Self-Force in Scalar-Tensor Theories of Gravity: Perturbative Approach **Beyond Linear Order**

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22nd September 2023

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1 Why we desire high accuracy scalar-tensor self-force waveforms



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- 2 Derive field equations and equations of motion up to second-order



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- 2 Derive field equations and equations of motion up to second-order
- **3** How easy will implementation be?

1st post-adiabatic (PA) waveforms in GR



[Wardell et al. PRL. 130.24 (2023): 241402]

1st post-adiabatic (PA) waveforms in GR



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But, to test GR we need waveforms to compare to...

Same methods

Same accuracy

New theories

New waveforms

Perturbation theory and self-force



[Source: NASA website]

Perturbation theory and self-force



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Perturbation theory and self-force



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$$\varepsilon \sim \frac{\mu}{M}$$

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Require:

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$$g_{ab} = g_{ab}^{(0)} + \varepsilon^1 h_{ab}^{(1)} + \varepsilon^2 h_{ab}^{(2)} + \mathcal{O}(\varepsilon^3),$$

$$\varphi = \varepsilon^1 \varphi^{(1)} + \varepsilon^2 \varphi^{(2)} + \mathcal{O}(\varepsilon^3),$$

Require:

0PA:
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 $F_{(1)\text{diss}}^{a}[h_{ab}^{(1)}, \varphi^{(1)}] + F_{(1)\text{cons}}^{a}[h_{ab}^{(1)}, \varphi^{(1)}] + F_{(2)\text{diss}}^{a}[h_{ab}^{(2)}, \varphi^{(2)}]$

 $S[\mathbf{g}_{ab}, \Psi] = S_0[\mathbf{g}_{ab}] + S_m[\mathbf{g}_{ab}, \Psi],$

[Detweiler & Whiting, PRD 67, 024025 (2003)], [Detweiler, PRD 85, 044048 (2012)], [Upton & Pound, PRD 103, 124016 (2021)]

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Scalar-Tensor Self-Force

 $S[\mathbf{g}_{ab}, \Psi] = S_0[\mathbf{g}_{ab}] + S_m[\mathbf{g}_{ab}, \Psi],$

$$S_0[\mathbf{g}_{ab}] = \int \frac{\sqrt{-\mathbf{g}}}{16\pi} R \ d^4x$$

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Scalar-Tensor Self-Force

 $\delta_{\mathbf{g}}S_0 \Rightarrow G_{ab}[\mathbf{g}_{cd}]$

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$$\delta_{\mathbf{g}}S_0 \Rightarrow G_{ab}[\mathbf{g}_{cd}]$$

$$\delta_{\tilde{g}}S_{\mathrm{m}} \Rightarrow T_{ab}$$

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 $\begin{array}{l} \text{Perturbative expansion:} \\ \delta G_{ab}[h_{cd}^{(1)}] = T_{ab}^{(1)} \\ \delta G_{ab}[h_{cd}^{(2)}] = T_{ab}^{(2)}[h_{cd}^{(1)\mathcal{R}}] - \delta^2 G_{ab}[h_{cd}^{(1)}, h_{cd}^{(1)}] \end{array}$

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Equation of motion from the action

$$\delta_{x^{\mu}}S \Rightarrow \quad \tilde{u}^{b}\tilde{\nabla}_{b}\tilde{u}^{a} = \tilde{a}^{a} = 0$$

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Perturbative expansion:

$$a_{(1)}^{a} = a_{\text{grav}(1)}^{a} [h_{ab}^{\mathcal{R}(1)}]$$
$$a_{(2)}^{a} = a_{\text{grav}(2)}^{a} [h_{ab}^{\mathcal{R}(1)}, h_{ab}^{\mathcal{R}(2)}]$$

[Mino & Tanaka, PRD 55, 3457 (1997)], [Quinn & Wald, PRD 56, 3381 (1997)], [Pound, PRL 109, 051101 (2012)], [Gralla, PRD 85, 124011 (2012)]

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Scalar-Tensor Self-Force

Scalar-tensor theories of gravity



Scalar-tensor action

 $S[\mathbf{g}_{ab},\varphi,\Psi] = S_0[\mathbf{g}_{ab},\varphi] + \alpha S_c[\mathbf{g}_{ab},\varphi] + S_m[\mathbf{g}_{ab},\varphi,\Psi],$

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$$S_0[\mathbf{g}_{ab},\varphi] = \int \frac{\sqrt{-\mathbf{g}}}{16\pi} \left(R - \frac{1}{2}\partial_\mu \varphi \partial^\mu \varphi\right) d^4x,$$

Dimensionless non-trivial coupling perturbation

$$\zeta := \frac{\alpha}{M^n} = \varepsilon^n \frac{\alpha}{\mu^n}.$$

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$$\frac{\alpha}{\mu^n} = \mathcal{O}(1), \ n \ge 2$$

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Hence, $\zeta \sim \varepsilon^2$

$$\Rightarrow \alpha S_{\rm c}[\mathbf{g}_{ab},\varphi] \approx \mathcal{O}(\varepsilon^3)$$

The Bigger the Bolder



[Keck, Caltech. Getty Images. Andriy_A / Shutterstock.]

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Point scalar charge action

$$S_{\rm m} = -\int_{\gamma} m[\varphi] \sqrt{\mathbf{g}_{ab} \mathbf{u}^a \mathbf{u}^b} d\tau$$

[Damour & Esposito-Farese, CQG 9, 2093 (1992)]

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$$\widetilde{g}_{ab} := g_{ab}^{(0)} + h_{ab}^{(1)\mathcal{R}} + h_{ab}^{(2)\mathcal{R}}...$$
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$$\tilde{g}_{ab} := g_{ab}^{(0)} + h_{ab}^{(1)\mathcal{R}} + h_{ab}^{(2)\mathcal{R}} \dots$$
 and $\tilde{\varphi} := \varphi^{\mathcal{R}}$

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Tensor field equation from the action

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$$G_{ab}[\mathbf{g}_{cd}] = T_{ab}[\tilde{g}_{cd}, \tilde{\varphi}] + T_{ab}^{\mathrm{scal}}[\varphi, \varphi] + \mathcal{O}(\varepsilon^{3})$$

Perturbative tensor field equation

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Equation of motion

$$\delta_{x^{\mu}}S \Rightarrow \quad m[\tilde{\varphi}]\tilde{a}^{a} = m'[\tilde{\varphi}](\tilde{g}^{ab} + \tilde{u}^{a}\tilde{u}^{b})\partial_{b}\tilde{\varphi}$$

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- Two-timescale approximation compatible

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- But, no part more challenging than GR







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- We derived field equations and equations of motion up to second order (up to our assumptions)



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- We derived field equations and equations of motion up to second order (up to our assumptions)
- **③** Implementation builds on GR calculation (and *no more* difficult)