Oligochromatic extreme mass-ratio inspirals (E-EMRIs)

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- Stellar-mass object spiraling into $10^4 - 10^6 \, M_\odot$.
- This range of masses corresponds to relaxed nuclei.
- With LISA z $\sim$ 1, 4.
  [Amaro-Seoane 2018, Babak et al +Amaro-Seoane 2017, Amaro-Seoane et al 2007]
- Rates are very low: 10^{-5}, 10^{-6} per year.
  (stellar-mass BHs and MW)
- Take into account the impact of asymmetry between pro- and retrograde orbits in the location of the LSO helps, if MBH is Kerr.
  [Amaro-Seoane, Sopuerta & Freitag 2013]
- In any case, we don't expect EMRIs at the Galactic Centre, right?

Not really.

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"the ones we have been talking about all along", (using Bernard's words).

Humour me...
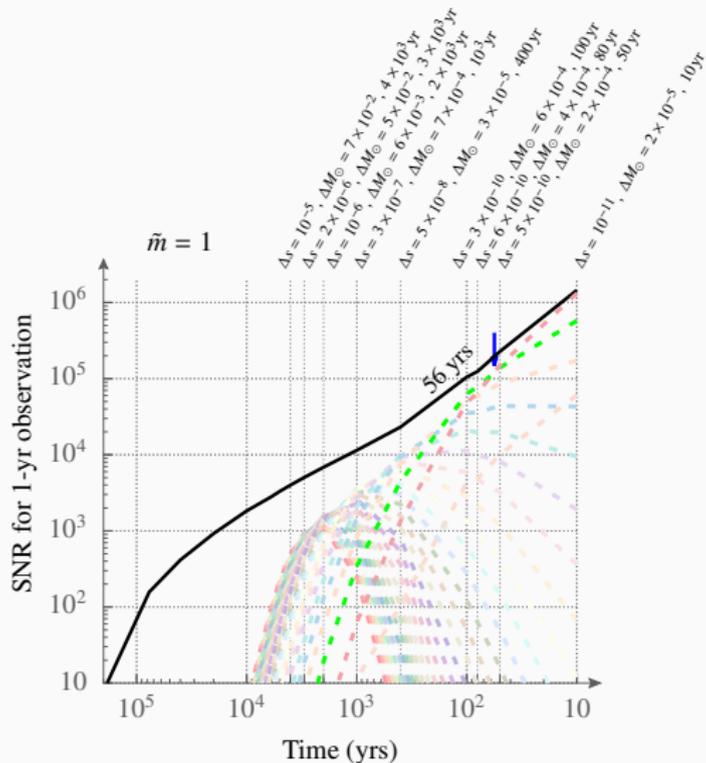
185000 yr before plunge, an E-EMRI would be already on band with
SNR> 10. Waveforms à la Barack and Cutler.
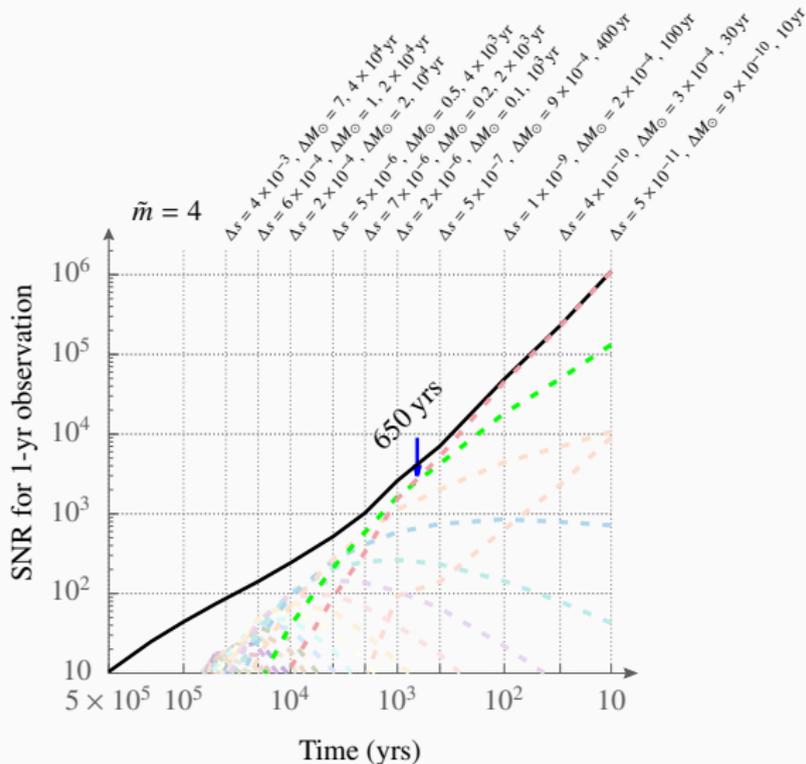
500000 yr before plunge, an E-EMRI would be already on band with

SNR> 10 *[PAS, Lin & Tzanavaris (TBS)]*

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Fine... and the event rate?

■ The event rate in phase-space can be calculated as follows

$$\dot{\Gamma}_{\mathrm{EMRI}} \simeq \int_{a_{\min}}^{a_{\mathrm{crit}}} \frac{dn_{\mathrm{bh}}(a)}{T_{\mathrm{rlx}}(a) \ln\left(\theta_{\mathrm{lc}}^{-2}\right)}$$

■ We need to determine **four quantities**

1. The loss-cone angle
2. The number of bh
3. The relaxation time as a function of the radius
4. The critical radius a_{crit}

(1) The loss-cone angle

■ It can be approximated as

$$\theta_{lc} \simeq \frac{1}{\sqrt{J_{max}/J_{lc}}}$$

$$J_{lc} \simeq \frac{4\,G}{c}\,M_{BH}, \; J_{max}^2 = GM_{BH}a$$

[Alexander & Livio 2001]

■ So that

$$\theta_{lc}^2 \simeq \sqrt{\frac{8\,R_S}{a}}$$

- Assuming the power-law distribution,

$$n_{\mathrm{bh}}(a) \sim f_{\mathrm{sub}} \cdot N_0 \left(\frac{a}{R_0} \right)^{3-\gamma}$$

- Differentiating,

$$dn_{\mathrm{bh}}(a) = f_{\mathrm{sub}} (3 - \gamma) \frac{N_0}{R_0} \left(\frac{a}{R_0} \right)^{2-\gamma} da$$

■ Relaxation due to the most massive stellar species, **stellar-mass black holes**

$$T_{\mathrm{rlx}}(a) = T_0 \left(\frac{a}{R_0} \right)^{\gamma - 3/2}$$

■ With

$$T_0 \simeq \frac{4.26}{(3 - \gamma)(1 + \gamma)^{3/2}} \frac{\sqrt{R_0^3 (G M_{\mathrm{BH}})^{-1}}}{\ln(\Lambda) \, N_0} \left(\frac{M_{\mathrm{BH}}}{m_{\mathrm{bh}}} \right)^2$$

■ From its definition, **it is the threshold between stellar dynamics and the GW-dominated regime**

$$T_{\mathrm{rlx,\,peri}} = C\,T_{\mathrm{GW}}(a,\,e)$$

■ And

$$\frac{8\,GM_{\mathrm{BH}}}{c^2} = (1-e)a\mathcal{W}(\iota,\,\mathrm{s})$$

$\mathcal{W}(\iota,\,\mathrm{s})$ takes into account the asymmetry between pro- and retrograde orbits for the location of the LSO for a Kerr MBH
[Amaro-Seoane et al 2013].

The function depends on the spin of the MBH a and the inclination of the orbit ι.

- EMRI orbits have $e \sim 1$, **hence**

$$T_{\mathrm{GW}}(a, e) \sim \sqrt{2} \, \frac{24}{85} \frac{c^5}{G^3} \frac{a^4 \, (1-e)^{7/2}}{m_{\mathrm{bh}} \, M_{\mathrm{BH}}^2}$$

- So that we obtain

$$a_{\mathrm{crit}} = R_0 \left[\frac{20480}{1207} (3-\gamma)(1+\gamma)^{3/2} C \mathcal{W}(\iota, \mathrm{s})^{5/2} N_0 \ln(\Lambda) \left(\frac{M_{\mathrm{BH}}}{m_{\mathrm{bh}}} \right)^{-1} \right]^{\frac{1}{\gamma-3}}$$
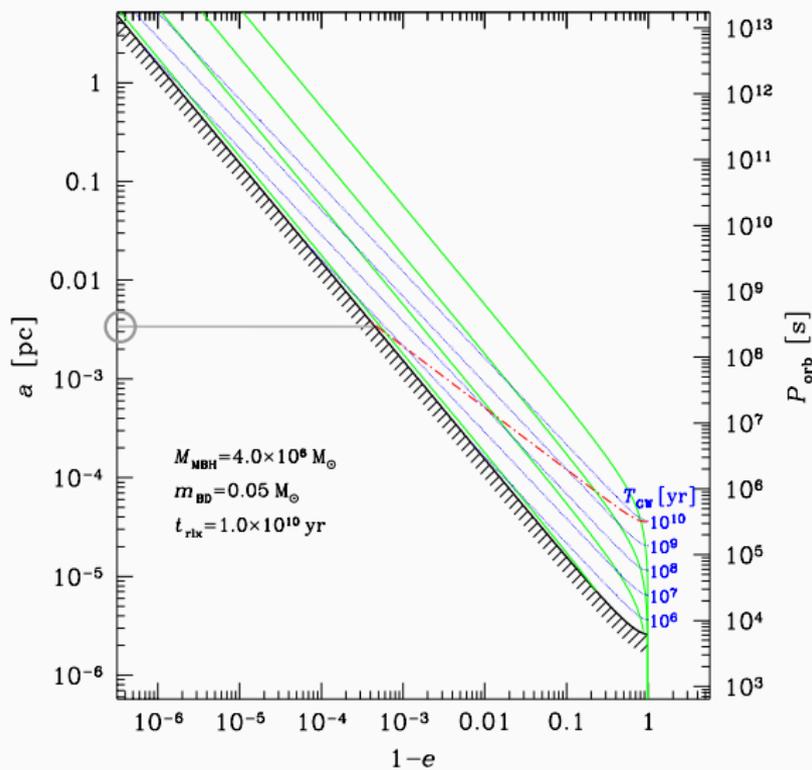
Definition of a_{crit}, at a fixed t_{rlx} for illustration.

■ The integral can be solved analytically (*)

$$
\dot{\Gamma} \sim 1.92 \times 10^{-6} \, \mathrm{yrs}^{-1} \tilde{N}_0 \, \tilde{\Lambda} \, \tilde{R}_0^{-2} \, \tilde{m}^2 \times
$$

$$
\left\{ 1.6 \times 10^{-1} \tilde{R}_0^{1/2} \tilde{N}_0^{-1/2} \tilde{\Lambda}^{-1/2} \tilde{m}^{1/2} \, \mathcal{W}(\iota, \mathrm{s})^{-5/4} \times \right.
$$

$$
\left[\ln \left(9138 \, \tilde{R}_0 \, \tilde{N}_0^{-1} \tilde{\Lambda}^{-1} \tilde{m} \, \mathcal{W}(\iota, \mathrm{s})^{-5/2} \right) - 2 \right] -
$$

$$
\left. 4 \times 10^{-2} \tilde{R}_0^{1/2} \times \left[\ln \left(618 \, \tilde{R}_0 \right) - 2 \right] \right\},
$$

(*) If you distrust computer algebra systems.

■ with the following notation,

$$\tilde{\Lambda} := \left(\frac{\ln(\Lambda)}{13} \right), \; \tilde{N}_0 := \left(\frac{N_0}{12000} \right)$$

$$\tilde{R}_0 := \left(\frac{R_{\mathrm{h}}}{1\mathrm{pc}} \right), \; \tilde{m} := \left(\frac{m}{10 \, M_\odot} \right).$$

■ The advantage is that $\dot{\Gamma}$ contains **all physical information, including the relaxation time and critical radius, embedded**
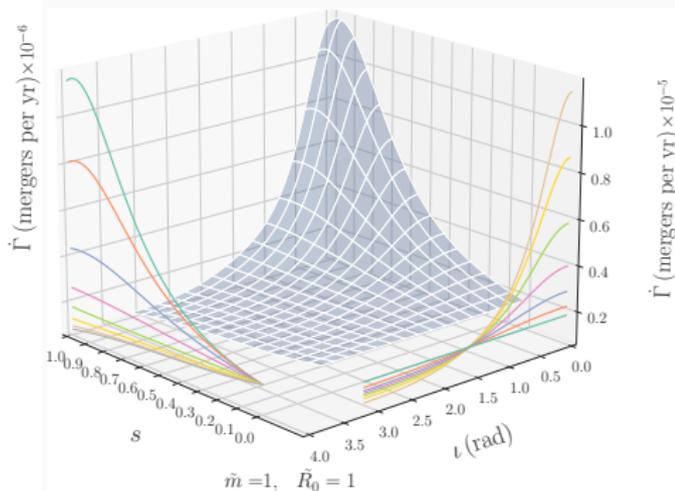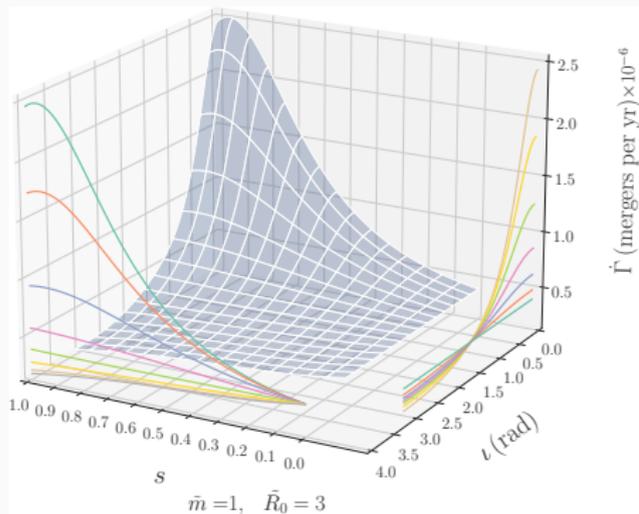
Assume a *Alexander &Hopman 2009, Preto & AS 2009* exponent based on *Peebles 1972* power-law solution, $m_{\rm bh} = 10\,M_\odot$. The event rate depends on the inclination of the orbit (ι) and the spin of the MBH (s). **The values for $\tilde{m} = 4$ are somewhat larger.**

So much fuss for this? We knew it.

$\tilde{m} = 4$

$\Delta s = 4 \times 10^{-3}, \Delta M_\odot = 7, 4 \times 10^5 \text{yr}$
$\Delta s = 6 \times 10^{-4}, \Delta M_\odot = 1, 2 \times 10^4 \text{yr}$
$\Delta s = 2 \times 10^{-4}, \Delta M_\odot = 2, 10^4 \text{yr}$
$\Delta s = 5 \times 10^{-6}, \Delta M_\odot = 0.5, 4 \times 10^3 \text{yr}$
$\Delta s = 7 \times 10^{-6}, \Delta M_\odot = 0.2, 2 \times 10^3 \text{yr}$
$\Delta s = 2 \times 10^{-6}, \Delta M_\odot = 0.1, 10^3 \text{yr}$
$\Delta s = 5 \times 10^{-7}, \Delta M_\odot = 9 \times 10^{-4}, 400 \text{yr}$
$\Delta s = 1 \times 10^{-9}, \Delta M_\odot = 2 \times 10^{-4}, 100 \text{yr}$
$\Delta s = 4 \times 10^{-10}, \Delta M_\odot = 3 \times 10^{-4}, 30 \text{yr}$
$\Delta s = 5 \times 10^{-11}, \Delta M_\odot = 9 \times 10^{-10}, 10 \text{yr}$

650 yrs

SNR for 1-yr observation

Time (yrs)
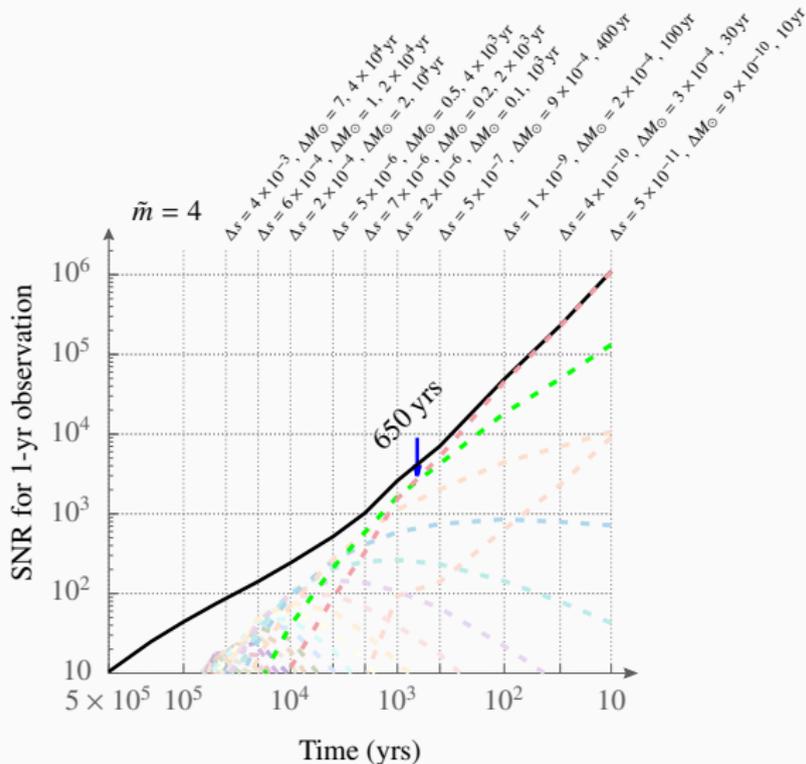
500000 yr before plunge, E-EMRI on band with SNR> 10

- E-EMRIs **spend a long time on band**
- The lifetime with SNR>10 in LISA is of $T \sim 10^5 \, \mathrm{yr}^{-1}$
  and the event rate $\dot{\Gamma} \cong 10^{-6} \, \mathrm{yr}^{-1}$
- Therefore... **How many of these in band??**

- From the continuity equation of the events **we can derive the relative occupation fractions of the line density** $g = dN/da$
- Taking into account the eccentricity of the sources when integrating N, the inclinations and spins **we find the final numbers**
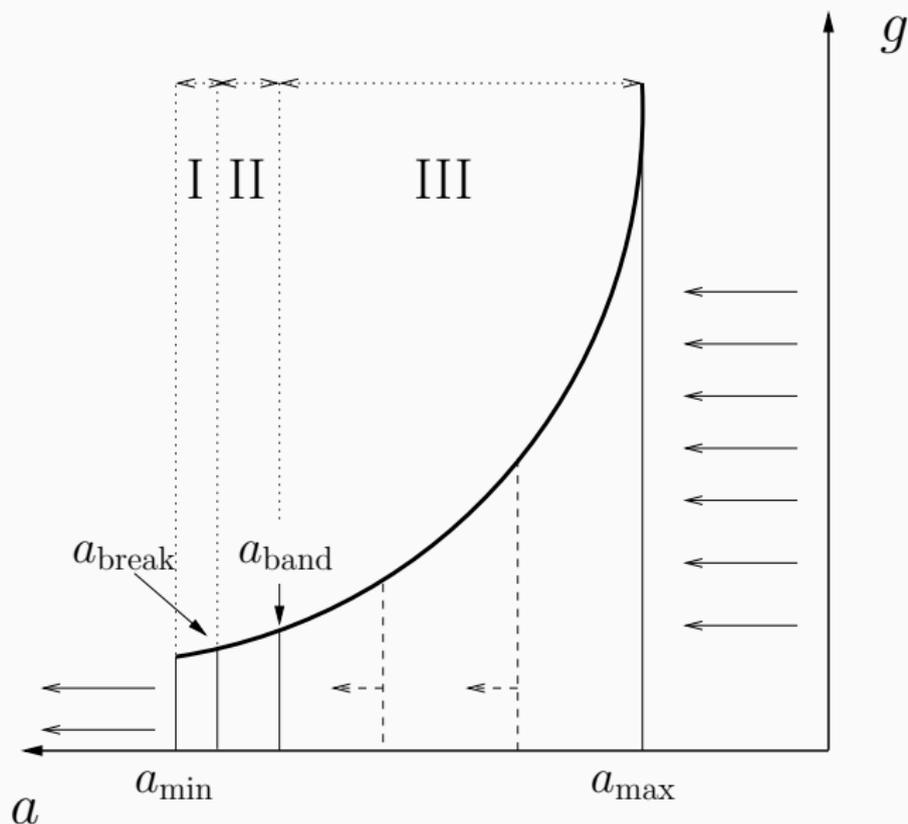
Three equations, three unknowns

$$\frac{N_{\mathrm{II}}}{N_{\mathrm{III}}} = \frac{a_{\mathrm{band}}^{1/2} - a_{\mathrm{thr}}^{1/2}}{a_{\mathrm{crit}}^{1/2} - a_{\mathrm{band}}^{1/2}}$$

$$\frac{N_{\mathrm{I}}}{N_{\mathrm{II}} + N_{\mathrm{III}}} = \frac{1}{8} \times \frac{1 - (a_{\mathrm{min}}/a_{\mathrm{thr}})^4}{(a_{\mathrm{crit}}/a_{\mathrm{thr}})^{1/2} - 1}$$

$$N_{\mathrm{I}} + N_{\mathrm{II}} = \dot{\Gamma} \times T(a_{\mathrm{crit}}, e)$$

$$N_{\mathrm{I}} = \dot{\Gamma} \times T(a_{\mathrm{crit}}, e) \times \Omega_1$$
$$N_{\mathrm{II}} = \dot{\Gamma} \times T(a_{\mathrm{crit}}, e) \times \Omega_2$$
$$N_{\mathrm{III}} = \dot{\Gamma} \times T(a_{\mathrm{crit}}, e) \times \Omega_3, \tag{1}$$

where we have introduced the weighting functions Ω_1, Ω_2 and Ω_3

$$\Omega_1 \equiv \frac{(\sqrt{a_{\mathrm{thr}}} - \sqrt{a_{\mathrm{crit}}})(a_{\mathrm{min}}^4 - a_{\mathrm{thr}}^4)}{a_{\mathrm{thr}}^4(8\sqrt{a}(\sqrt{a_{\mathrm{crit}}/a_{\mathrm{thr}}} - 1) + \sqrt{a_{\mathrm{crit}}}) + a_{\mathrm{thr}}^{9/2}(7 - 8\sqrt{a_{\mathrm{crit}}/a_{\mathrm{thr}}}) + a_{\mathrm{min}}^4(\sqrt{a_{\mathrm{thr}}} - \sqrt{a_{\mathrm{crit}}})}$$

$$\Omega_2 \equiv \frac{8a_{\mathrm{thr}}^4(\sqrt{a} - \sqrt{a_{\mathrm{thr}}})(\sqrt{a_{\mathrm{crit}}/a_{\mathrm{thr}}} - 1)}{a_{\mathrm{thr}}^4(8\sqrt{a}(\sqrt{a_{\mathrm{crit}}/a_{\mathrm{thr}}} - 1) + \sqrt{a_{\mathrm{crit}}}) + a_{\mathrm{thr}}^{9/2}(7 - 8\sqrt{a_{\mathrm{crit}}/a_{\mathrm{thr}}}) + a_{\mathrm{min}}^4(\sqrt{a_{\mathrm{thr}}} - \sqrt{a_{\mathrm{crit}}})}$$

$$\Omega_3 \equiv \frac{8a_{\mathrm{thr}}^4(\sqrt{a} - \sqrt{a_{\mathrm{crit}}})(1 - \sqrt{a_{\mathrm{crit}}/a_{\mathrm{thr}}})}{(a_{\mathrm{thr}}^4(8\sqrt{a}(\sqrt{a_{\mathrm{crit}}/a_{\mathrm{thr}}} - 1) + \sqrt{a_{\mathrm{crit}}}) + a_{\mathrm{thr}}^{9/2}(7 - 8\sqrt{a_{\mathrm{crit}}/a_{\mathrm{thr}}}) + a_{\mathrm{min}}^4(\sqrt{a_{\mathrm{thr}}} - \sqrt{a_{\mathrm{crit}}})}$$

$\tilde{m} = 1, \quad \tilde{R}_0 = 3, \quad \tilde{T}(a_{\text{crit}}, e) = 185000 \text{ yrs}$

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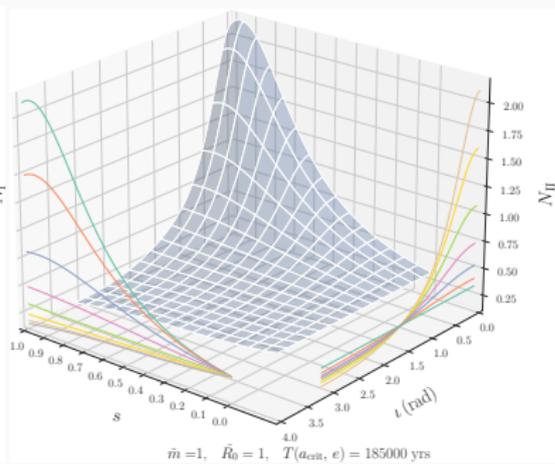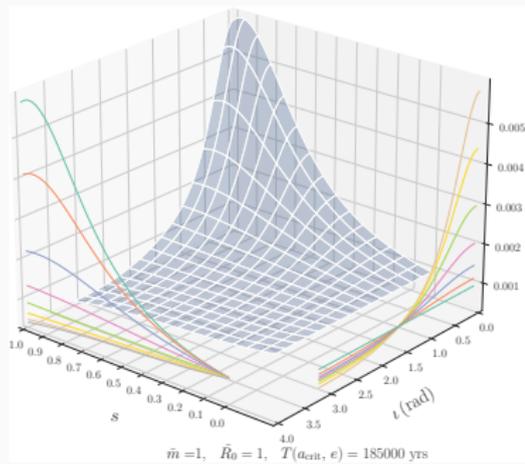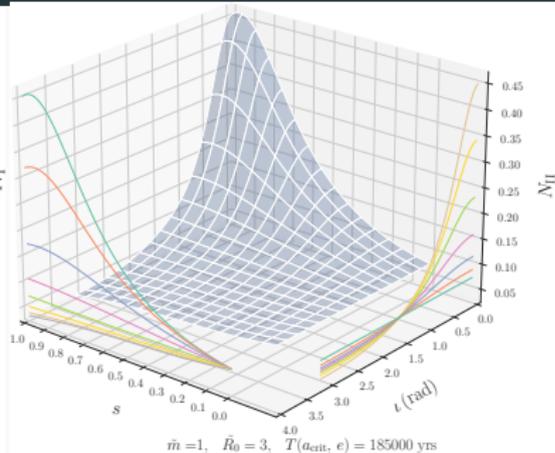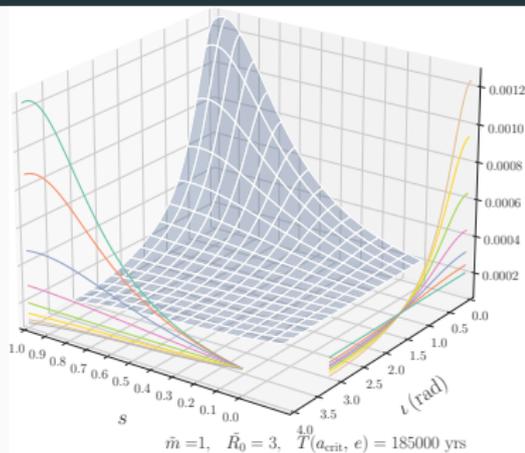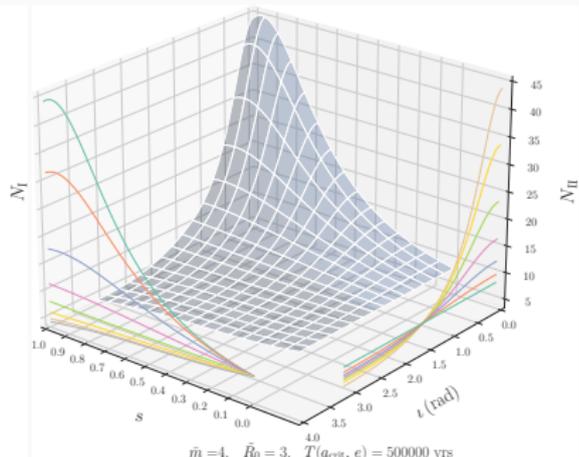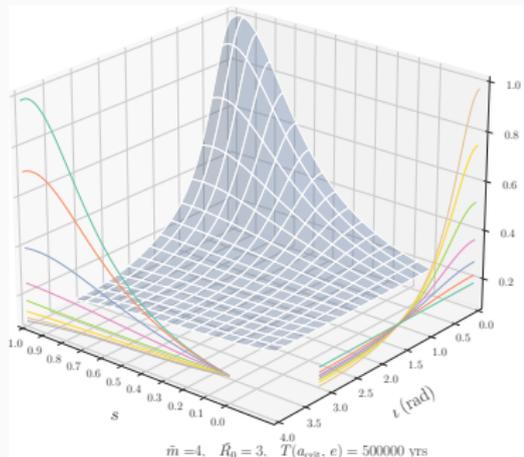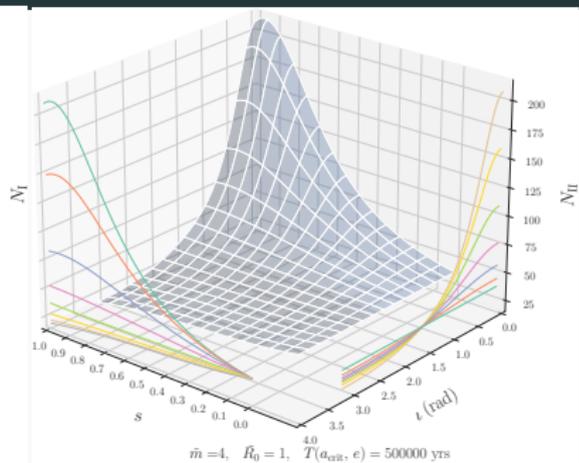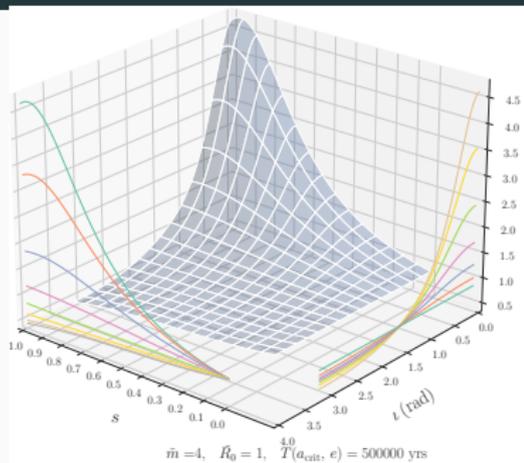
A forest of E-EMRIs:
fore- and background population

$\tilde{m} = 1$

Total SNR for 1-yr, from SgrA*

10 kpc · 10^2 kpc · 10^3 kpc · 10^4 kpc · 10^5 kpc

Time (yrs)
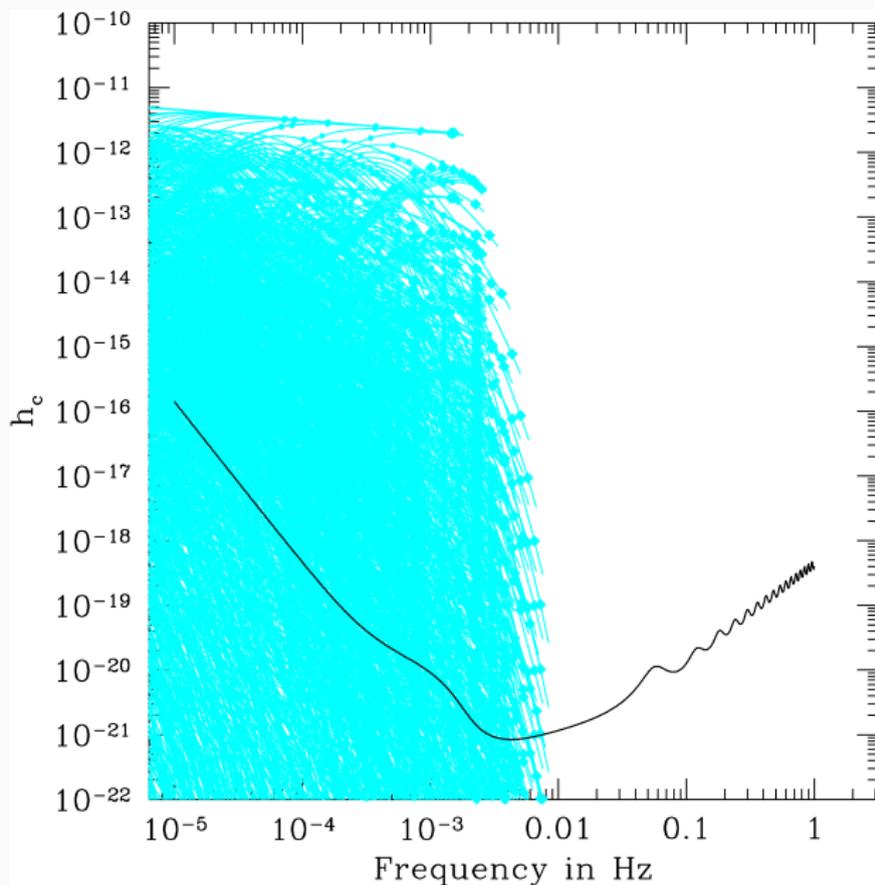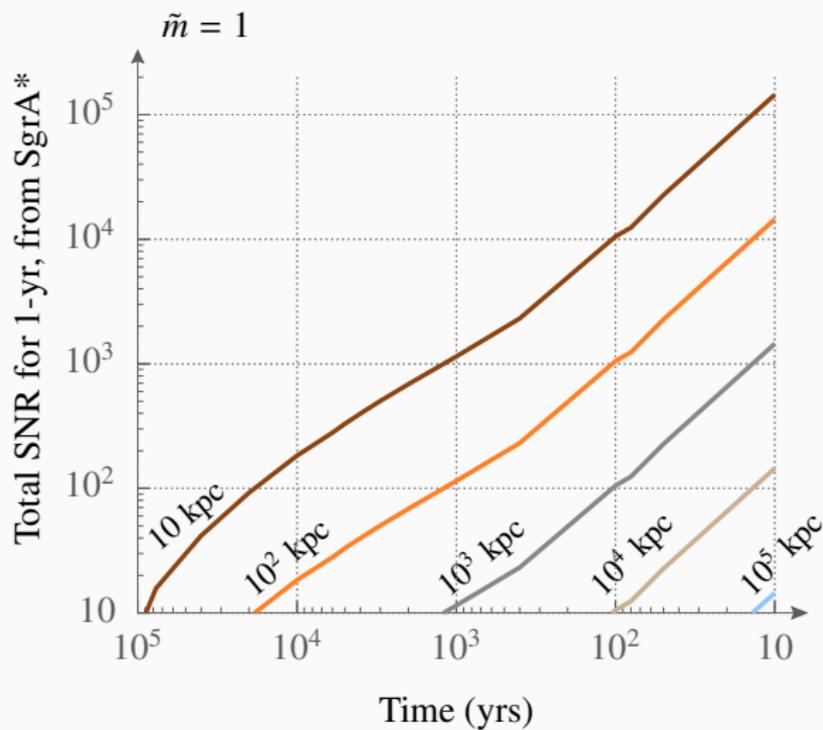
[PAS, Yiren & TzanavarisTBS]

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- The forest hence **extends beyond 0.1 Gpc**

So, how does the fore- and background look like?

I don't know (yet).

So…

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  Early-inspiral EMRIs

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- **I am working on it...**

Mono- and oligochromatic EMRIs

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EXTRA MATERIAL

- I am assuming that different mass "species" contribute to relaxation individually...
  ... and that the total amount of relaxation in the system can be added up linearly from them.
  Can I do that?
  "Yes."
- The distribution function of mass and velocity is $f(m, v)$,
  and a moment of the change of velocities is of the form

$$< dv^2 >= \int dv^2 f(m, v) \, dm \, dv.$$

- And this can be envisaged as

$$< dv^2 >= \sum_m n(m) \left(\int dv^2 f(v) dv \right),$$

with $n(m)$ the density of stars of mass m.

Harmonics

The strain amplitude in the n-th harmonic at a given distance D, normalized to the typical values of this work is

$$
\begin{aligned}
h_n &= g(n, e) \frac{G^2 \, M_{\mathrm{BH}} m_{\mathrm{CO}}}{D \, a \, c^4} \\
&\simeq 8 \times 10^{-23} g(n, e) \left(\frac{D}{500 \, \mathrm{Mpc}} \right)^{-1} \left(\frac{a}{10^{-5} \, \mathrm{pc}} \right)^{-1} \\
&\quad \left(\frac{M_{\mathrm{BH}}}{10^3 \, M_\odot} \right) \left(\frac{m_{\mathrm{CO}}}{10 \, M_\odot} \right).
\end{aligned}
$$

In this expression M_{BH} is the mass of the IMBH, m_{CO} is the mass of the compact object (CO), and $g(n, e)$ is a function of the harmonic number n and the eccentricity e *[Peters & Matthews 1963]*. We consider the RMS amplitude averaged over the two GW polarizations and all directions.