

Supersymmetric Dark Matter

From Historical Foundations to Future Challenges

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Conference in memory of Veniamin Sergeyevich Berezinsky
GSSI, L'Aquila – October 2, 2024

Supersymmetry

- Profound symmetry between fermions and bosons
- Extends spacetime symmetries
- Its local version contains a link to Gravity (SUGRA)
- Requires the introduction of new particles (superpartners + extended Higgs sector)
- Unfortunately, is broken

Dark Matter

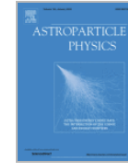
- Overwhelming evidence
 - Rotational curves of spiral galaxies
 - Galaxy clusters dynamics
 - Gravitational lensing
 - Hydrodynamical equilibrium of hot gas in galaxy clusters
 - Large scale structure of the Universe
 - Energy budget of the Universe
 - Structure formation
- However: DM evidence is purely gravitational
- A new particle?

A mutual opportunity?



ELSEVIER

Astroparticle Physics
Volume 5, Issue 1, June 1996, Pages 1-26



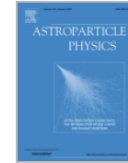
Neutralino dark matter in supersymmetric models with non-universal scalar mass terms

V. Berezhinsky^a✉, A. Bottino^{b c}✉, J. Ellis^d✉, N. Fornengo^{e c}✉, G. Mignola^{b c}✉,
S. Scopel^{f g}✉



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Astroparticle Physics
Volume 5, Issues 3–4, October 1996, Pages 333-352



Searching for relic neutralinos using neutrino telescopes

V. Berezhinsky^a✉, A. Bottino^{b c}✉, J. Ellis^d✉, N. Fornengo^{c e}✉, G. Mignola^{c d}✉,
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Minimal SUSY extension of the SM

Supersymmetry: Fermions \longleftrightarrow Bosons

Super-multiplets	Super-field	Bosonic fields	Fermionic partners	SU(3)	SU(2)	U(1)
gluon/gluino	\hat{V}_8	$g_{s=1}$	$\tilde{g}_{s=1/2}$	8	1	0
gauge boson/ gaugino	\hat{V}	W^\pm, W^0	$\tilde{W}^\pm, \tilde{W}^0$	1	3	0
	\hat{V}'	$B_{s=1}$	$\tilde{B}_{s=1/2}$	1	1	0
slepton/ lepton	\hat{L}	$(\tilde{\nu}_L, \tilde{e}_L^-)$	$(\nu, e^-)_L$	1	2	-1
	\hat{E}^c	$\tilde{e}_R^+_{s=0}$	$e_L^c_{s=1/2}$	1	1	2
squark/ quark	\hat{Q}	$(\tilde{u}_L, \tilde{d}_L)$	$(u, d)_L$	3	2	1/3
	\hat{U}^c	\tilde{u}_R^*	u_L^c	$\bar{3}$	1	-4/3
	\hat{D}^c	$\tilde{d}_R^*_{s=0}$	$d_L^c_{s=1/2}$	$\bar{3}$	1	2/3
Higgs boson/ higgsino	\hat{H}_d	(H_d^0, H_d^-)	$(\tilde{H}_d^0, \tilde{H}_d^-)$	1	2	-1
	\hat{H}_u	(H_u^+, H_u^0)	$(\tilde{H}_u^+, \tilde{H}_u^0)$	1	2	1

$s=0$

$s=1/2$

Particle content

Normal particles/fields		Supersymmetric partners			
		Interaction eigenstates		Mass eigenstates	
Symbol	Name	Symbol	Name	Symbol	Name
$q = d, c, b, u, s, t$	quark	\tilde{q}_L, \tilde{q}_R	squark	\tilde{q}_1, \tilde{q}_2	squark
$l = e, \mu, \tau$	lepton	\tilde{l}_L, \tilde{l}_R	slepton	\tilde{l}_1, \tilde{l}_2	slepton
$\nu = \nu_e, \nu_\mu, \nu_\tau$	neutrino	$\tilde{\nu}$	sneutrino	$\tilde{\nu}$	sneutrino
g	gluon	\tilde{g}	gluino	\tilde{g}	gluino
W^\pm	W -boson	\tilde{W}^\pm	wino	} $\tilde{\chi}_{1,2}^\pm$	chargino
H^-	Higgs boson	\tilde{H}_1^-	higgsino		
H^+	Higgs boson	\tilde{H}_2^+	higgsino		
B	B -field	\tilde{B}	bino		
W^3	W^3 -field	\tilde{W}^3	wino	} $\tilde{\chi}_{1,2,3,4}^0$	neutralino
h H_1^0 scalar	Higgs boson	\tilde{H}_1^0	higgsino		
H H_2^0 scalar	Higgs boson	\tilde{H}_2^0	higgsino		
A H_3^0 pseudoscalar	Higgs boson				

Neutral SUSY particles: sneutrinos, neutralinos [gravitinos]

Two required ingredients

- Supersymmetry breaking

- “Super”-particles cannot have the same mass as their SM partners (not observed)
- Large number of additional parameters (120+)

- R- parity

- For a generic SUSY models, a sort of (B-L) symmetry need to be enforced to prevent too-fast proton decay

$$R = (-1)^{3(B-L)+2S}$$

- This implies that the LSP is stable: if neutral, being massive and stable, it can be the DM

$$A + B \longrightarrow \tilde{X} + \tilde{Y}$$

$$\tilde{X} \longrightarrow \tilde{Y} + A + B$$

SUSY DM candidates

- Sneutrinos

$$\tilde{\nu}_e, \tilde{\nu}_\mu, \tilde{\nu}_\tau \longrightarrow \tilde{\nu}_1, \tilde{\nu}_2, \tilde{\nu}_3$$

The lightest of the three

- Neutralinos

$$\tilde{\gamma}, \tilde{Z}, \tilde{H}_1, \tilde{H}_2 \longrightarrow \chi_i = a_{1_i} \tilde{\gamma} + a_{2_i} \tilde{Z} + a_{3_i} \tilde{H}_1 + a_{4_i} \tilde{H}_2 \quad (i = 1, 2, 3, 4)$$

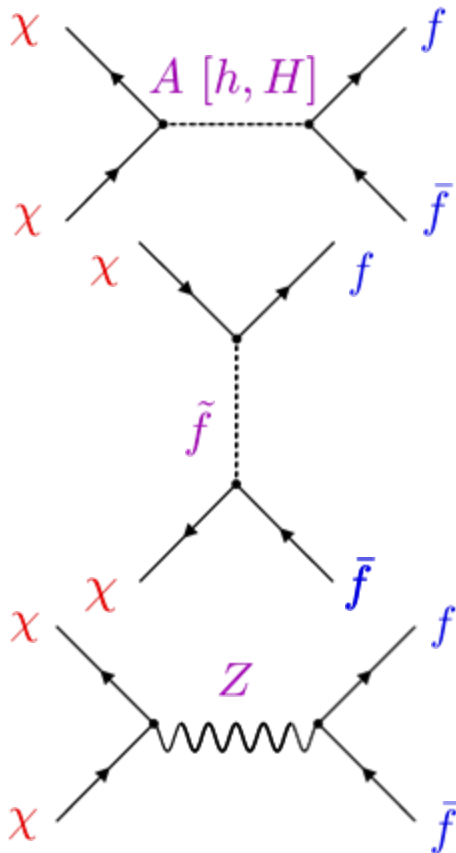
The lightest of the four

- Gravitino

$$\tilde{g}$$

Neutralino cosmology

$$\Omega h^2 \sim \langle \sigma_{\text{ann}} v \rangle^{-1}$$



* $\chi\chi \rightarrow f\bar{f}$	Z, h, H, A \tilde{f}_L, \tilde{f}_R	s channel t and u channels
* $\chi\chi \rightarrow hh, hH, HH, AA$	h, H $\chi_i (i = 1, 2, 3, 4)$	s channels t and u channels
* $\chi\chi \rightarrow hA, HA$	Z, A $\chi_i (i = 1, 2, 3, 4)$	s channel t and u channels
* $\chi\chi \rightarrow H^+H^-$	Z, h, H $\chi_j^+ (j = 1, 2)$	s channel t and u channels
* $\chi\chi \rightarrow ZZ$	h, H $\chi_i (i = 1, 2, 3, 4)$	s channel t and u channels
* $\chi\chi \rightarrow W^+W^-$	Z, h, H $\chi_j^+ (j = 1, 2)$	s channel t and u channels
* $\chi\chi \rightarrow hZ, HZ$	Z, A $\chi_i (i = 1, 2, 3, 4)$	s channel t and u channel
* $\chi\chi \rightarrow AZ$	h, H $\chi_i (i = 1, 2, 3, 4)$	s channel t and u channels
* $\chi\chi \rightarrow W^\pm H^\mp$	h, H, A $\chi_j^+ (j = 1, 2)$	s channel t and u channels

which one is open depends on the neutralino mass

Higgs sector: 2HDB

$$V_0 = (M_{H_1}^2 + \mu^2) |H_1|^2 + (M_{H_2}^2 + \mu^2) |H_2|^2 - B\mu (H_1 H_2 + \text{h.c.}) \\ + \text{quartic D terms.}$$

$$v_u^2 + v_d^2 = (246 \text{ GeV})^2$$

$$\tan \beta = v_u/v_d$$

$$\sin 2\beta = \frac{-2B\mu}{M_{H_1}^2 + M_{H_2}^2 + 2\mu^2}$$

$$M_Z^2 = 2 \frac{M_{H_1}^2 - M_{H_2}^2 \tan^2 \beta}{\tan^2 \beta - 1} - 2\mu^2$$

$$M_A^2 = M_{H_1}^2 + M_{H_2}^2 + 2\mu^2 > 0$$

+ radiative corrections

Parameters:

$\tan \beta$

μ

$M_{H_1}^2$

$M_{H_2}^2$

Soft SUSY breaking

$$\begin{aligned} \mathcal{V}_{soft} = & \sum_i m_i^2 |\phi_i|^2 \\ & + \left\{ \left[A_{ab}^l h'_{ab} \tilde{L}_a H_1 \tilde{R}_b + A_{ab}^d h_{ab}^d \tilde{Q}_a H_1 \tilde{D}_b + A_{ab}^u h_{ab}^u \tilde{Q}_a H_2 \tilde{U}_b + \text{h.c.} \right] - B\mu H_1 H_2 + \text{h.c.} \right\} \\ & + \sum_i M_i (\lambda_i \lambda_i + \bar{\lambda}_i \bar{\lambda}_i) \end{aligned}$$

Parameters:

m_i	Scalar masses
M_i	Gaugino masses
A_{ab}	Trilinear couplings

SUSY Frameworks

- The number of free parameters is exceedingly large (124)
- **Effective MSSM**: a bunch of relevant phenomenological parameters are taken as independent at the EW scale
- **SUGRA-inspired models**: EW-scale parameters are obtained through RGE-evolution from a very limited set of parameters defined at the GUT or Planck scale. RGE-evolutions determines EWSB and all the low-energy phenomenology

mSUGRA/CMSSM

- High-scale parameters unification

$$M_i(M_{GUT}) = m_{1/2}$$

At low-energy, implies

$$M_1 : M_2 : M_3 = \alpha_1 : \alpha_2 : \alpha_3$$

$$M_{H_1}^2(M_{GUT}) = M_{H_2}^2(M_{GUT}) = m_0^2$$

$$M_{\tilde{L}}^2(M_{GUT}) = M_{\tilde{E}}^2(M_{GUT}) = m_0^2 \mathbf{1}$$

$$M_{\tilde{Q}}^2(M_{GUT}) = M_{\tilde{U}}^2(M_{GUT}) = M_{\tilde{D}}^2(M_{GUT}) = m_0^2 \mathbf{1}$$

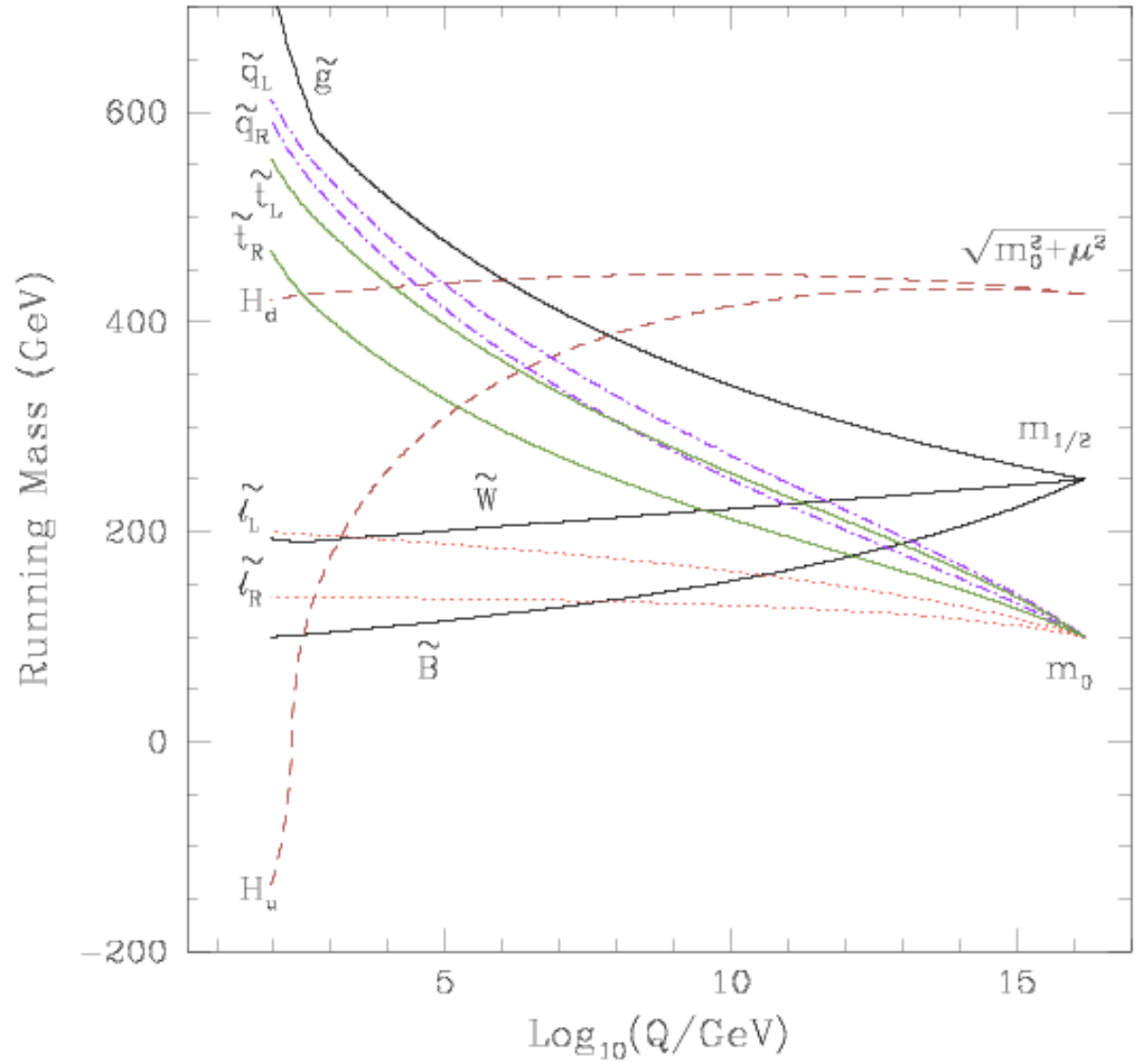
$$A_U(M_{GUT}) = A_D(M_{GUT}) = A_E(M_{GUT}) = A_0 \mathbf{1}$$

- Low-energy parameters

$$|\mu|$$

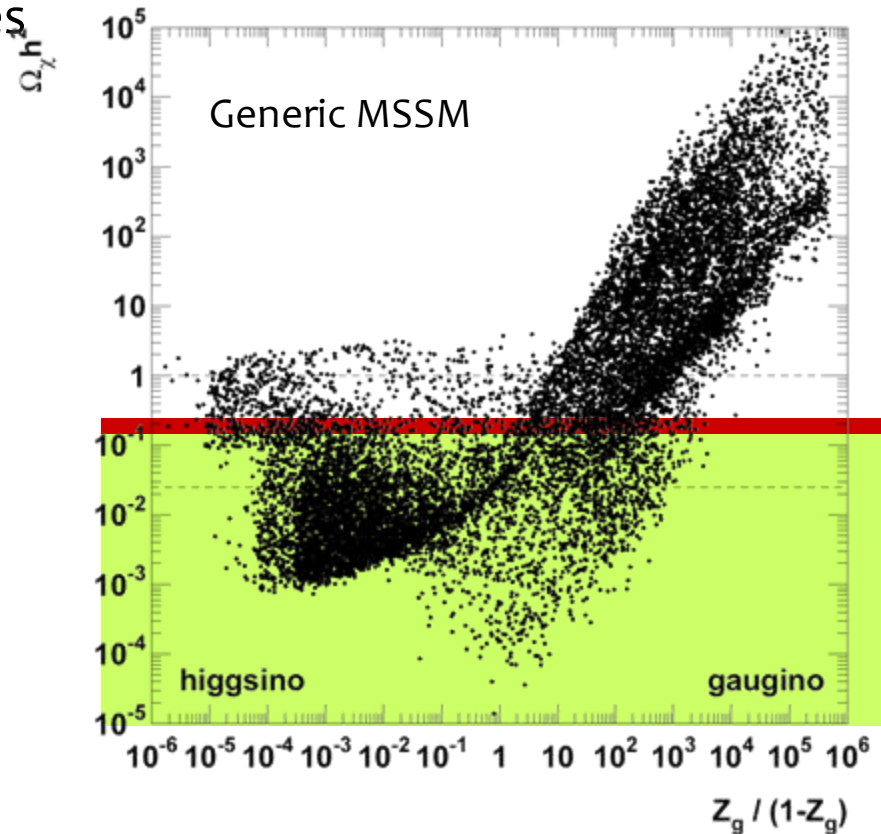
$$\tan \beta$$

RGE evolution



The most popular: mSUGRA/CMSSM

- Consequences for neutralino DM:
 - Neutralino is mostly a gaugino
 - Exceedingly large relic abundance
 - Suppressed DM detection rates



Caveats on mass universality

- Strict unification of mass parameters at the GUT scale (universality) is not at the same level of motivation as gauge coupling unification
- Threshold effects or evolution from Planck/string scale to the GUT scale can easily spoil mass universality
- Non-universality can alter neutralino DM phenomenology quite significantly through:
 - Change in the low-energy Higgs phenomenology (Higgs masses and couplings to matter and DM)
 - Change in the low-scale neutralino characteristics (flip from pure gaugino to a gaugino-higgsino mix)

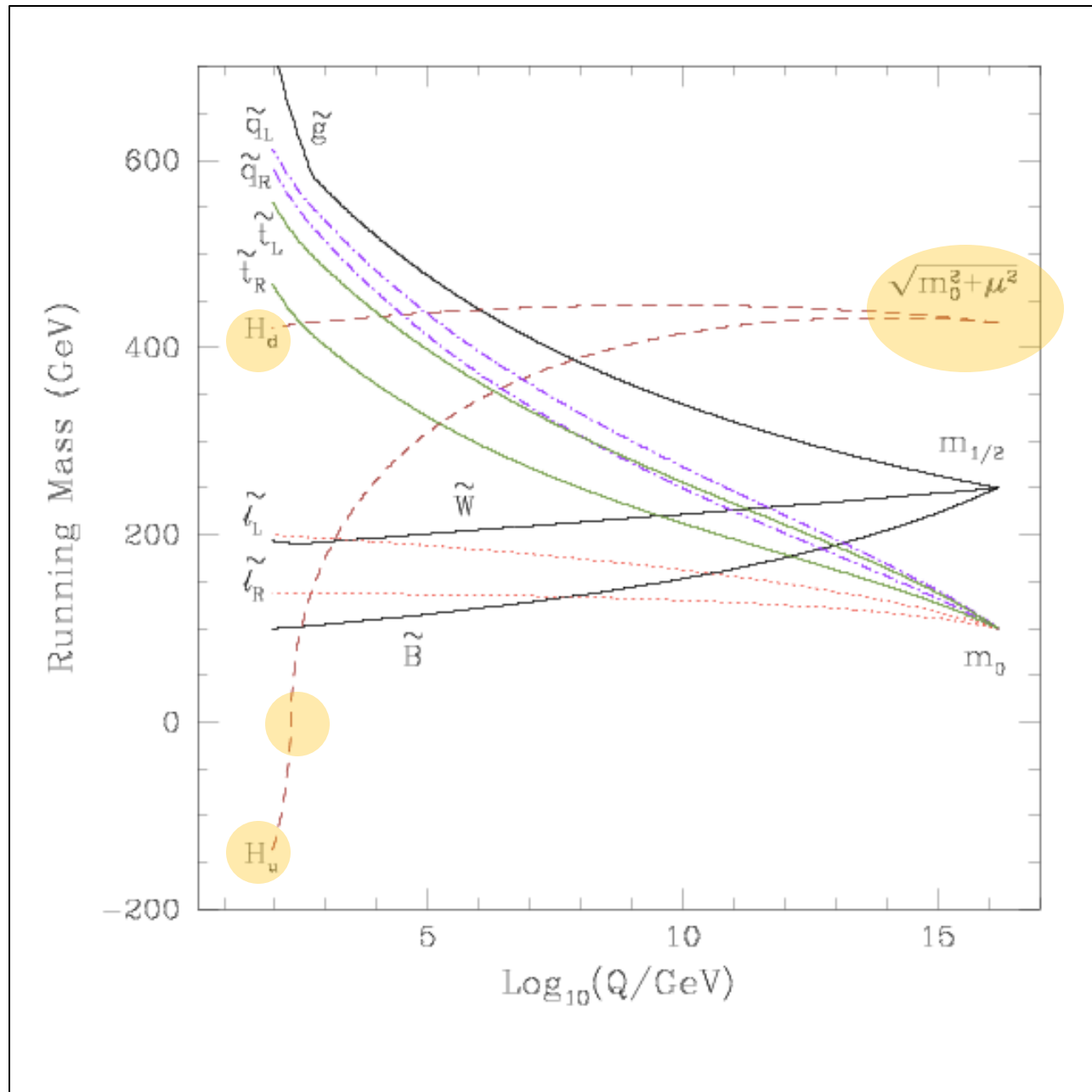
Non-universal-Higgs models

$$M_{H_i}^2 (M_{GUT}) = m_0^2 (1 + \delta_i)$$

Split in the Higgs soft masses, while maintaining universality in the slepton and squark sectors, as well as in the gaugino sector

This adds only 2 parameters to the model, but changes significantly the phenomenology, even with moderate non-universality, at the level of a few percent change

RGE evolution



Higgs sector

RGE evolution can be summarized in a dependence of the low energy parameters from the GUT-scale ones:

$$M_{H_i}^2 = a_i m_{1/2}^2 + b_i m_0^2 + c_i A_0^2 m_0^2 + d_i A_0 m_0 m_{1/2}$$

$$\mu^2 = J_1 m_{1/2}^2 + J_2 m_0^2 + J_3 A_0^2 m_0^2 + J_4 A_0 m_0 m_{1/2} - \frac{M_Z^2}{2}$$

$$M_A^2 = K_1 m_{1/2}^2 + K_2 m_0^2 + K_3 A_0^2 m_0^2 + K_4 A_0 m_0 m_{1/2} - M_Z^2$$

Higgs sector

Let's simplify and take $A_0 = 0$

$$M_{H_i}^2 = a_i m_{1/2}^2 + \textcircled{b_i} m_0^2$$

$$\mu^2 = J_1 m_{1/2}^2 + \textcircled{J_2} m_0^2$$

$$- \frac{M_Z^2}{2}$$

$$M_A^2 = K_1 m_{1/2}^2 + \textcircled{K_2} m_0^2$$

$$- M_Z^2$$

↑
These parameters in front of m_0 acquire dependence from the non-universality parameters δ and alter the low-energy Higgs phenomenology, including EWSB.

In particular, they can lower the A-higgs mass (relevant for neutralino relic abundance and detection rates) and the μ -parameter (relevant also for gaugino/higgsino mixing)

Neutralino sector

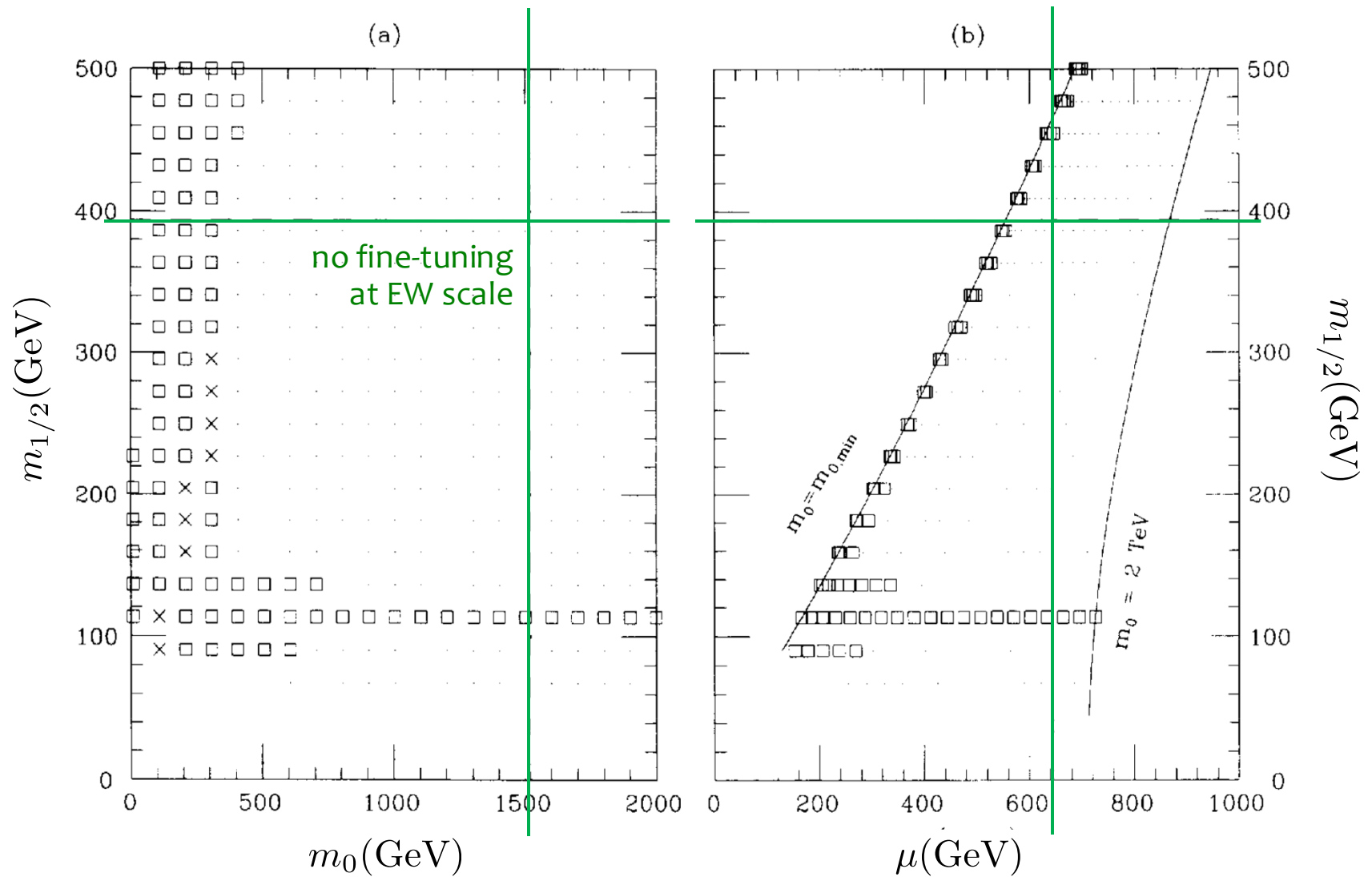
$$M_N \equiv \begin{pmatrix} \begin{array}{cc} \text{gaugino} & \\ \hline M_1 & 0 \\ 0 & M_2 \\ \hline \end{array} & \begin{array}{cc} -\frac{1}{2}g'v_d & \frac{1}{2}g'v_u \\ \frac{1}{2}gv_d & -\frac{1}{2}gv_u \end{array} \\ \begin{array}{cc} -\frac{1}{2}g'v_d & \frac{1}{2}gv_d \\ \frac{1}{2}g'v_u & -\frac{1}{2}gv_u \end{array} & \begin{array}{cc} \hline 0 & -\mu \\ -\mu & 0 \\ \hline \text{higgsino} \end{array} \end{pmatrix}$$

With gaugino-mass unification:

$$M_1 : M_2 : M_3 = \alpha_1 : \alpha_2 : \alpha_3$$

$$M_1 = \frac{5}{3} \tan^2 \theta_W M_2 \simeq 0.5 M_2$$

mSUGRA/CMSSM



$$\tan \beta = 8; \quad \delta_1 = 0; \quad \delta_2 = 0$$

No fine-tuning

The EW scale is obtained without excessive level of cancellations

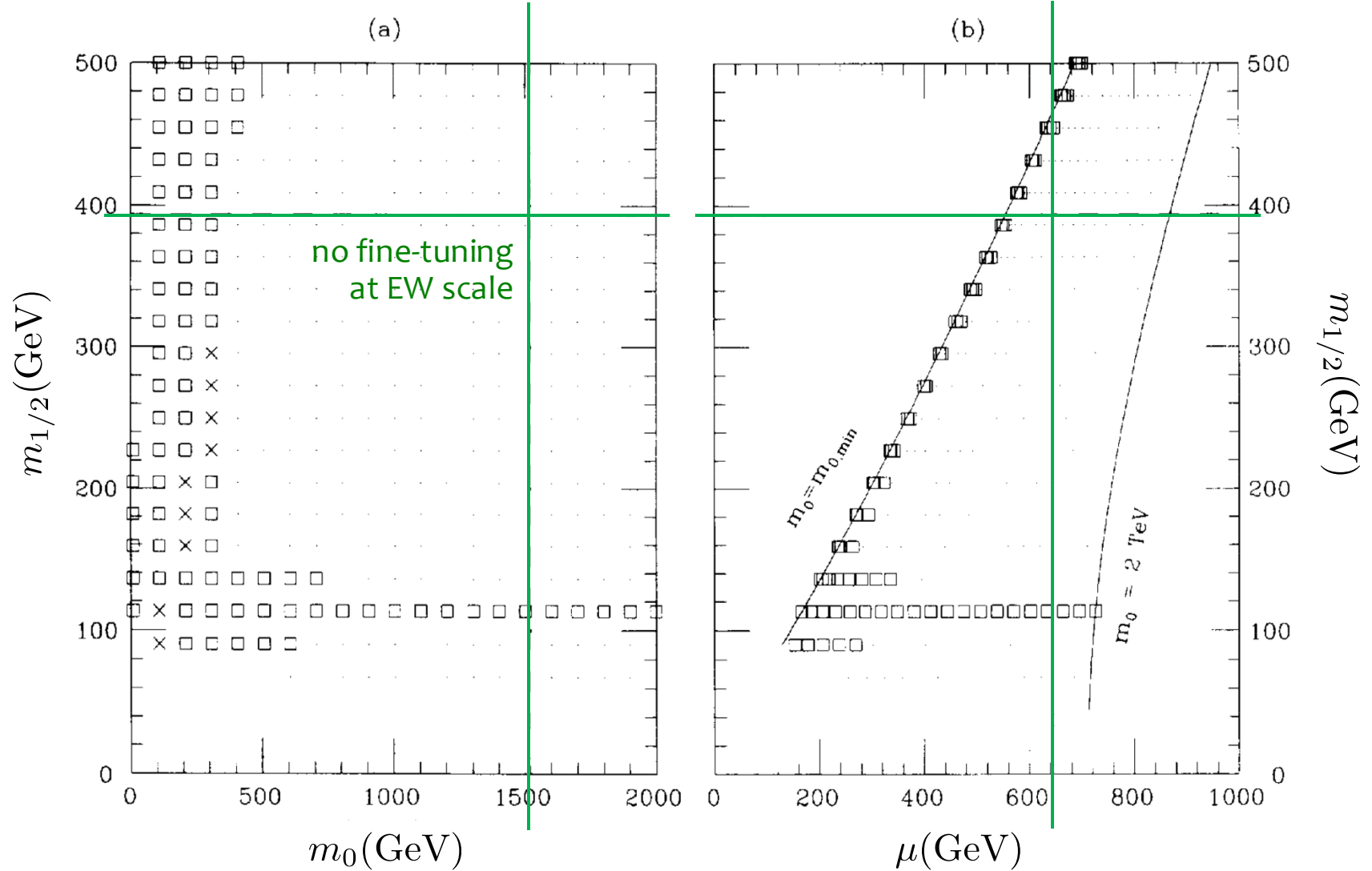
$$M_Z^2 = 2 \left(J_1 m_{1/2}^2 + J_2 m_0^2 + J_3 A_0^2 m_0^2 + J_4 A_0 m_0 m_{1/2} - \mu^2 \right)$$

$$\left| \frac{\Delta M_Z^2}{M_Z^2} \right| < \eta_f \left| \frac{\Delta x_i^2}{x_i^2} \right|$$

For $A_0 = 0$

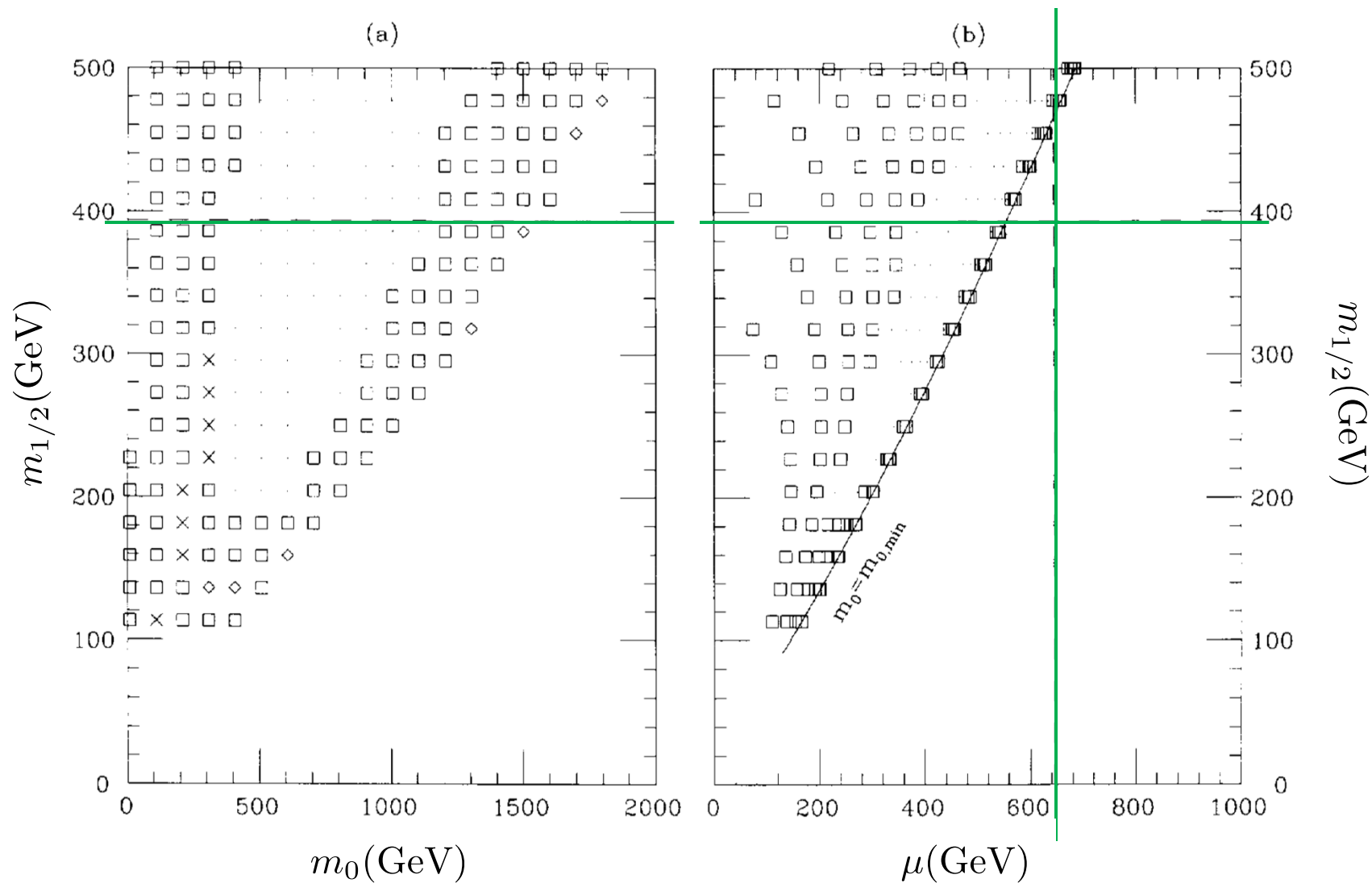
$$m_{1/2}^2 < \frac{\eta_f}{2|J_1|} M_Z^2, \quad m_0^2 < \frac{\eta_f}{2|J_2|} M_Z^2, \quad \mu^2 < \frac{\eta_f}{2} M_Z^2 \simeq (640 \text{ GeV})^2$$

mSUGRA/CMSSM



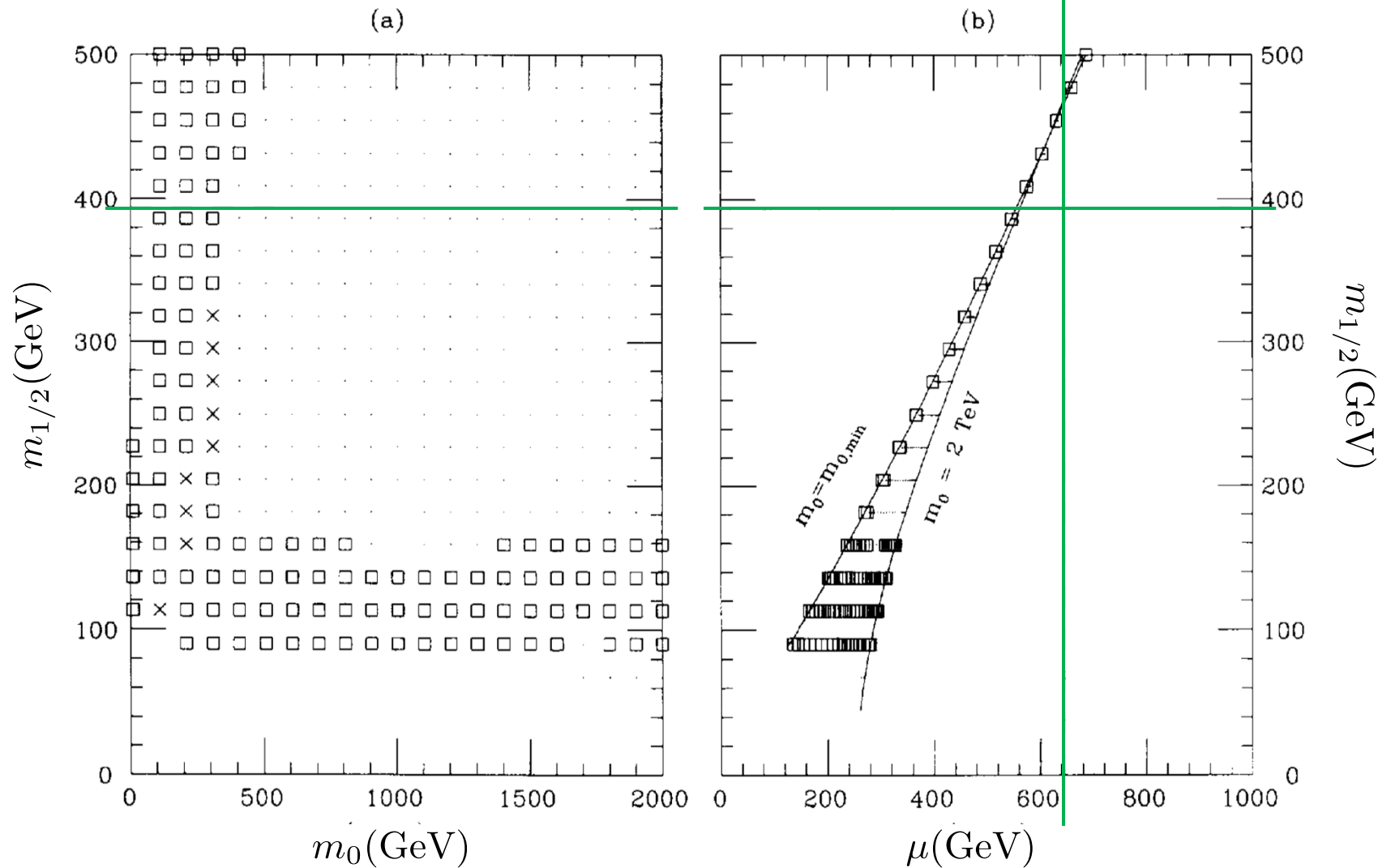
$$\tan \beta = 8; \quad \delta_1 = 0; \quad \delta_2 = 0$$

Non-universal model



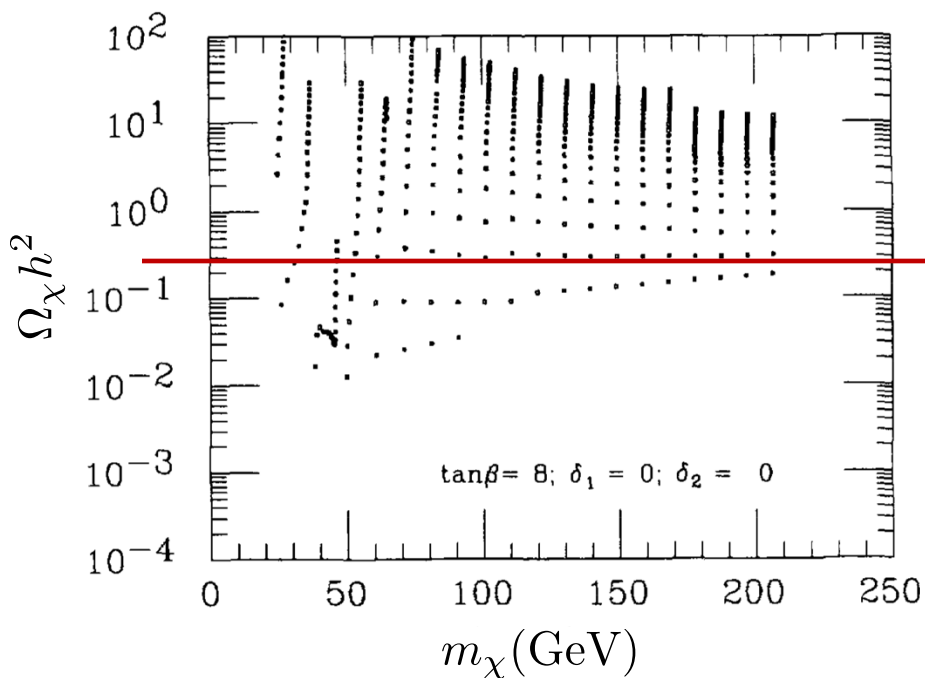
$$\tan \beta = 8; \quad \delta_1 = -0.2; \quad \delta_2 = 0.4$$

Non-universal model

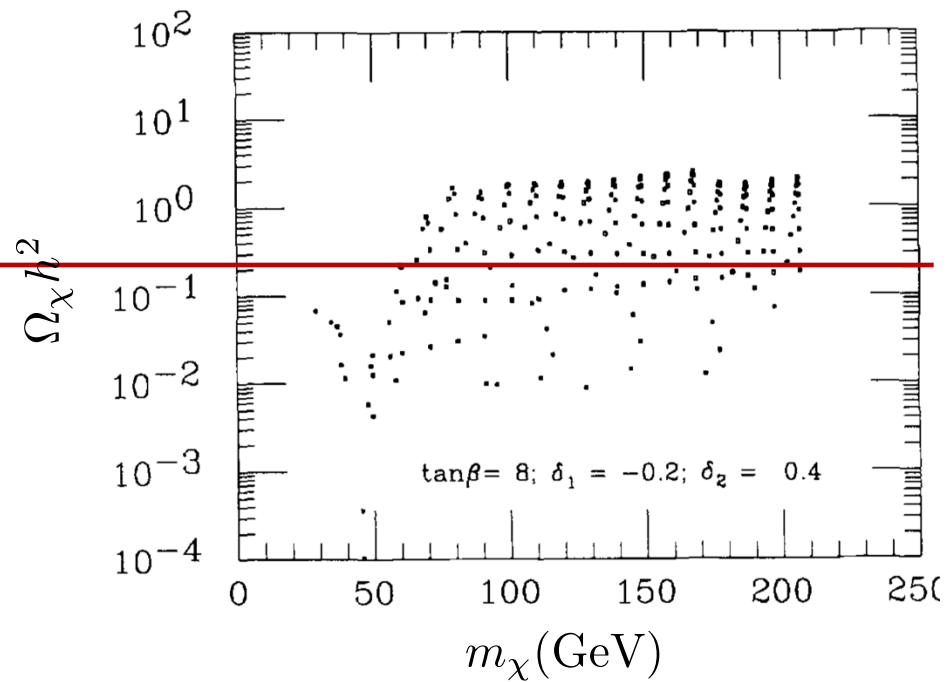


$$\tan \beta = 8; \quad \delta_1 = -0.8; \quad \delta_2 = 0.2$$

Neutralino relic abundance



mSUGRA/CMSSM



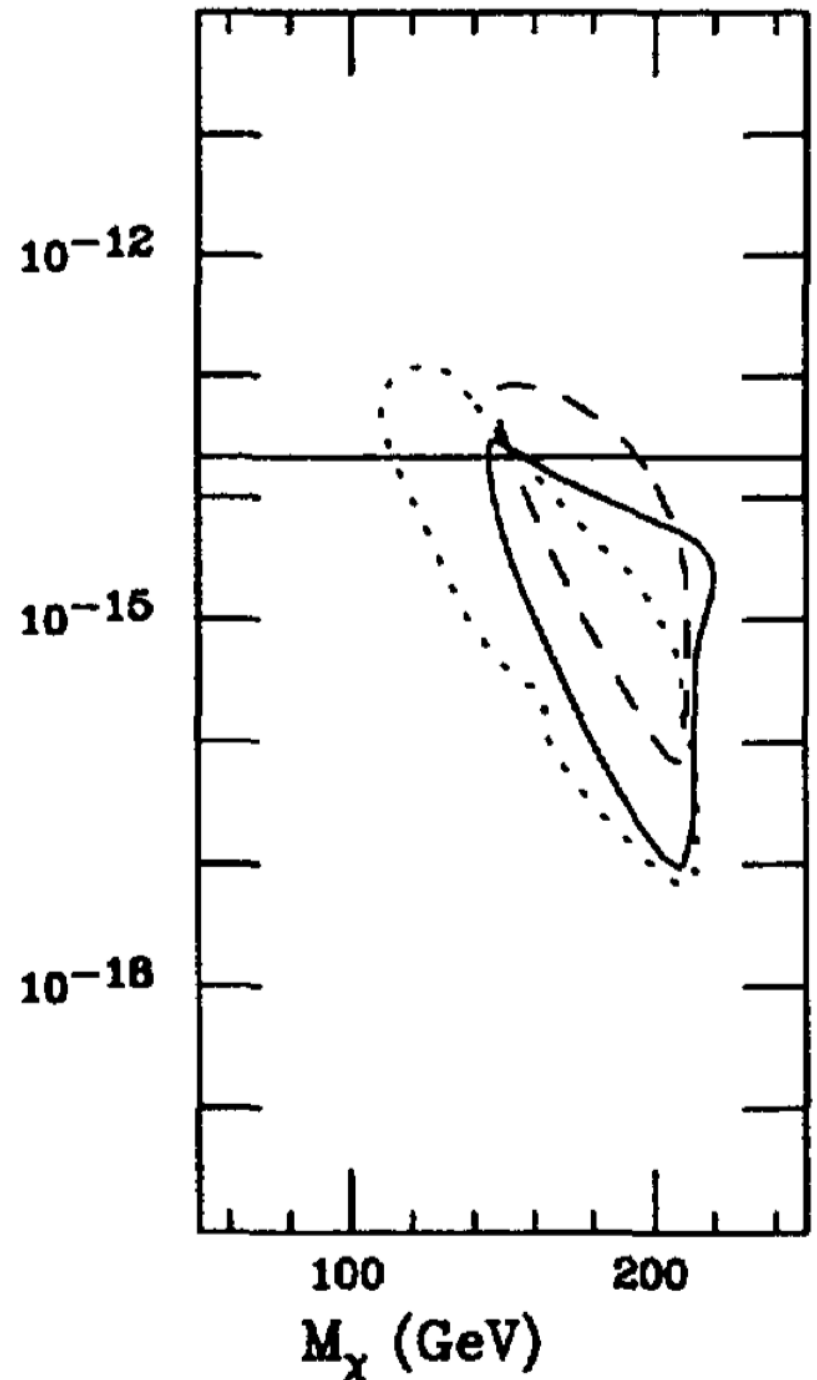
Non-universal model

Neutrino flux from the Earth (indirect detection)

Solid: mSUGRA/CMSSM

Dashed, dotted: Non-universal models

Flux ($\text{cm}^{-2} \text{sec}^{-1}$)



Gaugino non-universality

PHYSICAL REVIEW D

Light relic neutralinos

A. Bottino, N. Fornengo, and S. Scopel
Phys. Rev. D **67**, 063519 – Published 28 March 2003

- By allowing non-universality in the gaugino sector, neutralino lighter than the “canonical” ones become possible, by evading the bound induced by the non-observation of the chargino at accelerators

	Neutralino		Chargino
$M_N \equiv$	$\begin{pmatrix} M_1 & 0 & -\frac{1}{2}g'v_d & \frac{1}{2}g'v_u \\ 0 & M_2 & \frac{1}{2}gv_d & -\frac{1}{2}gv_u \\ -\frac{1}{2}g'v_d & \frac{1}{2}gv_d & 0 & -\mu \\ \frac{1}{2}g'v_u & -\frac{1}{2}gv_u & -\mu & 0 \end{pmatrix}$	$M_C \equiv$	$\begin{pmatrix} M_2 & \frac{1}{\sqrt{2}}gv_u \\ \frac{1}{\sqrt{2}}gv_d & \mu \end{pmatrix}$

SUGRA inspired

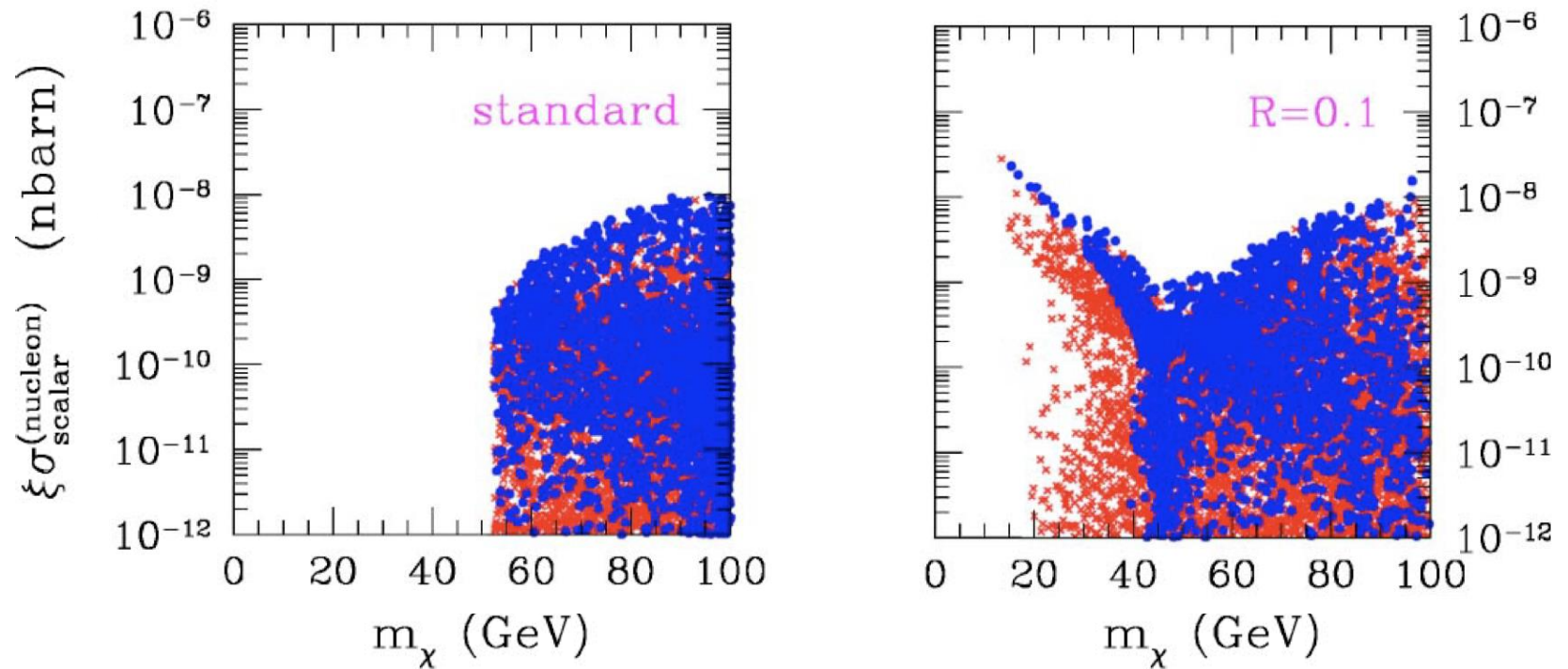
$$M_1 : M_2 : M_3 = \alpha_1 : \alpha_2 : \alpha_3$$

$$M_1 = \frac{5}{3} \tan^2 \theta_W M_2 \simeq 0.5 M_2$$

Gaugino non-universal

$$M_1 = R \times M_2$$

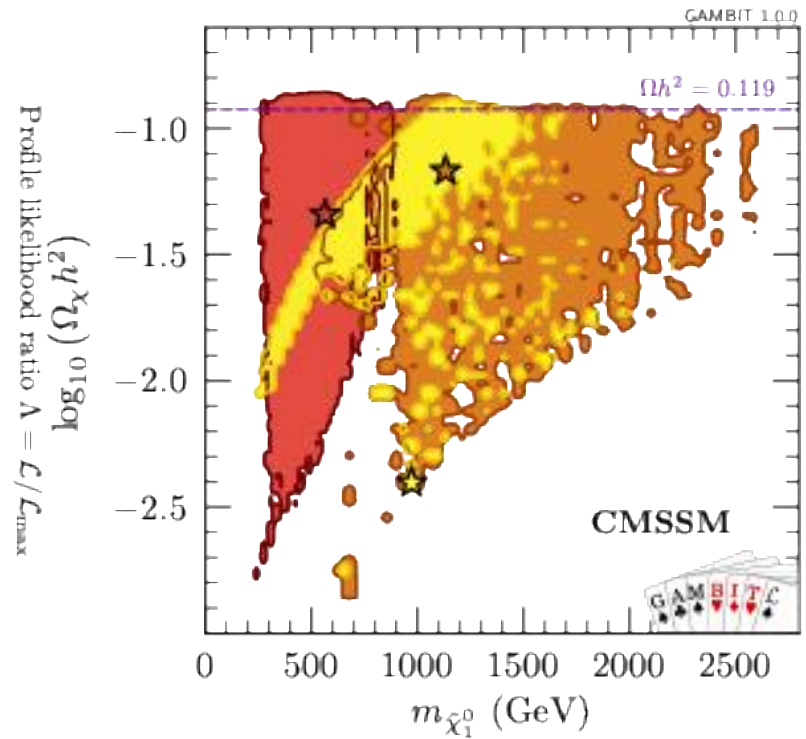
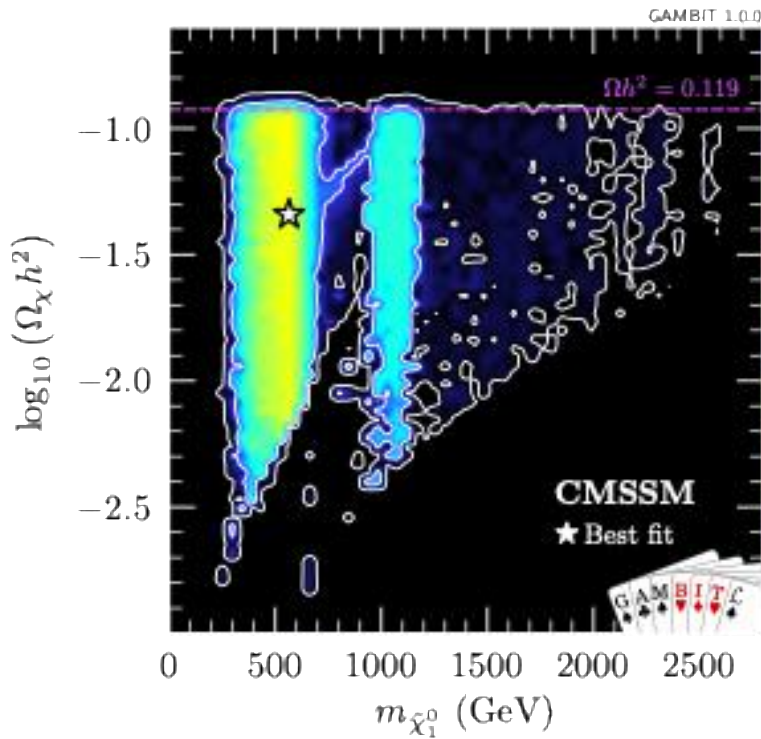
DM direct detection



Fast forward 25+ years

- The Higgs boson has been discovered (this fixes the Higgs mass and its coupling to the SM particles): what about 2HDM?
- Plenty of new data from the LHC
- No evidence of SUSY has been found
 - Several SUSY frameworks have been investigated since 1996:
 - mSUGRA/CMSSM
 - NUHM1, NUHM2
 - NMSSM
 - pMSSM
 - AMSB models
 - GMSB models
 - ...
- Dark Matter remains an unsolved outstanding problem
- Significant improvements in sensitivity for both direct detection and indirect detection searches

Constrained MSSM

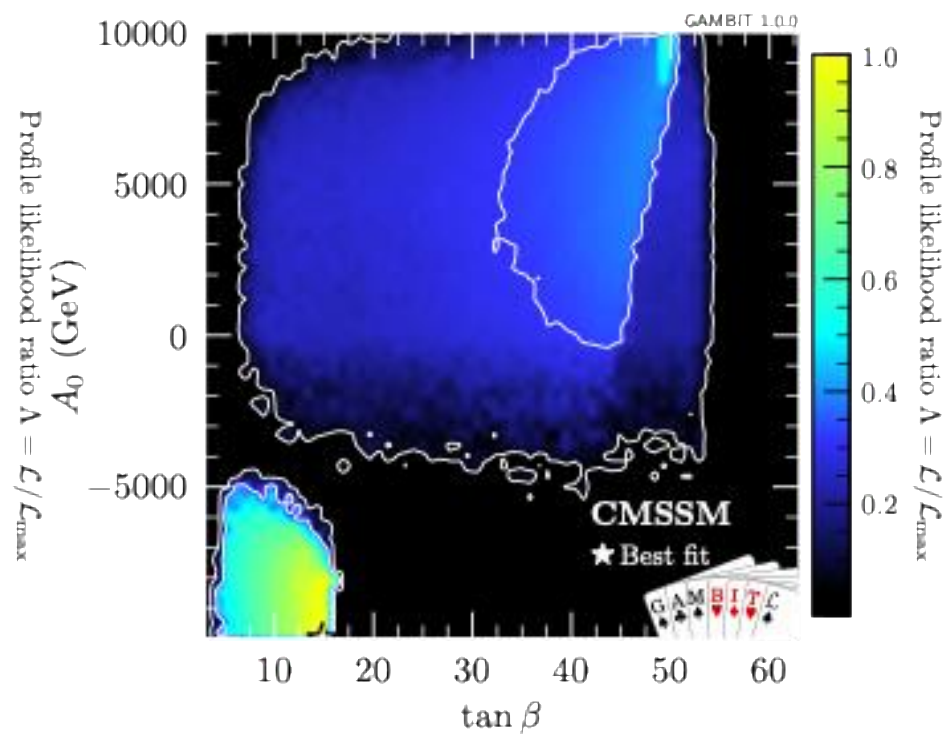
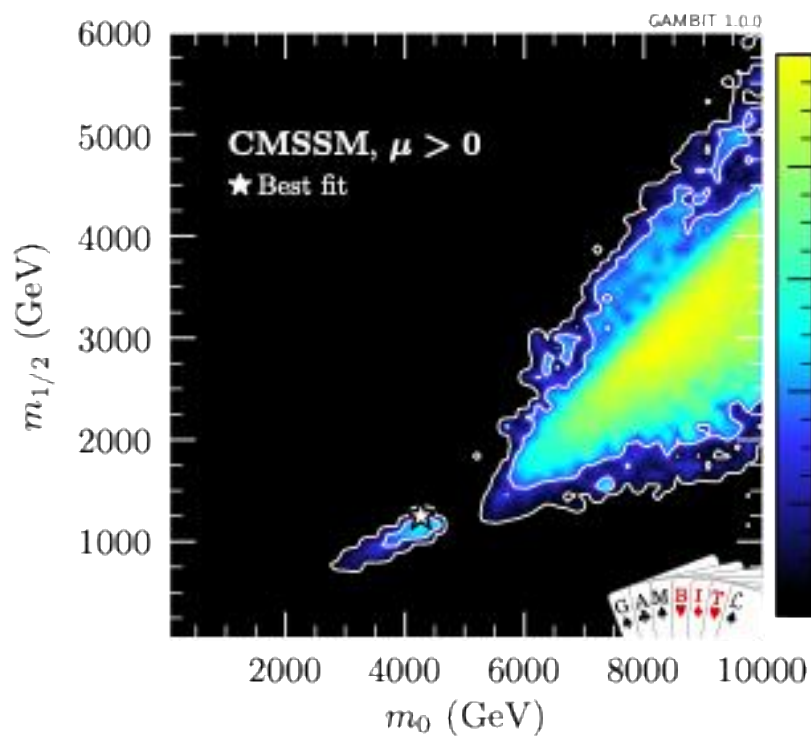


■ \tilde{t}_1 co-annihilation

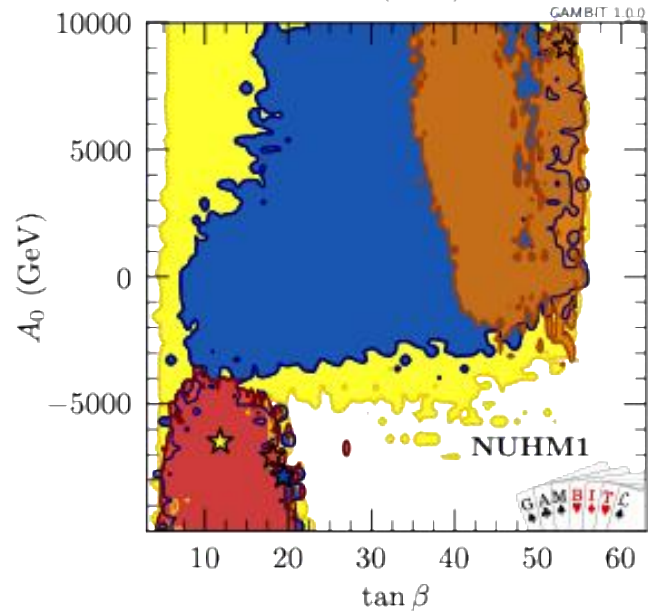
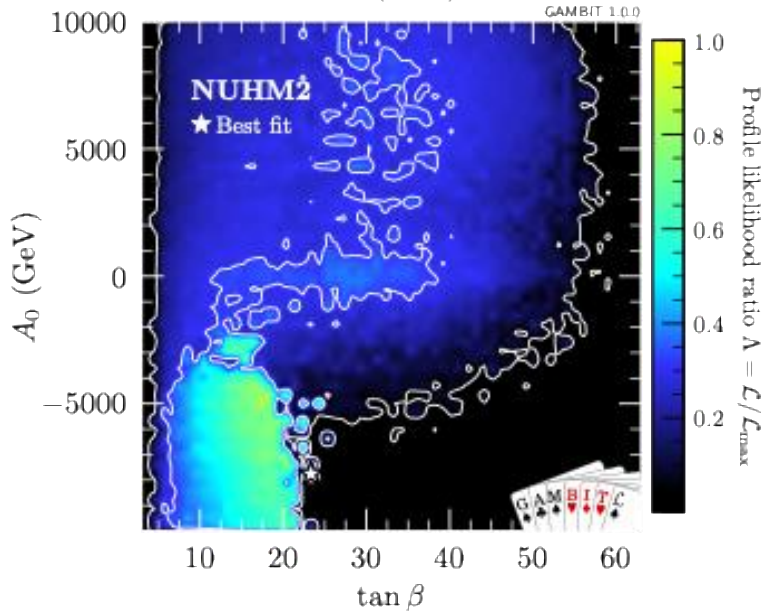
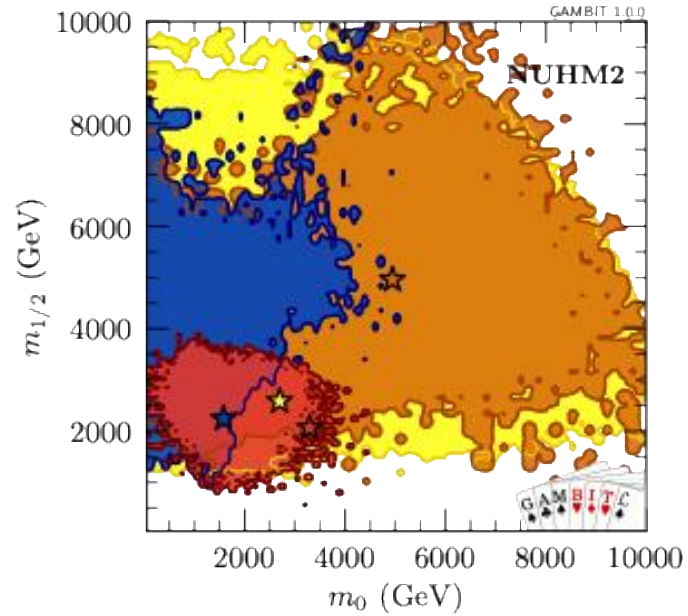
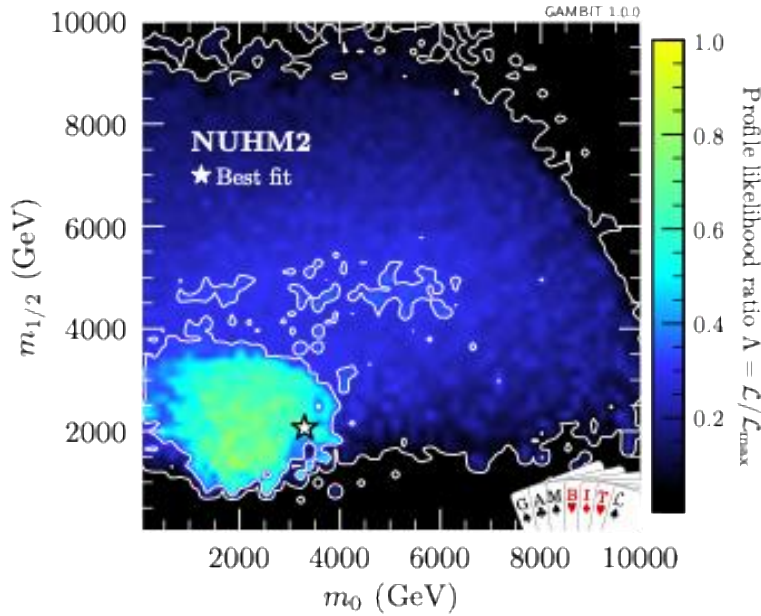
■ A/H funnel

■ $\tilde{\chi}_1^\pm$ co-annihilation

Constrained MSSM

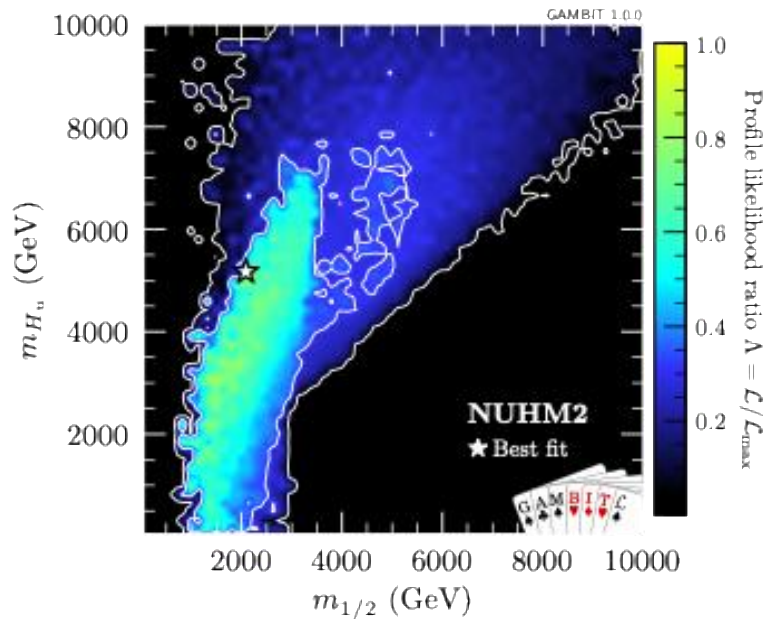
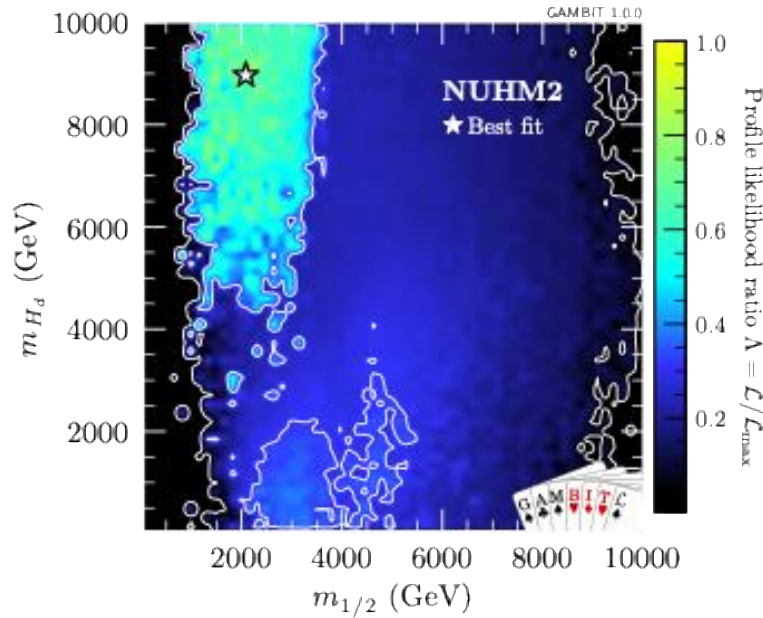


Non Universal Scalars (NUHM2)



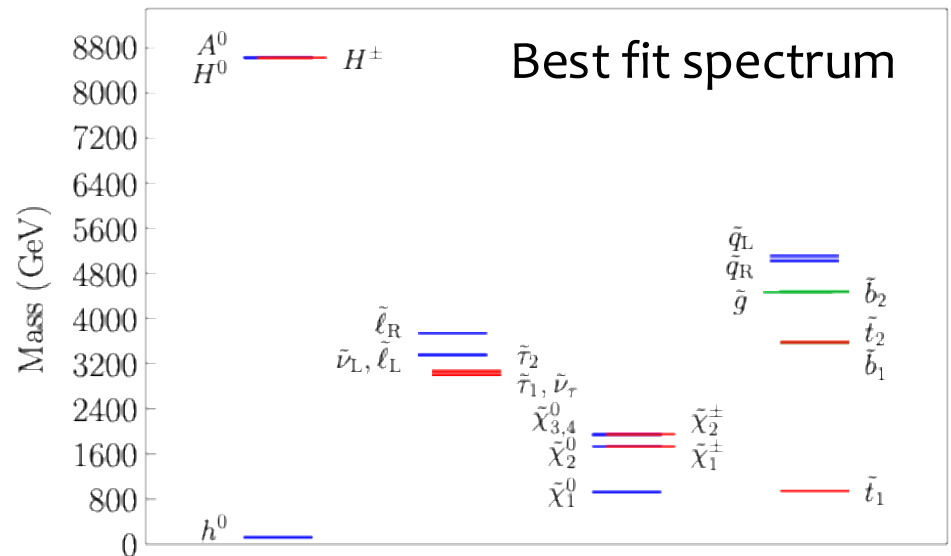
- \tilde{t}_1 co-annihilation
- A/H funnel
- $\tilde{\chi}_1^\pm$ co-annihilation
- $\tilde{\tau}_1$ co-annihilation

Non Universal Scalars (NUHM2)

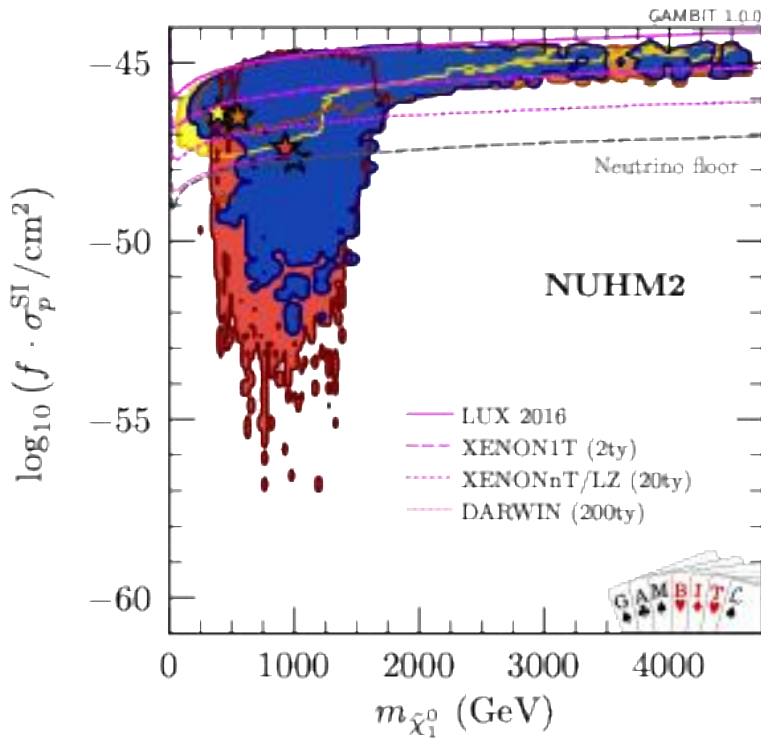


Models bounded by:

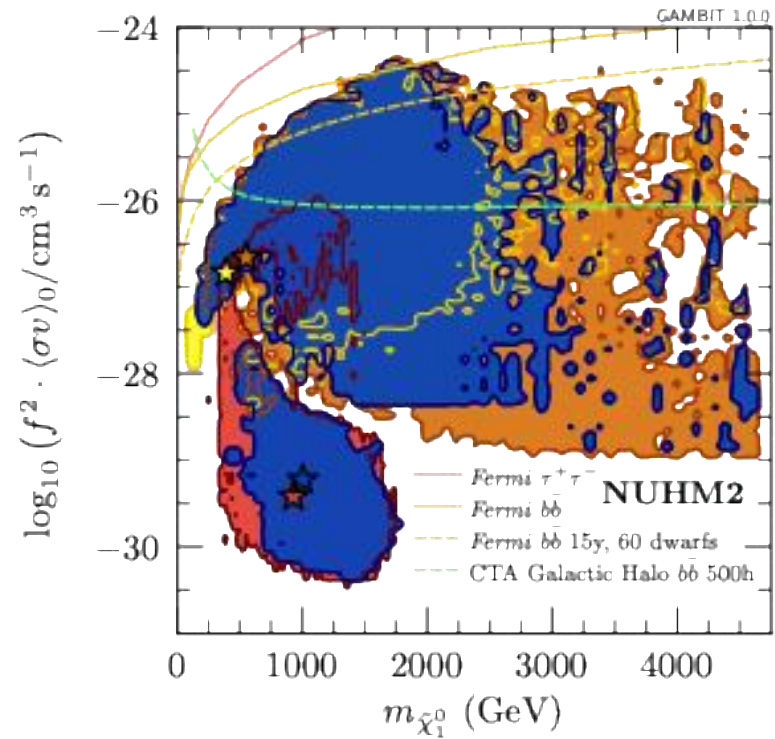
- Direct and indirect dark matter searches
- Electroweak precision and flavour observables
- Direct searches for supersymmetry at LEP and Runs I and II of the LHC
- Constraints from Higgs observables



Non Universal Scalars (NUHM2)



Direct DM searches



Indirect DM searches

Conclusions

Venya, as well as for cosmic rays, neutrinos and other astroparticle physics topics, has been a pioneer also in the study of dark matter, bringing new bright ideas both to theory and phenomenology

Small scale clumps and minihalos: 1405.2204, 1308.6742, 1107.2751, 1102.3445, 1002.3445, 1002.3444, 0712.3499, astro-ph/0511494, astro-ph/0301551

Superheavy DM: 0810.3012, astro-ph/0604311

Neutralino stars: astro-ph/9610060

DM annihilation in the galactic center: hep-ph/9402215

Majoron DM: hep-ph/9309214

DM distribution in the galaxy and implications for SUSY: PLB294(1992)221

Gamma-ray line from DM annihilation: PLB274(1992)122

Signatures of broken R-parity SUSY: PLB286(1991)382

Cosmology of gravitinos: PLB261(1991)71

It has been a privilege to work with Venya and to have the opportunity to learn so many things from his deep knowledge and understading of Physics!