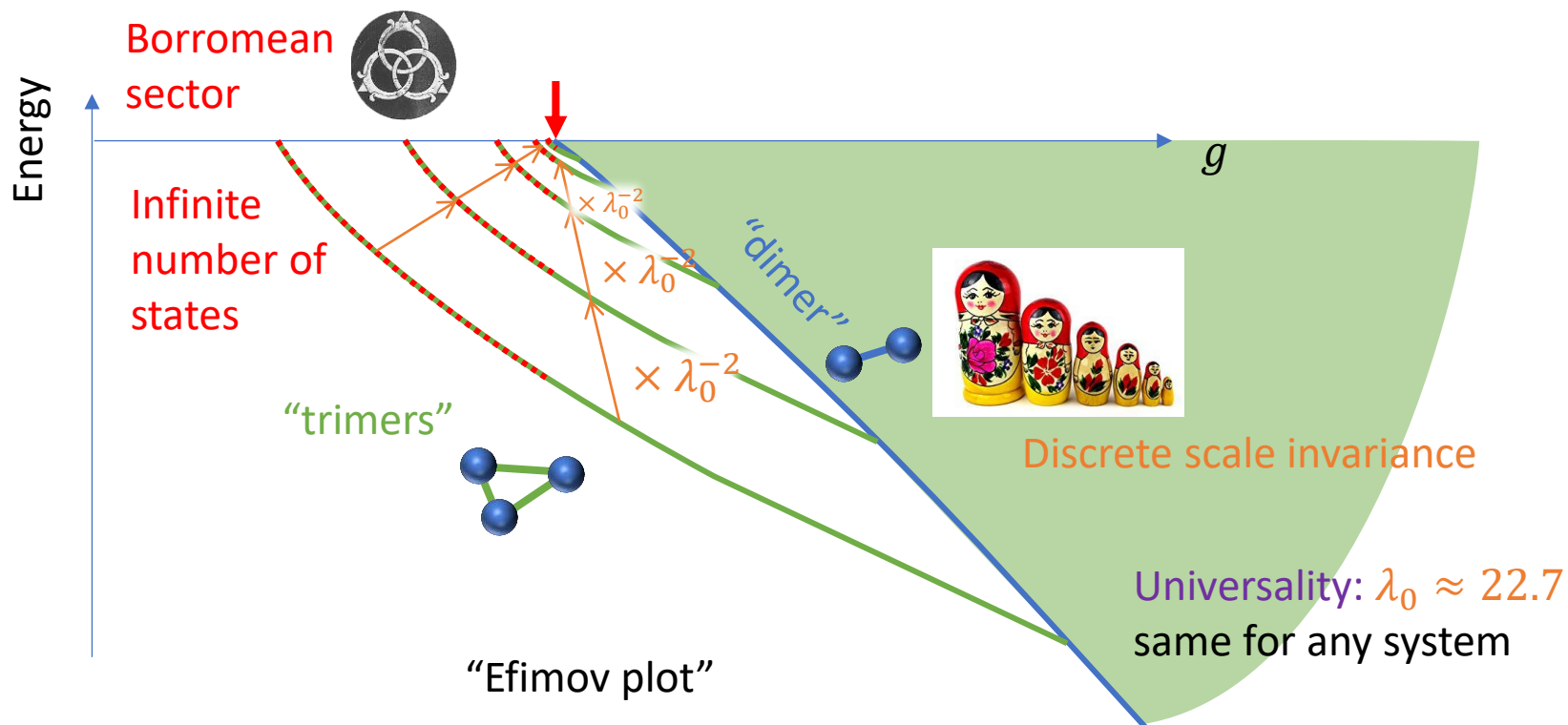
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Critical strengths g for binding two particles \Rightarrow “unitarity”



What is the origin of the Efimov effect?

- The quantum origin of the Efimov effect

Why universality and scale invariance?

At the two-body level:

Short-range interaction



At low-energy



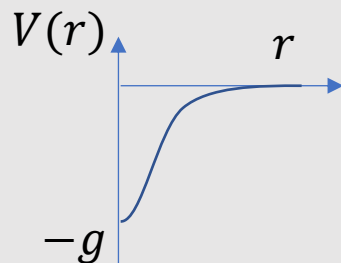
At “unitarity” ($a \rightarrow \infty$)

Single physical scale:
“Scattering length” a

No more length scale!
(Scale invariance)



Original interaction

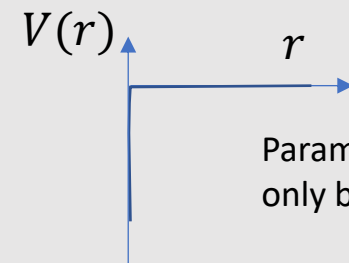


“Low-energy universality”
“Short-range universality”
“Zero-range universality”

The wave function is much larger than the range of interactions

can be replaced by

“Zero-range” interaction




Parametrised only by a

What is the origin of the Efimov effect?

- The quantum origin of the Efimov effect

Why infinitely many states?

At the three-body level:

Effective three-body potential $V(R) \propto -\frac{\hbar^2}{m} \frac{1}{R^2}$  Scales like the kinetic energy $K \propto \frac{\hbar^2}{m} \frac{d^2}{dR^2}$



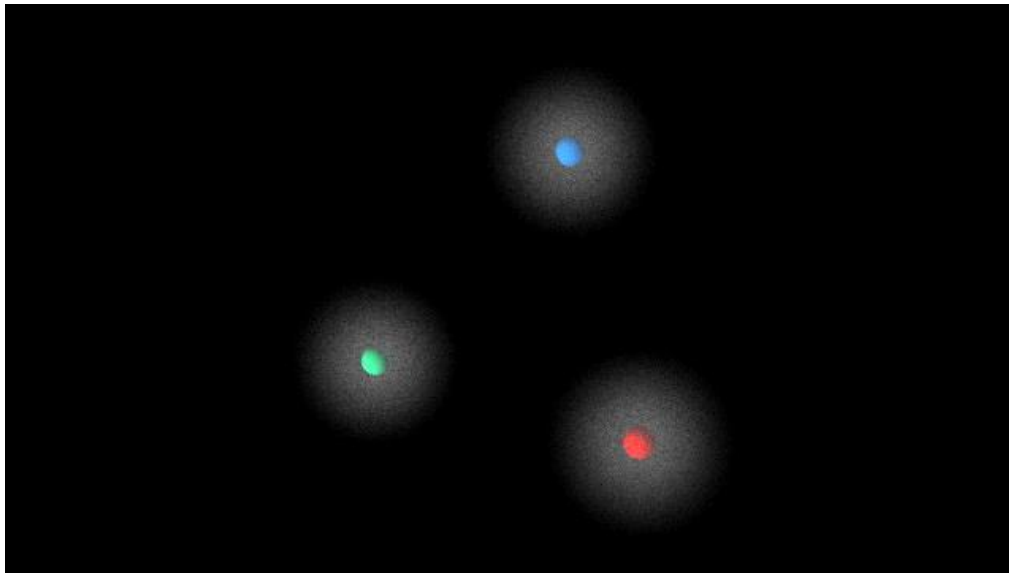
- Scale invariant equation
- Long-range attractive potential, may support **infinitely many states**

What is the origin of the Efimov effect?

- The quantum origin of the Efimov effect

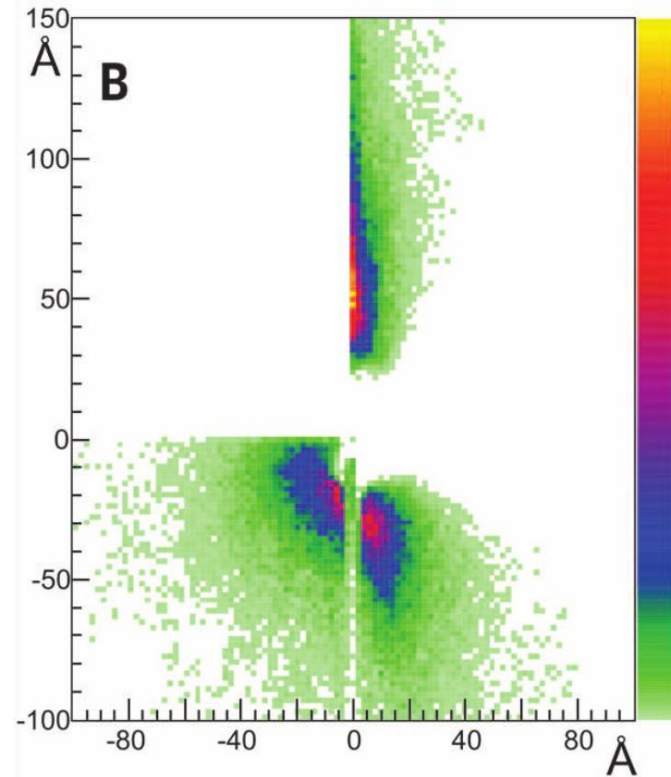
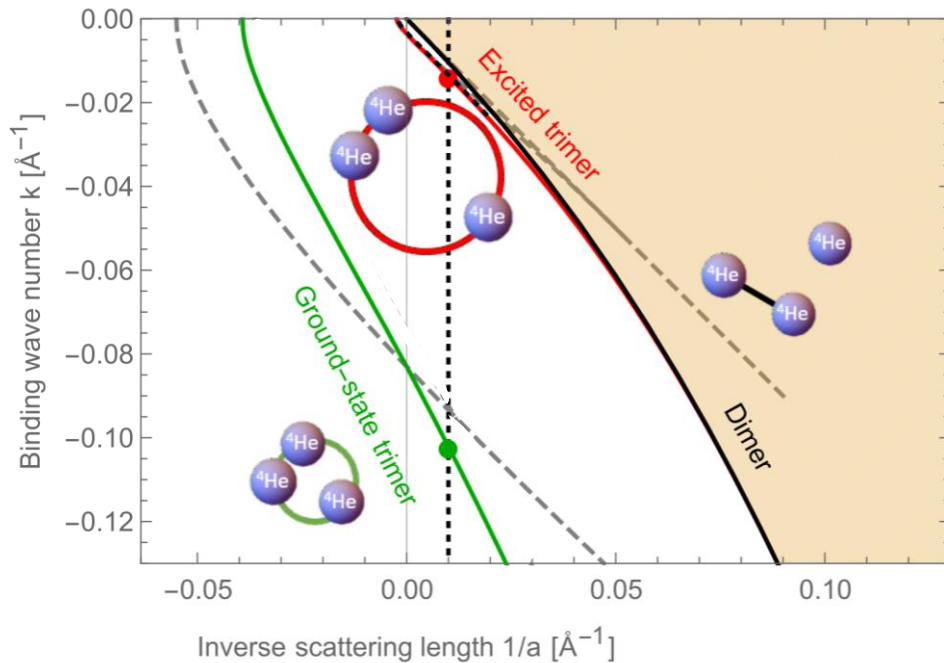
Why long range ? in spite of short-range two-body interactions?

The Efimov attraction may be viewed as an interaction between two particles mediated by a third particle



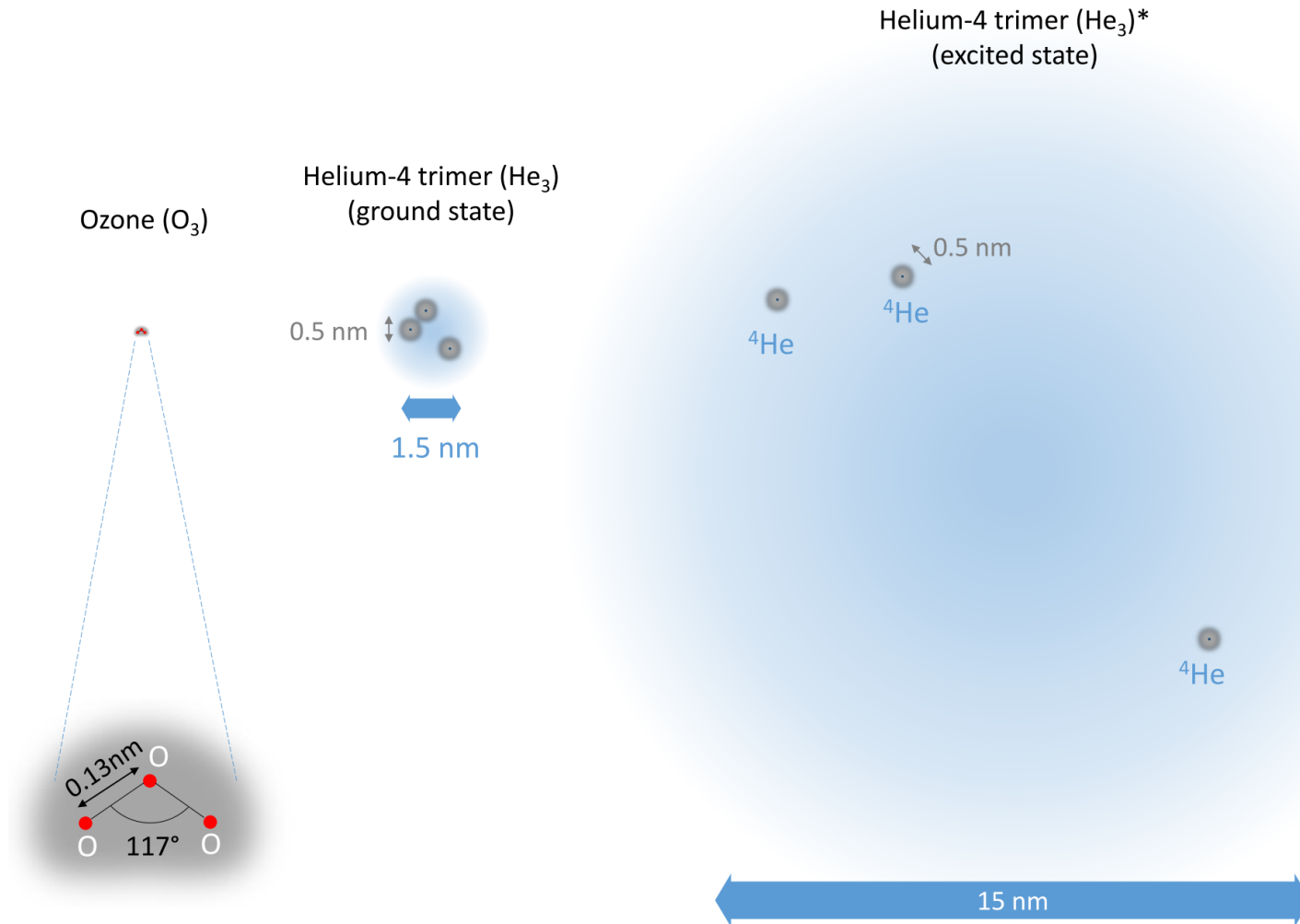
Physical examples

The helium triatomic molecules ${}^4\text{He}_3$

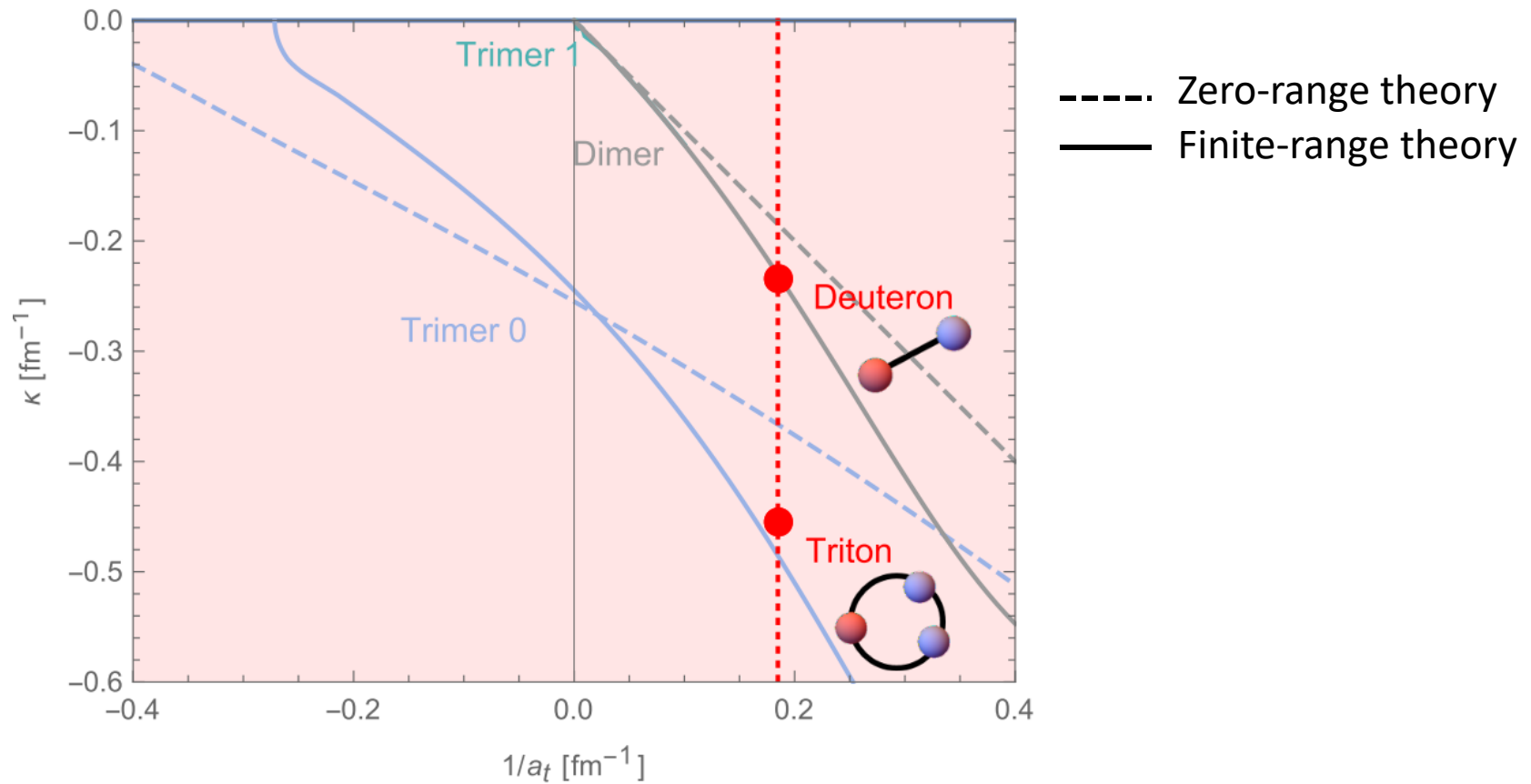


Kunitski, Science 348, 551 (2015)

The helium triatomic molecules ${}^4\text{He}_3$



The triton (2 neutrons + 1 proton)



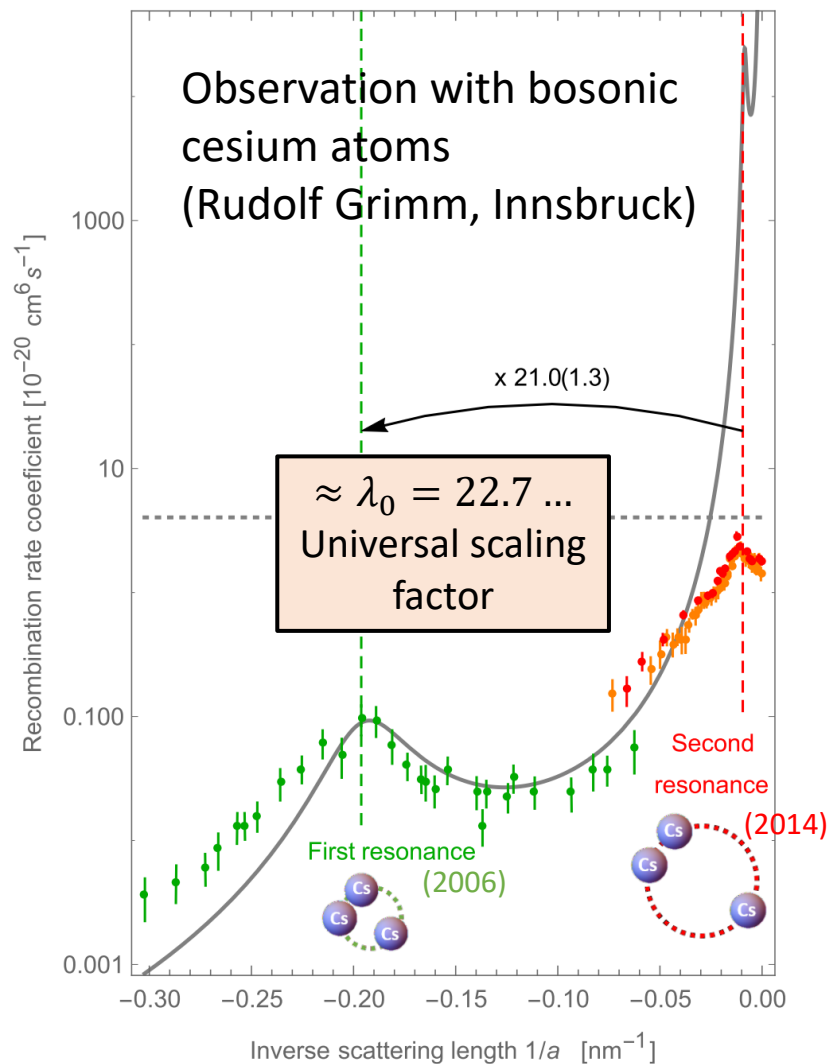
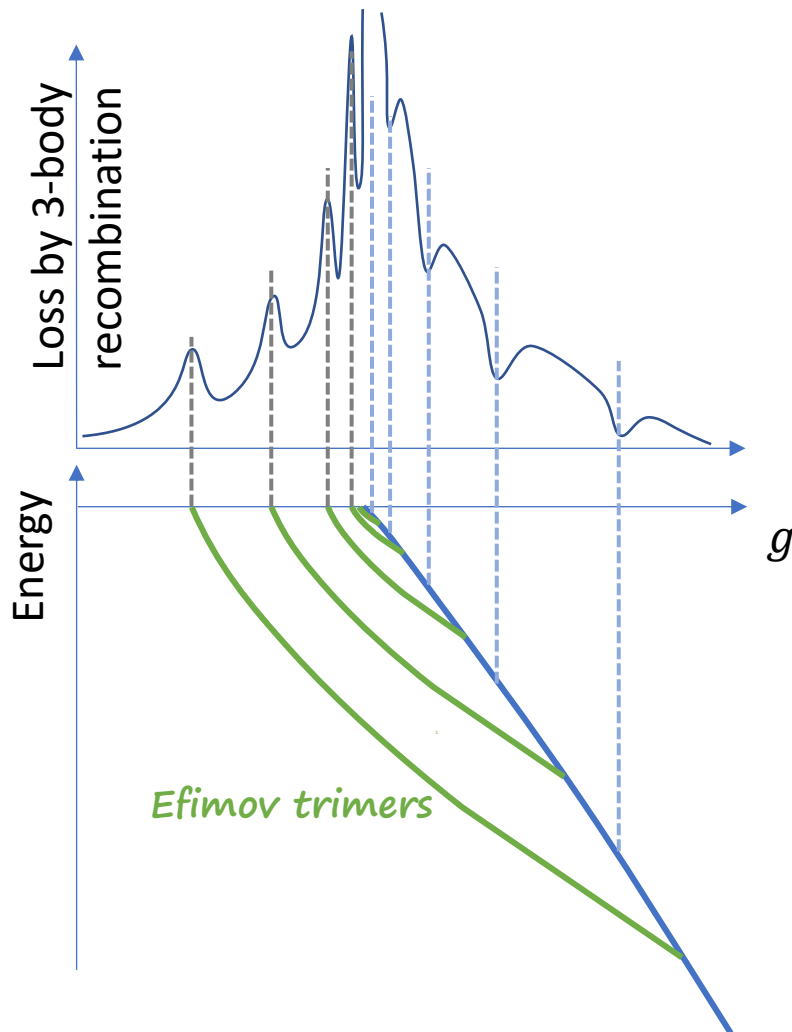
Observations in ultra-cold atomic gases

Ultra-cold atoms are gases of atoms cooled to extremely low temperatures ($< \mu\text{K}$) to reach the quantum degeneracy regime.

Magic tool: Feshbach resonances

By applying a magnetic field, it is possible to change the strength of interatomic interactions!

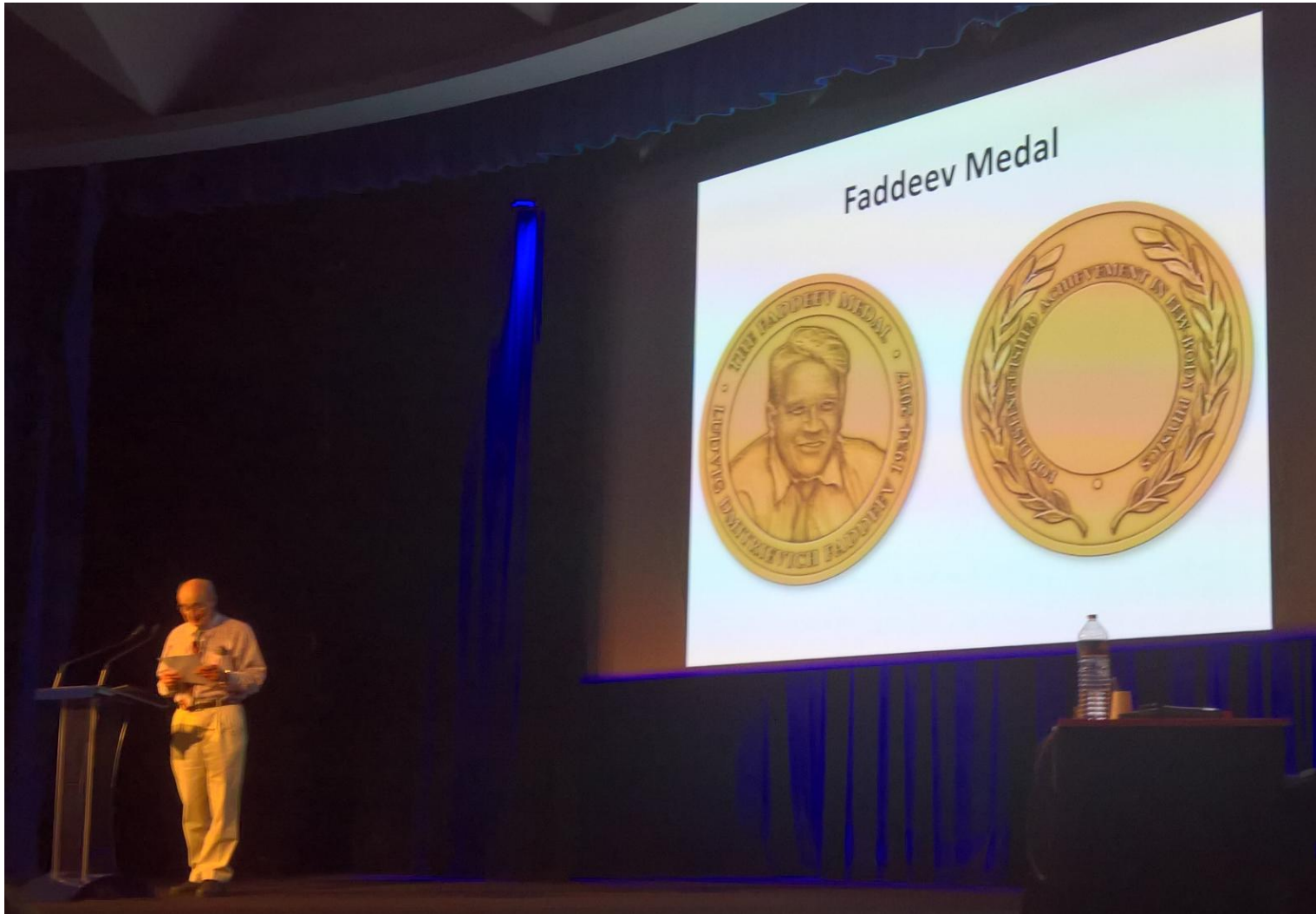
Observations in ultra-cold atomic gases



Vitaly Efimov and Rudolf Grimm receive the first Faddeev medal in Caen (July 11, 2018)

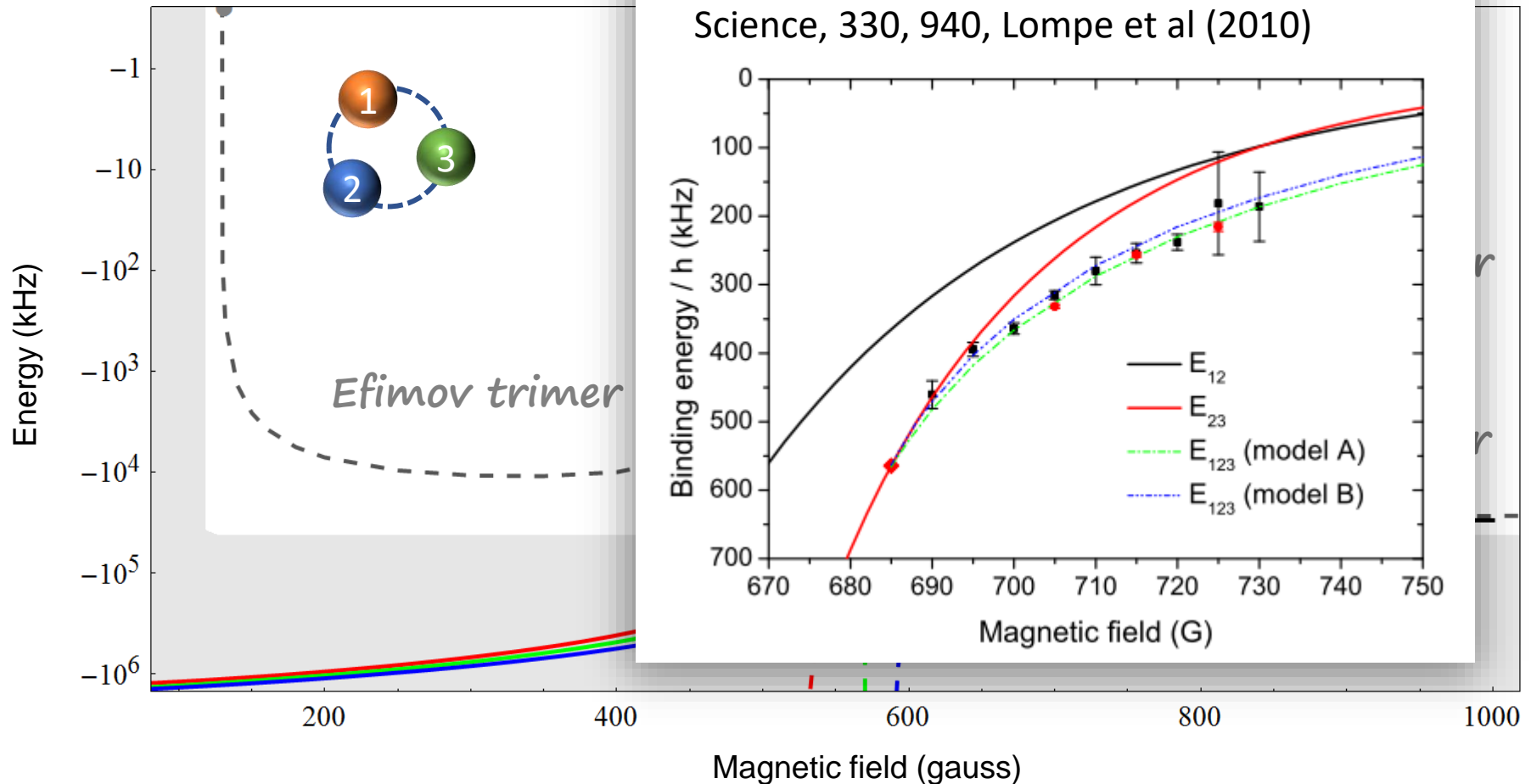


Vitaly Efimov's speech after receiving the prize



Observations in ultra-cold atomic gases

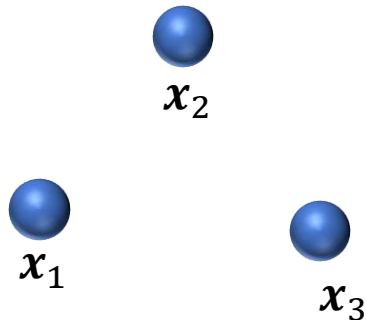
Observation with three distinguishable states of lithium atoms
(Heidelberg, Tokyo)



Mathematical description

Three-body equation

3 identical bosons interacting through the two-body potential $V(r)$



Bosonic symmetry:

$$\psi(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3) = \psi(\mathbf{x}_i, \mathbf{x}_j, \mathbf{x}_k)$$

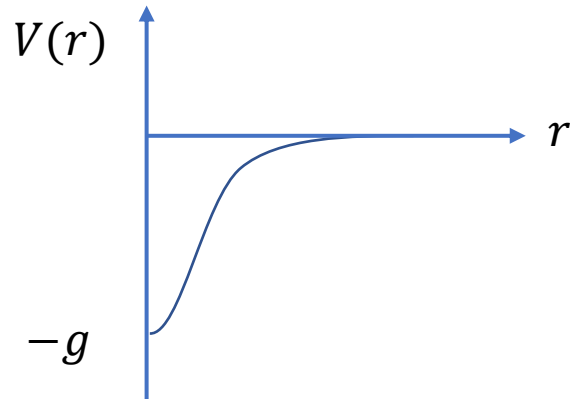
Schrödinger equation: $(\hat{H} - E)\psi(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3) = 0$

$$\text{Hamiltonian: } \hat{H} = -\frac{\hbar^2}{2m}(\Delta_1 + \Delta_2 + \Delta_3) + V(\mathbf{r}_{12}) + V(\mathbf{r}_{23}) + V(\mathbf{r}_{31})$$

With $\mathbf{r}_{ij} = \mathbf{x}_i - \mathbf{x}_j$

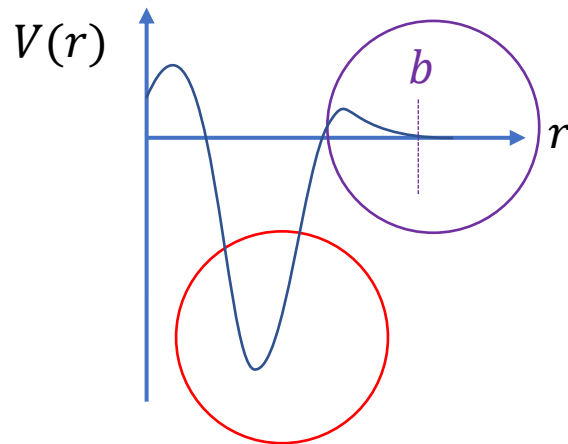
Conditions on the two-body interaction potentials $V(r)$

What is a **short-range resonant** interaction?



Conditions on the two-body interaction potentials $V(r)$

What is a **short-range resonant** interaction?



Should decay faster than $1/r^3$ to qualify as “**short-range**” and admit:

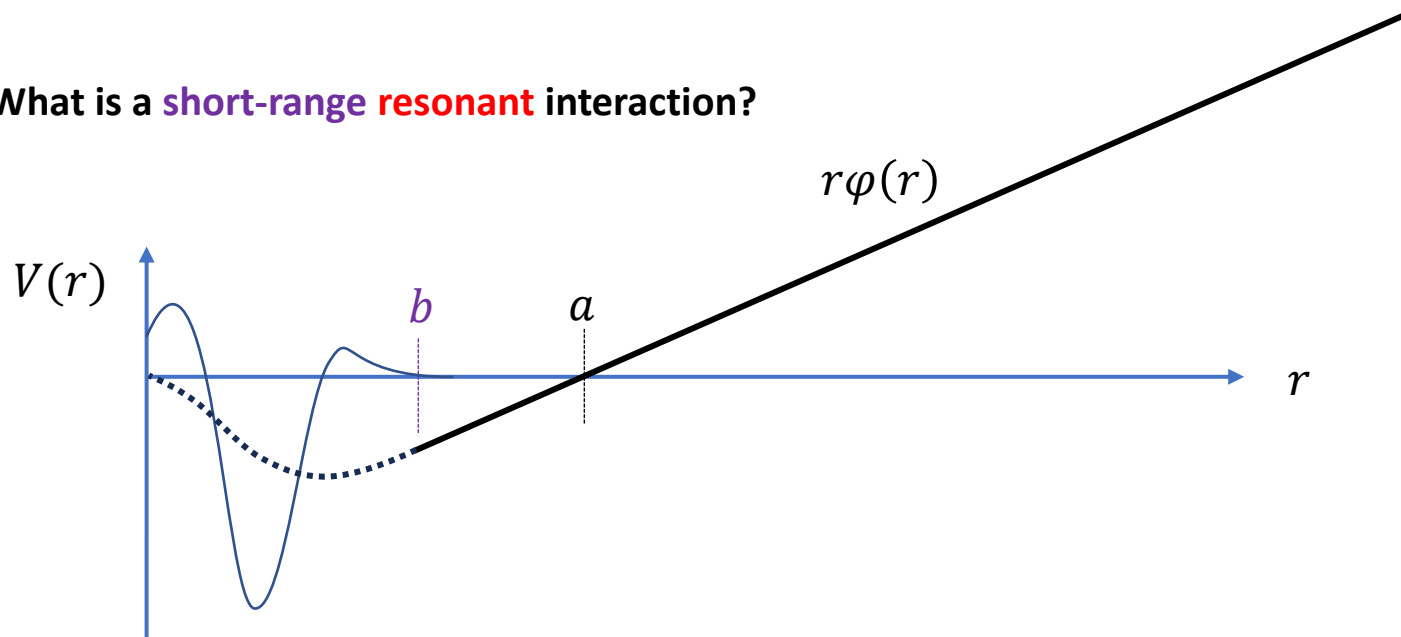
- **typical range b** : $V(r) \approx 0$ for $r \gg b$
- **scattering length a**

Should feature enough **attractive** part to approach the appearance of a two-body bound state.

- $|a| \gg b$

Conditions on the two-body interaction potentials $V(r)$

What is a **short-range resonant** interaction?



Two-body equation at zero energy:

$$\left(-\frac{1}{2}\Delta_1 - \frac{1}{2}\Delta_2 + V(r_{12}) - 0 \right) \varphi(\mathbf{x}_1, \mathbf{x}_2) = 0$$

Elimination of the centre of mass and s-wave:

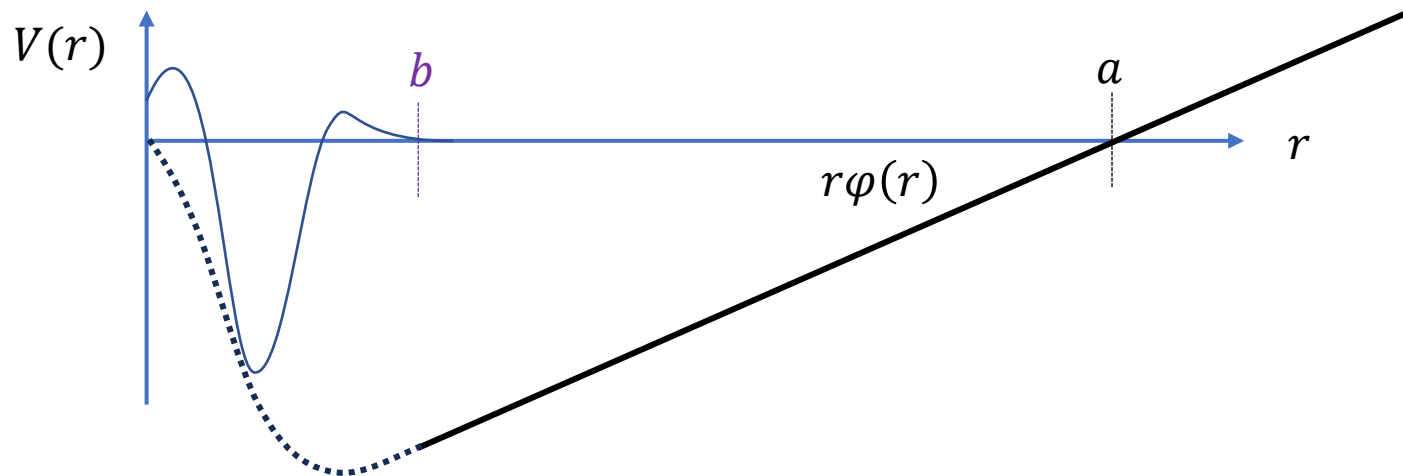
$$\left(-\frac{d^2}{dr^2} + \cancel{V(r)} \right) [r\varphi(r)] = 0$$

For $r \gg b$,

$$r\varphi(r) \propto r - a$$

Conditions on the two-body interaction potentials $V(r)$

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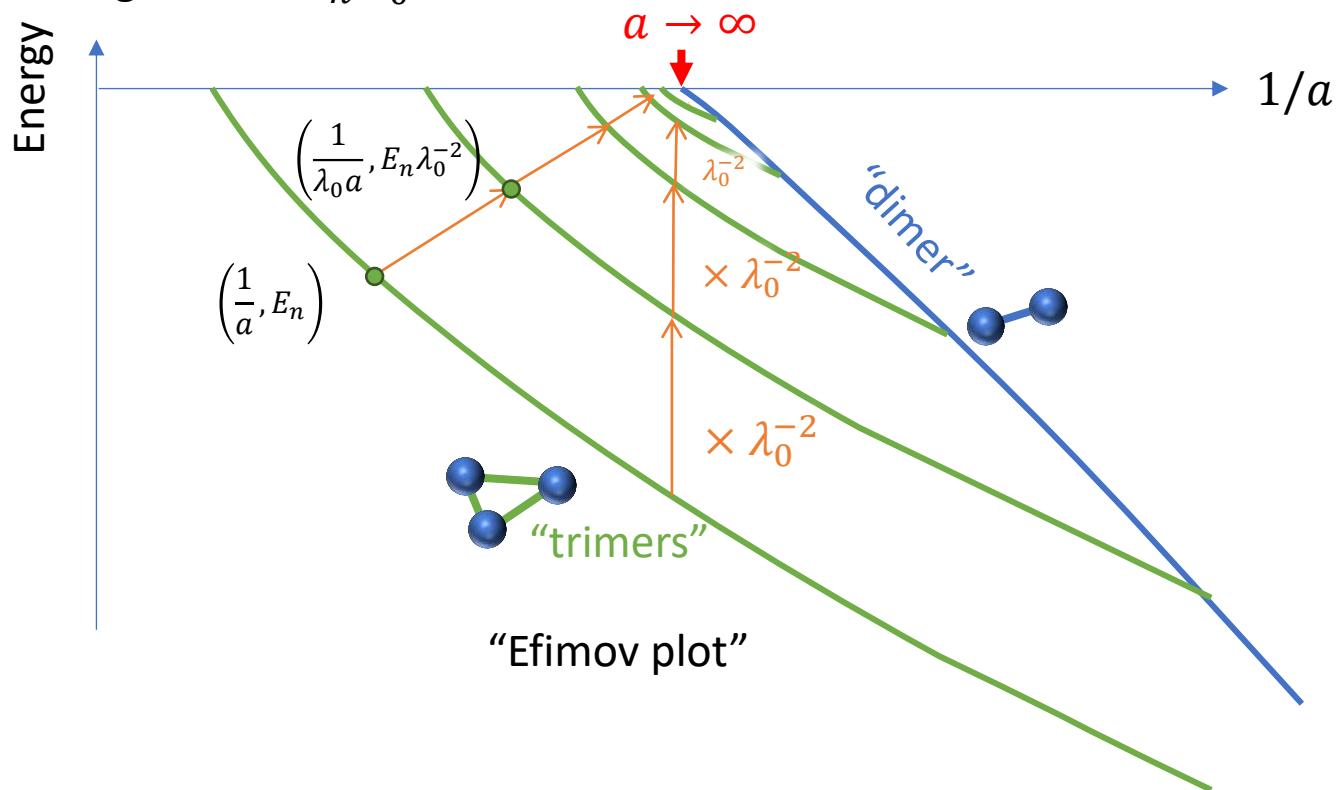
$$\left(-\frac{d}{dr} + V(r)\right)[r\varphi(r)] = 0$$

For $r \gg b$,

$$r\varphi(r) \propto r - a$$

Characterisation of the Efimov effect

- For $a \rightarrow \infty$, infinite sequence of discrete eigenvalues: $E_n \xrightarrow{n \rightarrow \infty} E_0 \lambda_0^{-2n}$
- **Discrete scale invariance:** for large n , if $\psi_n(r)$ is an eigenvector of $\hat{H}(a)$ with eigenvalue E_n , then $\psi_n(r/\lambda_0)$ is (almost) an eigenvector of $\hat{H}(a \lambda_0)$ with eigenvalue $E_n \lambda_0^{-2}$



How did Efimov “demonstrate” the effect?



Vitaly Efimov in 1977

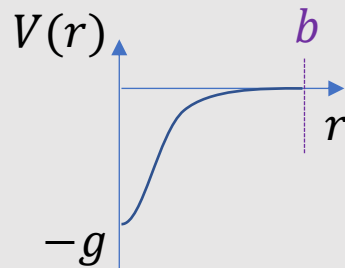
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V. Efimov, “Energy levels arising from resonant two-body forces in a three-body system.” **Physics Letters B**, **33**, 563 – 564, 1970.

How did Efimov “demonstrate” the effect?

(1) Zero-range approximation:

Original interaction

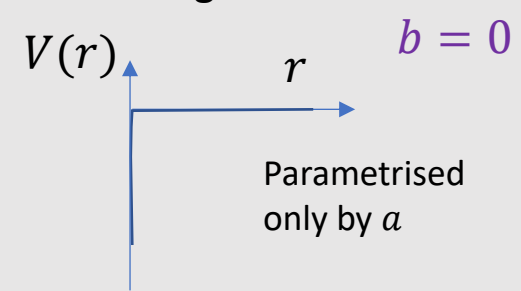


The wave function is much larger than the range of interactions

can be replaced by

Universality and scale invariance become exact!

“Zero-range” interaction



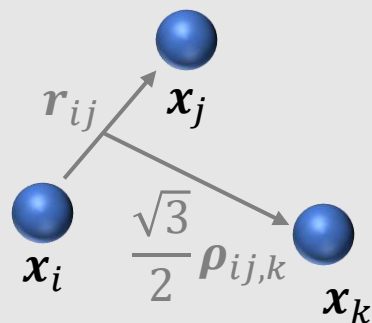
There are many ways to implement this:

- “Delta function” $V(r) \rightarrow g\delta^3(\mathbf{r})$ regularised by a cutoff (like renormalisation in quantum field theory)
- Zero-range boundary condition: $\psi \xrightarrow{r \rightarrow 0} \propto \frac{1}{r} - \frac{1}{a} \iff \left[\frac{d}{dr} \ln(r\psi) \right]_{\ell=0} \xrightarrow{r \rightarrow 0} -\frac{1}{a}$
- Pseudo-potential $V(r) = g\delta^3(\mathbf{r}) \frac{d}{dr}(r \cdot)$

How did Efimov “demonstrate” the effect?

However, the 3-boson problem with such zero-range interaction is not well posed: divergence of energies with large cutoff (“Thomas collapse”)

(2) Hyper-spherical adiabatic expansion (crucial insight!)



Change to hyper-spherical coordinates

$$\psi(\mathbf{r}_{ij}, \boldsymbol{\rho}_{ij,k}) = \psi(\rho, \alpha, \hat{r}_{ij}, \hat{\rho}_{ij,k})$$

$$\text{Hyper-radius } \rho = \sqrt{r_{ij}^2 + \rho_{ij,k}^2}$$

$$\text{Hyper-angle } \alpha = \arctan \frac{r_{ij}}{\rho_{ij,k}}$$

$$\rho^{-5/2} \left(-\frac{d^2}{d\rho^2} - \frac{1/4}{\rho^2} - \frac{\hat{\Lambda}_\Omega}{\rho^2} + \sum_{ij} V(\rho \sin \alpha_{ij}) - E \right) (\rho^{5/2} \psi) = 0$$

How did Efimov “demonstrate” the effect?

$$\rho^{-5/2} \left(-\frac{d^2}{d\rho^2} - \frac{1/4}{\rho^2} - \frac{\hat{\Lambda}_\Omega}{\rho^2} + \sum_{ij} V(\rho \sin \alpha_{ij}) - E \right) (\rho^{5/2} \psi) = 0$$

Expansion: $\psi(\rho, \Omega) = \frac{1}{\rho^{5/2}} \sum_{n=1}^{\infty} \overset{\text{Hyper-radial}}{f_n(\rho)} \overset{\text{Hyper-angular}}{\Phi_n(\Omega; \rho)}$

Hyper-angular equation at fixed ρ : $\left(\hat{\Lambda}_\Omega - \rho^2 \sum_{ij} V(\rho \sin \alpha_{ij}) + s_n^2(\rho) \right) \Phi_n(\Omega; \rho) = 0$

Coupled hyper-radial equations: $\left(-\frac{d^2}{d\rho^2} + \frac{s_n^2(\rho) - 1/4}{\rho^2} - E \right) f_n - \sum_p \left(2P_{np} \frac{df_p}{d\rho} + Q_{np} f_p \right) = 0$

For zero-range interactions and $a \rightarrow \infty$, the equations decouple!

How did Efimov “demonstrate” the effect?

$$\left(\hat{\Lambda}_\Omega - \rho^2 \sum_{ij} V(\rho \sin \alpha_{ij}) + s_n^2(\rho) \right) \Phi_n(\Omega; \rho) = 0 \quad \Rightarrow \quad s_n \cos\left(\frac{s_n \pi}{2}\right) + \frac{8}{\sqrt{3}} \sin\left(\frac{s_n \pi}{6}\right) = 0$$

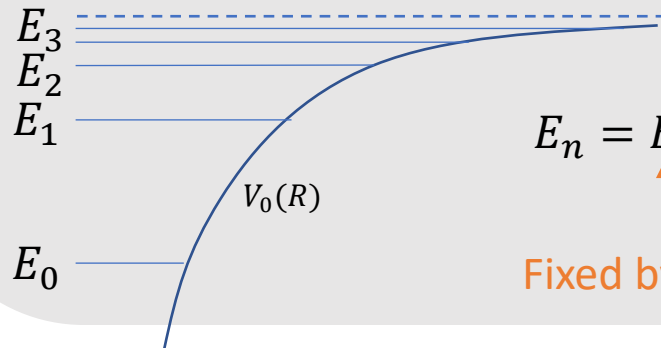
$$\left(-\frac{d^2}{d\rho^2} + \frac{s_n^2(\rho) - 1/4}{\rho^2} - E \right) f_n = 0$$

All s_n are real, except one:
 $s_0 = \pm i1.00624$

For $n = 0$, one gets the Efimov attractive potential $V_0(\rho) = -\frac{|s_0|^2 + \frac{1}{4}}{\rho^2}$

For small ρ , $f_0(\rho) = \alpha \rho^{i|s_0|} + \beta \rho^{-i|s_0|} \propto \cos(|s_0| \ln \Lambda \rho)$

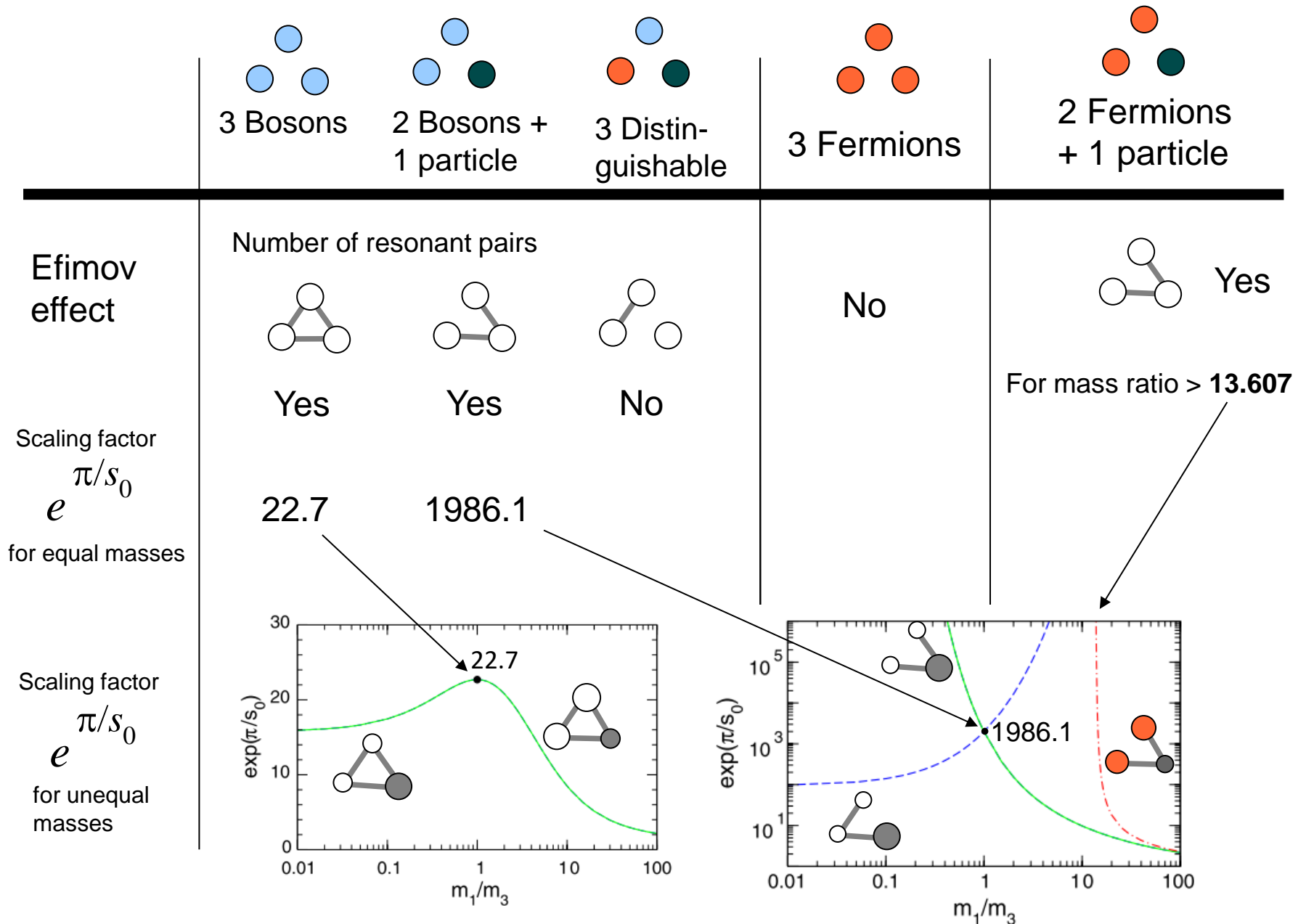
⋮



$$E_n = E_0 \underbrace{(\exp \pi / |s_0|)^{-2n}}_{\lambda_0 \approx 22.7 \text{ (scaling factor)}}$$

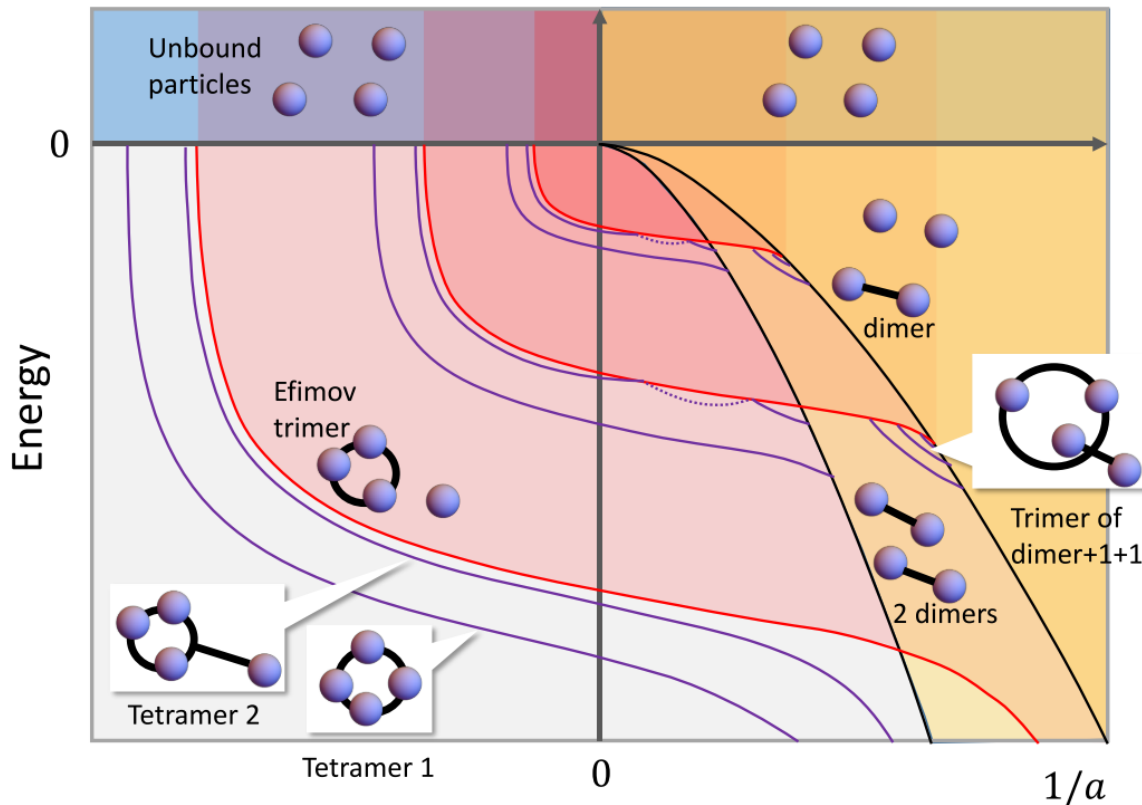
↑
Fixed by Λ

↪ Three-body parameter set by a short-range 3-body boundary condition



More than three particles

Tetramers of four identical bosons



No four-body Efimov effect

Two “universal tetramers”
attached to each Efimov
trimer

Controversy:

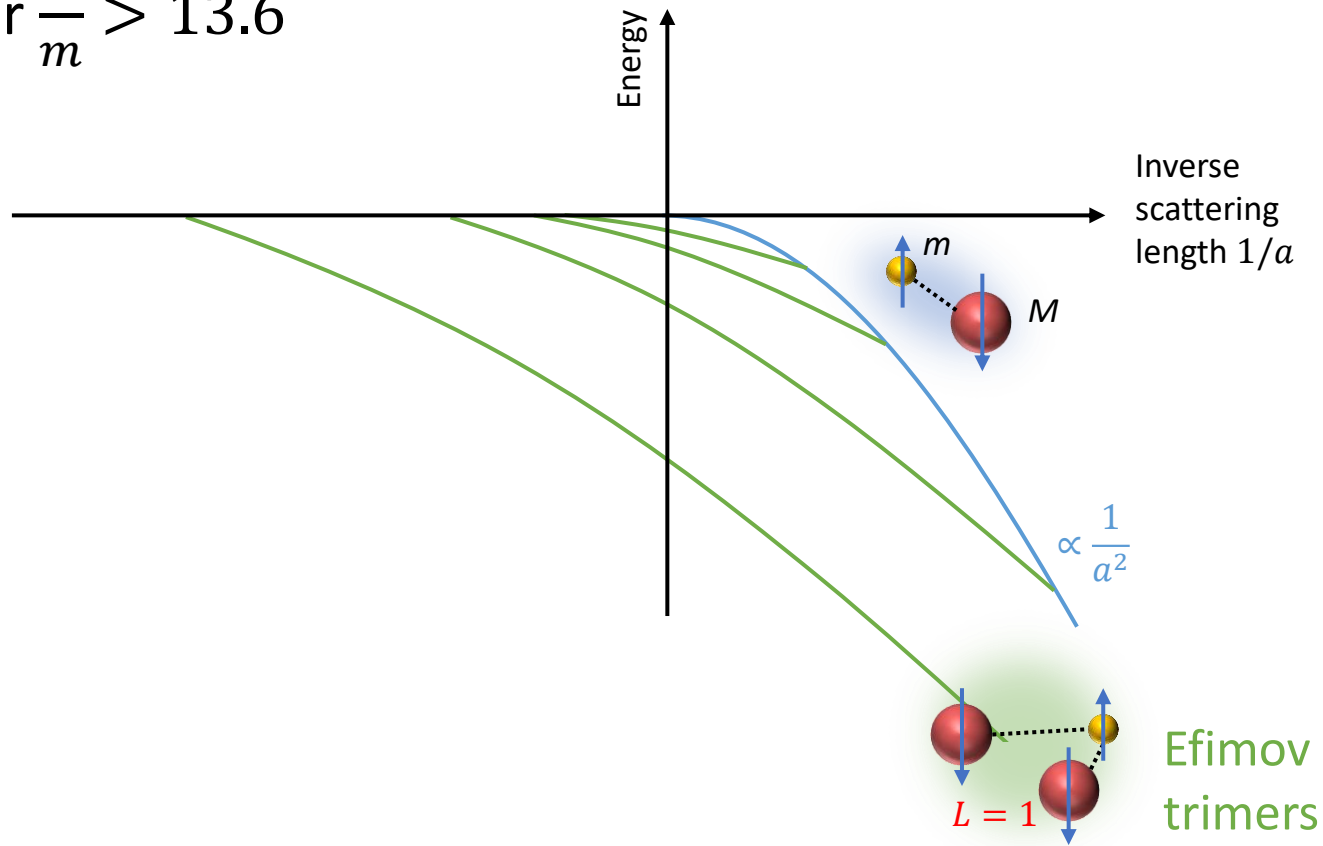
- 1) There is in general a need for a 4-body parameter
- 2) The universal states do not require any 4-body parameter

J. von Stecher, J. P. D’Incao, and C. H. Greene, *Nature Physics*, 5, 417–421, 2009.

A. Deltuva, *Europhysics Letters*, 95, 43002, 2011.

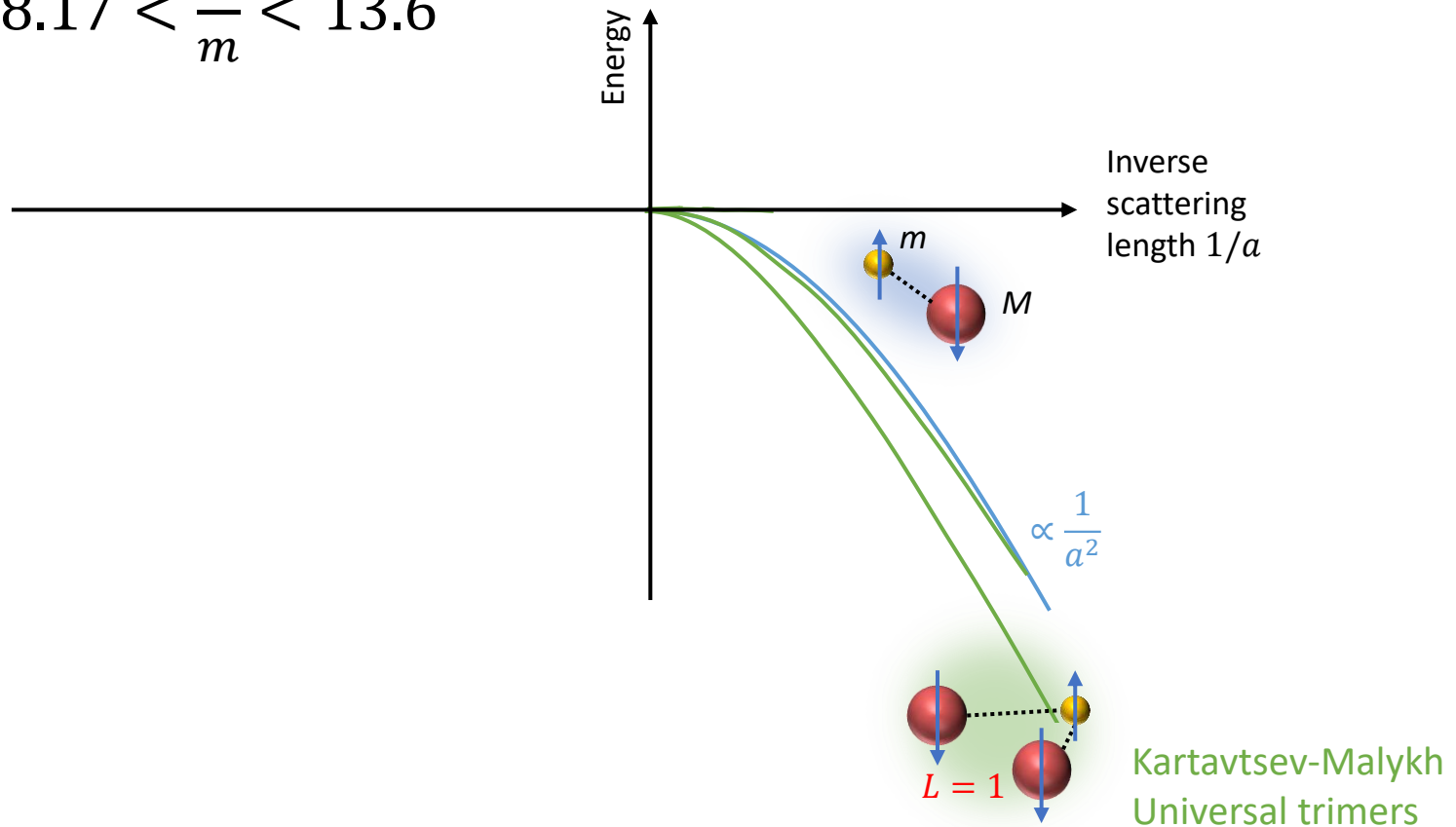
Tetramers of 3+1 fermions

For $\frac{M}{m} > 13.6$



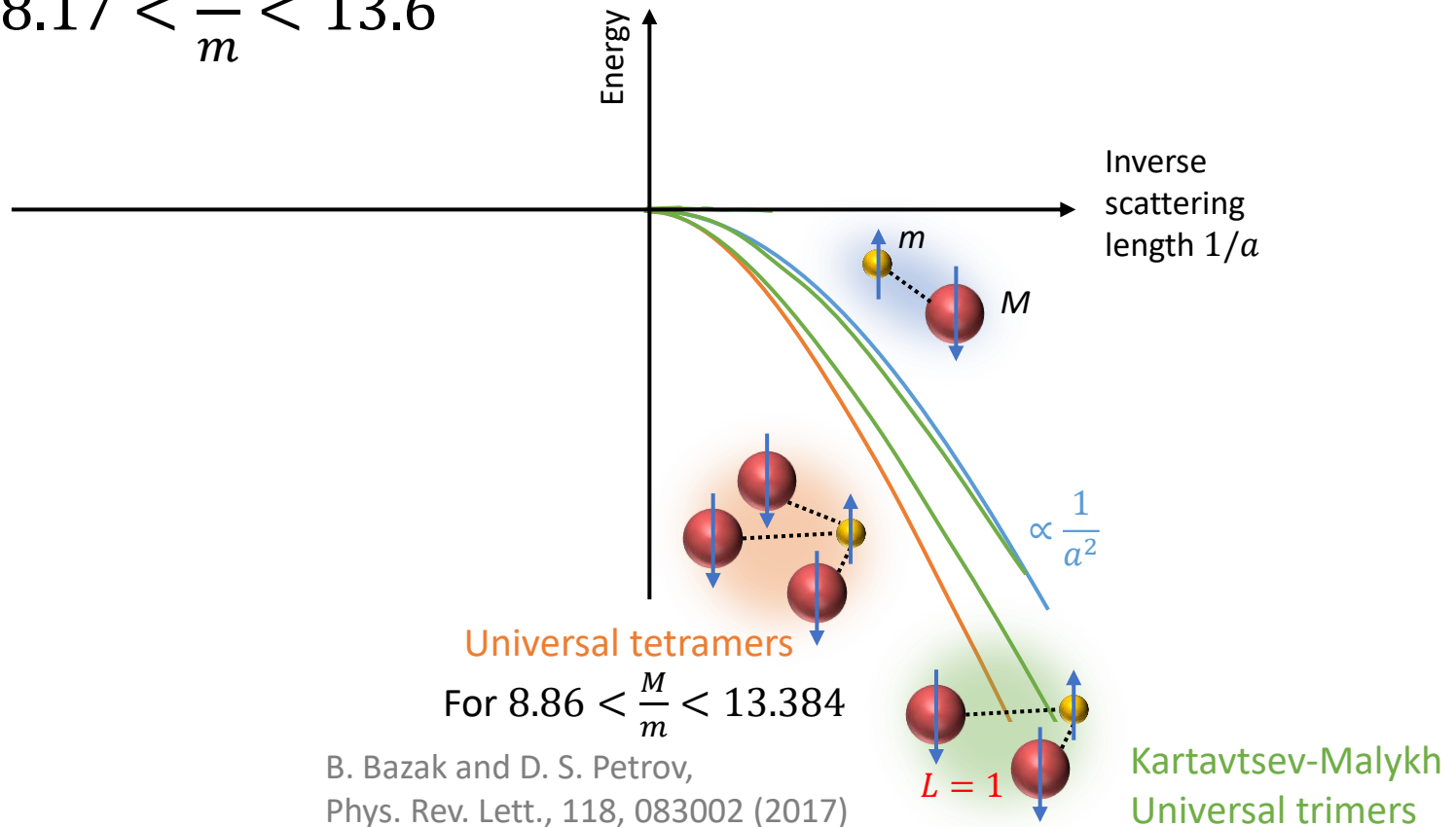
Tetramers of 3+1 fermions

For $8.17 < \frac{M}{m} < 13.6$



Tetramers of 3+1 fermions

For $8.17 < \frac{M}{m} < 13.6$

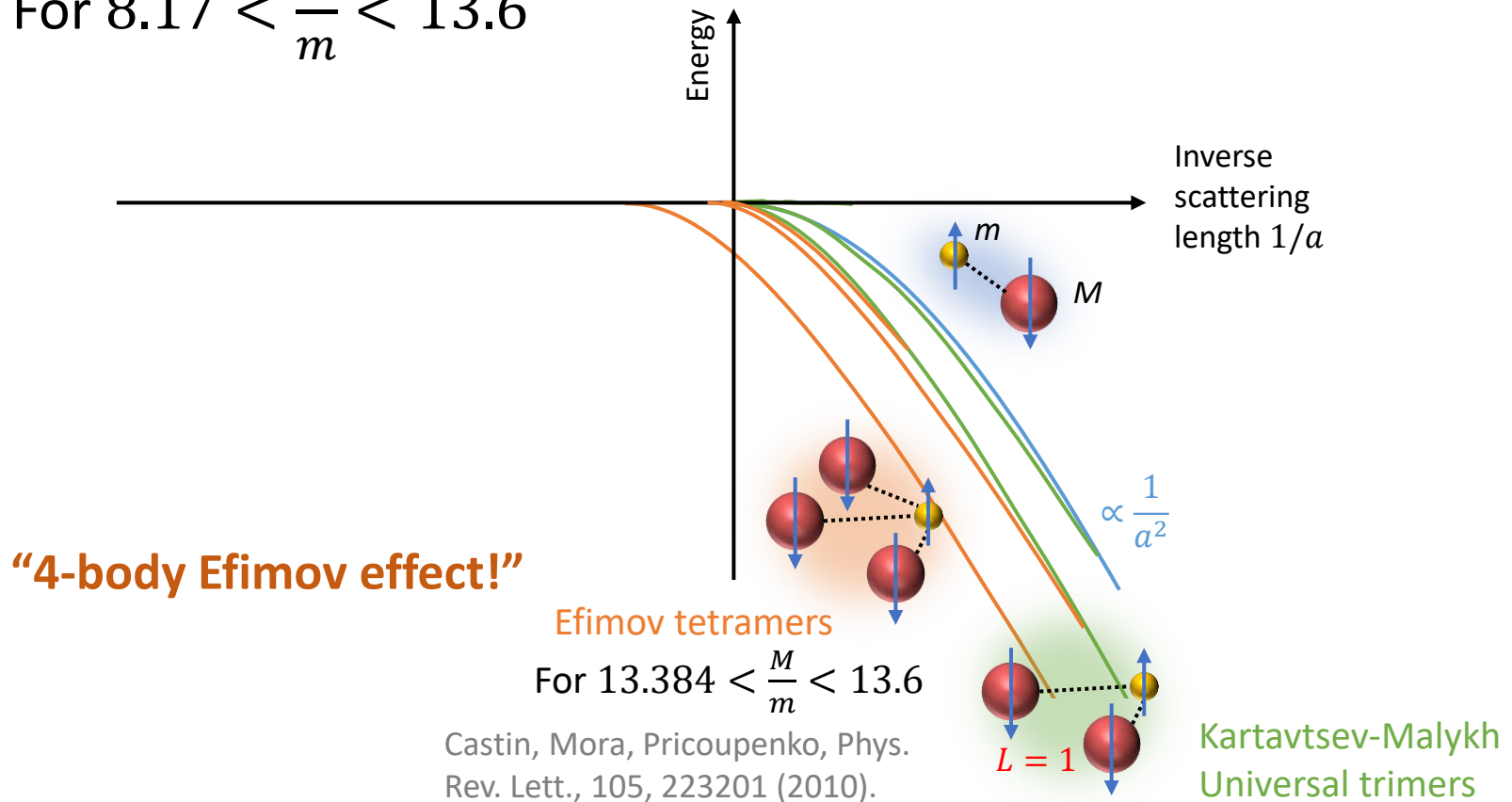


B. Bazak and D. S. Petrov,
Phys. Rev. Lett., 118, 083002 (2017)

Kartavtsev+Malykh
J Phys B 40, 1429 (2007)

Tetramers of 3+1 fermions

For $8.17 < \frac{M}{m} < 13.6$



Castin, Mora, Pricoupenko, Phys. Rev. Lett., 105, 223201 (2010).

Kartavtsev+Malykh
J Phys B 40, 1429 (2007)

Conclusion

Efimov physics has been a developing field of quantum physics, both theoretical and experimental, unveiling a whole collection of universal few-body states with remarkable mathematical properties and challenges.