Constraining cosmological models with the Effective Field Theory of Large-Scale Structures

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The Effective Field Theory of Large-Scale Structures (EFTofLSS)

Before EFTofLSS…

There are **two main** historical ways of using LSS data:

1. From the linear perturbation theory:

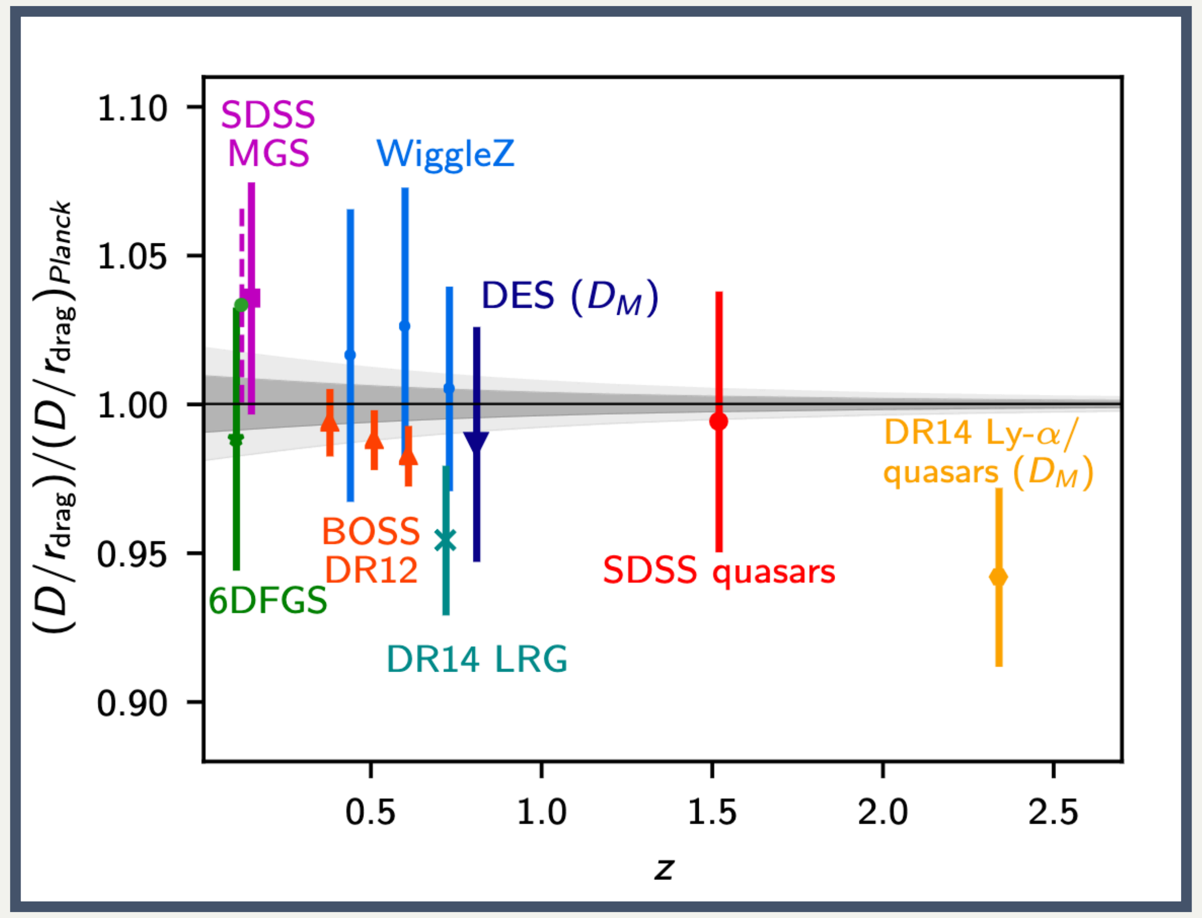
$$P_g(z,k,\mu) \simeq \left[b_1(z) + f\mu^2\right]^2 P_m(z,k)$$

Kaiser '87

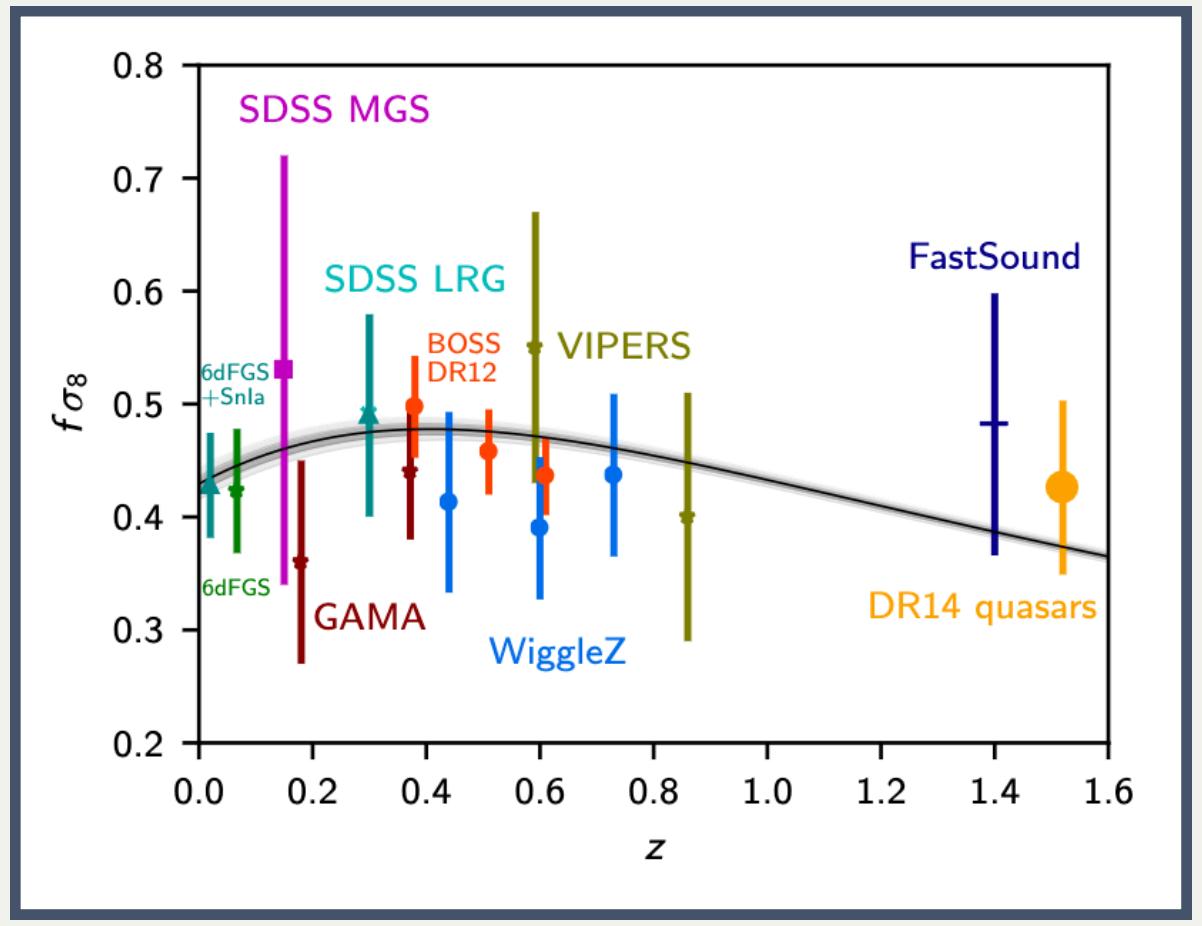
b_1: bias parameter, f: growth factor and $\mu = \hat{z} \cdot \hat{k}$

2. BAO angles + Redshift Space Distortion (RSD) information: BAO/$f\sigma_8$

LSS collaborations conventionally use the second method



Planck Collaboration [arXiv:1807.06209]



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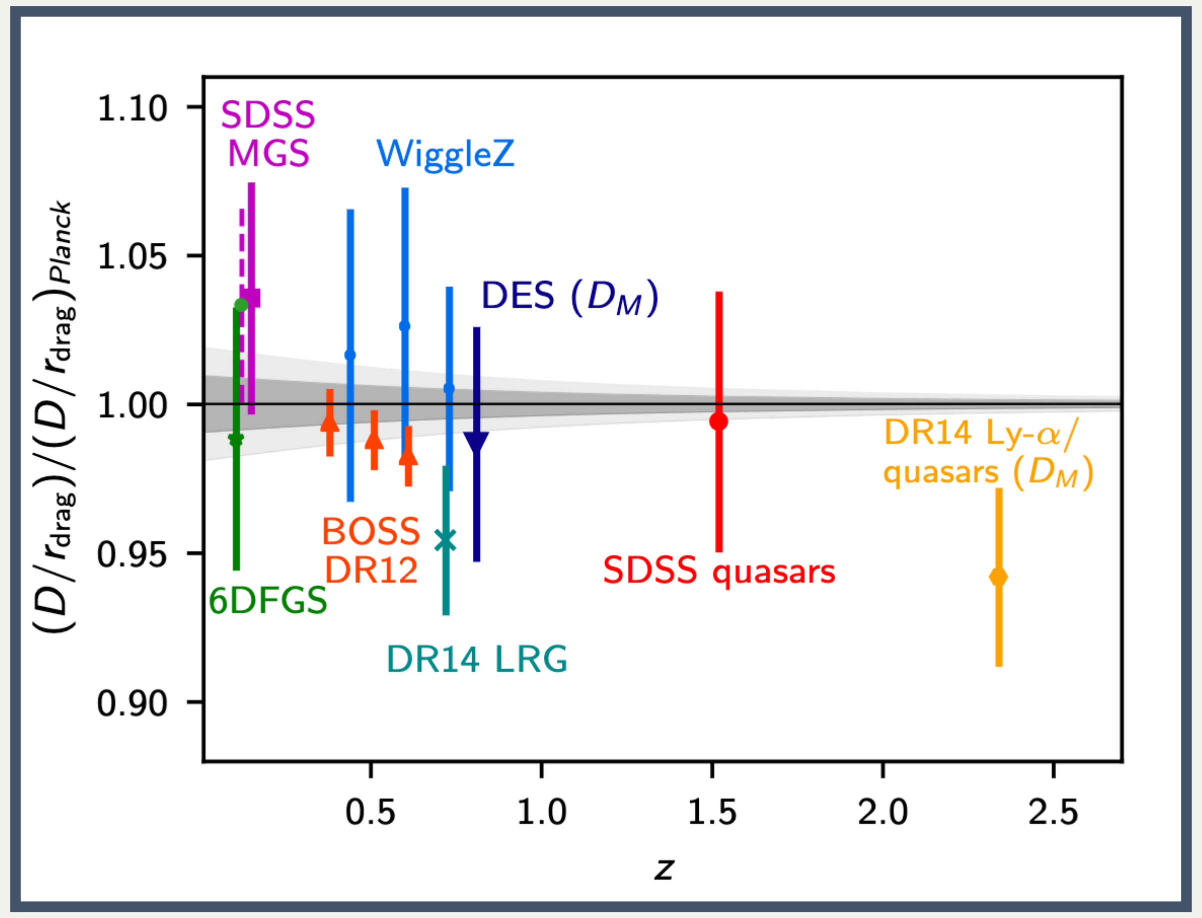

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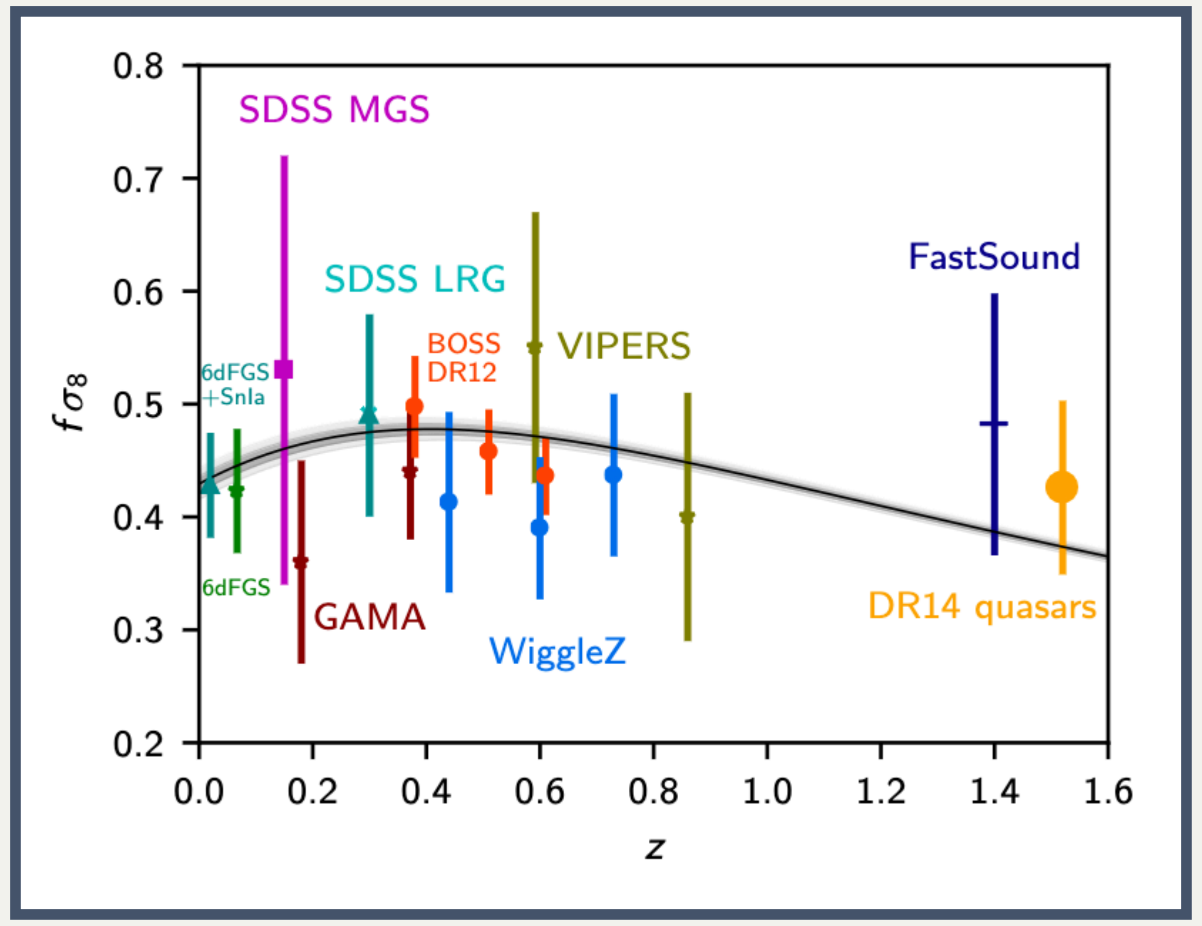

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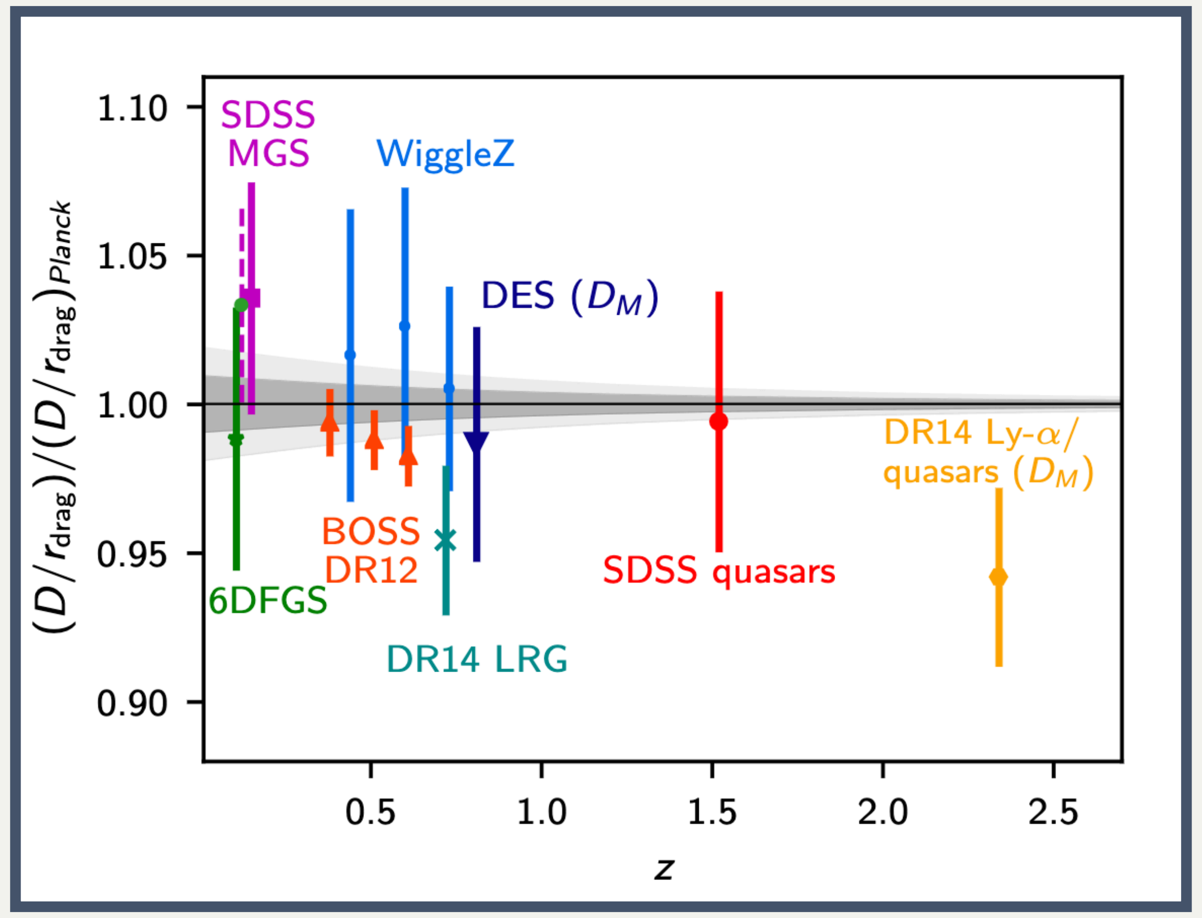

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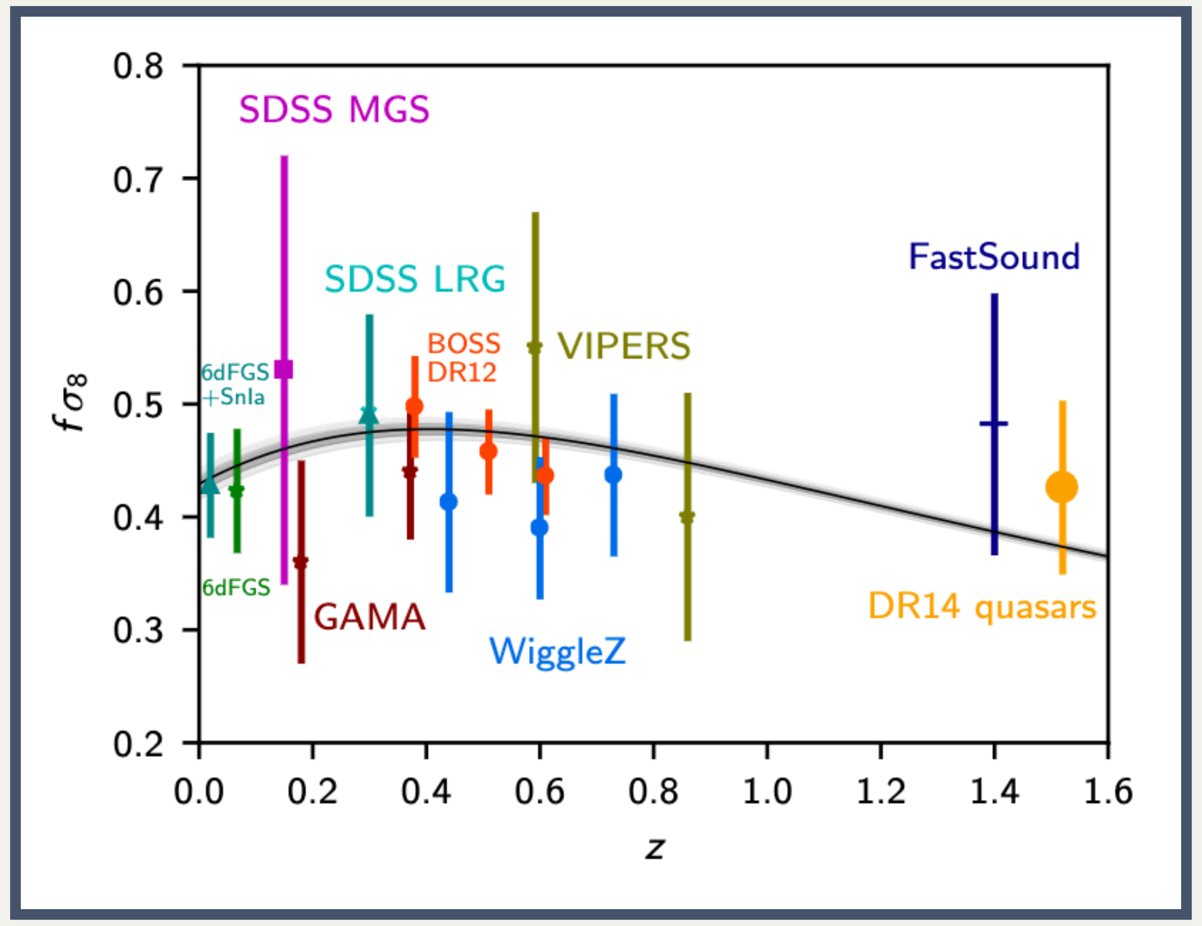

The effective field theory of large-scale structures (EFTofLSS)

Main ingredients

EFTofLSS

Carrasco++ [arXiv:1206.2926]

Baumann++ [arXiv:1004.2488]

The effective field theory of large-scale structures (EFTofLSS)

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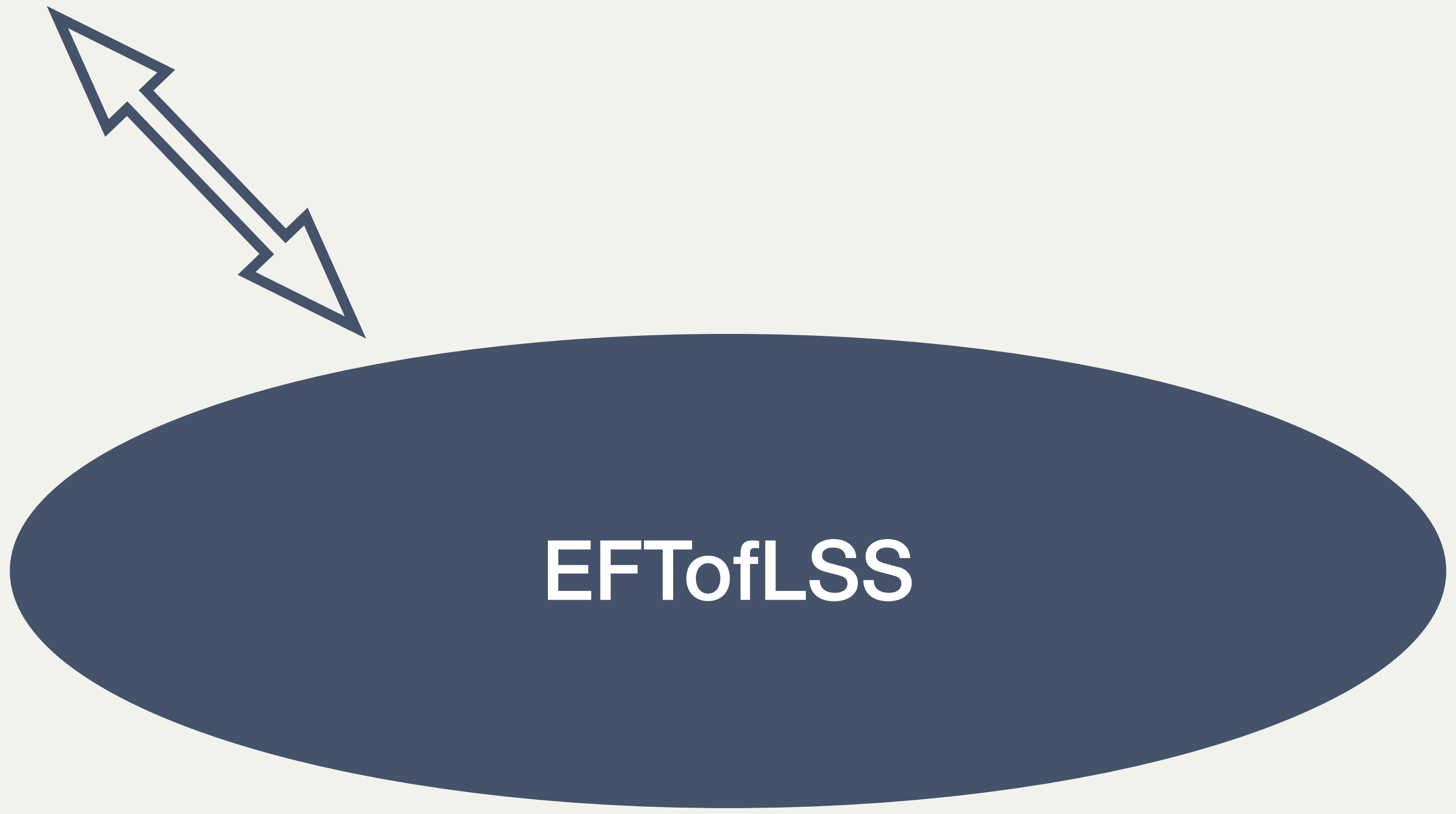
Solve cosmological equations only for **large-scale physics**

$$\delta(\boldsymbol{k}) = \delta_l(\boldsymbol{k}) + \delta_s(\boldsymbol{k})\,,$$
$$\delta_l(\boldsymbol{k}) = \delta(\|\boldsymbol{k}\| < \Lambda^{-1})$$



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Dark matter: the Vlasov system

$$\dot{\delta}_l + \theta_l = -\delta\theta_l - v^j\partial_j\delta_l\,,$$
$$\dot{\theta}_l + aH\theta_l + \nabla^2\psi_l = -v^j\partial_j\theta_l - \partial_i v_l^j \cdot \partial_j v_l^i - \boldsymbol{\partial_j}\left(\frac{\mathbf{1}}{\boldsymbol{\rho_l}}\boldsymbol{\partial_i}\left[\boldsymbol{\tau}^{ij}\right]_{\boldsymbol{\Lambda}}\right)$$

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EFTofLSS

$$\nabla^2\psi_l = \frac{3}{2}\Omega_m(a)(aH)^2\delta_l$$

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EFTofLSS

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Gravity: the Poisson equation

$$\vec{x} \to \vec{x} + \vec{a}$$
$$\vec{v} \to \vec{v} + \partial_t\vec{a}$$

Symmetries: Galilean invariance

Carrasco++ [arXiv:1206.2926]

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The effective field theory of large-scale structures (EFTofLSS)

Main steps

Step by Step…

1- **Solve** dark matter equations **perturbatively**: $\delta_l(\vec{x},t) = \delta_l^{(1)}(\vec{x},t) + \delta_l^{(2)}(\vec{x},t) + \ldots + \delta_l^{(n)}(\vec{x},t)$ *Bernardeau++ '01*

2- Obtain the **mildly non-linear matter power spectrum**:

Carrasco++ [arXiv:1206.2926]

$$P_m(k,\tau) = P_{11}(k,\tau) + P_{22}^{\Lambda}(k,\tau) + 2P_{13}^{\Lambda}(k,\tau) + 2P_{c^2_{\rm comb}}^{\Lambda}(k,\tau)$$

Senatore [arXiv:1406.7843]

Mirbabayi++ [arXiv:1412.5169]

3- Write down **all possible operators in the galaxy bias expansion**: $\delta_g = b_1\,\delta_l^{(1)} + b_2\,\delta_l^{(2)} + R_*^2\,\partial^2\delta_l^{(1)} + \ldots$

4- Take into account the **redshift-space distortion** (RSD) effect (to subtract the contribution of the peculiar velocity of the galaxies to the cosmological redshift) *Senatore++ [arXiv:1409.1225]*

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Tree-level — One-loop level — Counterterm

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Galaxy power spectrum

The **galaxy power spectrum** in the EFTofLSS framework:

$$P_g(k,\mu) \simeq \left[b_1 + f\mu^2\right]^2 P_m(k) = Z_1(\mu)^2 P_m(k)$$

We go from 1 to 10 free parameters

$$P_g(k,\mu) = Z_1(\mu)^2 P_{11}(k) + 2Z_1(\mu)P_{11}(k)\left(c_{\rm ct}\frac{k^2}{k_{\rm M}^2} + c_{r,1}\mu^2\frac{k^2}{k_{\rm R}^2} + c_{r,2}\mu^4\frac{k^2}{k_{\rm R}^2}\right)$$
$$+ 2\int\frac{d^3q}{(2\pi)^3}\, Z_2(\boldsymbol{q},\boldsymbol{k}-\boldsymbol{q},\mu)^2 P_{11}(|\boldsymbol{k}-\boldsymbol{q}|)P_{11}(q) + 6Z_1(\mu)P_{11}(k)\int\frac{d^3q}{(2\pi)^3}\, Z_3(\boldsymbol{q},-\boldsymbol{q},\boldsymbol{k},\mu)P_{11}(q)$$
$$+ \frac{1}{\bar{n}_g}\left(c_{\epsilon,0} + c_\epsilon^{\rm mono}\frac{k^2}{k_{\rm M}^2} + 3c_\epsilon^{\rm quad}\left(\mu^2 - \frac{1}{3}\right)\frac{k^2}{k_{\rm M}^2}\right),$$

Carrasco++ [arXiv:1206.2926] ; Baumann++ [arXiv:1004.2488]
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$$\begin{aligned}P_g(k,\mu) &= Z_1(\mu)^2 P_{11}(k) + 2 Z_1(\mu) P_{11}(k)\left(c_{ct}\frac{k^2}{k_{\rm M}^2} + c_{r,1}\mu^2\frac{k^2}{k_{\rm R}^2} + c_{r,2}\mu^4\frac{k^2}{k_{\rm R}^2}\right)\\ &+ 2\int\frac{d^3q}{(2\pi)^3} Z_2(\boldsymbol{q},\boldsymbol{k}-\boldsymbol{q},\mu)^2 P_{11}(|\boldsymbol{k}-\boldsymbol{q}|)P_{11}(q) + 6 Z_1(\mu)P_{11}(k)\int\frac{d^3q}{(2\pi)^3} Z_3(\boldsymbol{q},-\boldsymbol{q},\boldsymbol{k},\mu)P_{11}(q)\\ &+ \frac{1}{\bar{n}_g}\left(c_{\epsilon,0} + c_\epsilon^{\rm mono}\frac{k^2}{k_{\rm M}^2} + 3c_\epsilon^{\rm quad}\left(\mu^2 - \frac{1}{3}\right)\frac{k^2}{k_{\rm M}^2}\right),\end{aligned}$$

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2 renormalization scales

(green circle) Renormalization scale controlling the **spatial derivative expansion**, given by the typical size of a **virialized halo**

(red circle) Renormalization scale of the **velocity products** appearing in the redshift-space expansion

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$P_g(k,\mu)$ can be determined directly from $P_{11}(k) = P_m^{\rm lin}(k)$

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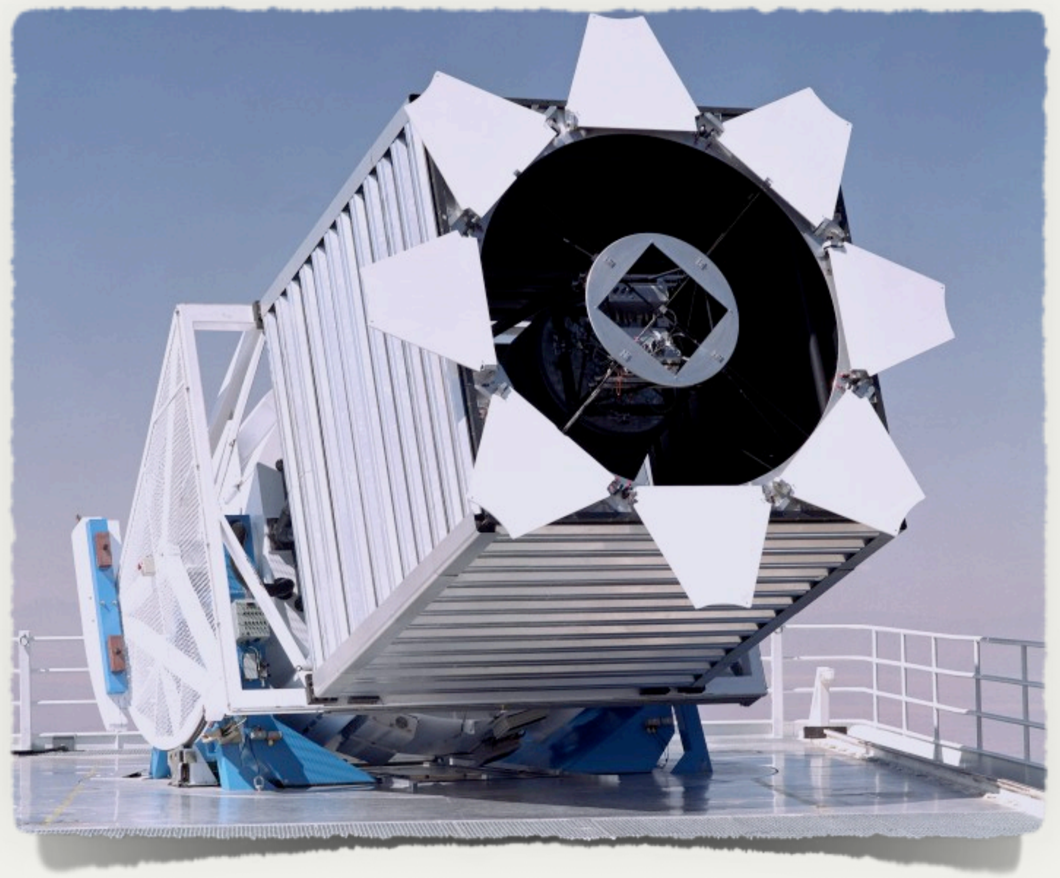
10 EFT parameters

4 parameters b_i $(i = 1,2,3,4)$ to describe the **galaxy bias** which arises from the one-loop contributions

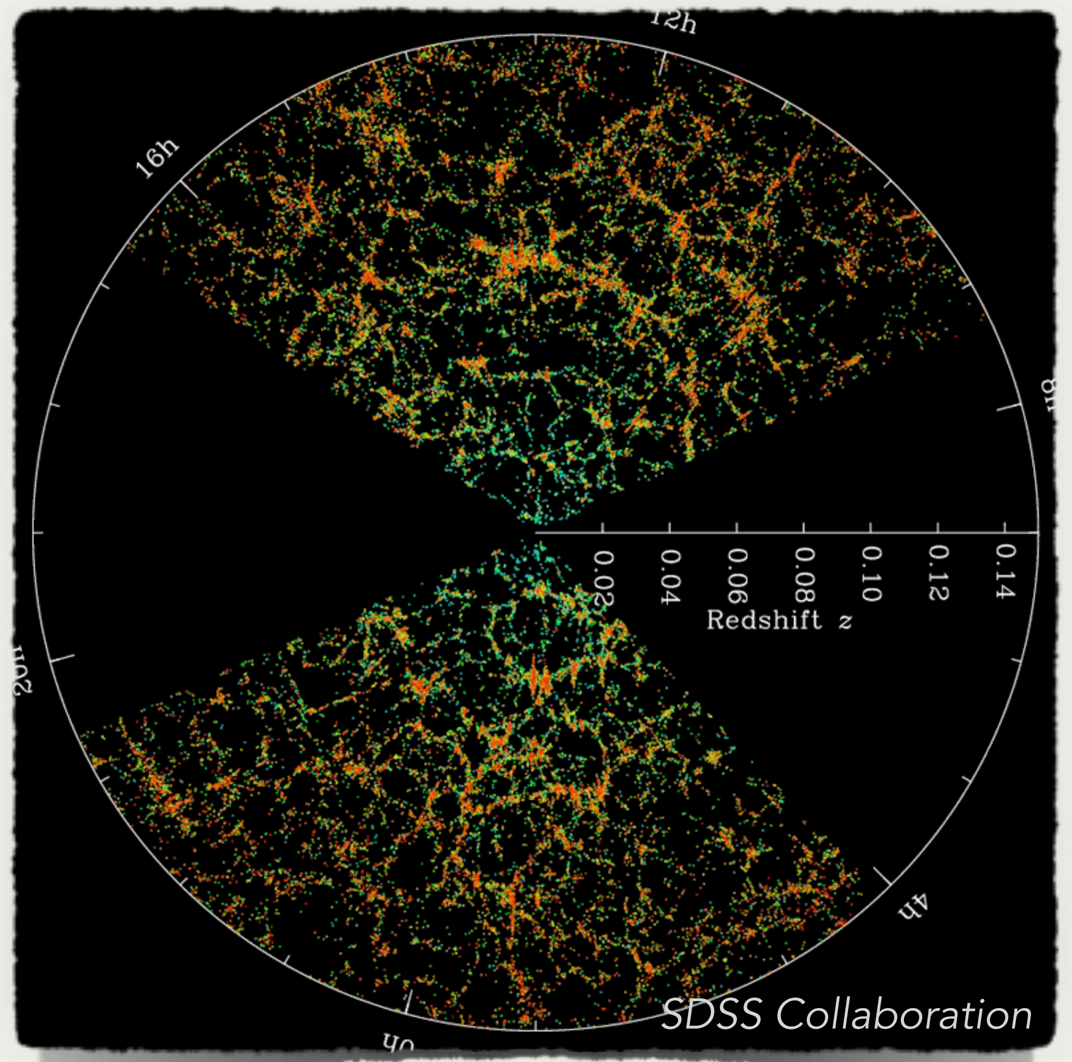
3 parameters corresponding to **counterterms** (c_{ct} linear combination of a higher derivative bias and the dark matter sound speed, while $c_{r,1}$ and $c_{r,2}$ are the redshift-space counterterms)

3 parameters which describe **stochastic** terms

The Sloan Digital Sky Survey (SDSS)



www.sdss.org



SDSS

2000 - 2025

BOSS DR12 LRG (Luminous Red Galaxies)

Galaxies (~ 1.5 million) selected in two redshift ranges:

→ LOWZ (SGC/NGC): $0.2 < z < 0.43$ ($z_{\rm eff} = 0.32$)

→ CMASS (SGC/NGC): $0.43 < z < 0.7$ ($z_{\rm eff} = 0.57$)

BOSS Collaboration [arXiv:1607.03155]

eBOSS DR16 QSO

Quasars (~ 300 000) selected in one redshift range:

$0.8 < z < 2.2$ ($z_{\rm eff} = 1.5$)

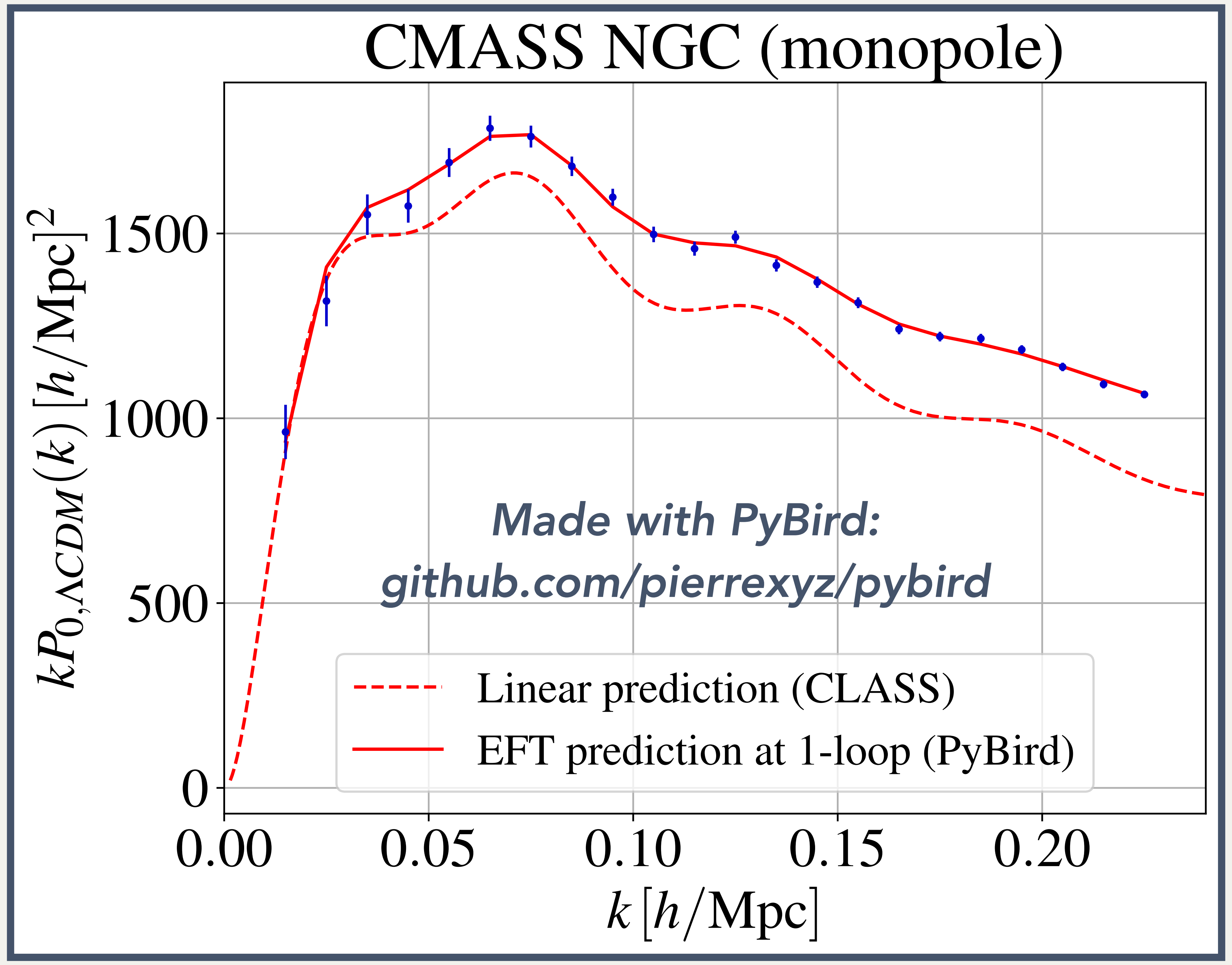
eBOSS Collaboration [arXiv:2007.08991]

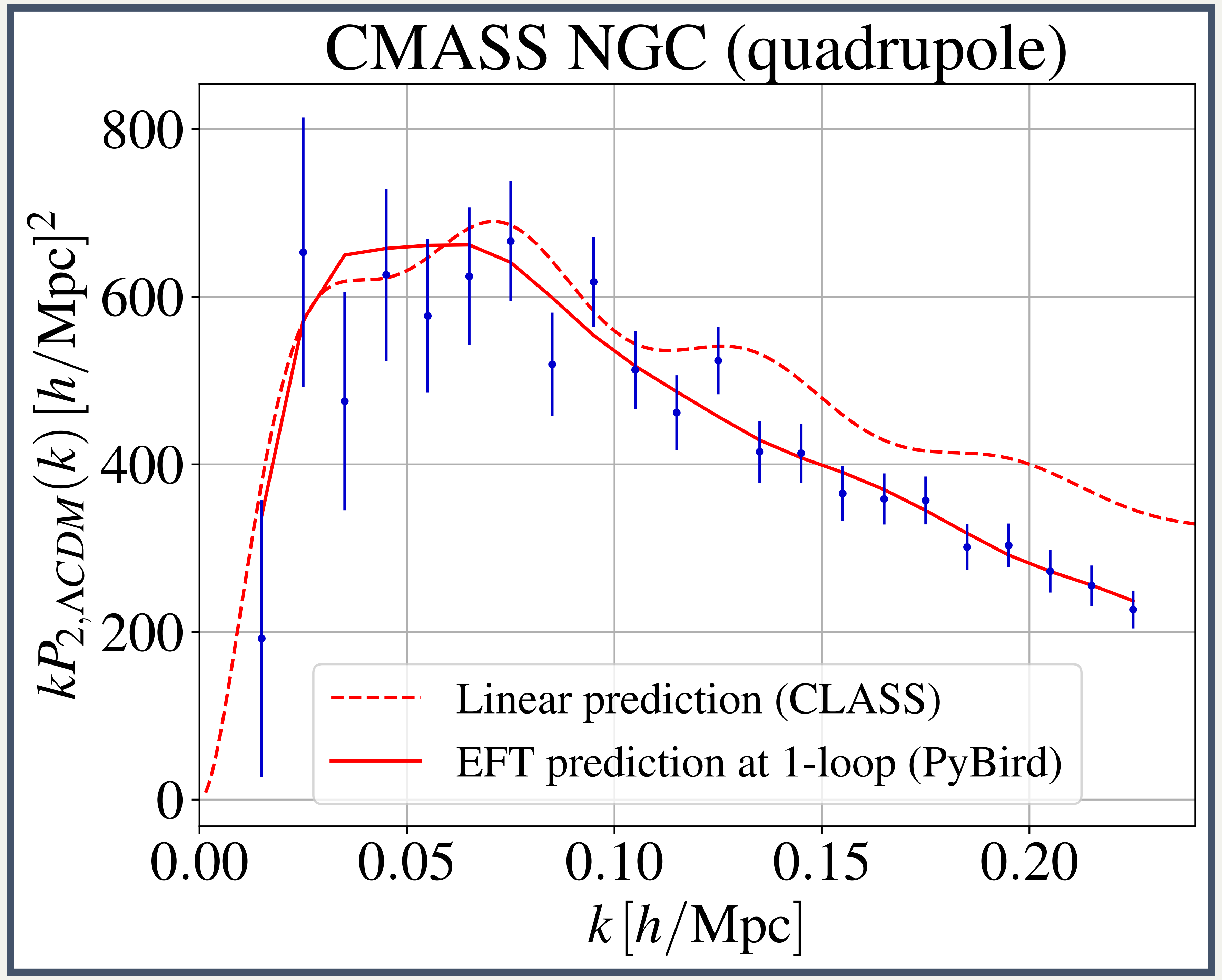
Applying EFTofLSS to SDSS data

Multipoles of the galaxy power spectrum, obtained through a **Legendre** polynomials ($\mathscr{L}_\ell$) decomposition:

$$P_g(z,k,\mu) = \sum_{\ell \text{ even}} \mathscr{L}_\ell(\mu) P_\ell(z,k)$$

→ the two main contributions to $P_g(z,k,\mu)$ are the **monopole** ($\ell = 0$) and the **quadrupole** ($\ell = 2$)





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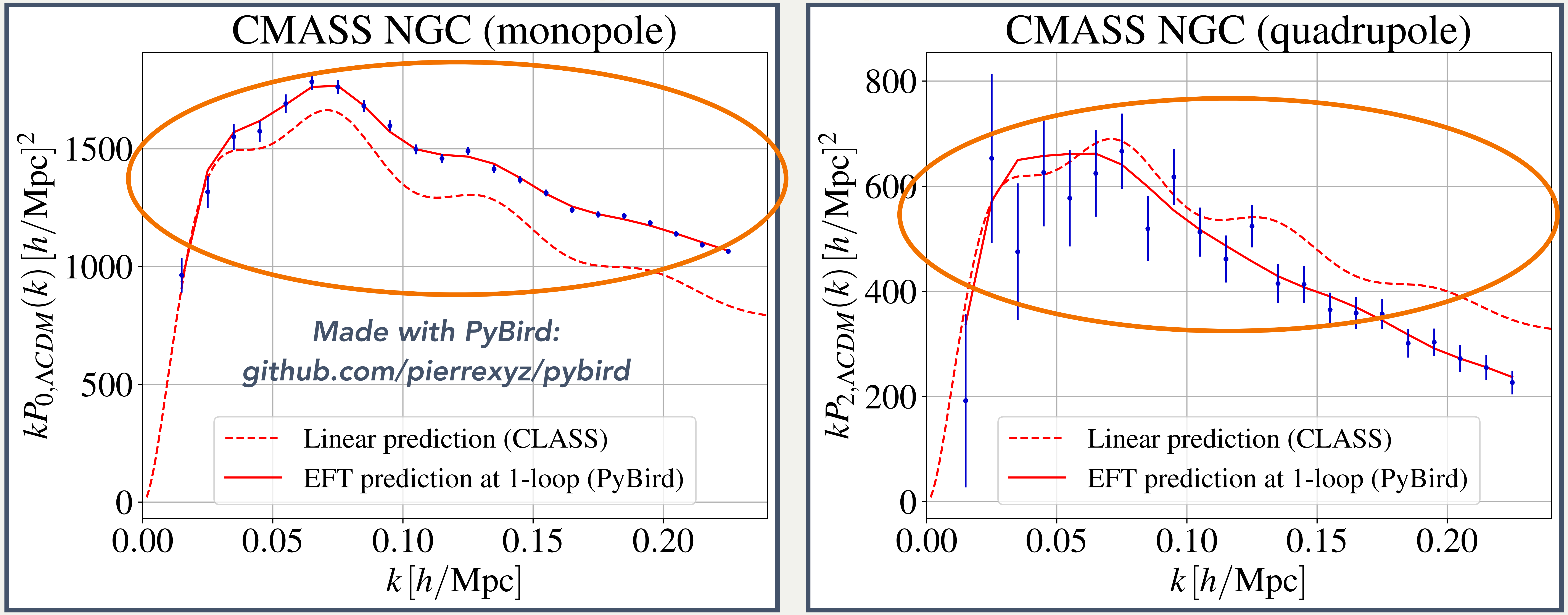
Improvement in precision!



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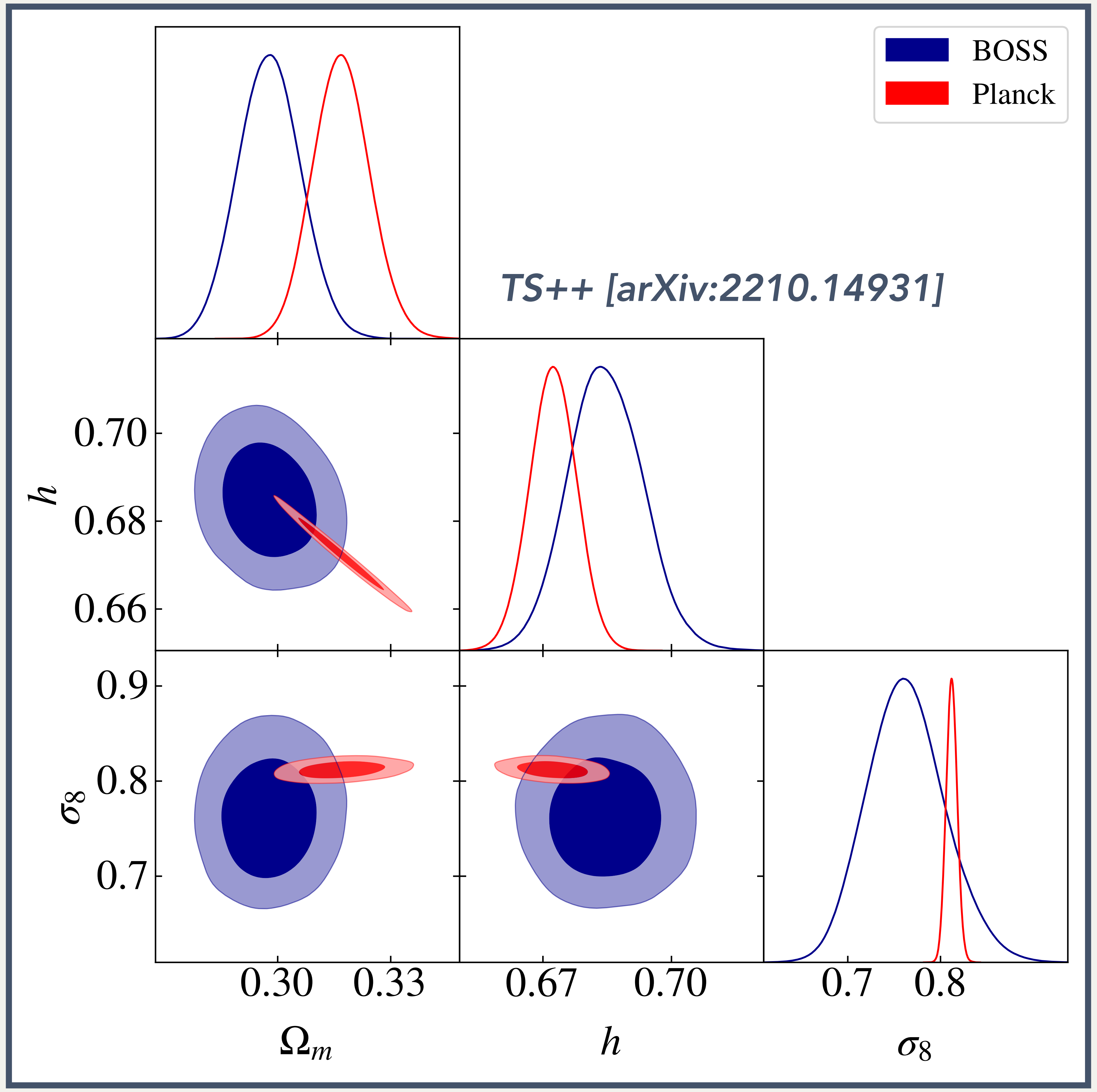
The effective field theory of large-scale structures (EFTofLSS)

Application to BOSS data



The EFTofLSS analysis of BOSS data allows to determine Ω_m and h at a **precision only** $10\,\%$ **and** $60\,\%$ worse than Planck

This is ~ 5.4 (for Ω_m) and ~ 3.2 (for h) times better than the BAO/$f\sigma_8$ analysis

See also D'Amico++ [arXiv:1909.05271] ; Philcox++ [arXiv:2002.04035]

On the consistency of EFTofLSS

Presentation of the problem

There are **several codes** in the literature with **different parametrizations**:

→ ***CLASS PT*** with the **EC** parametrization *Chudaykin++ [arXiv:2004.10607]*

→ ***PyBird*** with the **WC** parametrization *D'Amico++ [arXiv:2003.07956]*

(+ **Velocileptors** + **CLASS-OneLoop**)

Chen++ [arXiv:2005.00523] ; Linde++ [arXiv:2402.09778]

→ these codes use **different sets of priors** on EFT parameters

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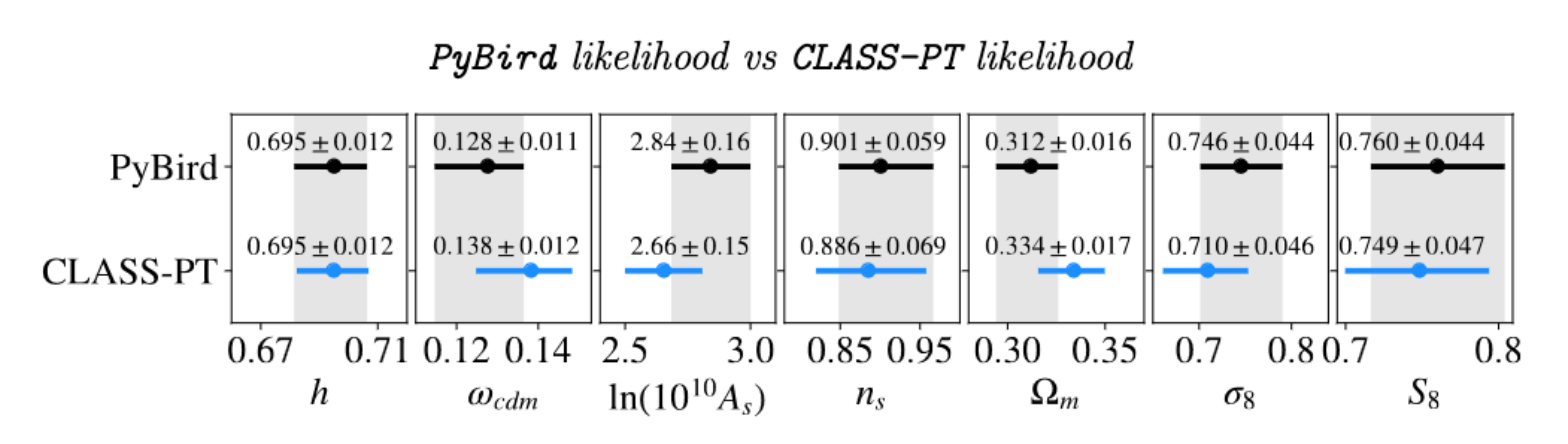

TS++ [arXiv:2208.05929]

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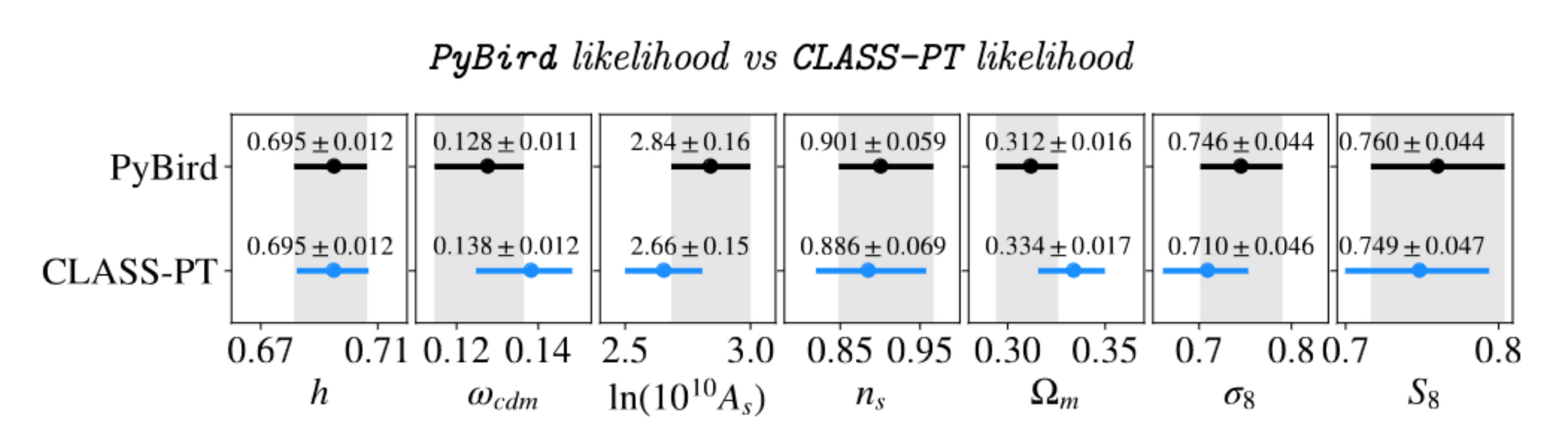

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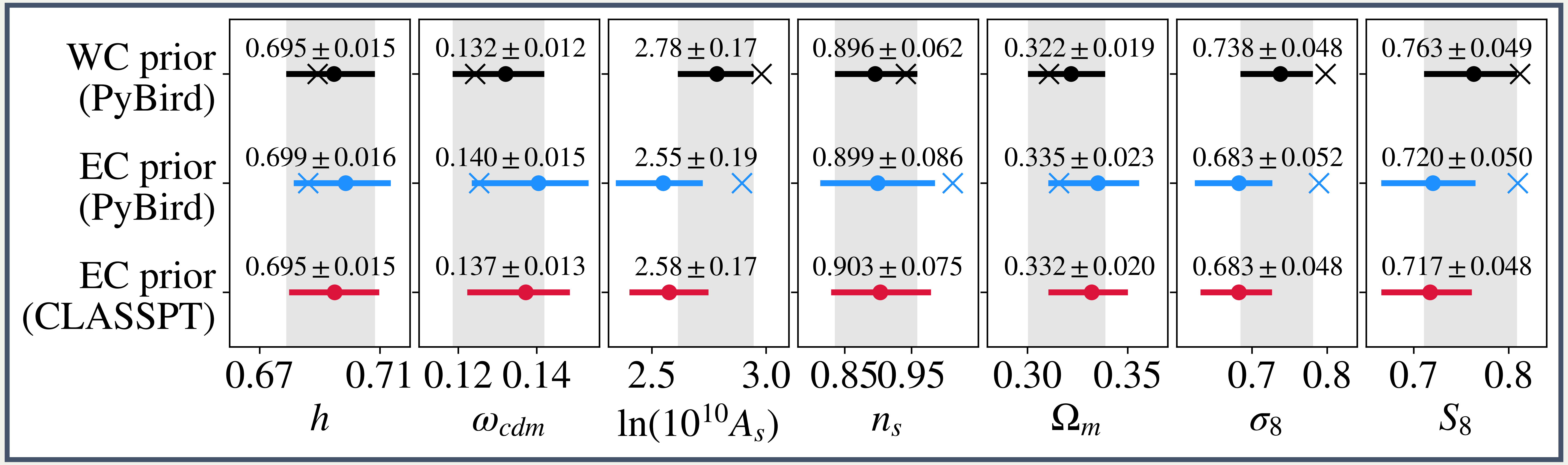
Data, theoretical **parametrizations** and **codes** are supposed to be **equivalent**: **what is going on?**

On the consistency of EFTofLSS

The EFT prior issue



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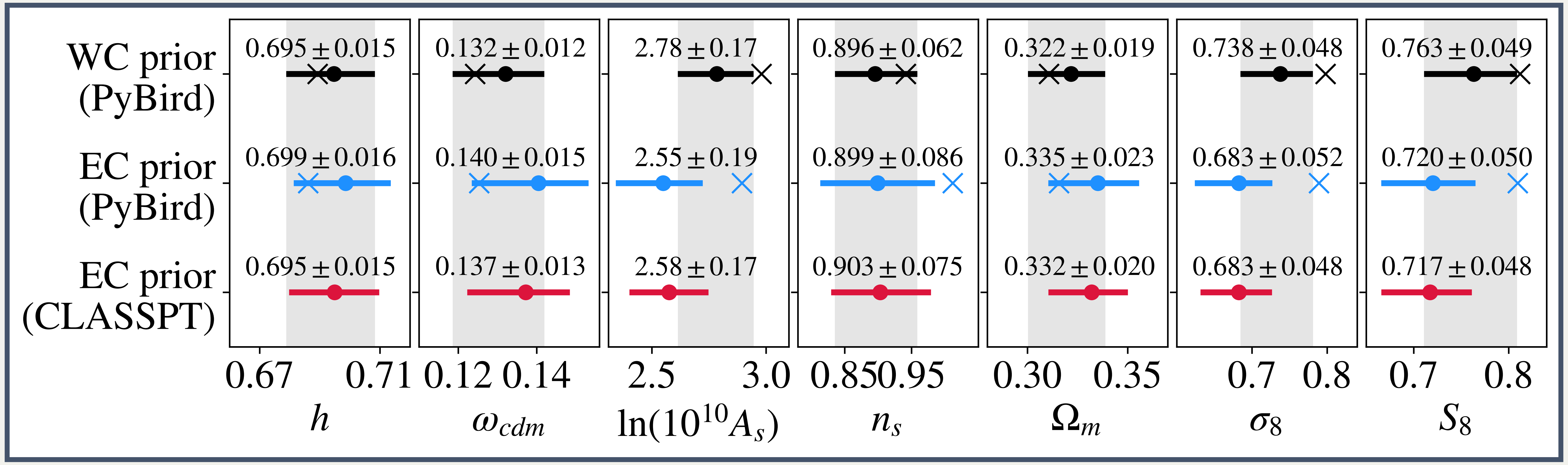


TS++ [arXiv:2208.05929]

Prior effects

→ ***The prior weight effect***: if the region allowed by the prior is far from the true value of a parameter, then the posterior and bestfit will be shifted from the true value

→ ***The prior volume effect***: a posterior depends on the volume enclosed by the priors $\Longrightarrow$ large parameter regions are emphasized compared to smaller regions

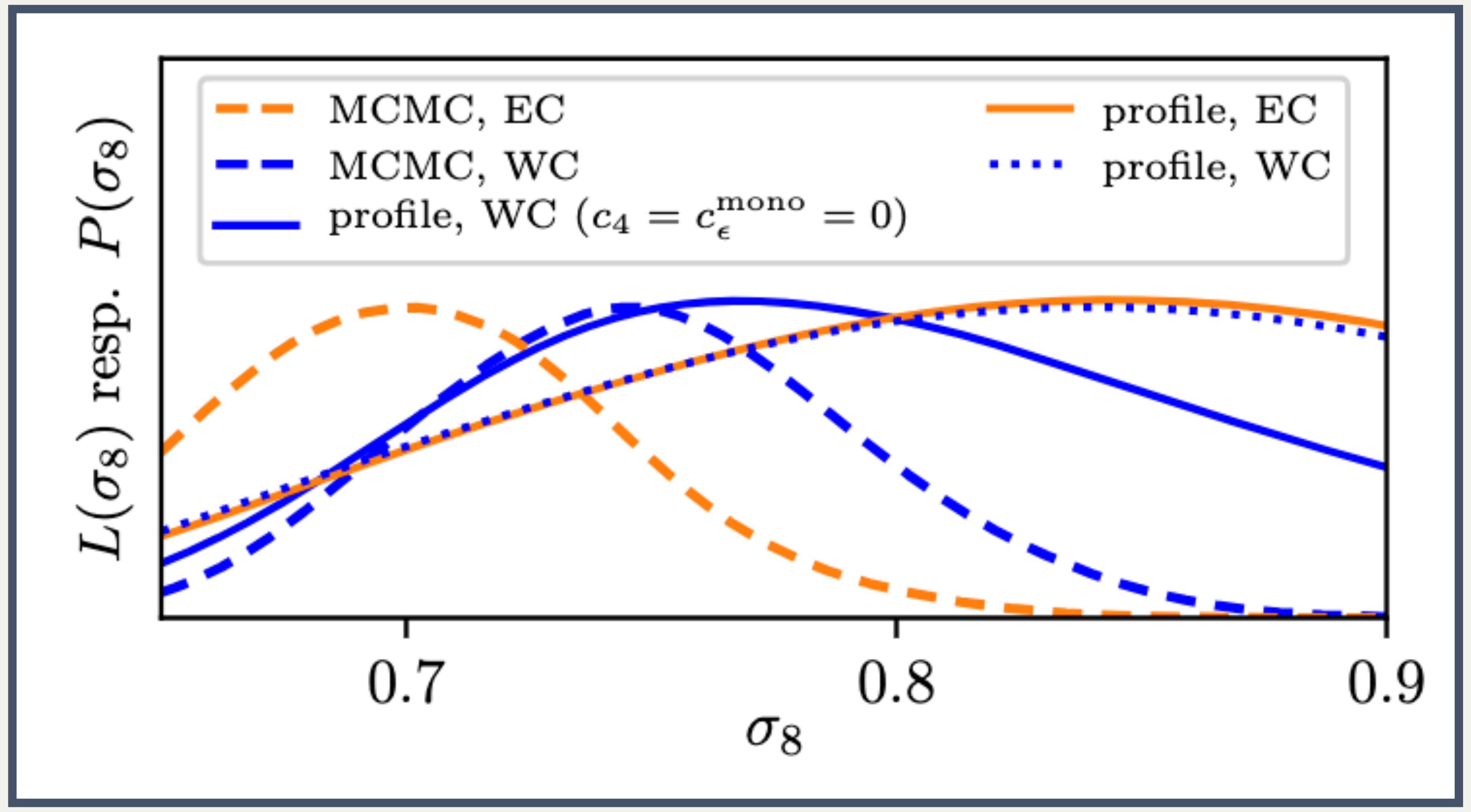
Bayes' theorem:
$P \propto \mathscr{L} \times p$

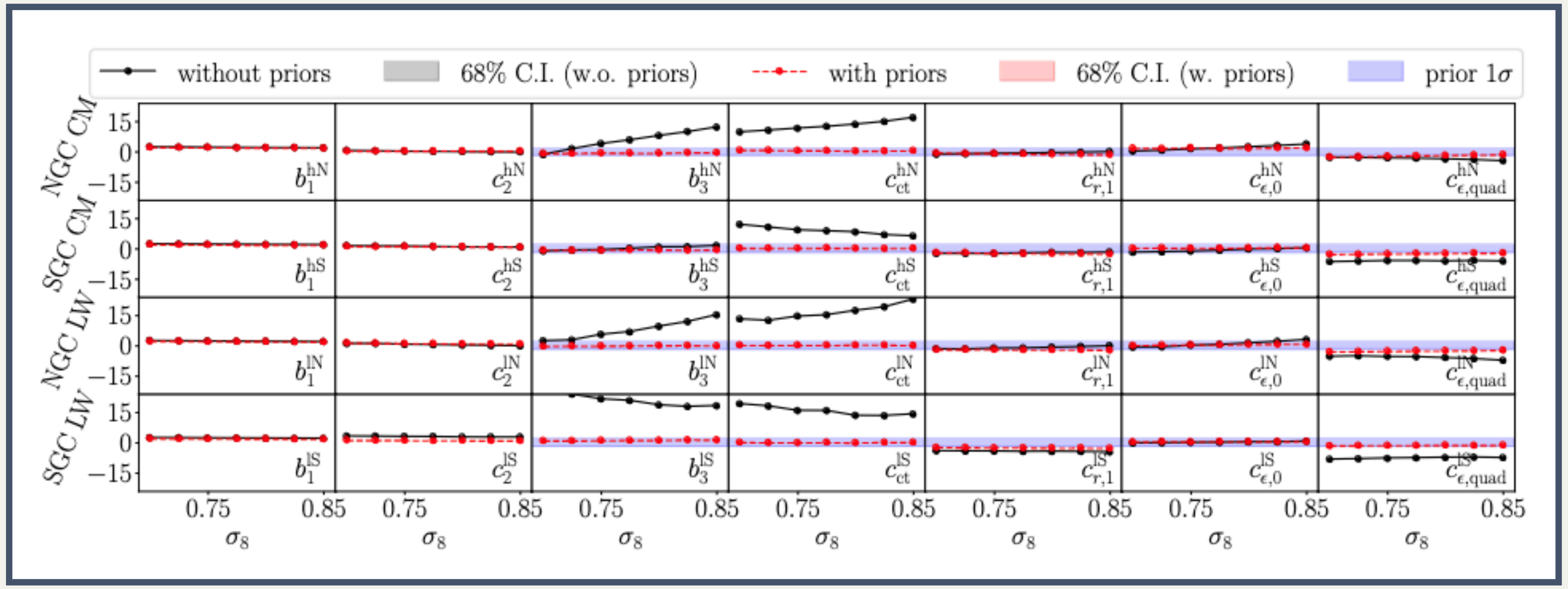
On the consistency of EFTofLSS

Profile likelihood

Advantage: frequentist analysis is **independent of priors** and therefore of projection effects

Disadvantage: the data prefers several EFT parameters to take on **extreme values**, possibly breaking the perturbative nature of the theory





Brinch, Herold, TS++ [arXiv:2309.04468]

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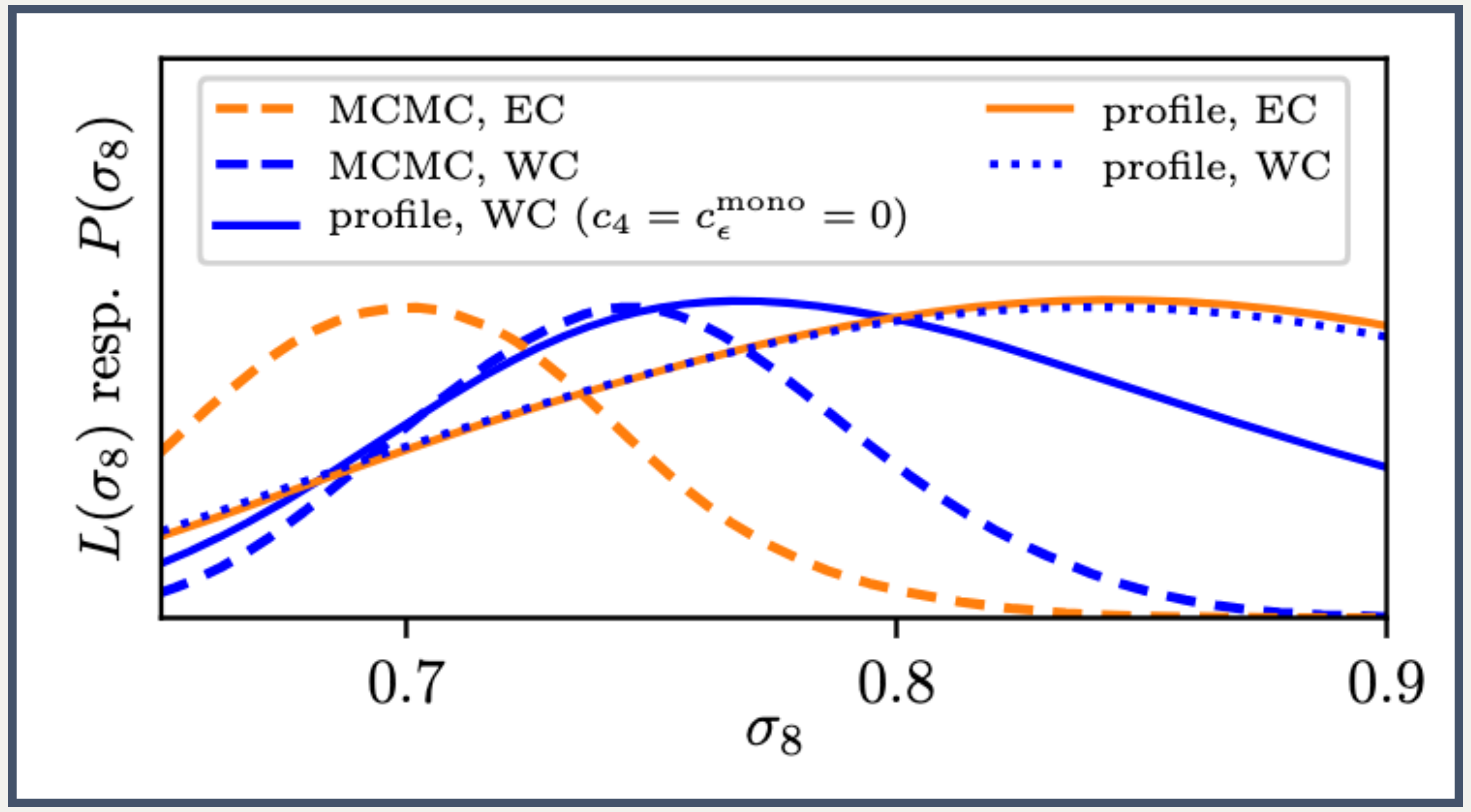
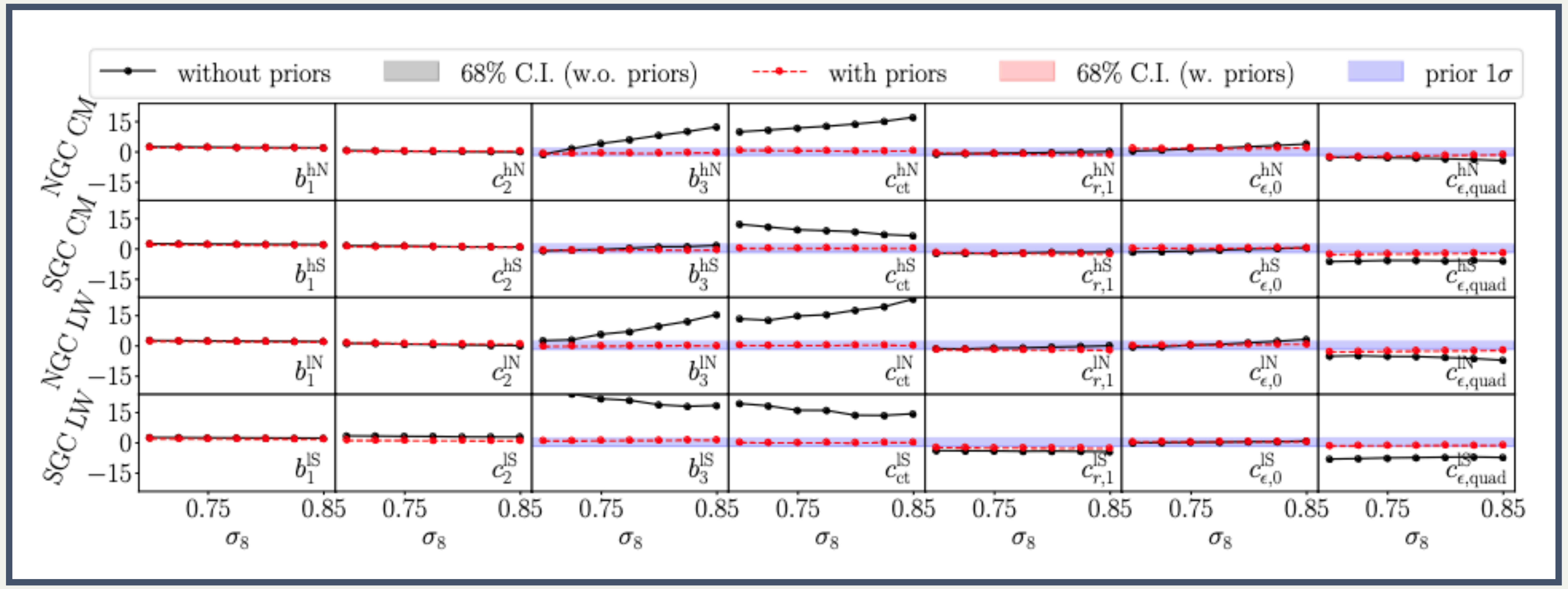


The low value of σ_8 is due to prior effects!

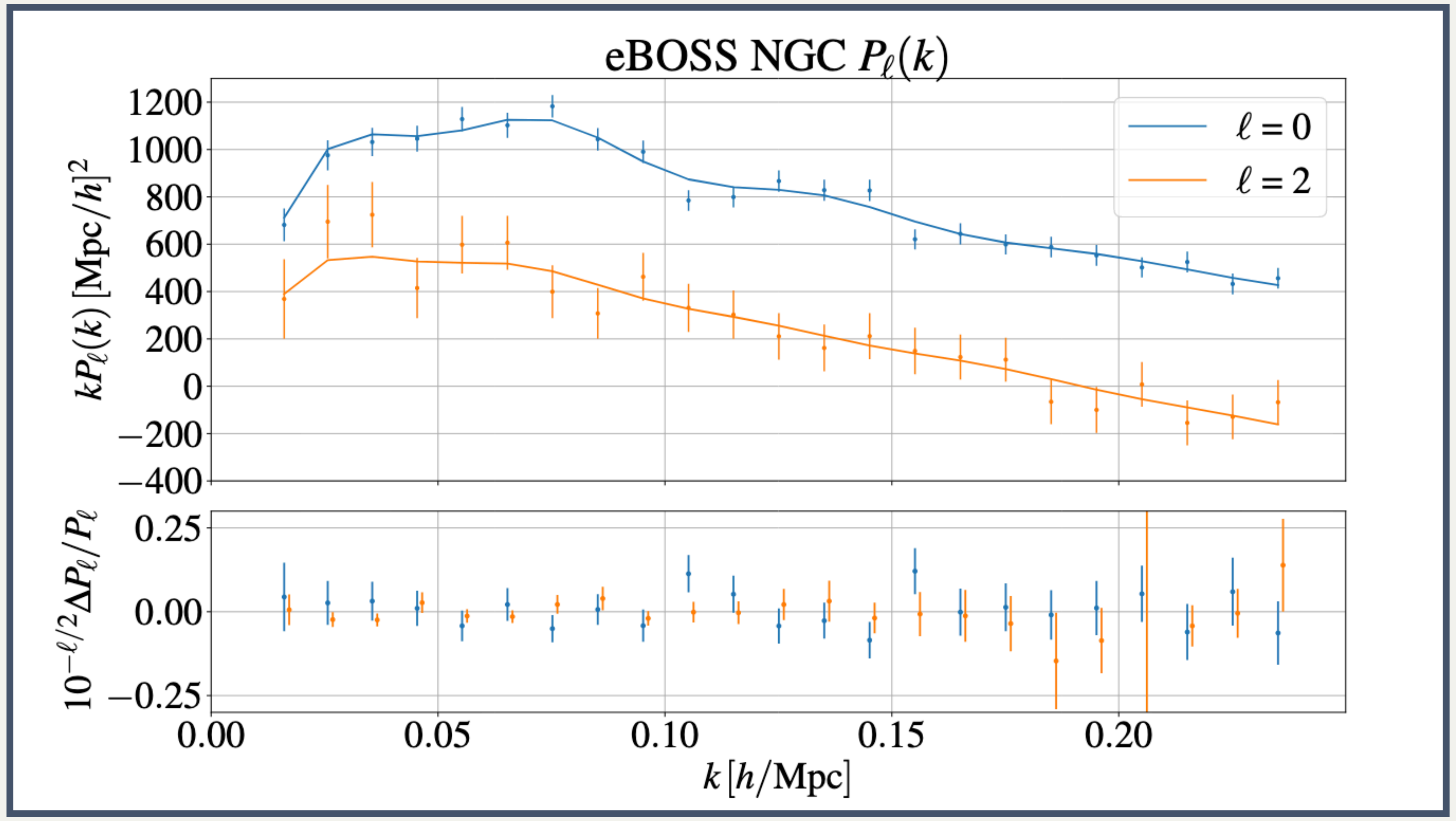


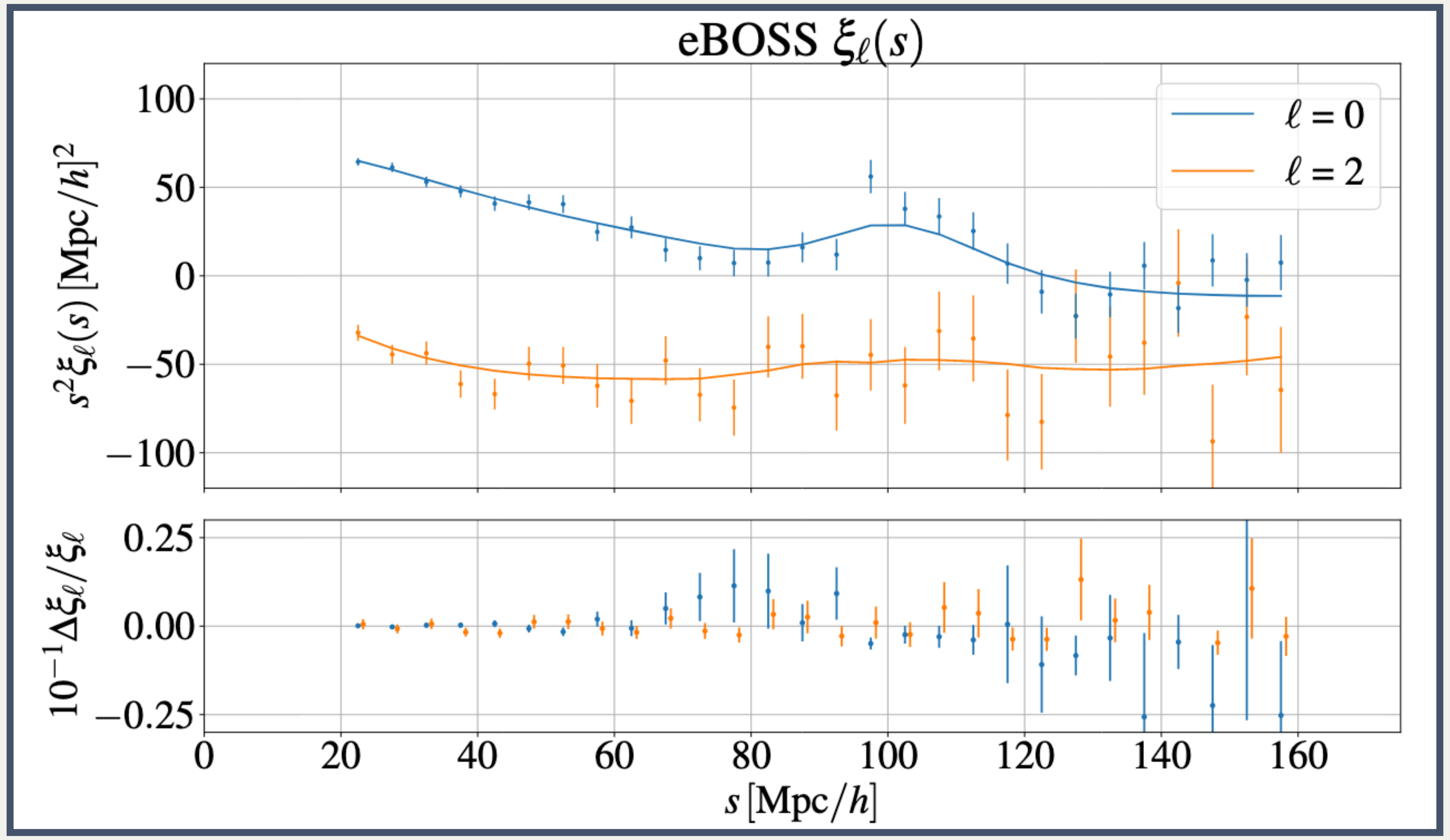
Brinch, Herold, TS++ [arXiv:2309.04468]

EFTofLSS applied to eBOSS QSO data

TS, P. Zhang and V. Poulin, JCAP **[arXiv:2210.14931]**

Cosmological inference from the EFTofLSS: the eBOSS QSO full-shape analysis





TS++ [arXiv:2210.14931]

See also Chudaykin++ [arXiv:2210.17044]

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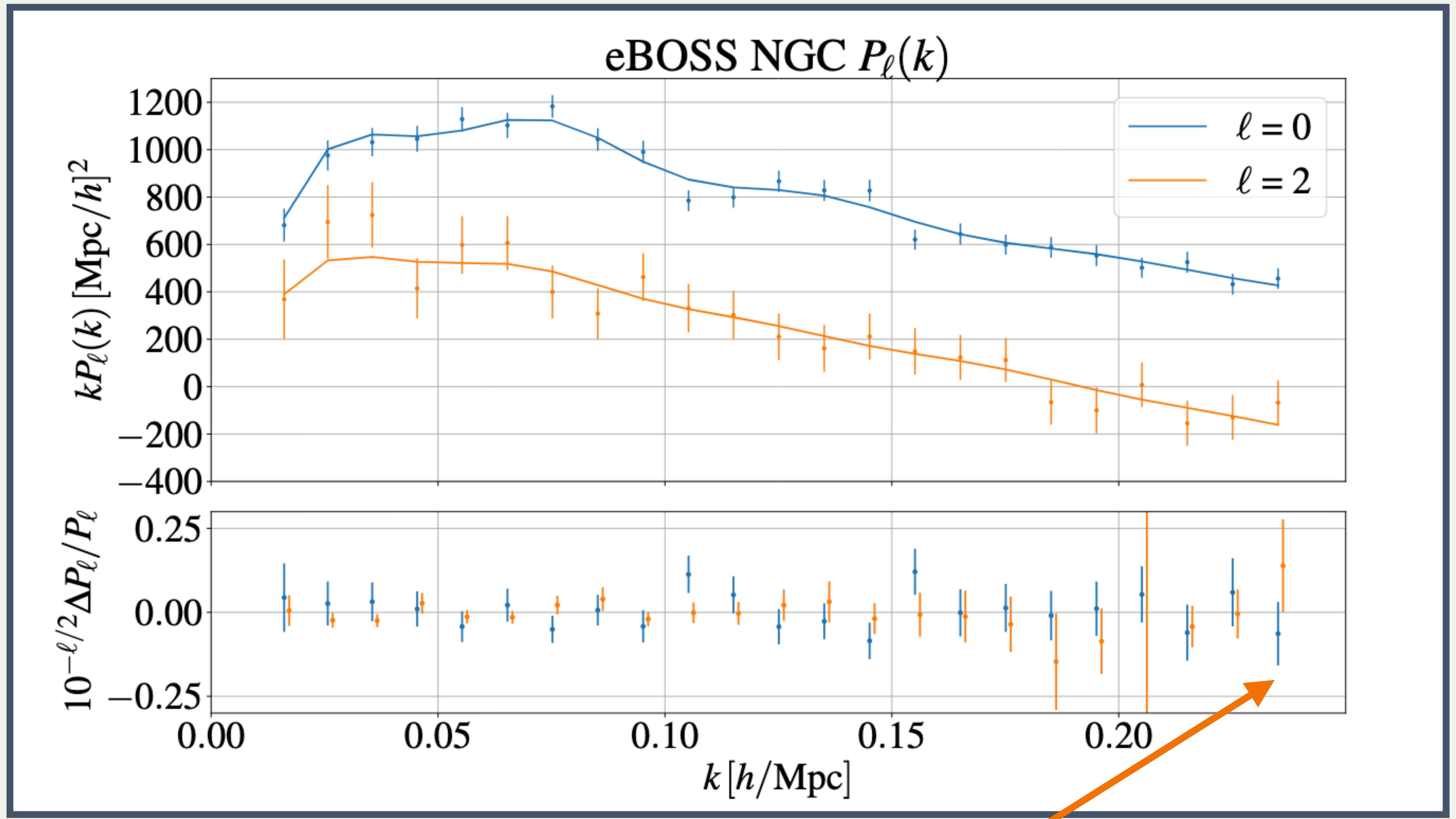

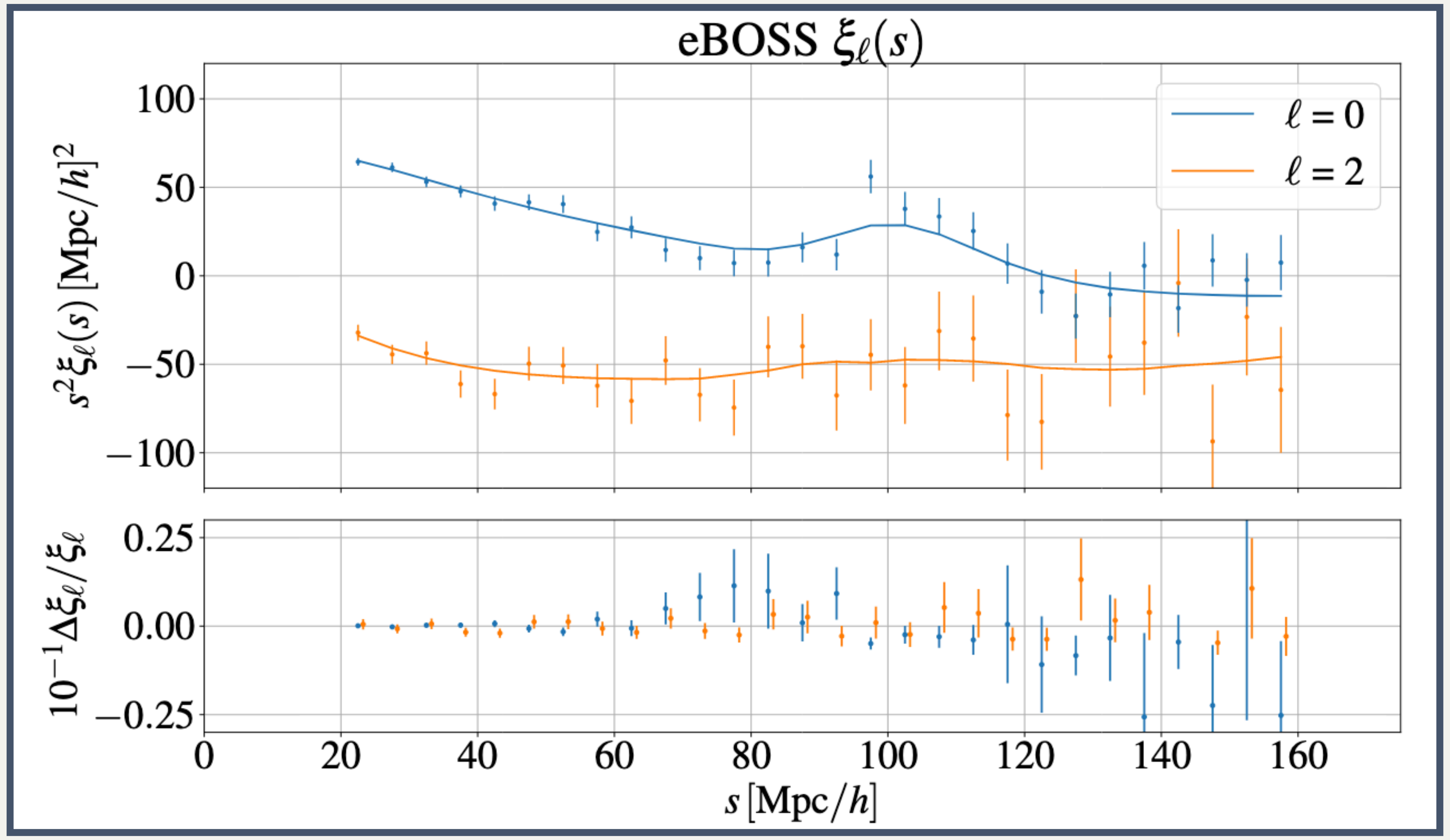

cut-off scale $k_{\mathrm{max}} = 0.24\,h\,.\,\mathrm{Mpc}^{-1}$

TS++ [arXiv:2210.14931]

See also Chudaykin++ [arXiv:2210.17044]

Determination of the cut-off scale $k_{\rm max}$ of the one-loop prediction

The next-to-next-to-leading order (NNLO) terms

At **one-loop order**, the galaxy power spectrum reads:

$$P_g(k,\mu) = Z_1(\mu)^2 P_{11}(k) + 2Z_1(\mu)P_{11}(k)\left(c_{\rm ct}\frac{k^2}{k_{\rm M}^2} + c_{r,1}\mu^2\frac{k^2}{k_{\rm M}^2} + c_{r,2}\mu^4\frac{k^2}{k_{\rm M}^2}\right)$$
$$+ 2\int \frac{d^3q}{(2\pi)^3} Z_2(\boldsymbol{q},\boldsymbol{k}-\boldsymbol{q},\mu)^2 P_{11}(|\boldsymbol{k}-\boldsymbol{q}|)P_{11}(q) + 6Z_1(\mu)P_{11}(k)\int \frac{d^3q}{(2\pi)^3} Z_3(\boldsymbol{q},-\boldsymbol{q},\boldsymbol{k},\mu)P_{11}(q)$$
$$+ \frac{1}{\bar{n}_g}\left(c_{\epsilon,0} + c_{\epsilon,1}\frac{k^2}{k_{\rm M}^2} + c_{\epsilon,2} f\mu^2\frac{k^2}{k_{\rm M}^2}\right),$$

One can add the **NNLO terms** (*i.e.*, the dominant two-loop terms):

$$P_{\rm NNLO}(k,\mu) = \frac{1}{4}b_1\left(c_{r,4}b_1 + c_{r,6}\mu^2\right)\mu^4\frac{k^4}{k_{\rm R}^4}P_{11}(k)$$

Zhang++ [arXiv:2110.07539]

If the contribution of $P_{\rm NNLO}(k,\mu)$ becomes **too large,** the one-loop prediction is **not accurate enough** → this determines the **cut-off scale** $k_{\rm max}$ of the prediction

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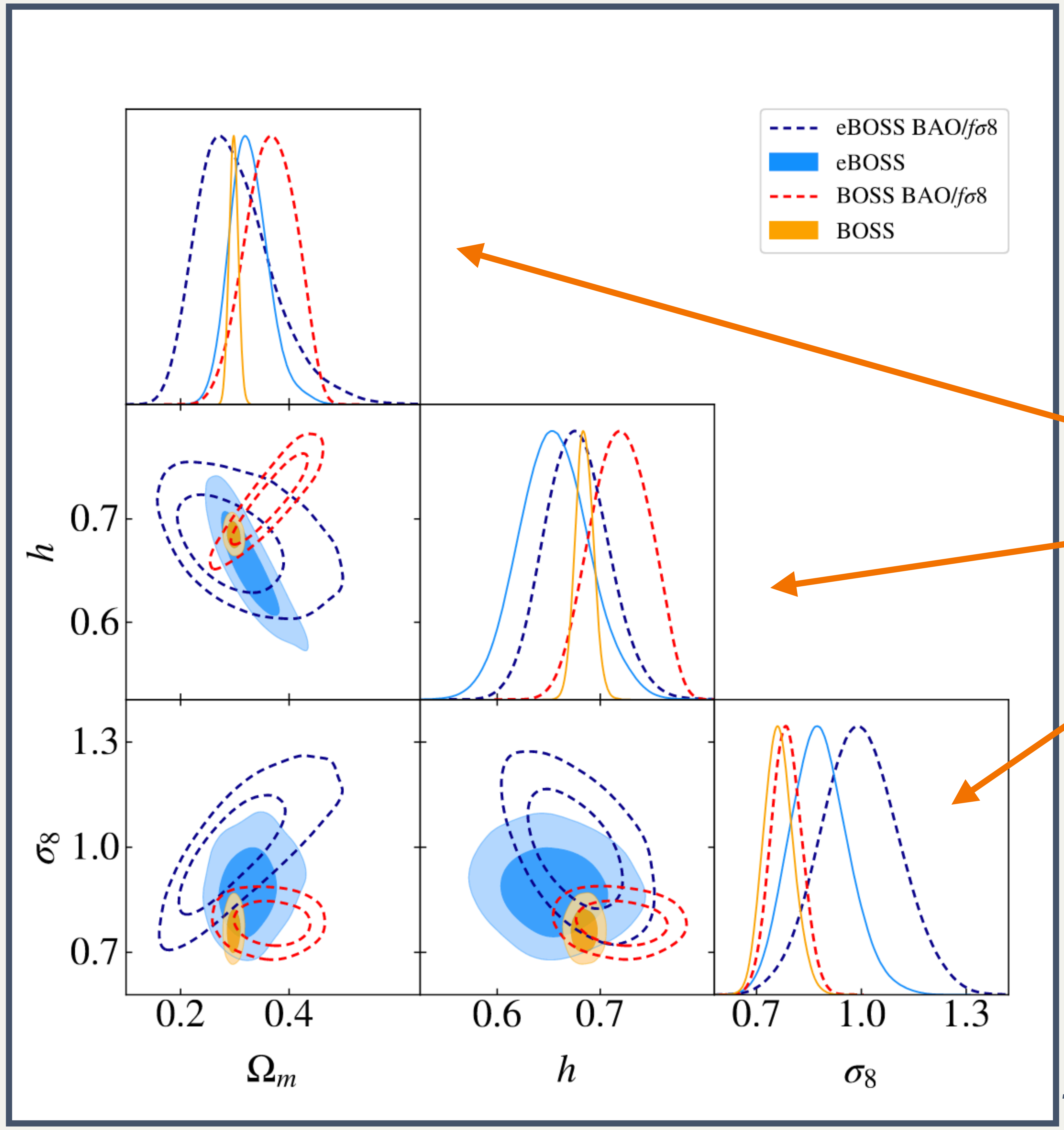
2 new EFT parameters

If the contribution of $P_{\rm NNLO}(k,\mu)$ becomes **too large,** the one-loop prediction is **not accurate enough** → this determines the **cut-off scale** $k_{\max}$ of the prediction

BAO/$f\sigma_8$ vs EFTofLSS

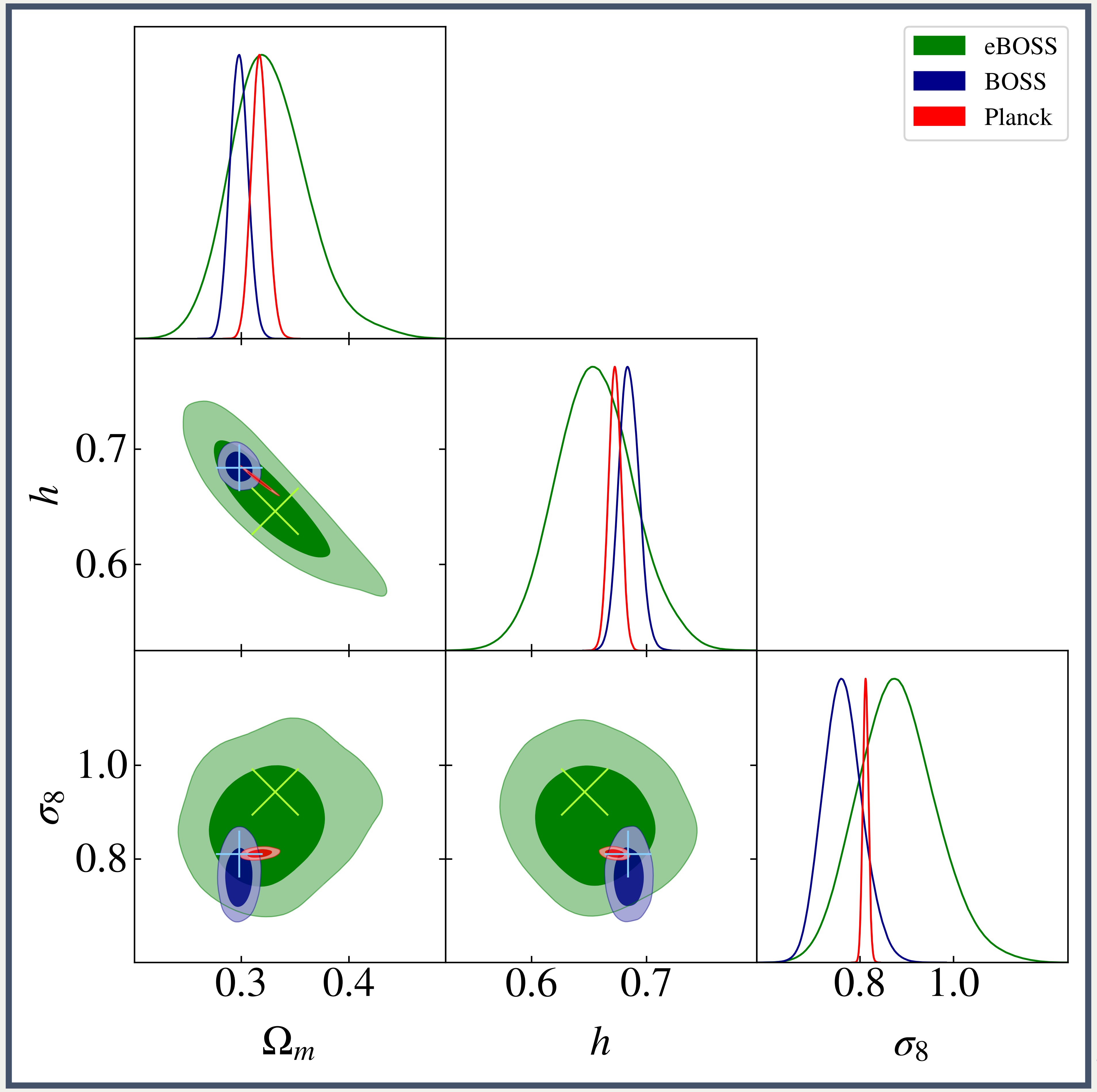


Additional information

TS++ [arXiv:2210.14931]

- For **eBOSS**, the error bars of Ω_m and σ_8 are reduced by a factor ~ 2.0 and ~ 1.3
- For **BOSS**, the error bars of Ω_m and h are reduced by a factor ~ 5.4 and ~ 3.2

LSS data vs Planck

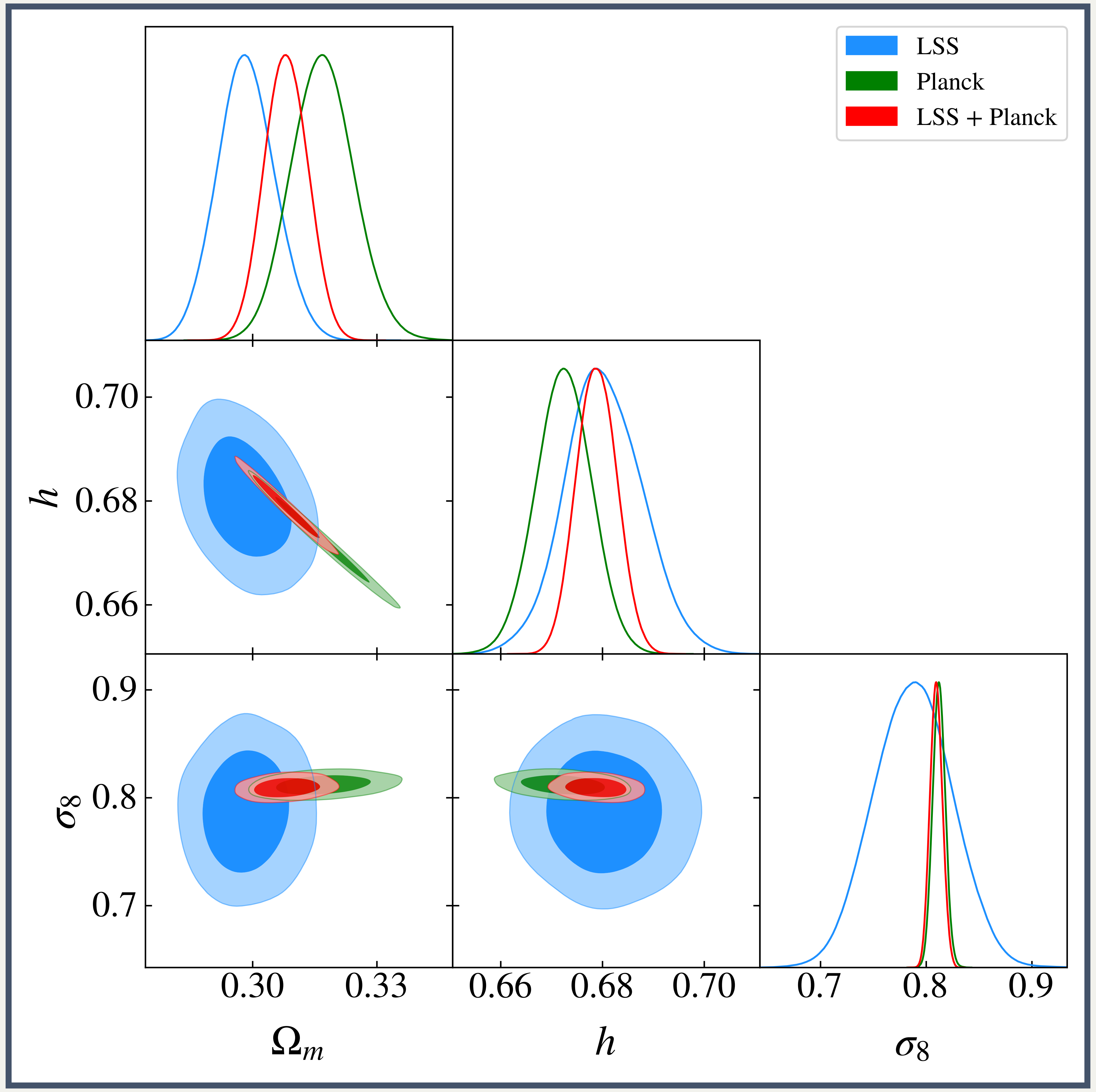


TS++ [arXiv:2210.14931]

- **eBOSS, BOSS and Planck are consistent** at $\lesssim 1.8\sigma$ on all cosmological parameters
- The h and σ_8 Planck values are **in-between** those of BOSS and eBOSS

→ **there is no tension between Planck and BOSS/eBOSS**

LSS data combined with Planck



LSS: eBOSS + BOSS + ext-BAO + Pantheon

(Uncalibrated Supernovae)

- The combination of eBOSS + BOSS allows to determine Ω_m and h at a **precision similar to Planck**
- **Compared to Planck alone**, the constraints on Ω_m and h are **improved by** $\sim 30\,\%$

TS++ [arXiv:2210.14931]

Extensions to ΛCDM: total neutrino mass $\sum m_\nu$

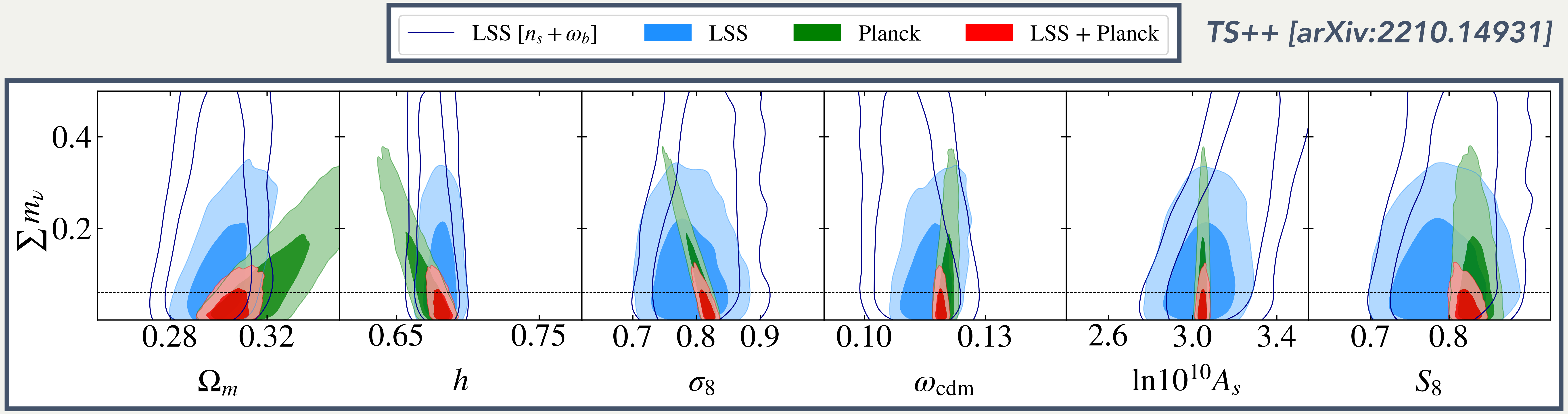
- The LSS constraint derived in this work is **only** $\sim 10\,\%$ **weaker than the Planck constraint** ($\sum m_\nu < 0.241 eV$)
- The EFT analysis **significantly improves the constraints** on $\sum m_\nu$ (by a factor of ~ 18) over the conventional BAO/$f\sigma_8$ analysis ($\sum m_\nu < 4.84 eV$)

Palanque-Delabrouille++ [arXiv:1911.09073]

- This analysis **disfavors the inverse hierarchy** at $\sim 2.2\sigma$ & is **competitive to the Lyman-α constraints**

LSS:
$\sum m_\nu < 0.274 eV$

LSS+Planck:
$\sum m_\nu < 0.093 eV$



TS++ [arXiv:2210.14931]

Conclusion

- The EFTofLSS is a novel method that provides an **accurate description of LSS data (up to mildly non linear scales) at a controlled precision**

- Constraints from LSS data are **competitive with CMB data** and their combination **improves over Planck alone**

- EFTofLSS allows to highlight that **there is no tension** between current BOSS/eBOSS data and Planck data (but not in tension with weak lensing neither)

- Data are consistent with ΛCDM at $\lesssim 1.3\sigma \rightarrow$ Strong constraints on canonical extensions to ΛCDM
  e.g. LSS+Planck: $\sum m_\nu < 0.093 e$V

- The same analysis can be applied to non-canonical extensions of ΛCDM
  → see *e.g.* **TS [arXiv:2310.16800] ; TS++ [arXiv:2310.16800] ; TS++ [arXiv:2203.07440] ; Schöneberg, Abellán, TS++ [arXiv:2306.12469] ; …**

Redshift uncertainties and EFTofLSS

« Highly ionized gas in the broadline region of quasars is subject to radiation-driven winds. **It is therefore likely that the measured redshifts** largely determined by these emission lines **are offset from the systemic redshift**. »

eBOSS Collaboration [arXiv:1508.04473]

Redshift uncertainties and EFTofLSS

Redshift Space Distorsion (RSD)

- **Observed redshift**: $1 + z_{\rm obs} = (1+z)(1+\delta z_{\rm pec})$
- **Comoving coordinate in redshift space**: $s(z) \simeq x + \frac{v \cdot \hat{n}}{\mathcal{H}}\hat{n}$
- **Relation, in Fourier space, between the overdensities in redshift space and real space**:

$$\delta_{g,r}(k) = \delta_g(k) + \int d^3x\, e^{-ik\cdot x}\left(e^{-ik\cdot \frac{v\cdot\hat{n}}{\mathcal{H}}\cdot\hat{n}} - 1\right)(1+\delta_g(x))$$

Kaiser '87

Redshift error

- **Observed redshift**: $1 + z_{\rm obs} = (1+z)(1+\delta z_{\rm pec} + \delta z_{\rm sys})$
- Taking this uncertainty into account is equivalent to carrying out the **transformation**:

$$v \cdot \hat{n} \rightarrow v \cdot \hat{n} + v_{\rm sys}$$

Redshift uncertainties and EFTofLSS

$$P^{\rm sys}_{g,r}(k,\mu) = P_{g,r}(k,\mu) - 2i\mu k \frac{\bar{v}_{\rm sys}}{\mathcal{H}}\sigma_0^2 - 2i\mu k \frac{\bar{v}_{\rm sys}}{\mathcal{H}}(b_1 + f\mu^2)^2 P_{11}(k) - \mu^2 k^2 \frac{\sigma^2_{v,{\rm sys}}}{\mathcal{H}^2}(\delta_D(k) + 3\sigma_0^2) - 2\mu^2 k^2 \frac{\sigma^2_{v,{\rm sys}}}{\mathcal{H}^2}(b_1 + f\mu^2)^2 P_{11}(k) + \dots$$

TS++ [arXiv:2210.14931]

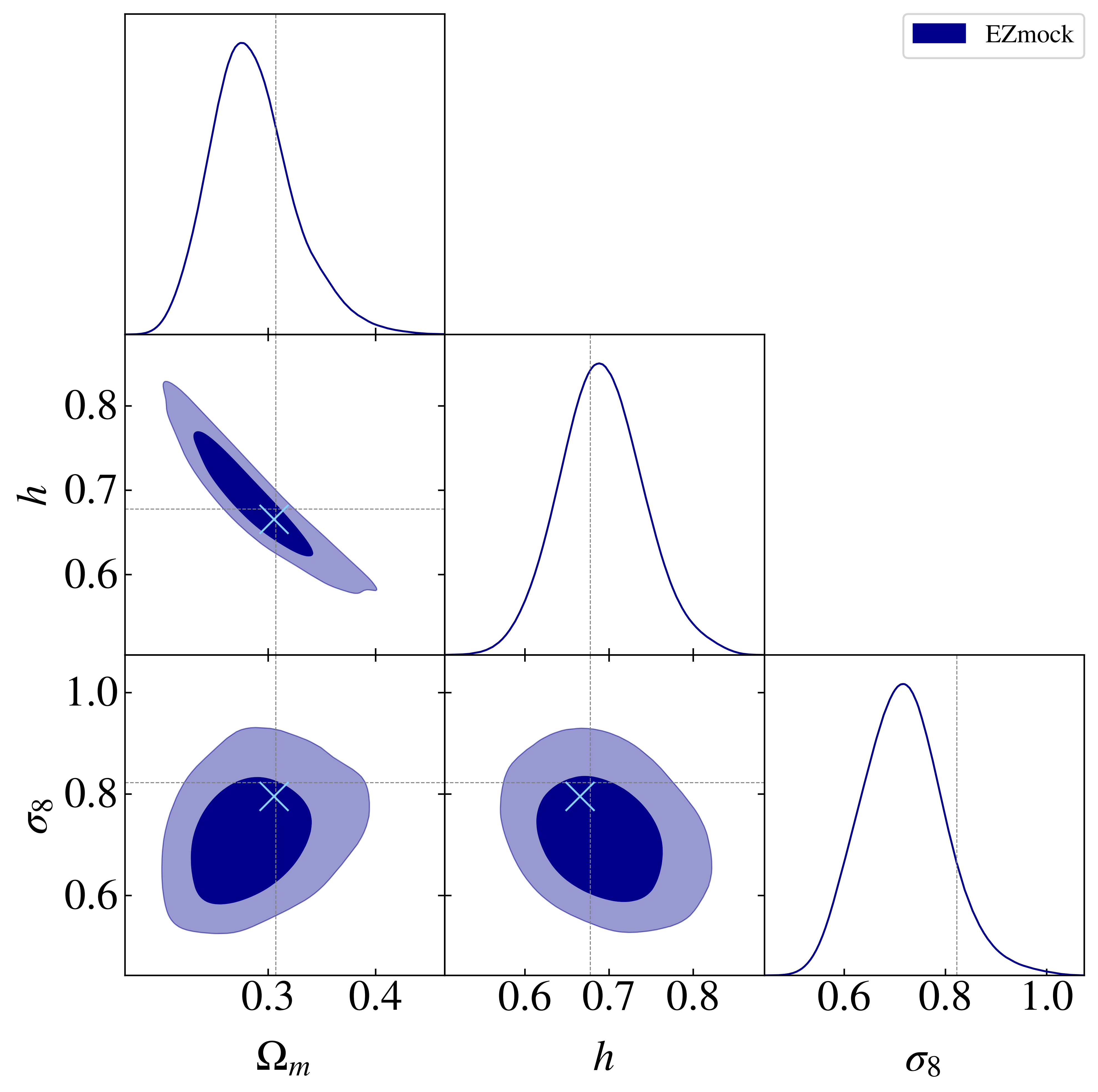
With $< v_{\rm sys} > \sim 0$ and $\sigma_{\rm v,sys} \sim 300$ km/s *eBOSS Collaboration [arXiv:1801.03062]*

- (green) **Purely imaginary**, and thus do not appear in the even multipoles. Significant only if the determination of the redshifts is biased on average $< v_{\rm sys} > \neq 0$.

- (red) The leading corrections to uncertainties in the redshift determination are **degenerate with EFT counterterms** going like $\sim \mu^2 k^2$ or $\sim \mu^2 k^2 P_{11}(k)$.

Determination of the cut-off scale $k_{\max}$ of the one-loop prediction

The EZmock



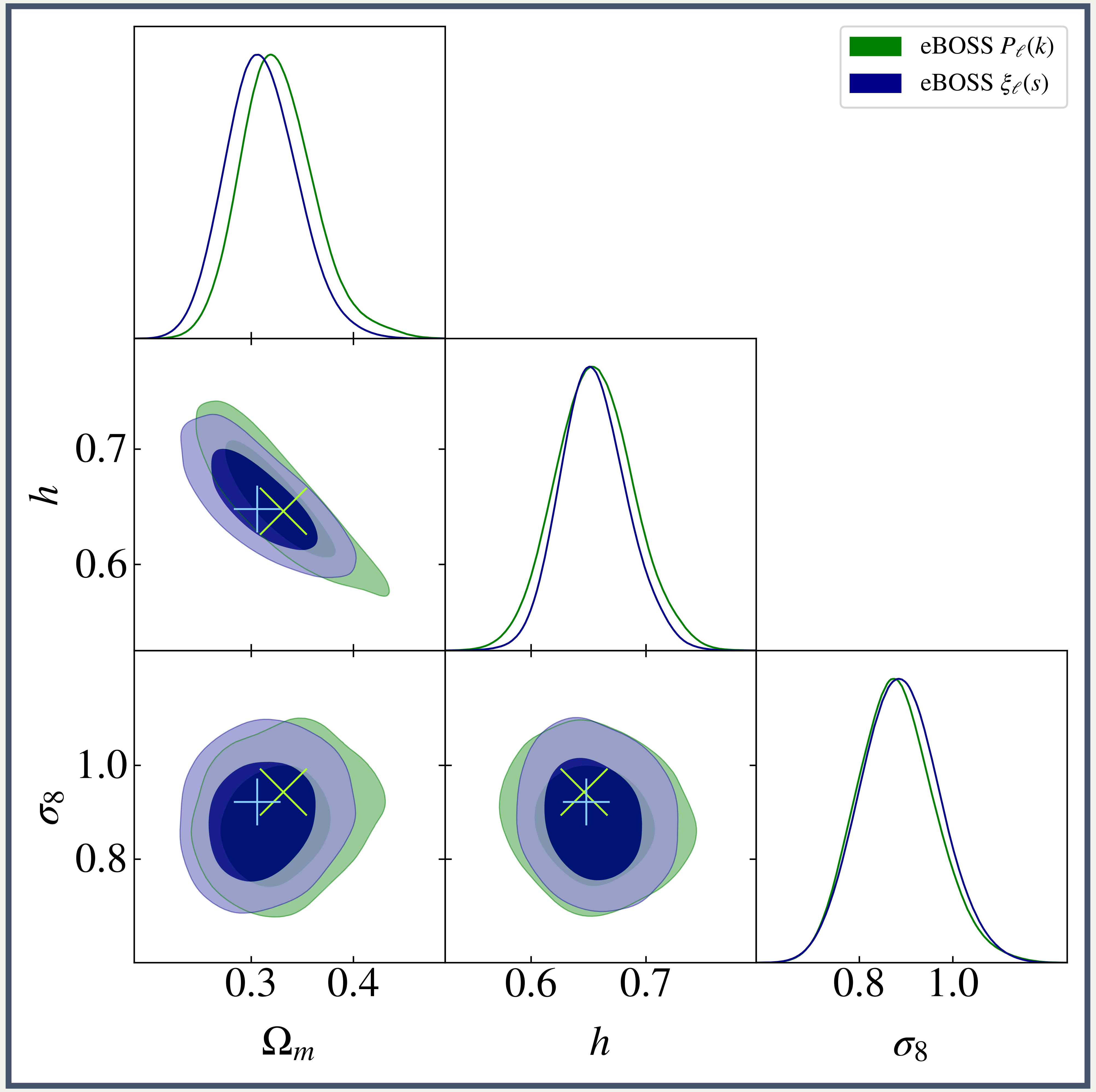
TS++ [arXiv:2210.14931]

- **EZmock**: mocks that are built to simulate eBOSS observational characteristics

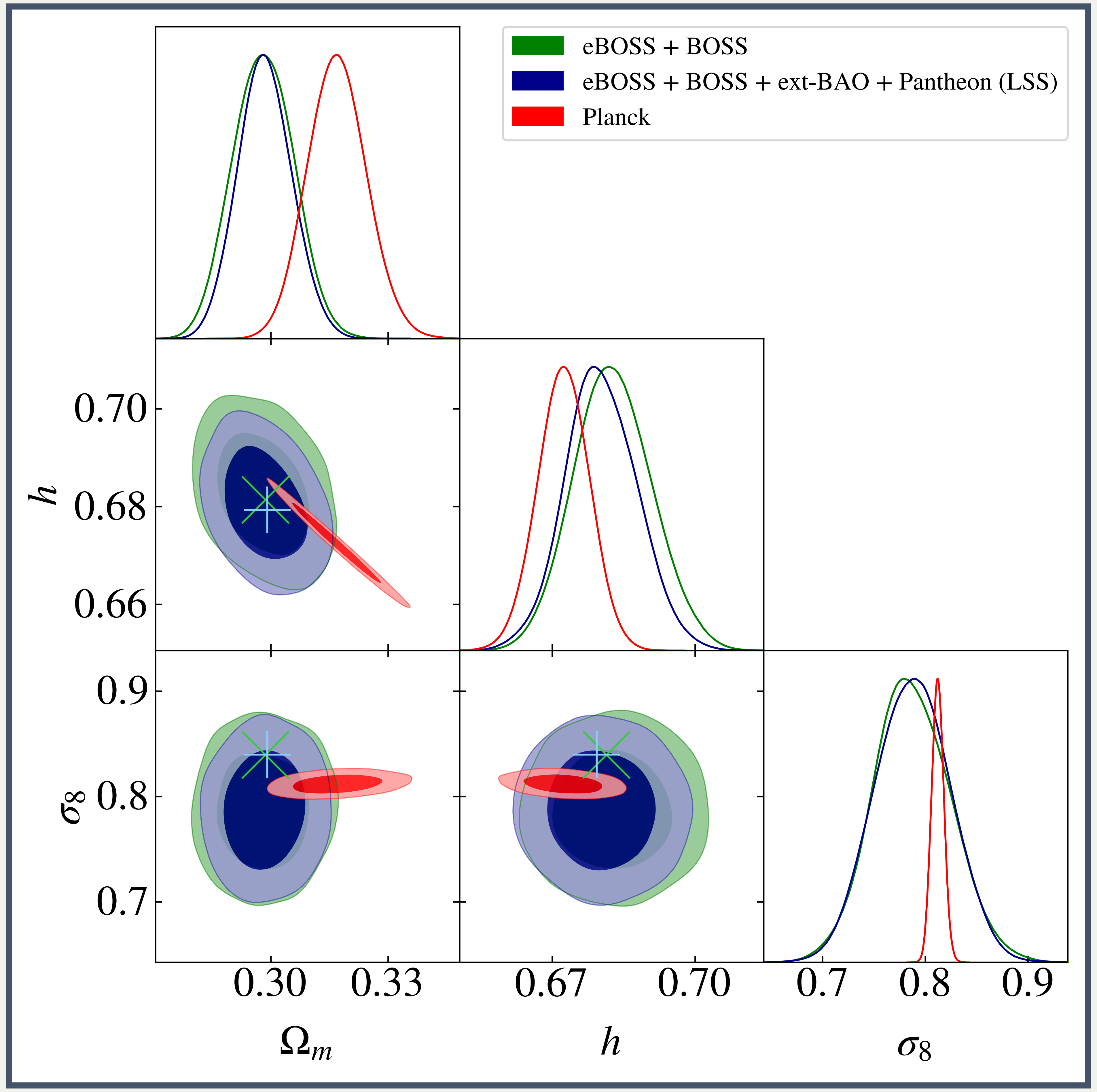
  Chuang++ [arXiv:1409.1124]

- Up to $k_{\max} = 0.24h\,\mathrm{Mpc}^{-1}$, the best-fit values of the cosmological parameters are shifted with respect to the truth of the simulations by $\lesssim 1/3 \cdot \sigma$

eBOSS $P_\ell(k)$ vs eBOSS $\xi_\ell(k)$



LSS data vs Planck



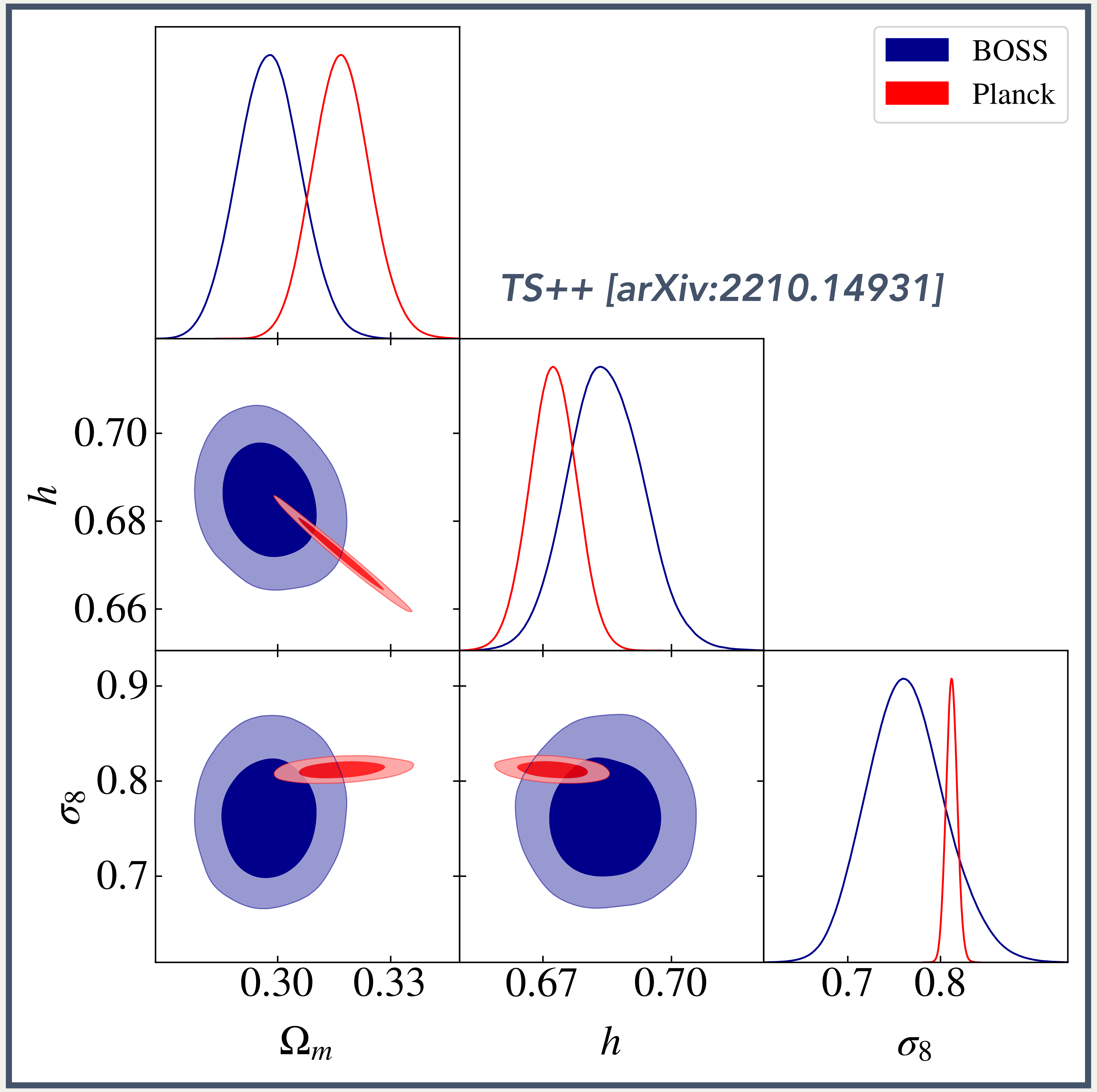
ext-BAO: 6dF & MGS (SDSS) data

TS++ [arXiv:2210.14931]

- The combination of eBOSS + BOSS allows to determine Ω_m and h at a **precision similar to Planck**
- The combination of LSS data remains consistent with Planck → **we can combine them!**

Applying EFTofLSS to BOSS data

Cosmological constraints



The EFTofLSS analysis of BOSS data allows to determine Ω_m and h at a **precision only** $10\,\%$ **and** $60\,\%$ worse than Planck

This is ~ 5.4 (for Ω_m) and ~ 3.2 (for h) times better than the BAO/$f\sigma_8$ analysis

See also D'Amico++ [arXiv:1909.05271] ; Philcox++ [arXiv:2002.04035]

On the consistency of EFTofLSS

Presentation of the problem

There are **several codes** in the literature with **different parametrizations**:

→ ***CLASS PT*** with the **EC** parametrization *Chudaykin++ [arXiv:2004.10607]*

→ ***PyBird*** with the **WC** parametrization *D'Amico++ [arXiv:2003.07956]*

(+ **Velocileptors** + **CLASS-OneLoop**)

Chen++ [arXiv:2005.00523] ; Linde++ [arXiv:2402.09778]

→ these codes use **different sets of priors** on EFT parameters

D'Amico++ [arXiv:1909.05271]
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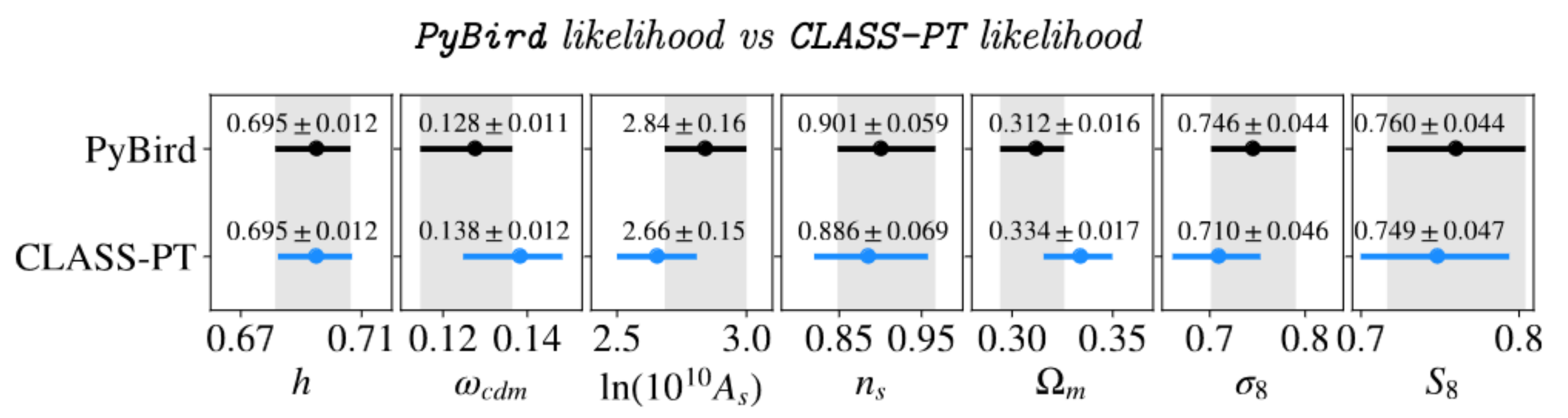

TS++ [arXiv:2208.05929]

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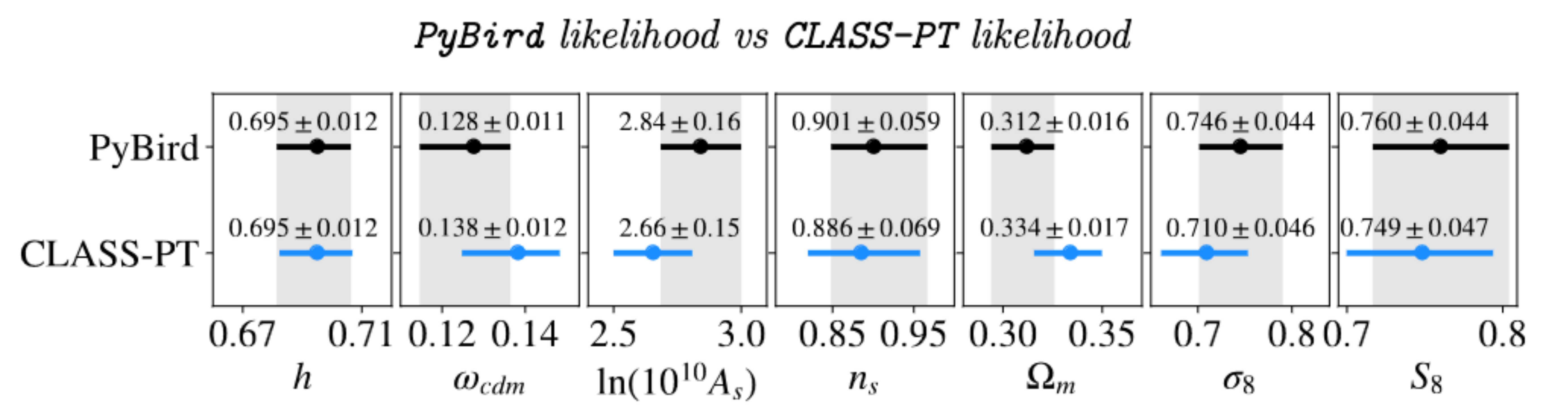
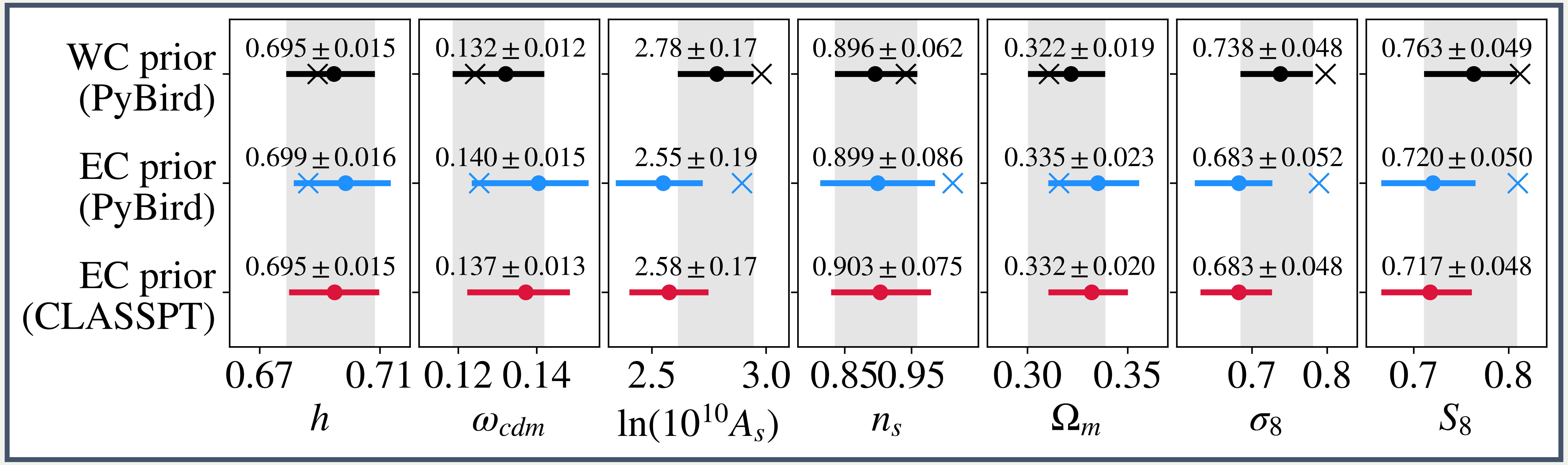


TS++ [arXiv:2208.05929]

Data, theoretical **parametrizations** and **codes** are supposed to be **equivalent**: **what is going on?**

On the consistency of EFTofLSS

The EFT prior issue



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On the consistency of EFTofLSS

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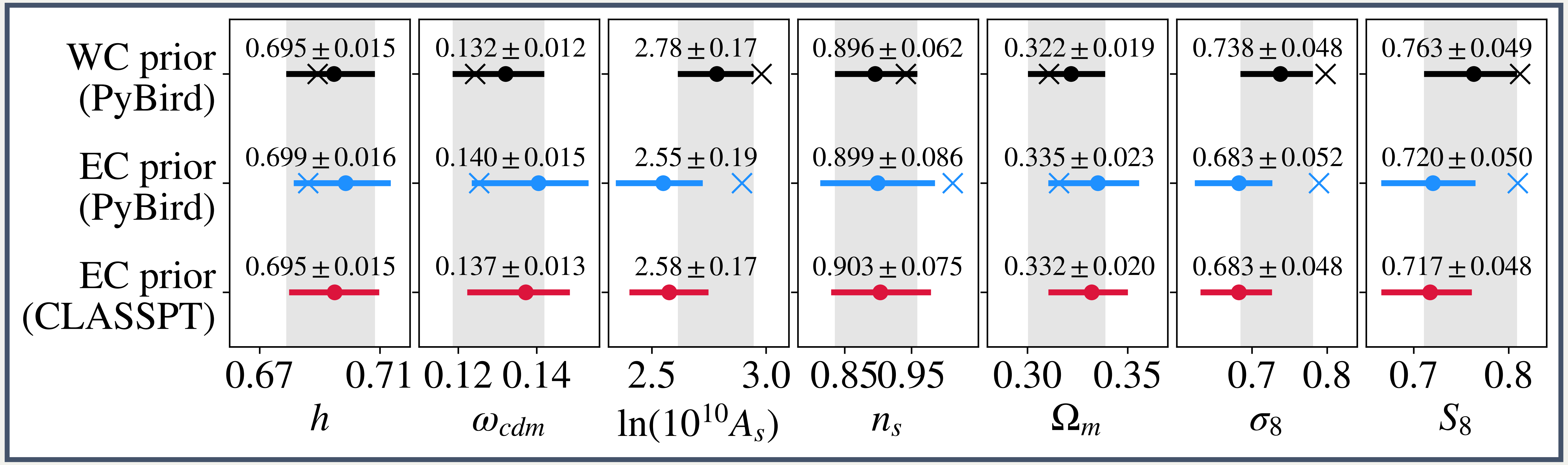


TS++ [arXiv:2208.05929]

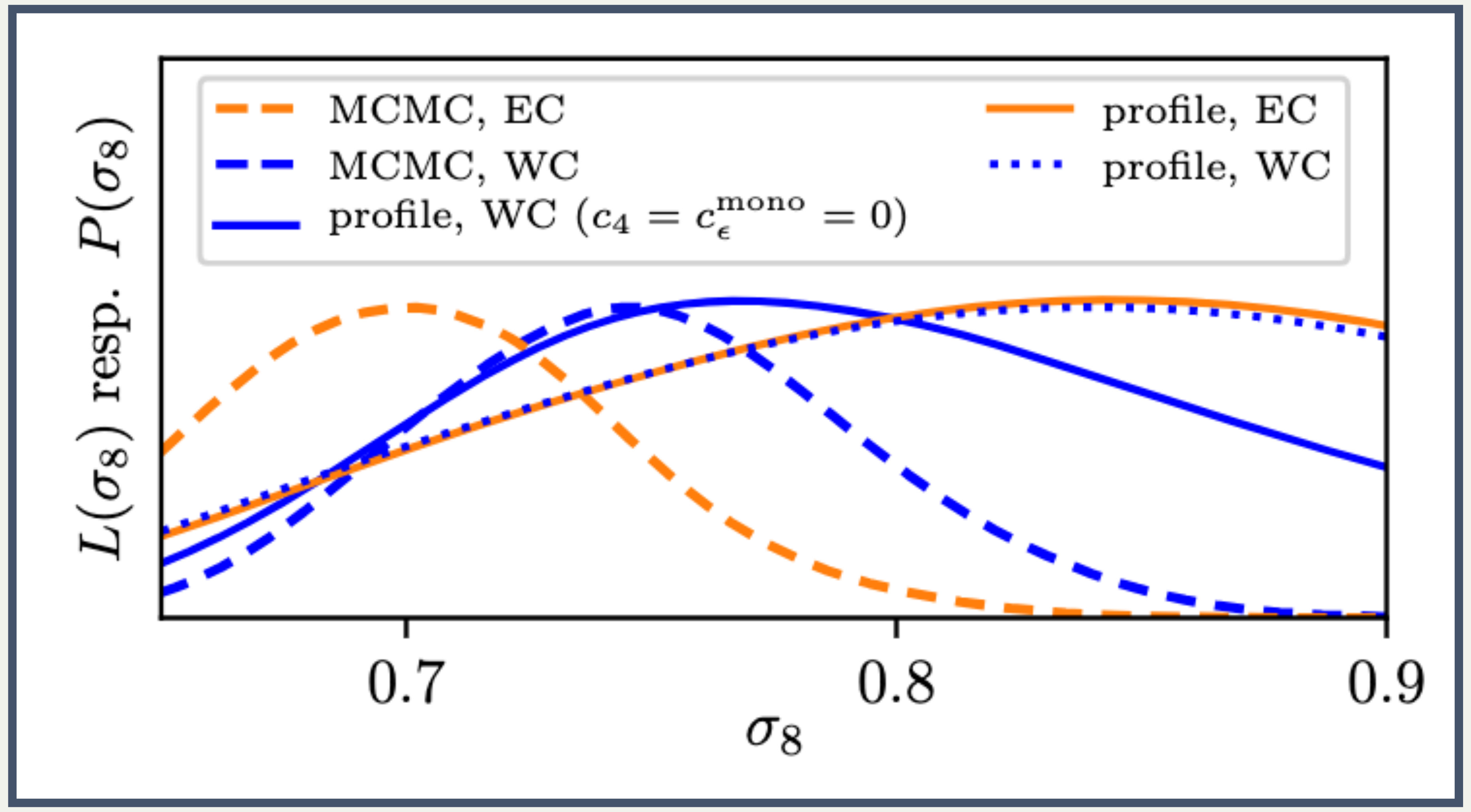
Prior effects

→ ***The prior weight effect***: if the region allowed by the prior is far from the true value of a parameter, then the posterior and bestfit will be shifted from the true value

→ ***The prior volume effect***: a posterior depends on the volume enclosed by the priors $\Longrightarrow$ large parameter regions are emphasized compared to smaller regions

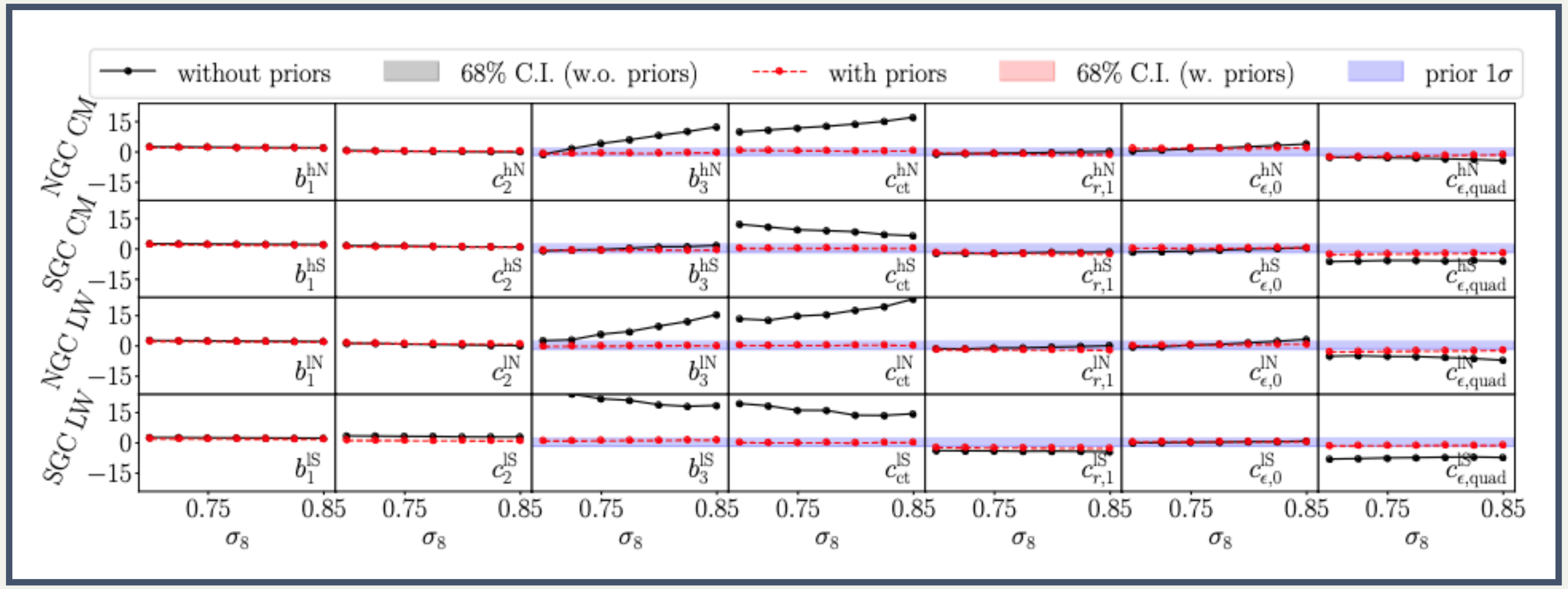
Bayes' theorem:
$P \propto \mathscr{L} \times p$

Profile likelihood



Advantage: frequentist analysis is **independent of priors** and therefore of projection effects

Disadvantage: the data prefers several EFT parameters to take on **extreme values**, possibly breaking the perturbative nature of the theory



Brinch, Herold, TS++ [arXiv:2309.04468]

On the consistency of EFTofLSS

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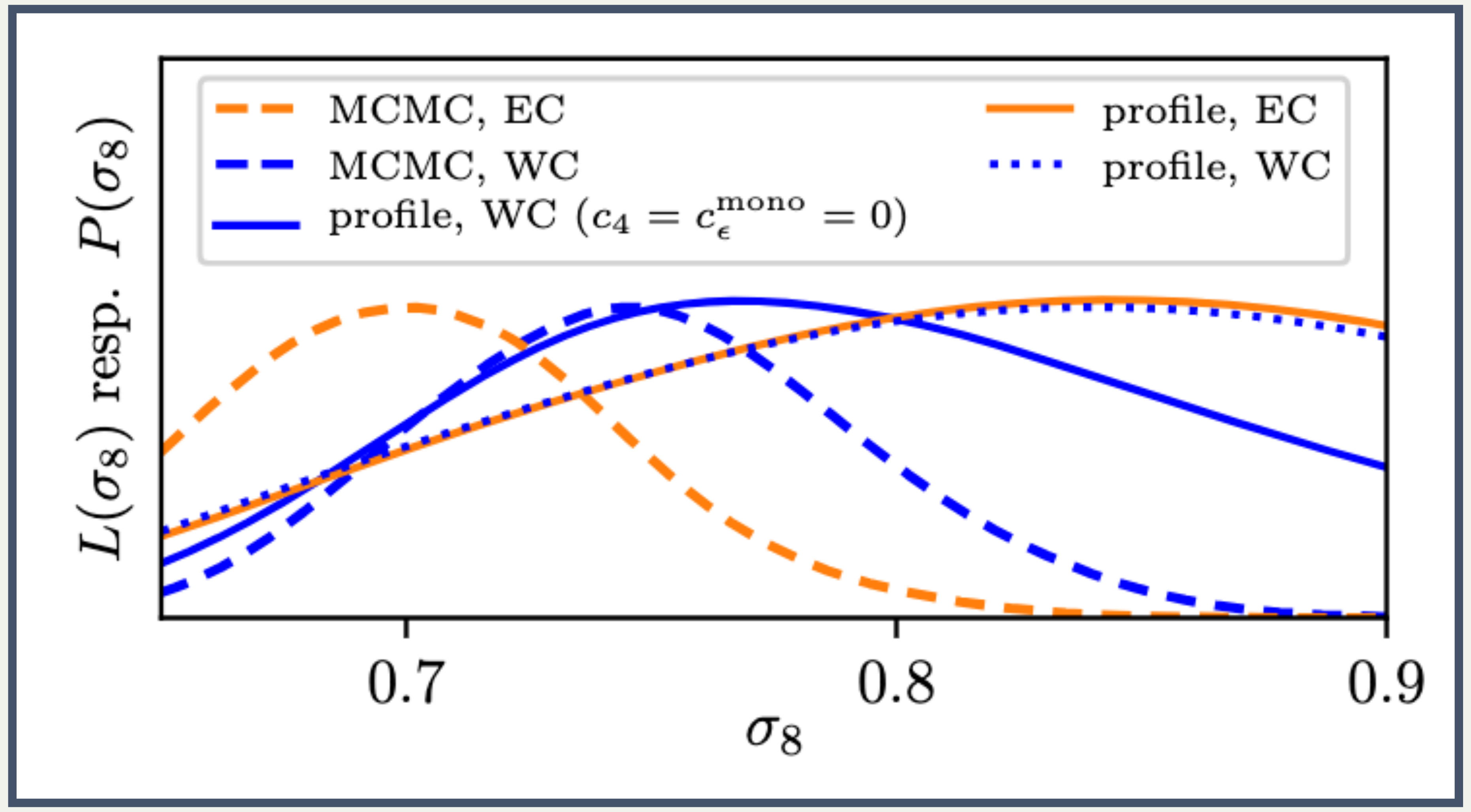

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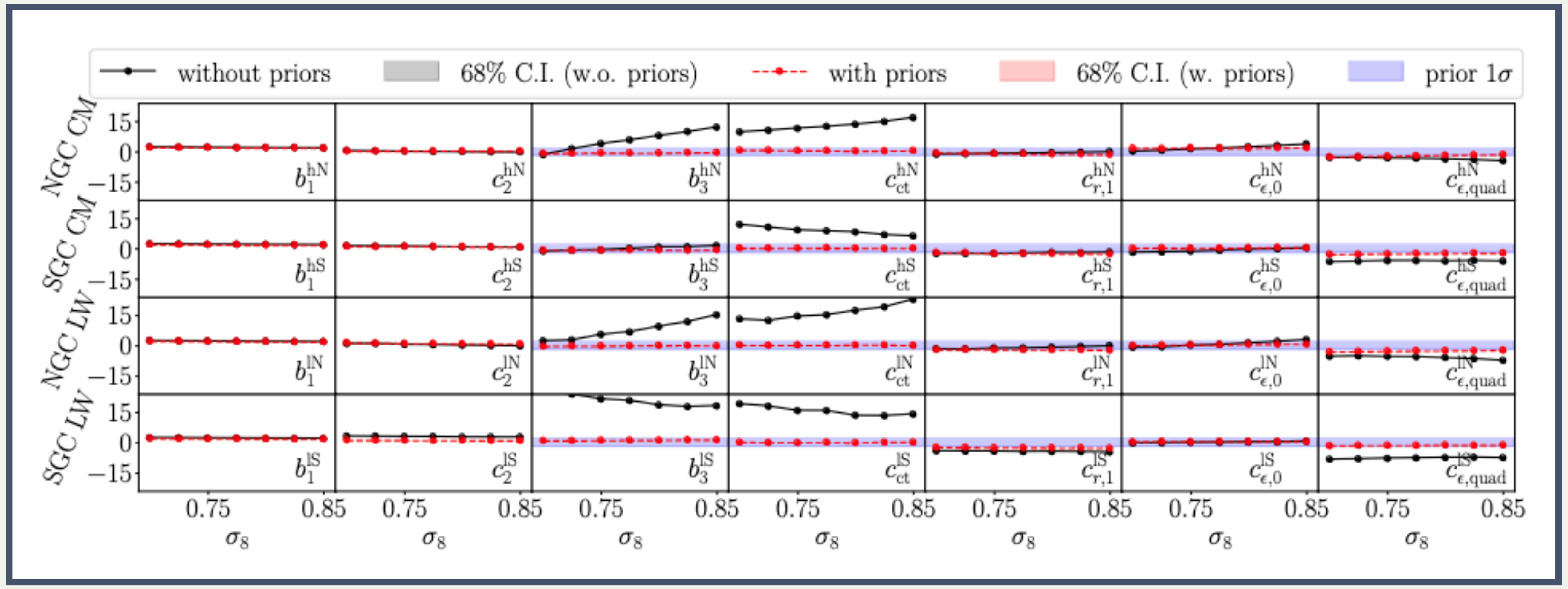
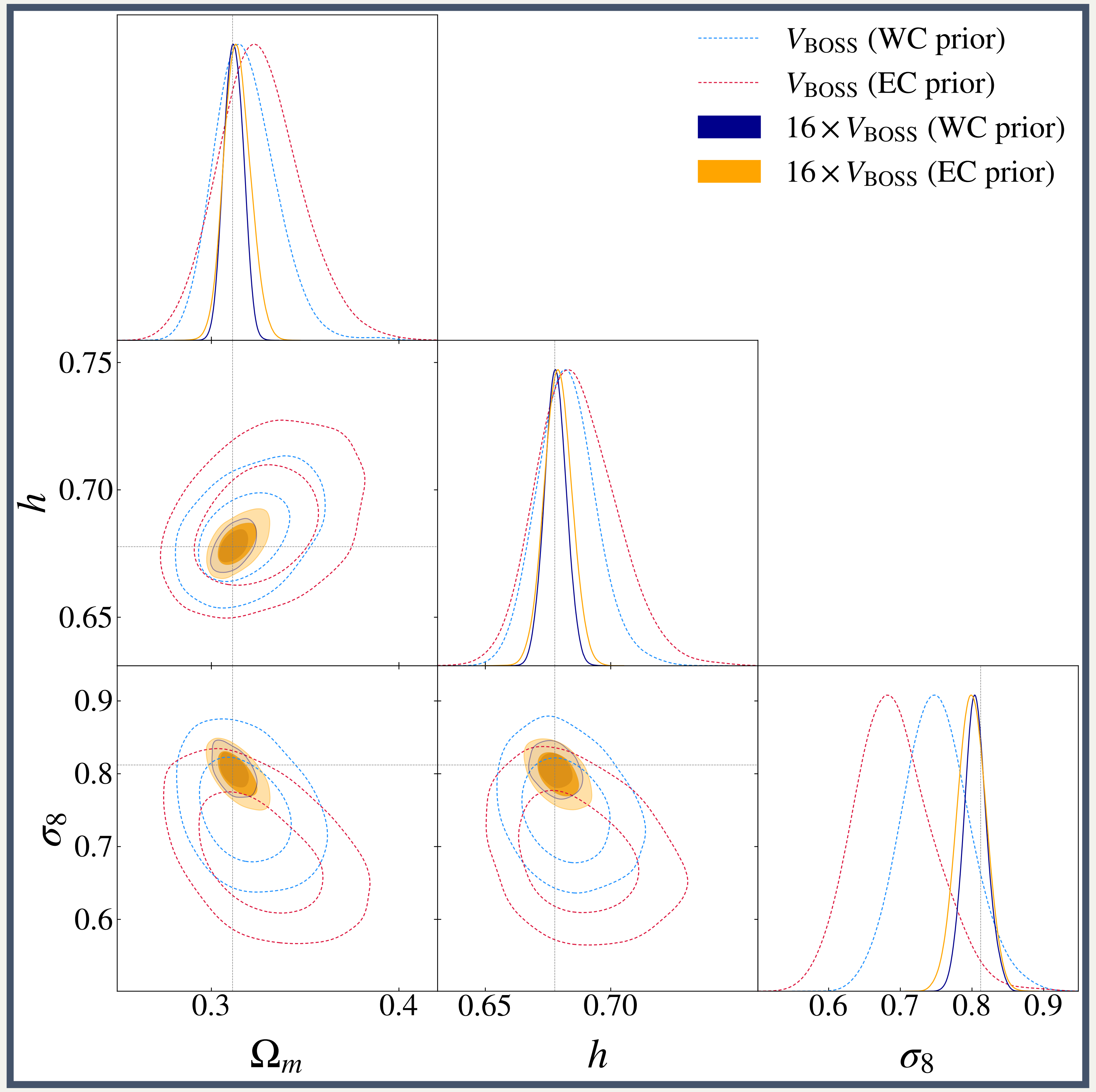


Brinch, Herold, TS++ [arXiv:2309.04468]

On the consistency of EFTofLSS

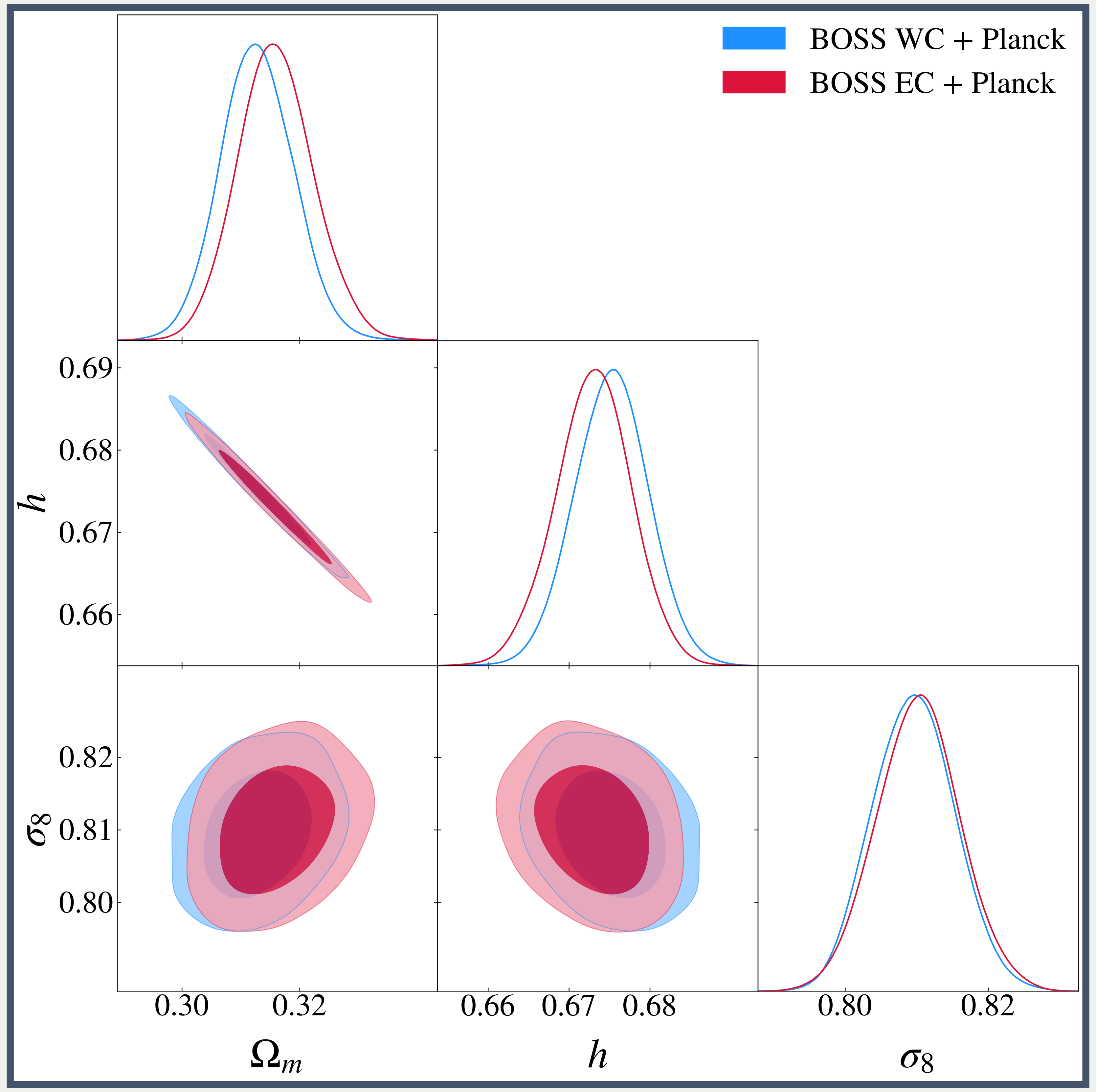
How to overcome this problem?



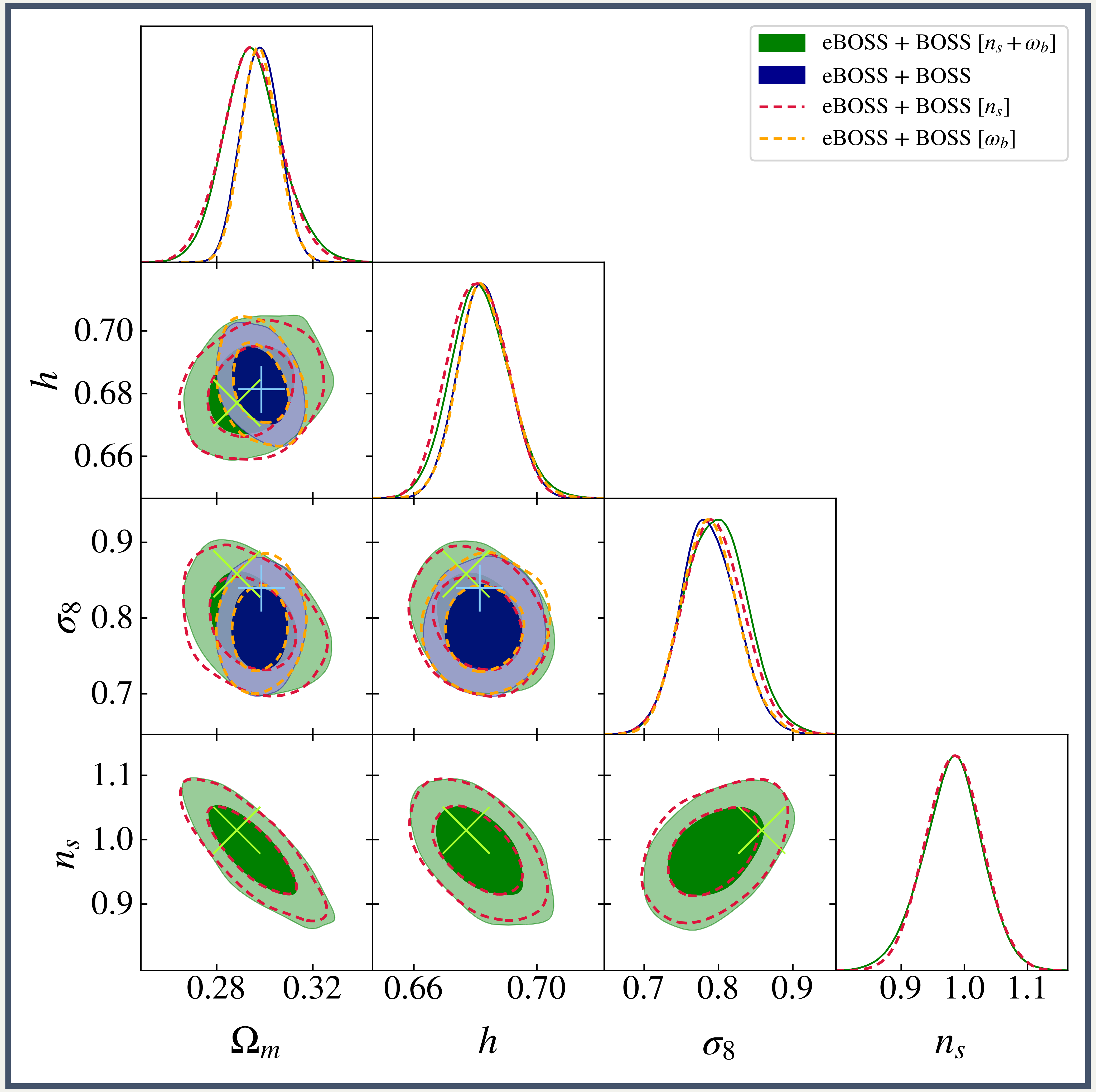
We find good consistency for:

- a larger volume of data (future experiments like DESI or EUCLID)
- a combination with Planck data

TS++ [arXiv:2208.05929]



Variation of n_s and ω_b



- We impose a uninformative large flat prior on n_s, while we impose a BBN Gaussian prior on ω_b
- The variation of ω_b within the BBN prior has a negligible impact on the cosmological results: we have a relative shift of $\lesssim 0.04\sigma$
- The variation of n_s within a uninformative large flat prior leads to a relative shift $\lesssim 0.4\sigma$

Extensions to ΛCDM: curvature density fraction Ω_k

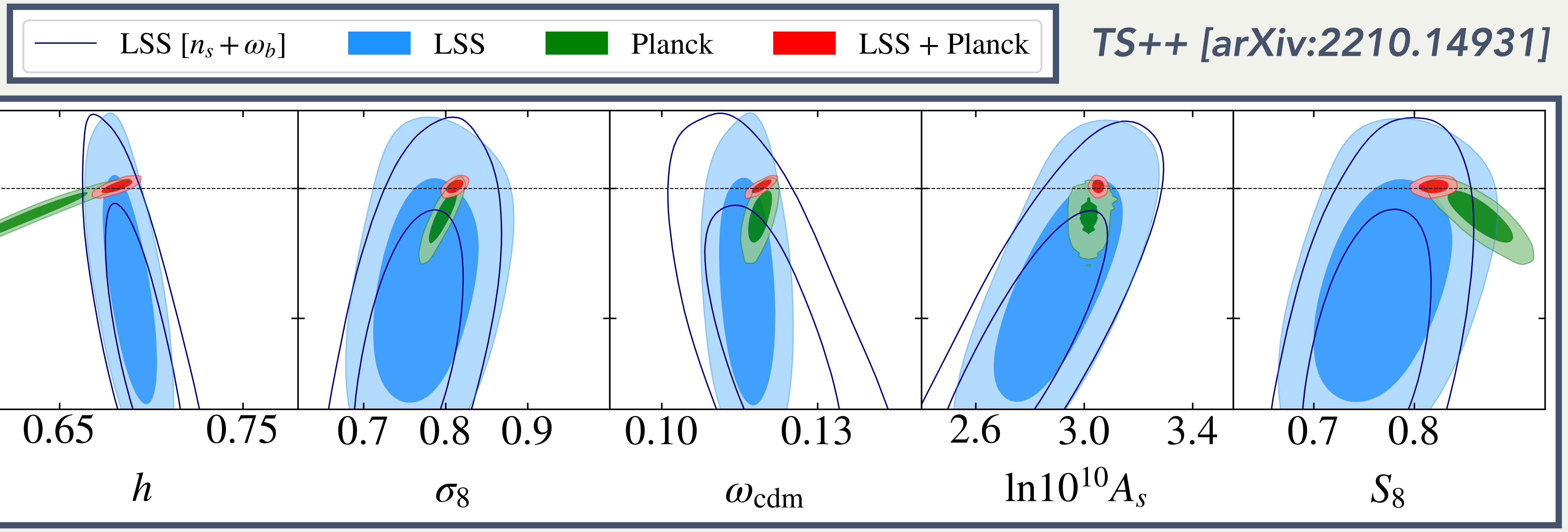
- With LSS data only, we find Ω_k **compatible with zero curvature at** 1.3σ
- The EFT analysis **significantly improves the constraints** on Ω_k by $\sim 50\,\%$ compared to the conventional BAO/$f\sigma_8$ analysis
- The combination of LSS and Planck leads to a **strong constraint** and excludes the (slightly favored) negative values of Ω_k

LSS:

$$\Omega_k = -0.039^{+0.028}_{-0.029}$$

LSS+Planck:

$$\Omega_k = 0.0008^{+0.0018}_{-0.0017}$$



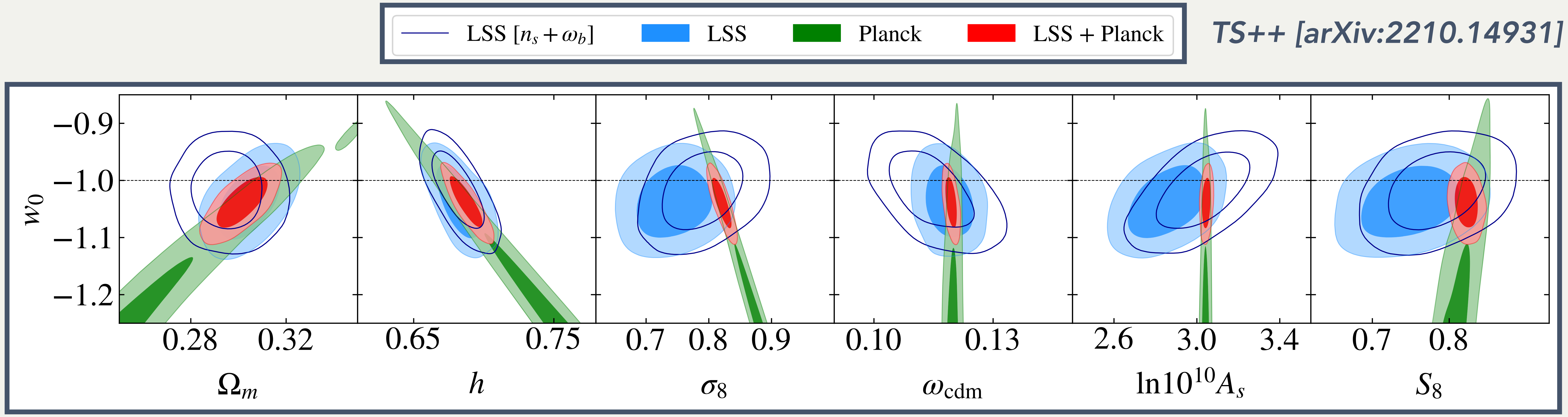
TS++ [arXiv:2210.14931]

Extensions to ΛCDM: dark energy equation of state w_0

- With the LSS data only, we find **no evidence for a universe with $w_0 \neq -1$**
- The EFT analysis **improves the constraints** on w_0 by $\sim 20\,\%$ compared to the conventional BAO/$f\sigma_8$ analysis
- The addition of LSS data select values of w_0 close to -1, located in the 2σ region reconstructed from Planck data

LSS:
$w_0 = -1.038 \pm 0.041$

LSS+Planck:
$w_0 = -1.039 \pm 0.029$



TS++ [arXiv:2210.14931]

Extensions to ΛCDM: effective number of relativistic species $N_{\rm eff}$

- The value of $\Delta N_{\rm eff}$ is **compatible with the standard model**
- Unlike EFTofLSS, **the conventional BAO/$f\sigma_8$ analysis is unable to constrain this parameter**
- The addition of the LSS data **improves** the results of Planck alone by $\sim 25\,\%$

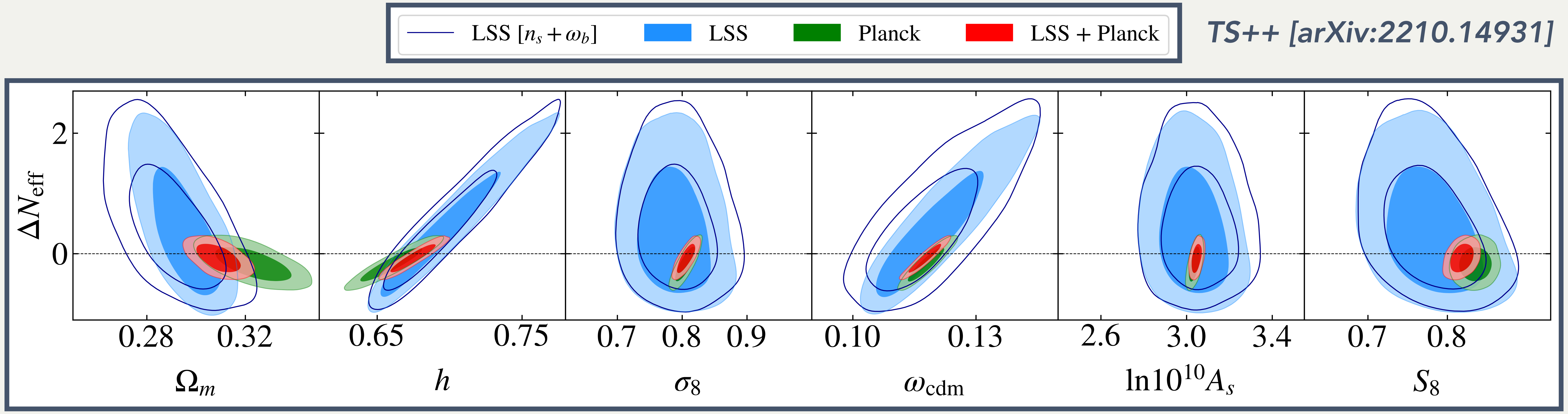
TS++ [arXiv:2210.14931]

LSS:

$$\Delta N_{\rm eff} = 0.40^{+0.44}_{-0.91}$$

LSS+Planck:

$$\Delta N_{\rm eff} = -0.07^{+0.15}_{-0.16}$$



Dark energy equation of state $w_0 \geq -1$

- One can see that this new prior shifts the 2D posteriors inferred from the LSS data in a non-negligible way, while it remains globally stable for the LSS + Planck
- For these analyses, $\Delta\chi^2 = 0$ with respect to ΛCDM, since we obtain best-fit values of $w_0 = -1$

LSS:
$w_0 < -0.932$

LSS+Planck:
$w_0 < -0.965$