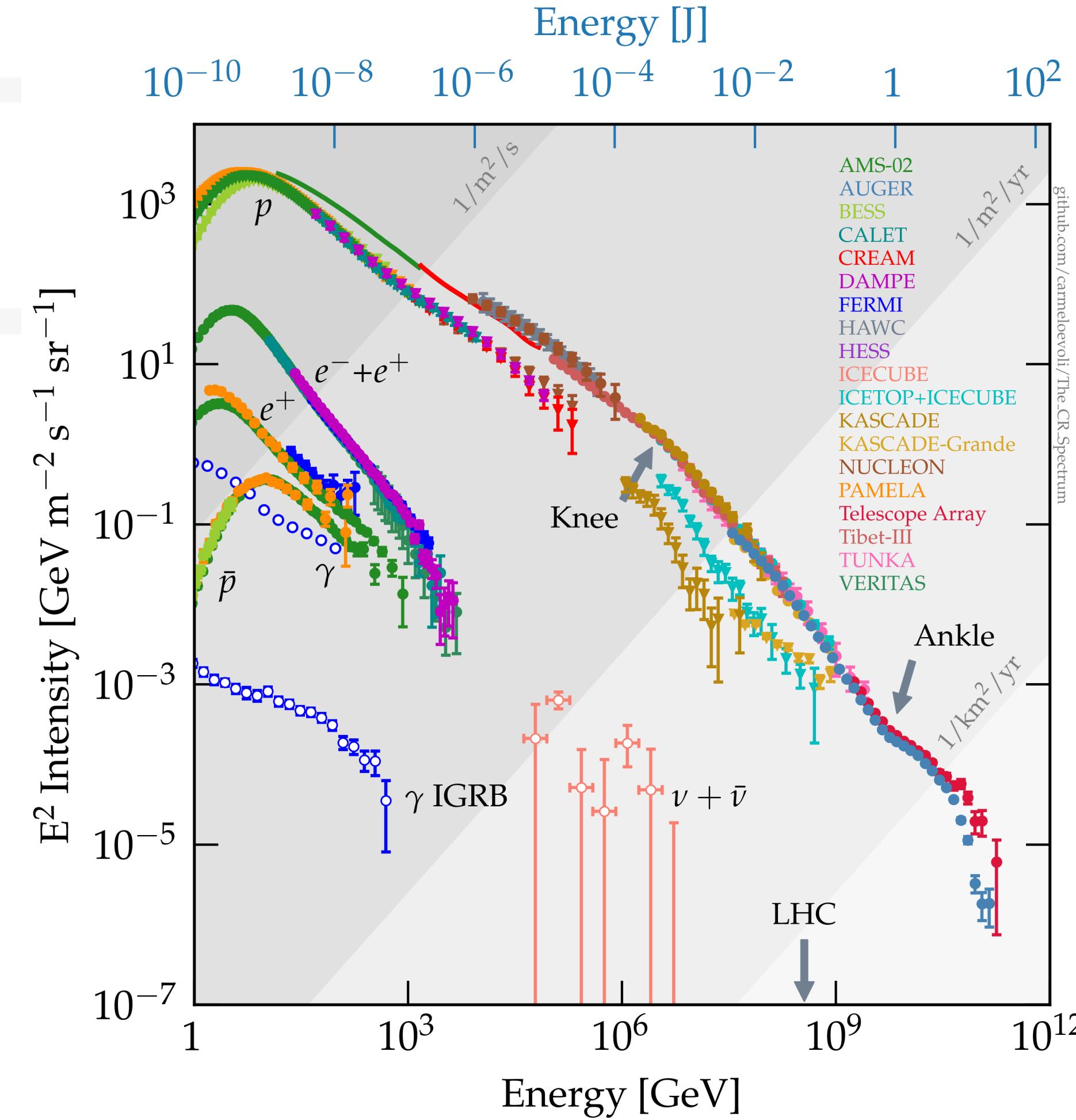


**GRAN SASSO
SCIENCE INSTITUTE**

CR transport in MHD turbulence

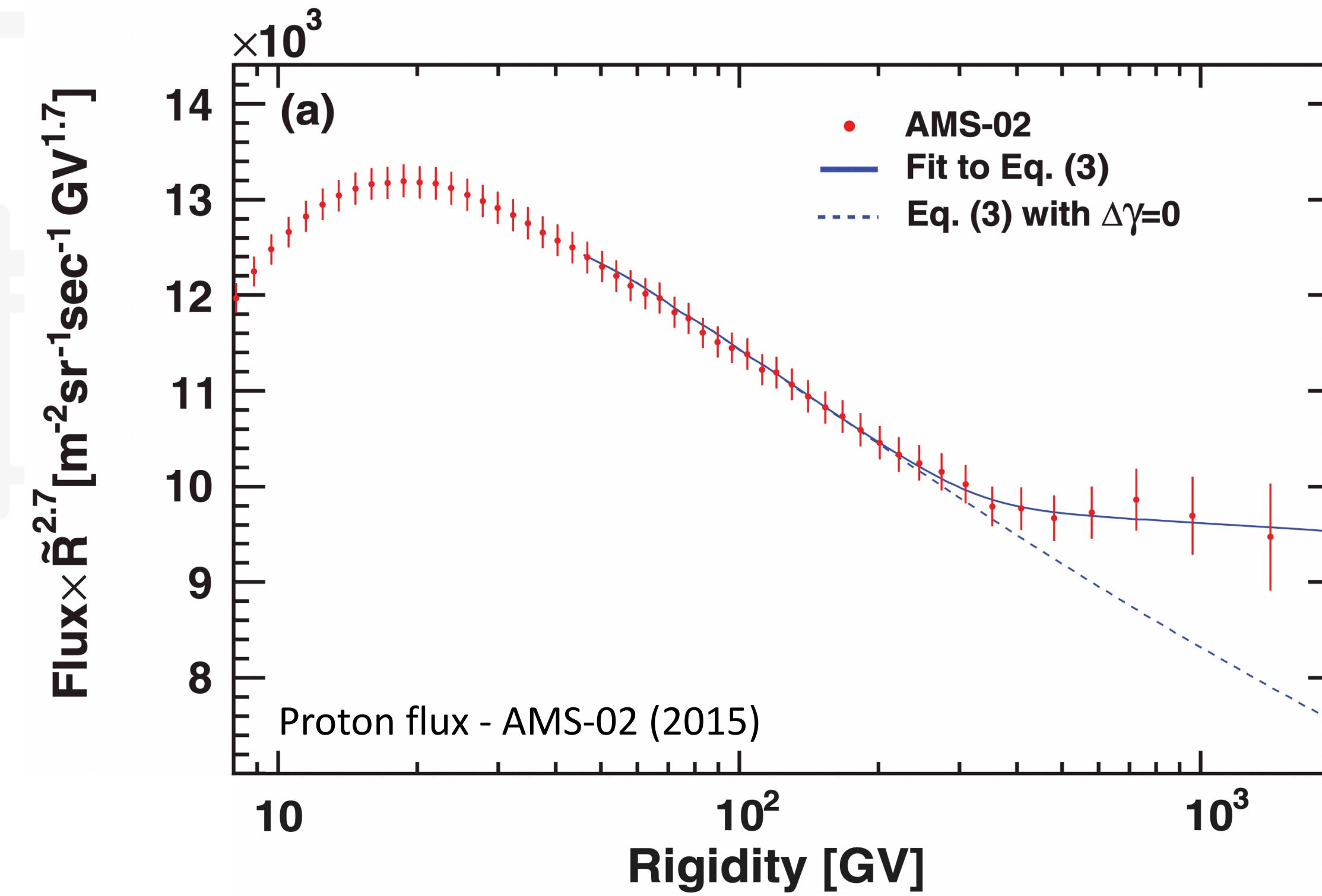
Presented by: Ottavio Fornieri

The general picture



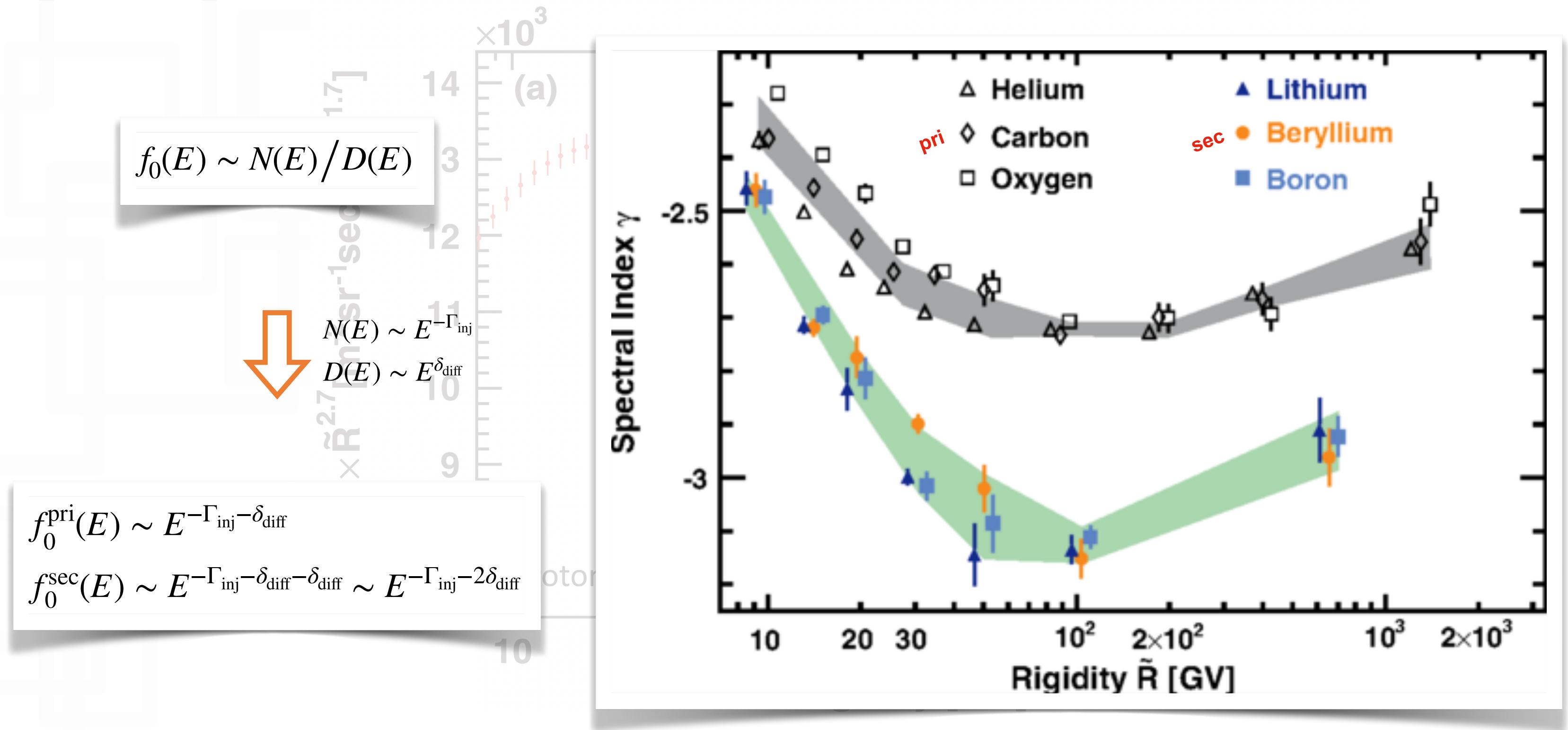
G S
S I

Anomalies in the details

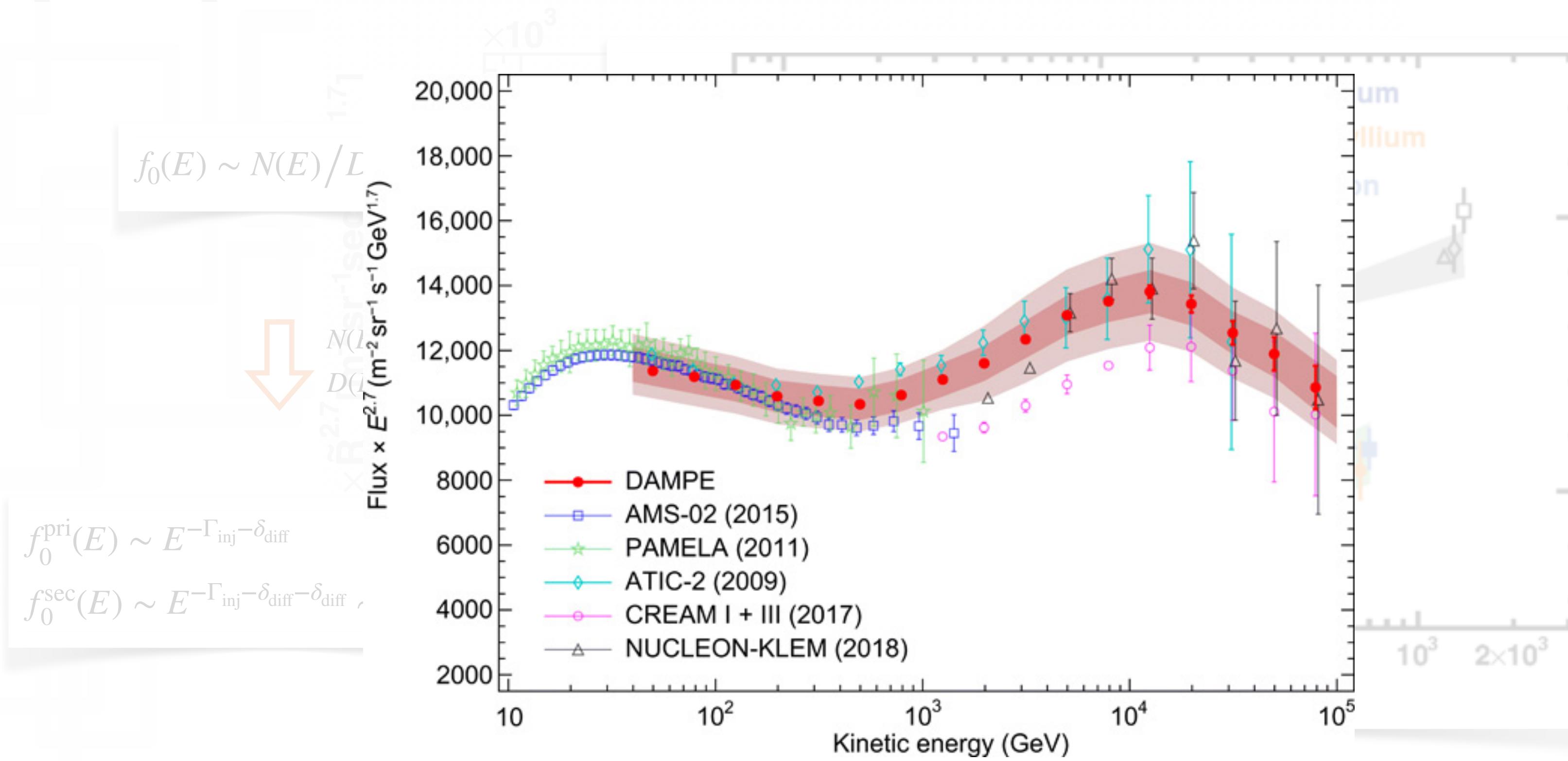


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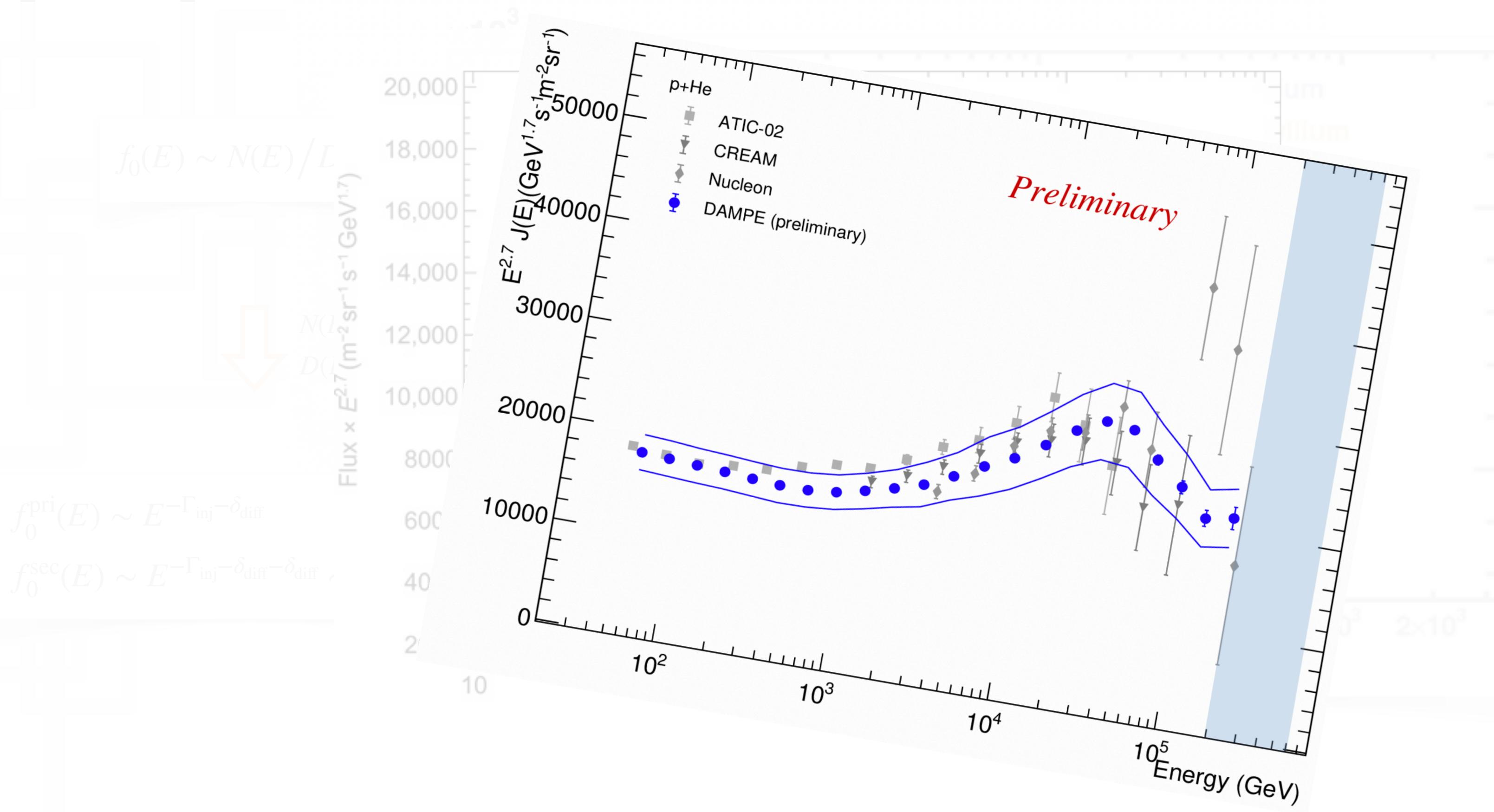
Anomalies in the details



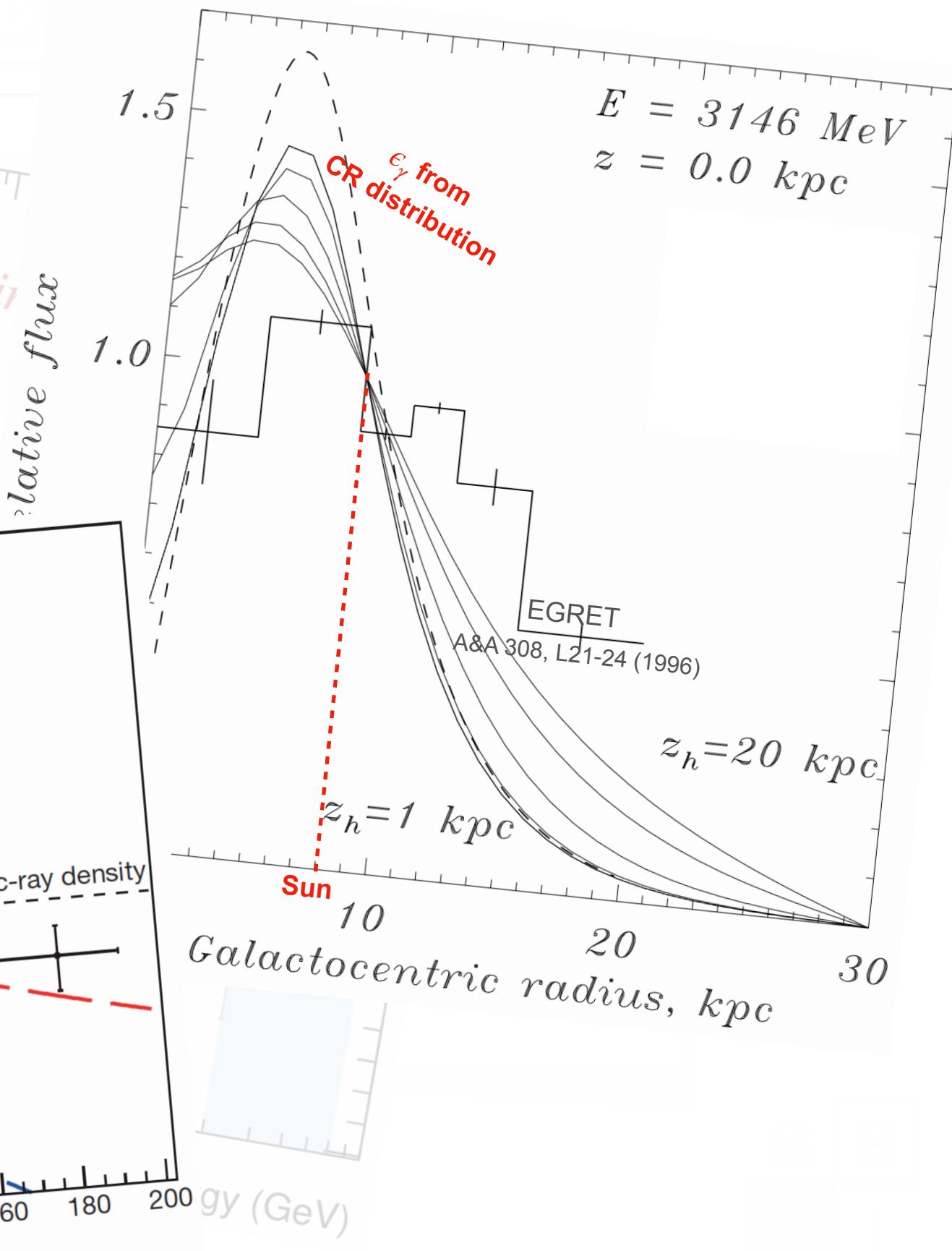
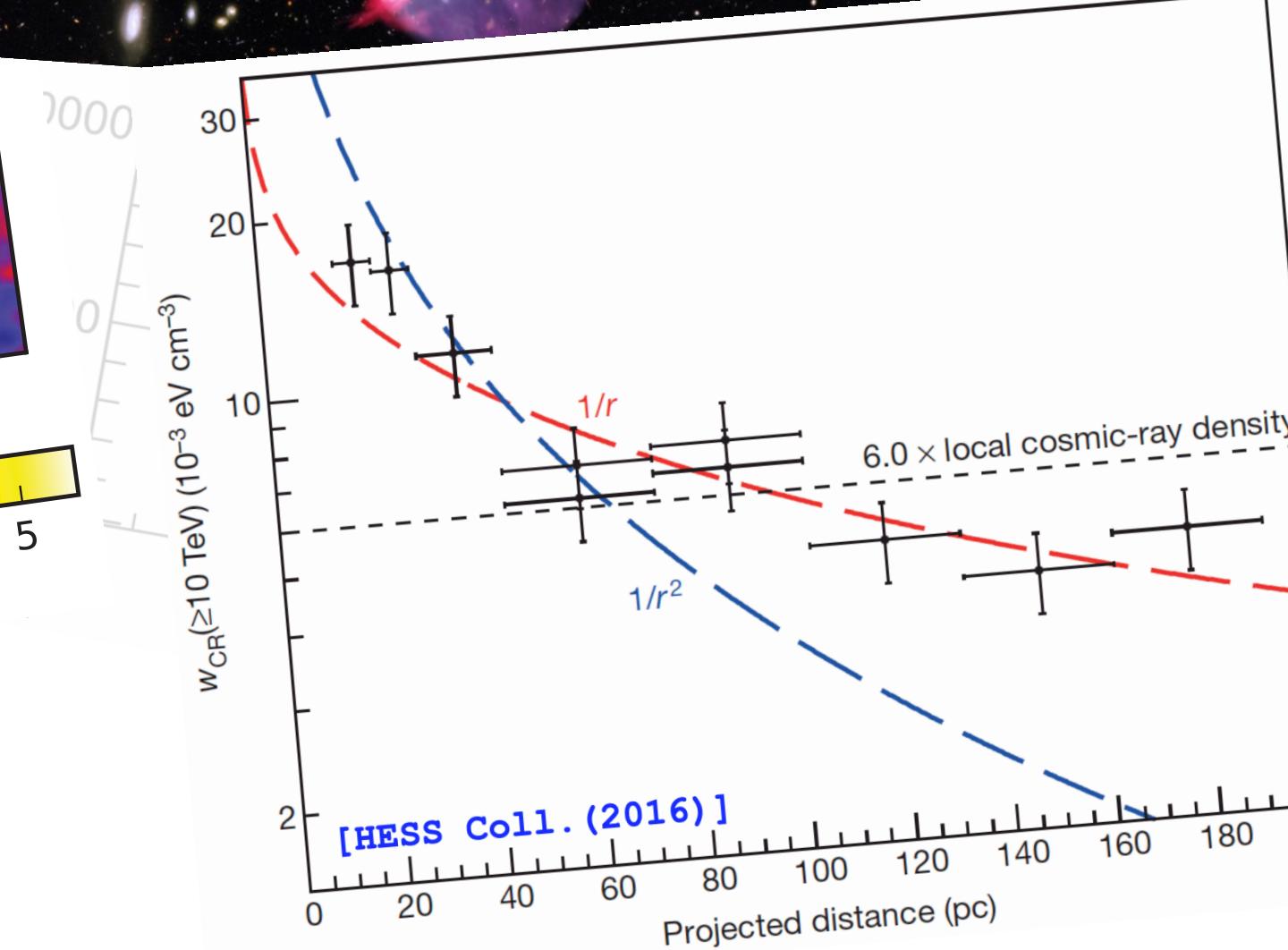
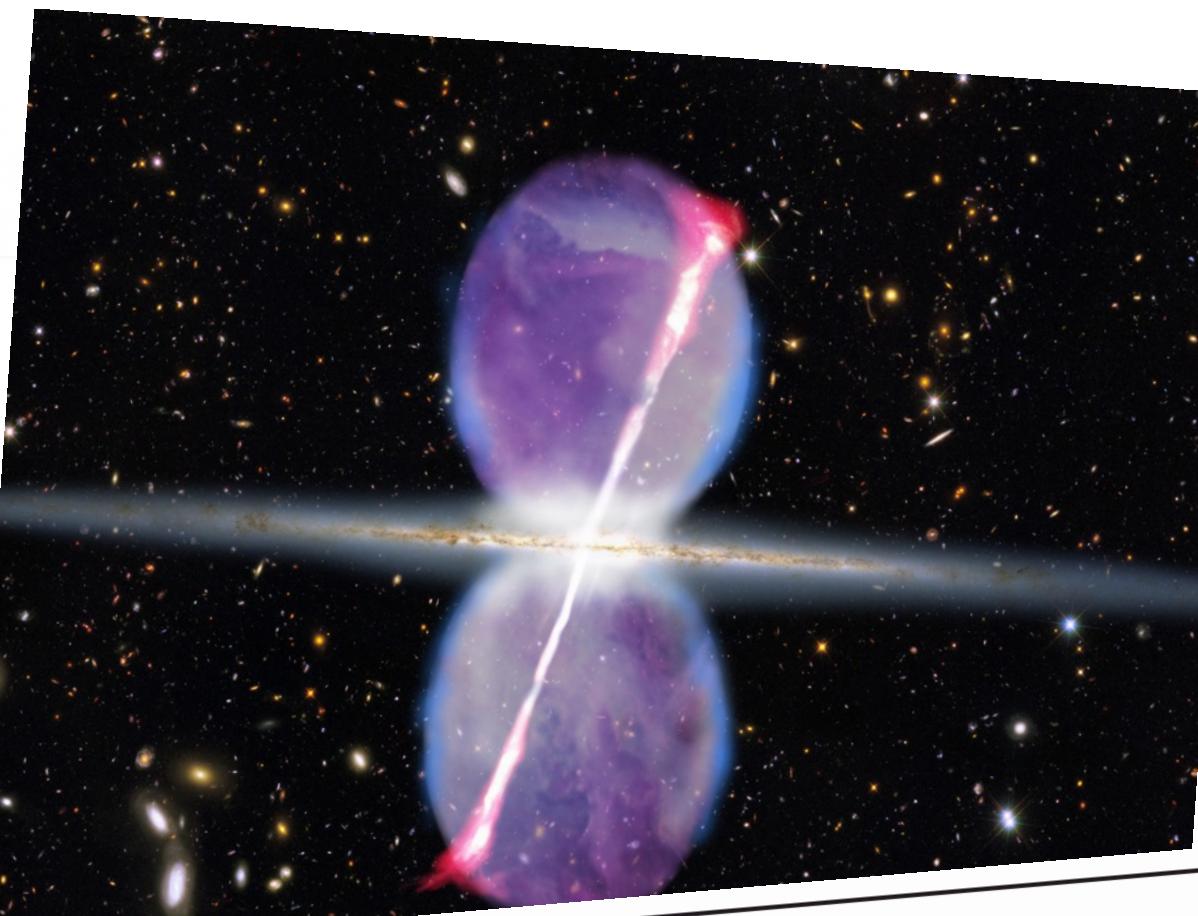
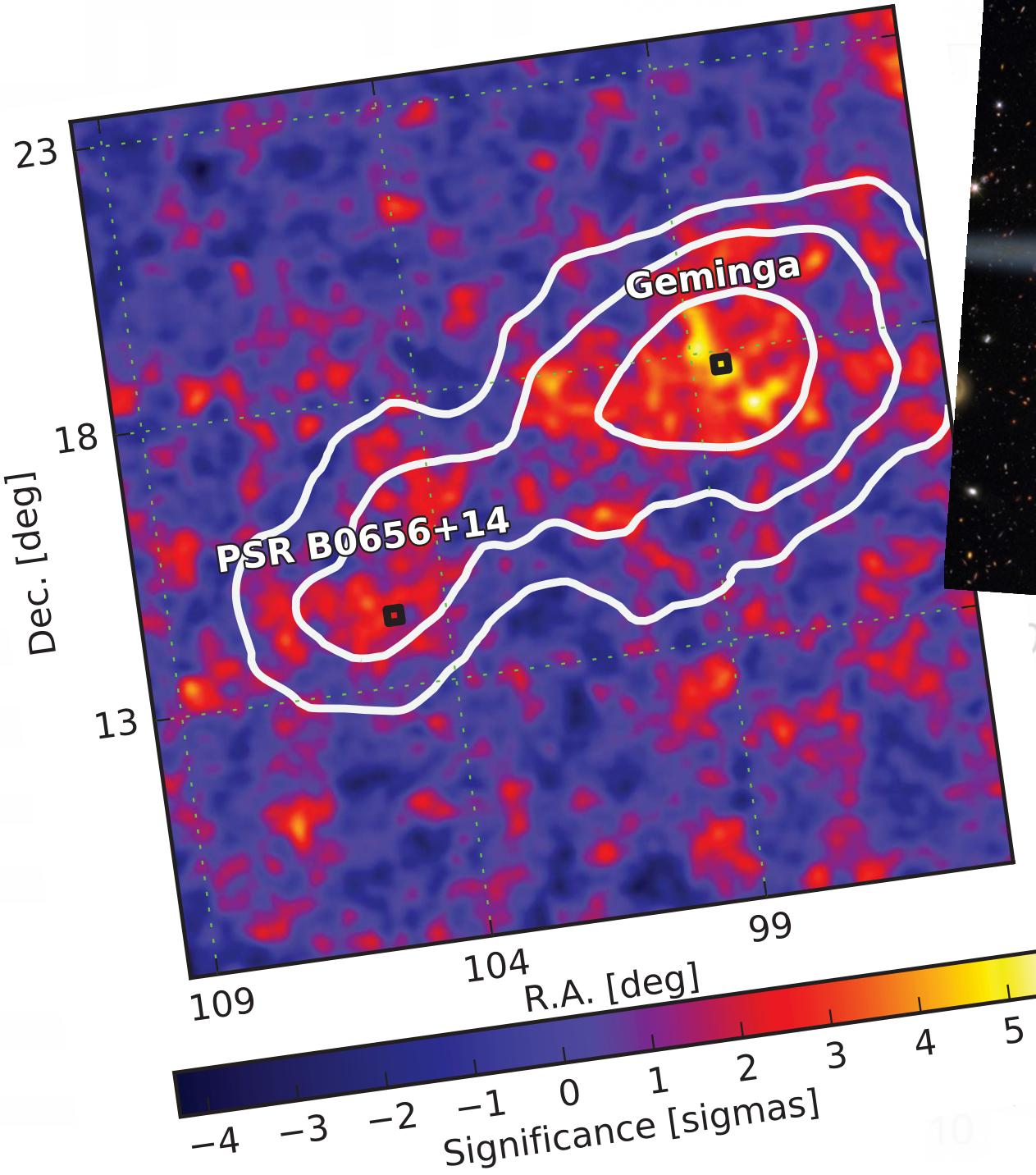
Anomalies in the details



Anomalies in the details



Anomalies in the details



CR spectrum at Earth =

**Acceleration at
the Sources**

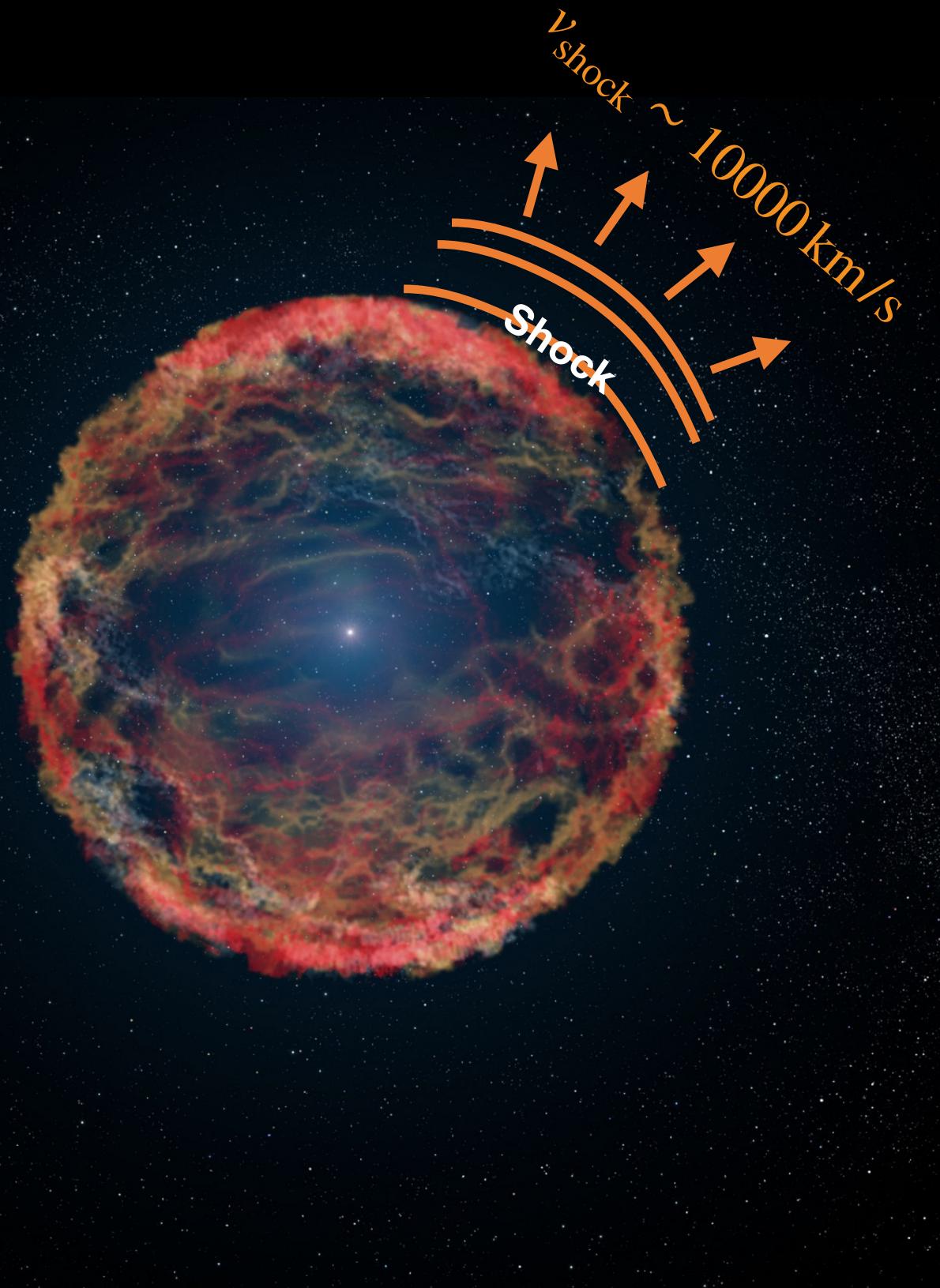


**Transport across
our Galaxy**

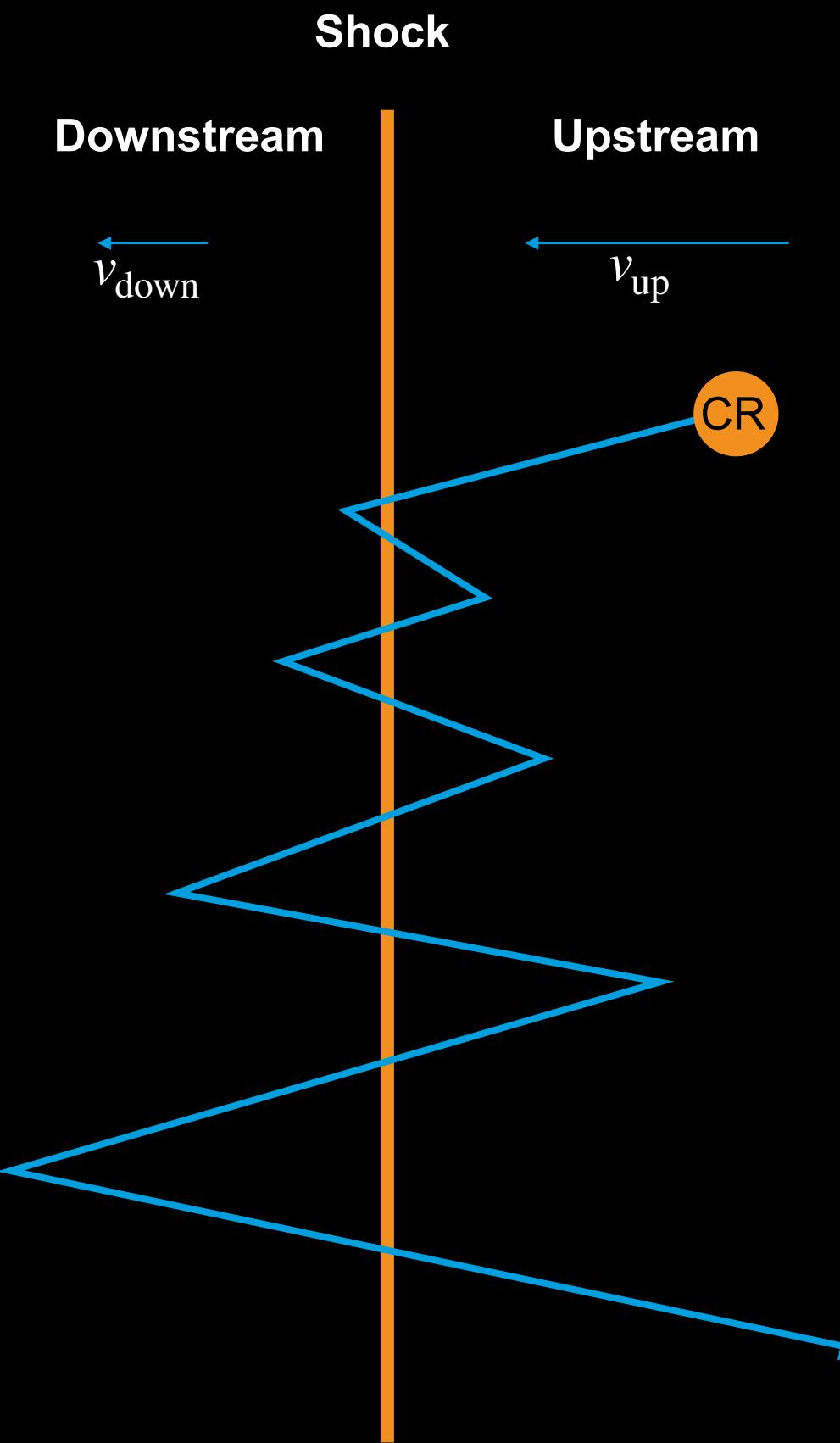
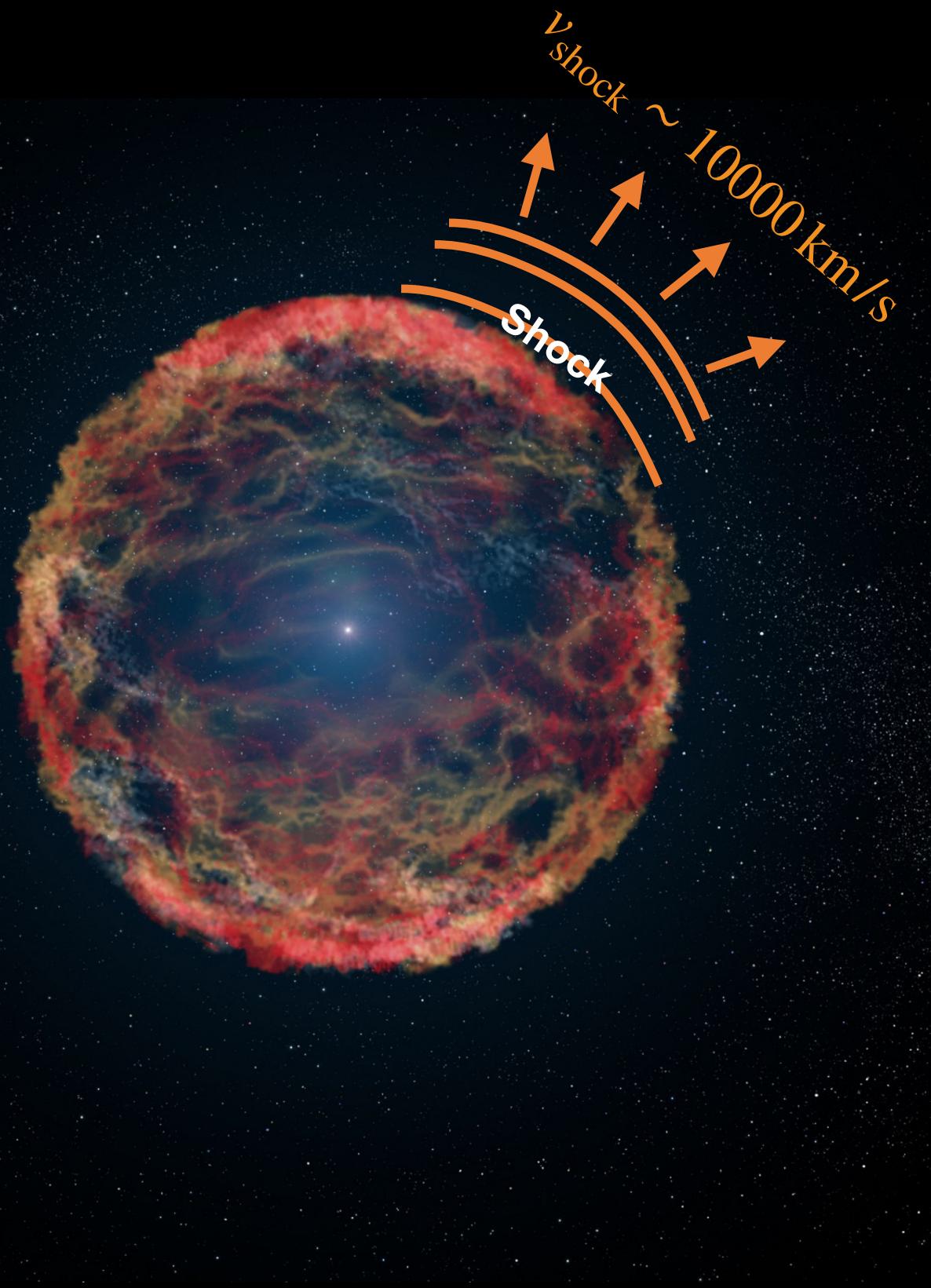


CR production up to high energies

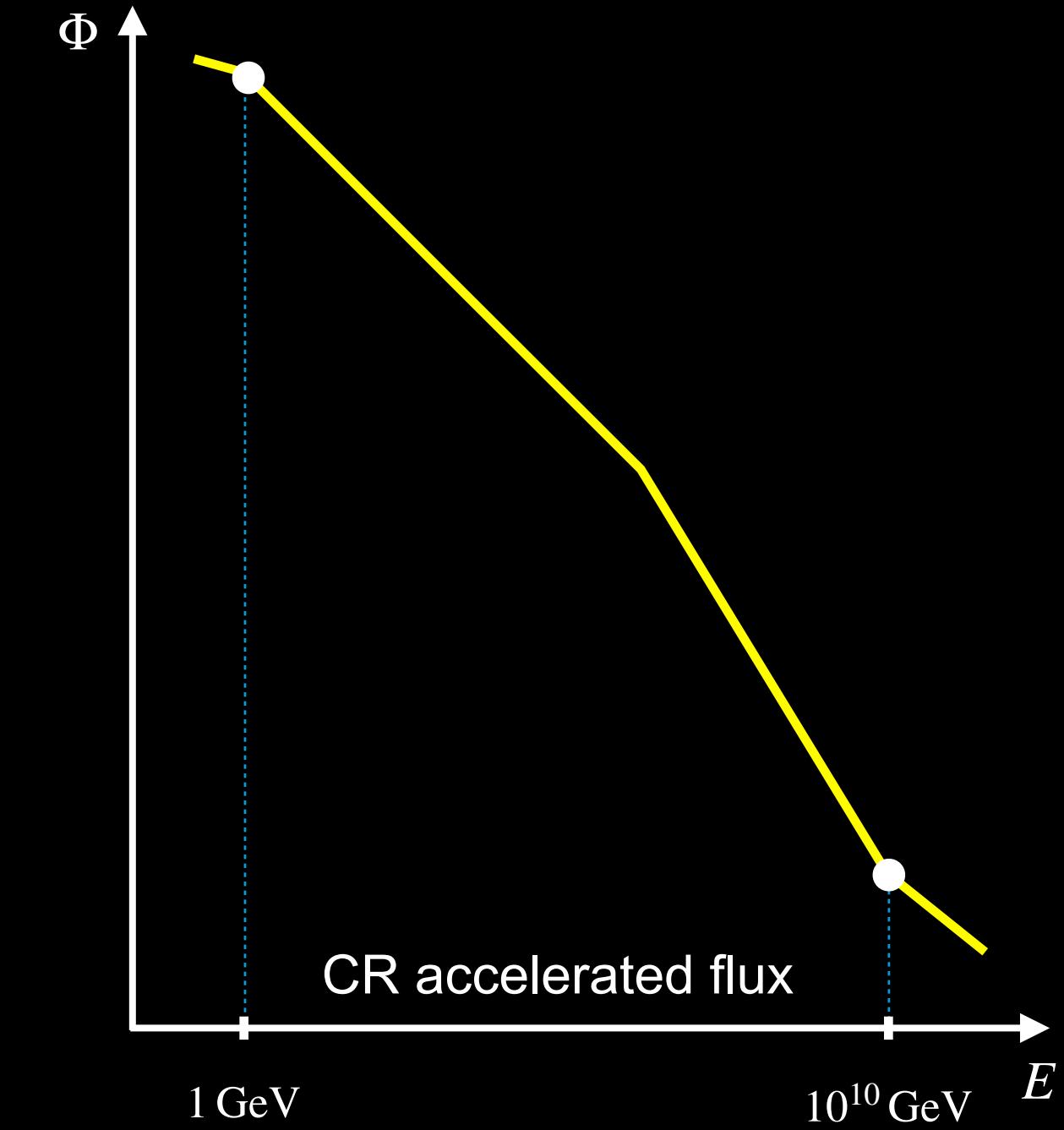
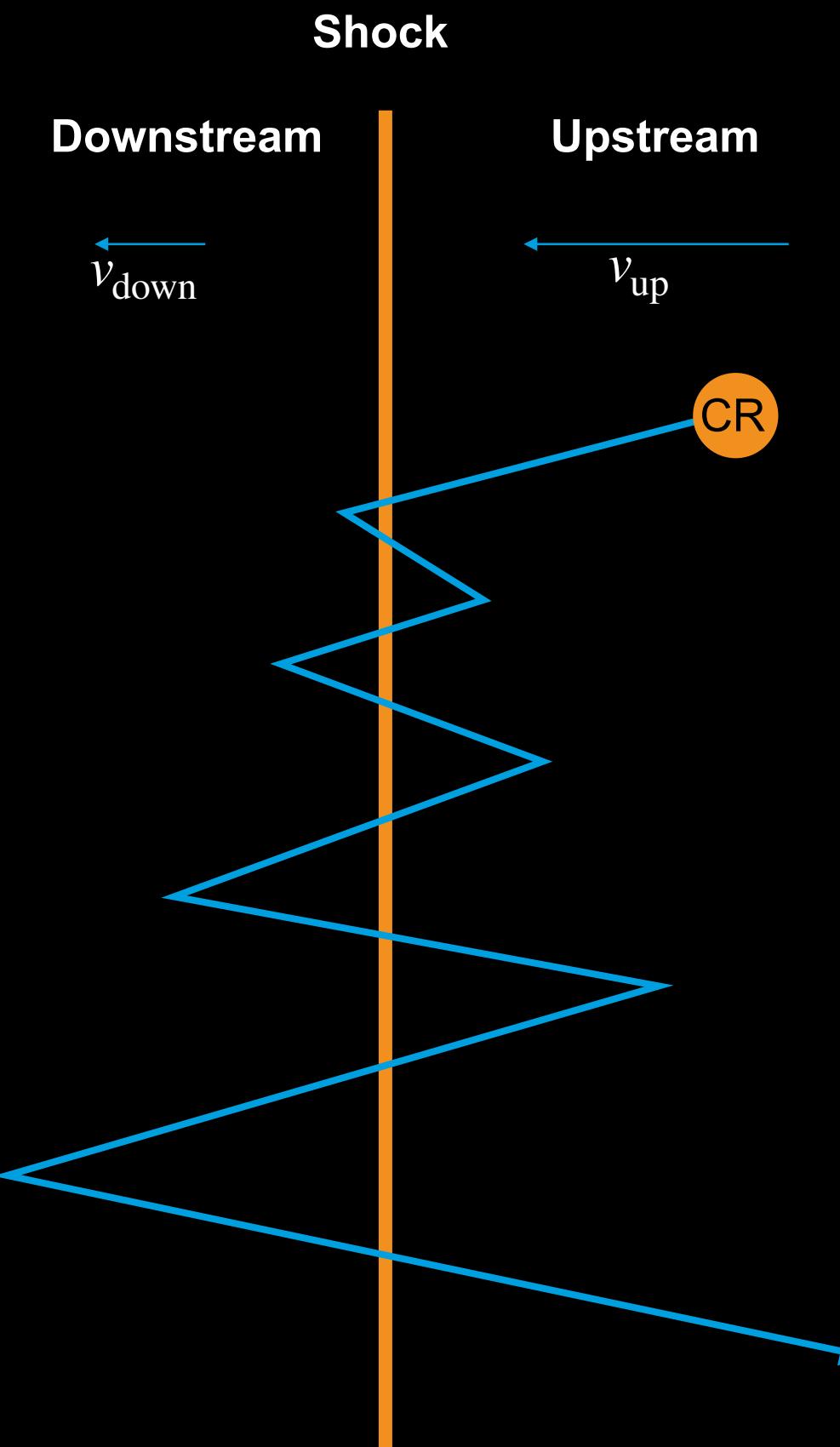
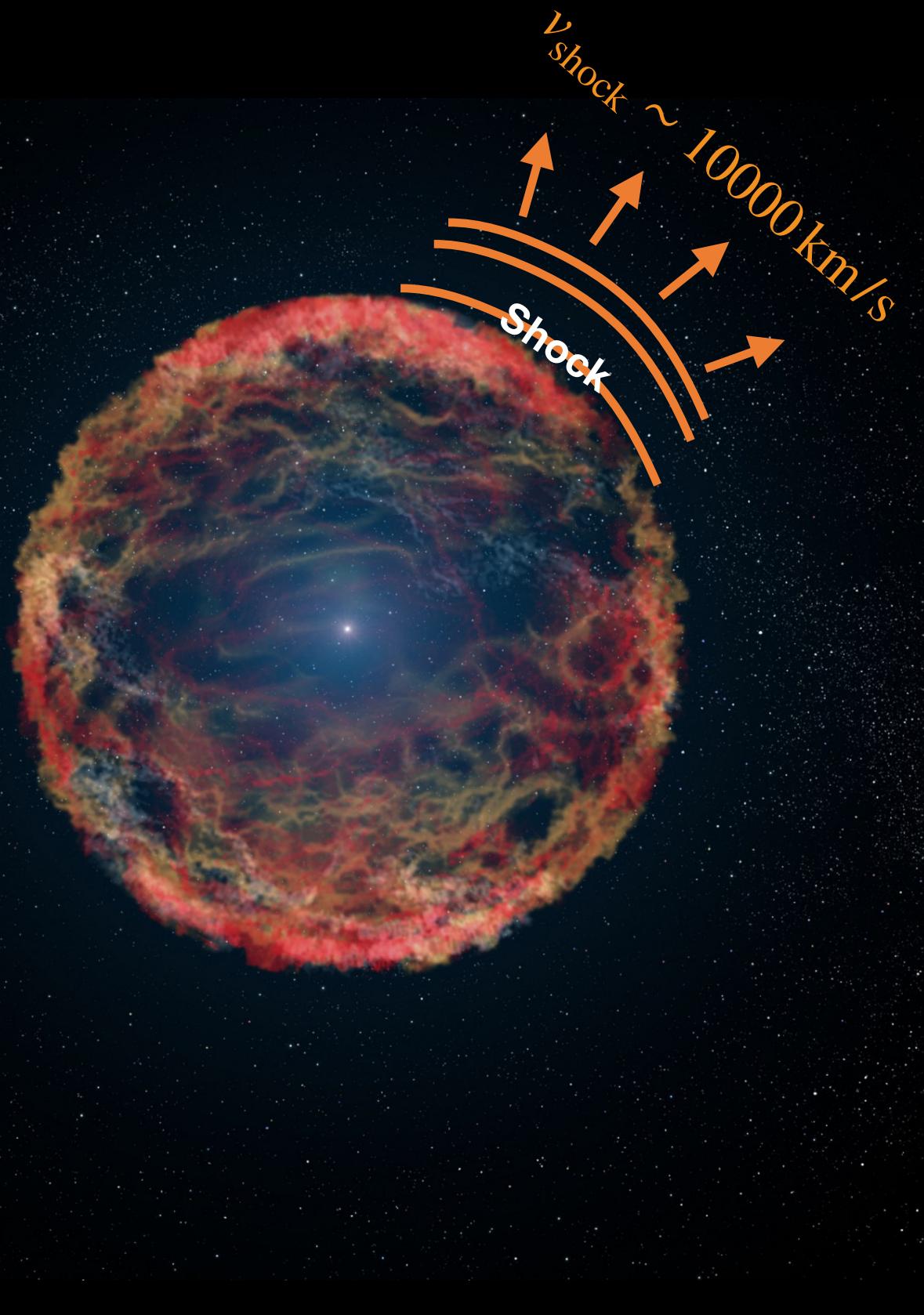
CR production up to high energies



CR production up to high energies



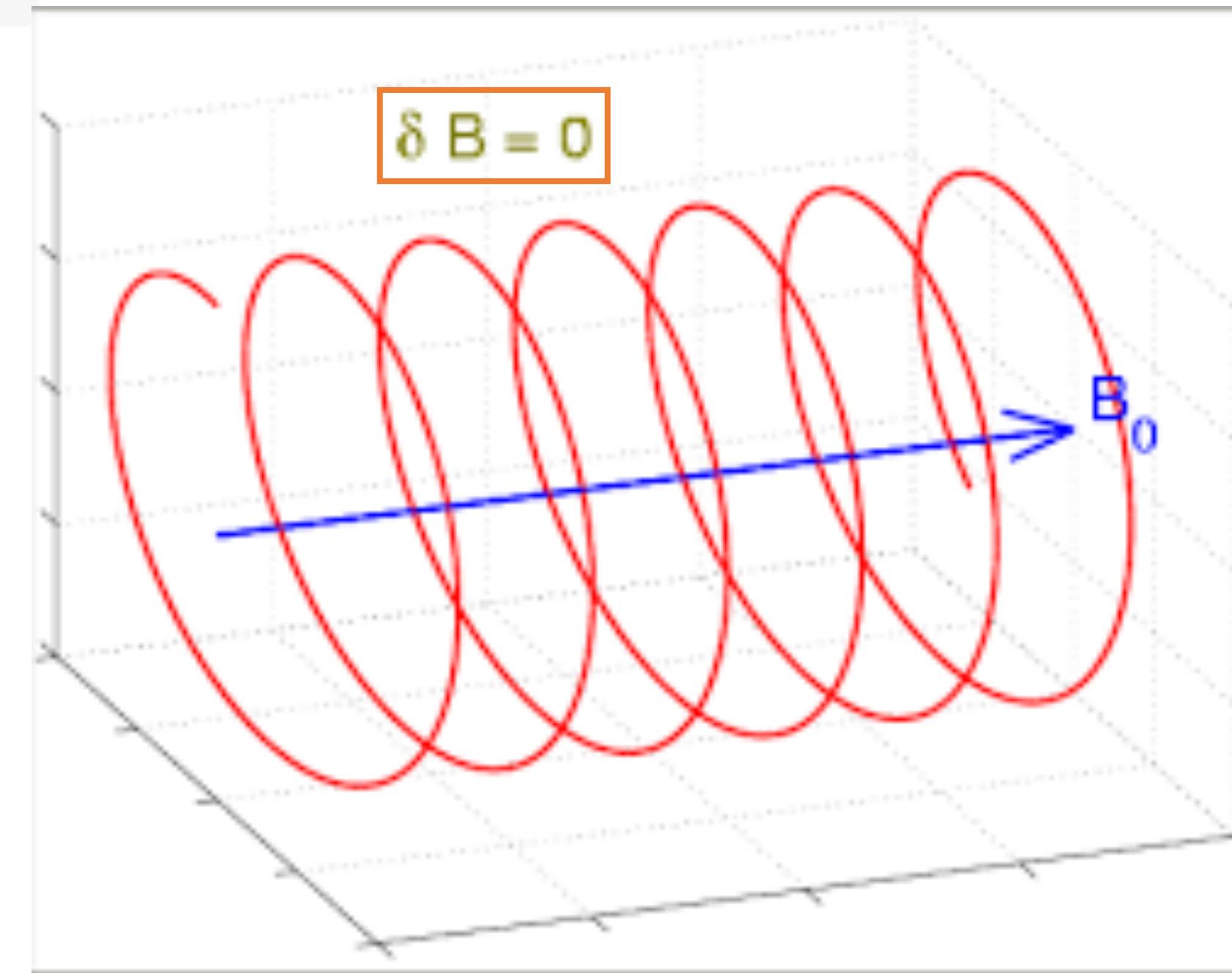
CR production up to high energies



Transport in the Galaxy

Gyro-motion of charged particles

$$m\gamma \frac{d\mathbf{v}}{dt} = \frac{q}{c} (\mathbf{v} \wedge \mathbf{B}_{\text{tot}})$$



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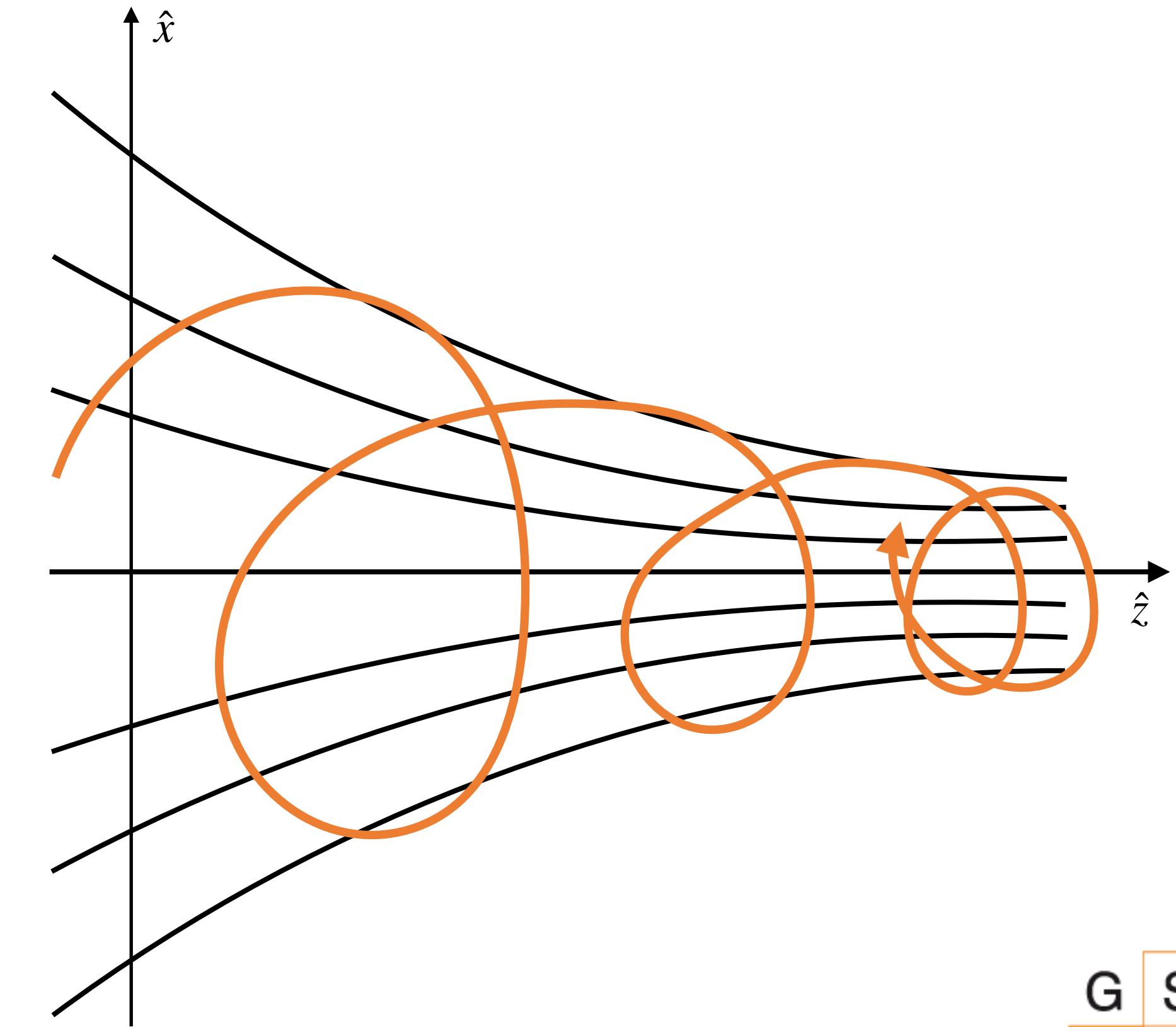
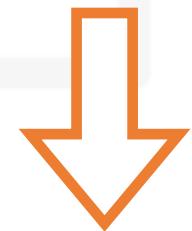
A little complication: magnetic bottles

$$v_0^2 = v_{\parallel,0}^2 + v_{\perp,0}^2$$
$$\underset{\forall z}{=} v_{\parallel}^2 + v_{\perp}^2$$



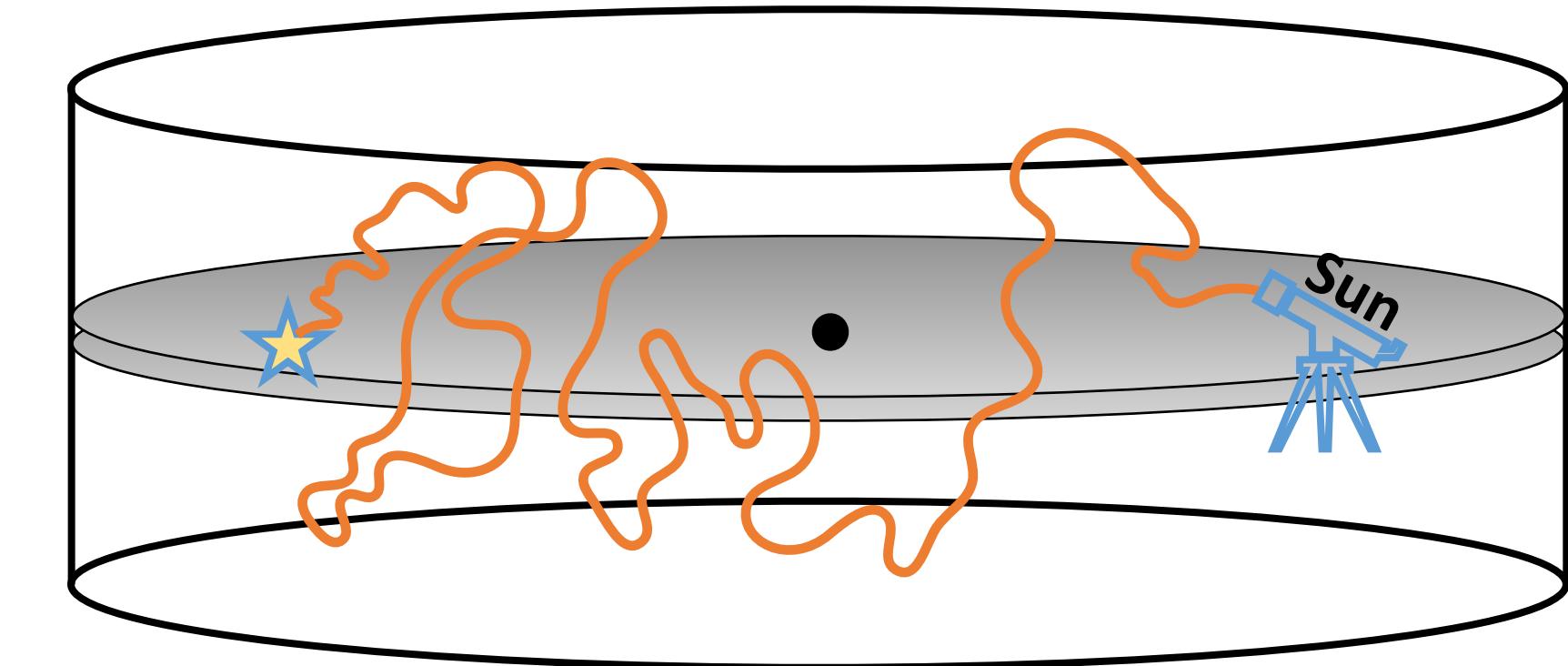
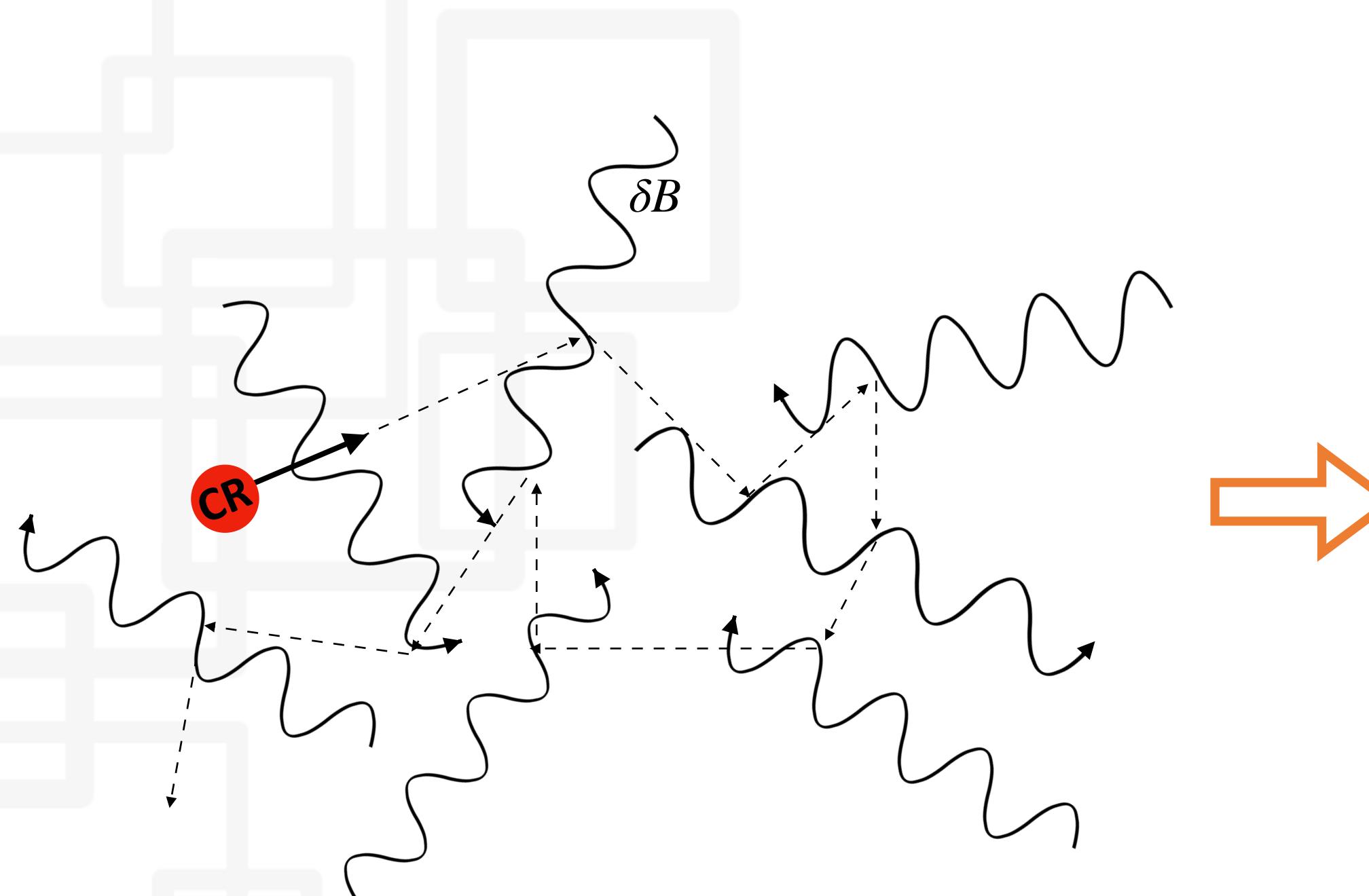
$$\frac{v_{\perp,0}^2}{B} = \underset{\forall z}{\frac{v_{\perp}^2}{B(z)}}$$

$$\Delta(v_{\parallel}^2) = v_{\perp,0}^2 (B_0 - B(z)) \cdot \frac{1}{B_0}$$



G S
S I

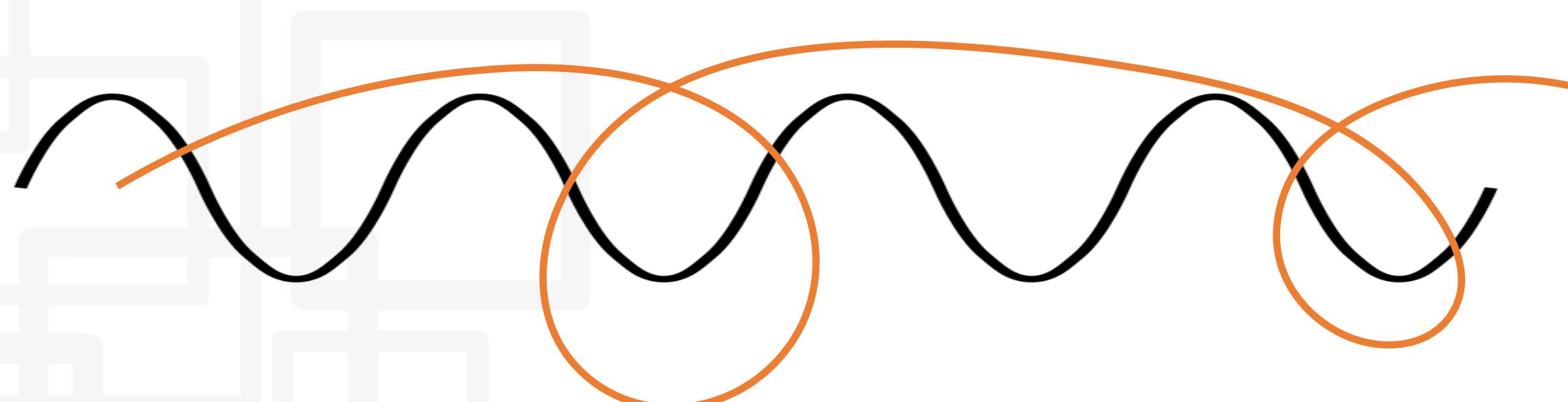
What is CR diffusion?



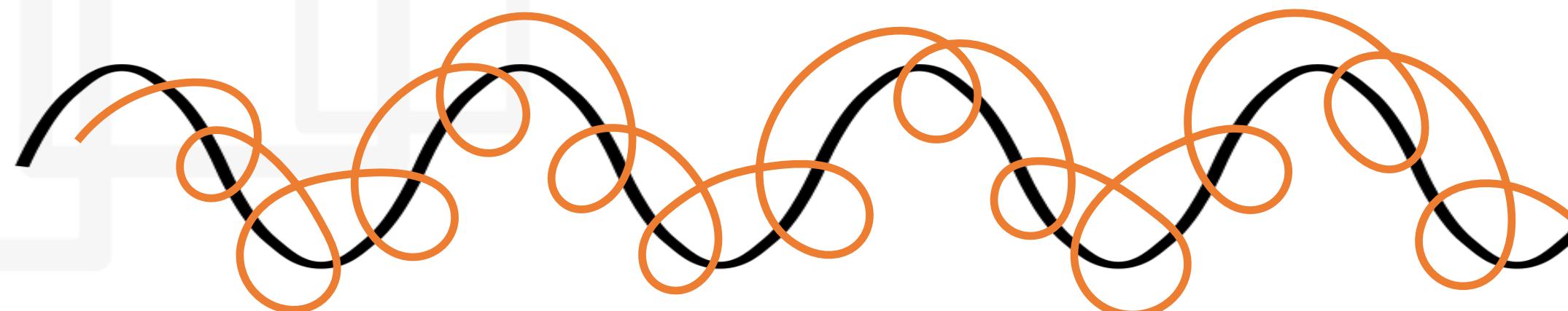
We call “CR diffusion” the random motion resulting from the scattering of charged particles against turbulent magnetic fluctuations.

G S
S I

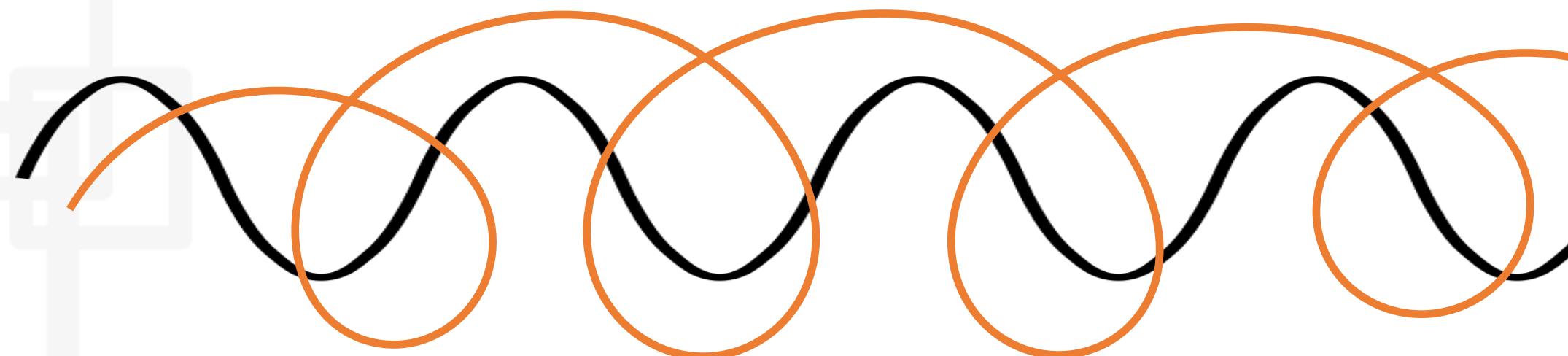
Pitch-angle scattering on B-fluctuations



$$r_L \gg \frac{1}{k_{\text{fluctuation}}} \rightarrow \mathbf{B}_0$$



$$r_L \ll \frac{1}{k_{\text{fluctuation}}} \rightarrow \delta\mathbf{B}$$



$$r_L \sim \frac{1}{k_{\text{fluctuation}}} \rightarrow \mathbf{B}_0 + \delta\mathbf{B}$$

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Generation of turbulence

Generation of turbulence in ISM



G S
S I

Generation of turbulence in ISM



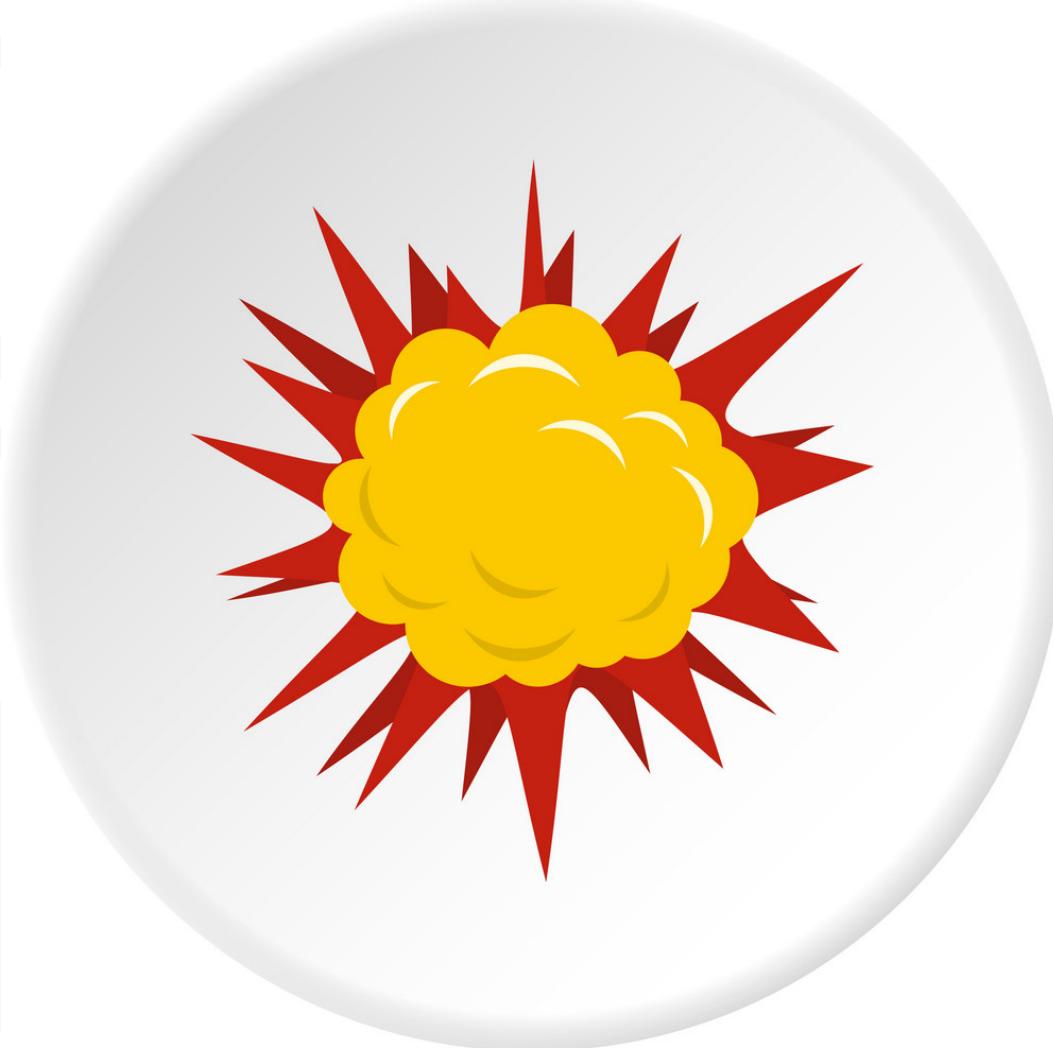
G S
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Turbulent cascade in the inertial range

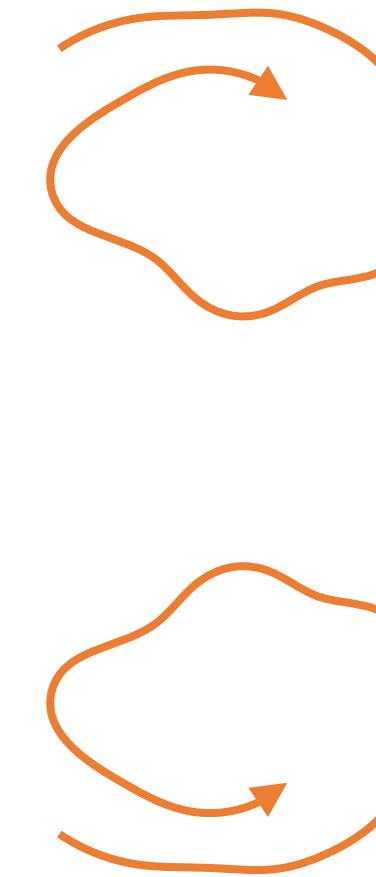


$$L_{\text{inj}} \sim \frac{1}{k_{\text{inj}}}$$

Turbulent cascade in the inertial range

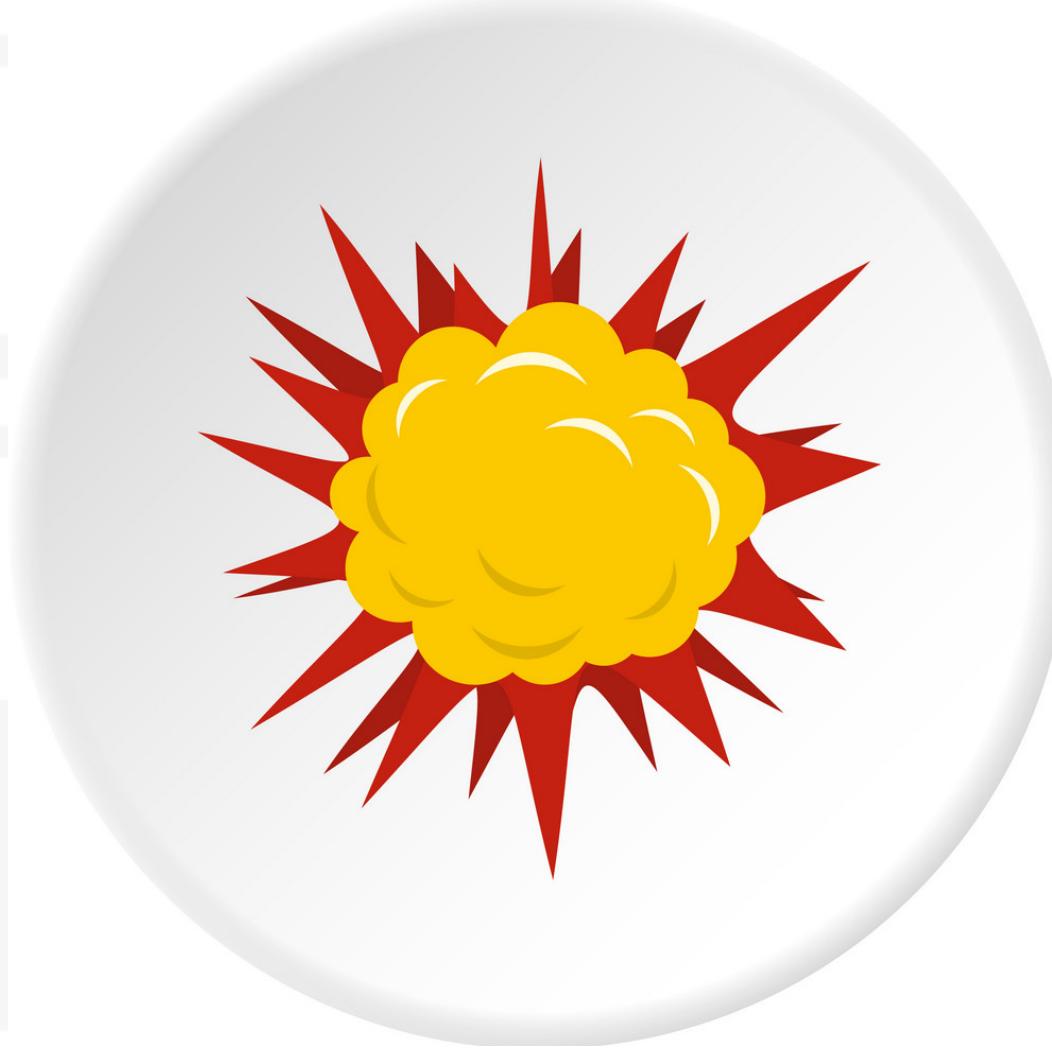


$$L_{\text{inj}} \sim \frac{1}{k_{\text{inj}}}$$



$$\ell_1 \sim \frac{1}{k_1}$$

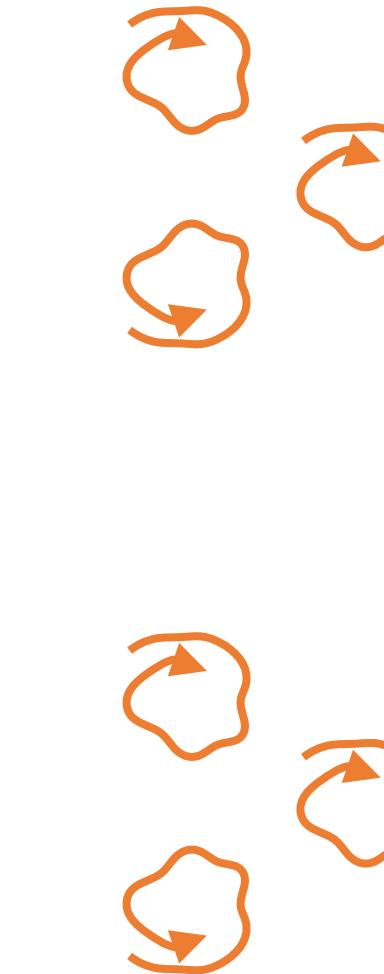
Turbulent cascade in the inertial range



$$L_{\text{inj}} \sim \frac{1}{k_{\text{inj}}}$$

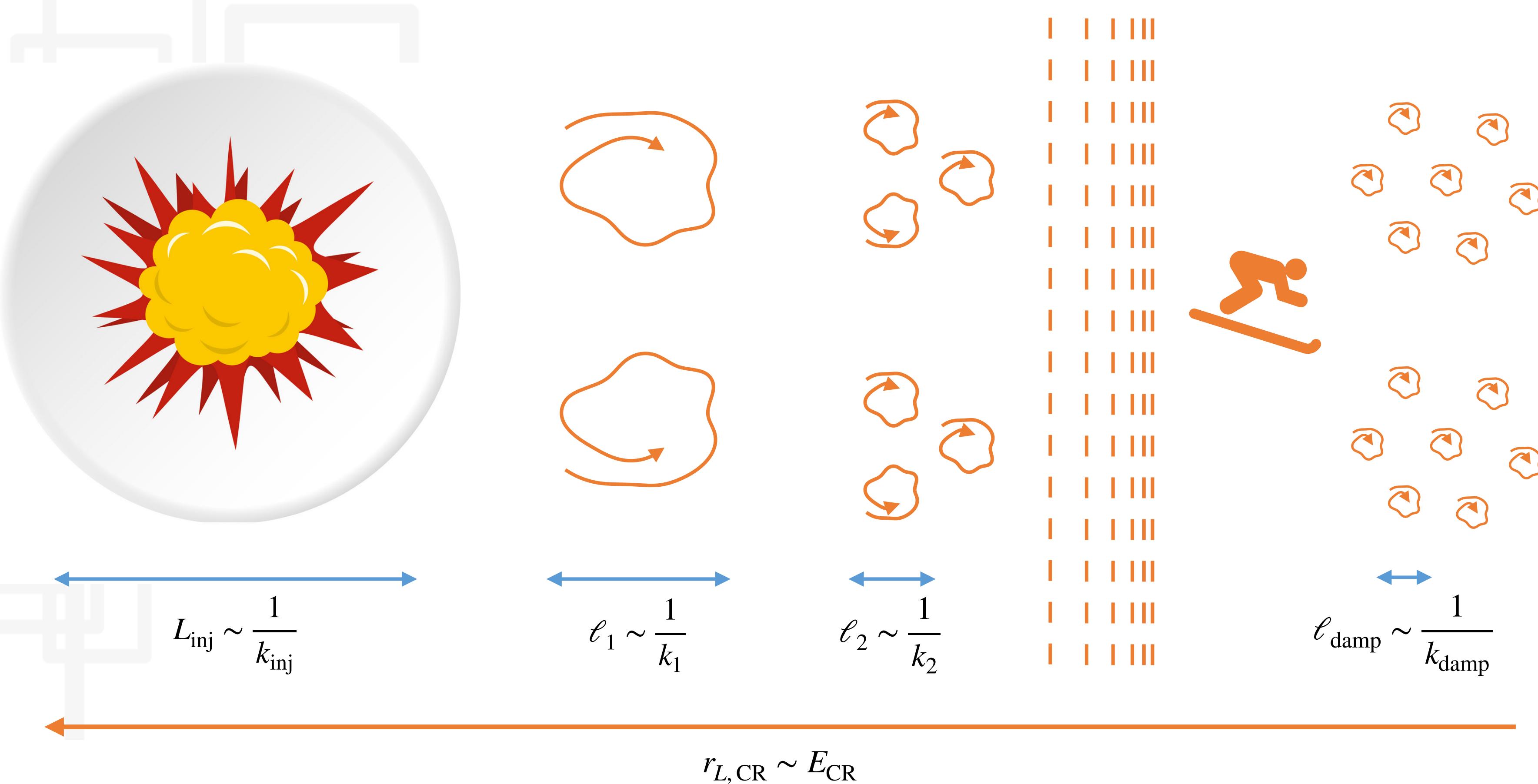


$$\ell_1 \sim \frac{1}{k_1}$$



$$\ell_2 \sim \frac{1}{k_2}$$

Turbulent cascade in the inertial range



MHD decomposition along the cascade

The diagram illustrates the process of decomposing the MHD equations. It starts with a 'Magnetized medium' leading to the 'MHD equations', which then lead to a system of three coupled equations involving δu_x , δu_y , and δu_z .

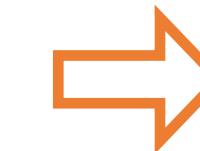
$$\begin{pmatrix} \omega^2 - k^2 v_A^2 - k_\perp^2 c_s^2 & 0 & -k_\perp k_\parallel c_s^2 \\ 0 & \omega^2 - k_\parallel^2 v_A^2 & 0 \\ -k_\perp k_\parallel c_s^2 & 0 & \omega^2 - k_\parallel^2 c_s^2 \end{pmatrix} \begin{pmatrix} \delta u_x \\ \delta u_y \\ \delta u_z \end{pmatrix} = 0$$

MHD decomposition along the cascade

Magnetized medium

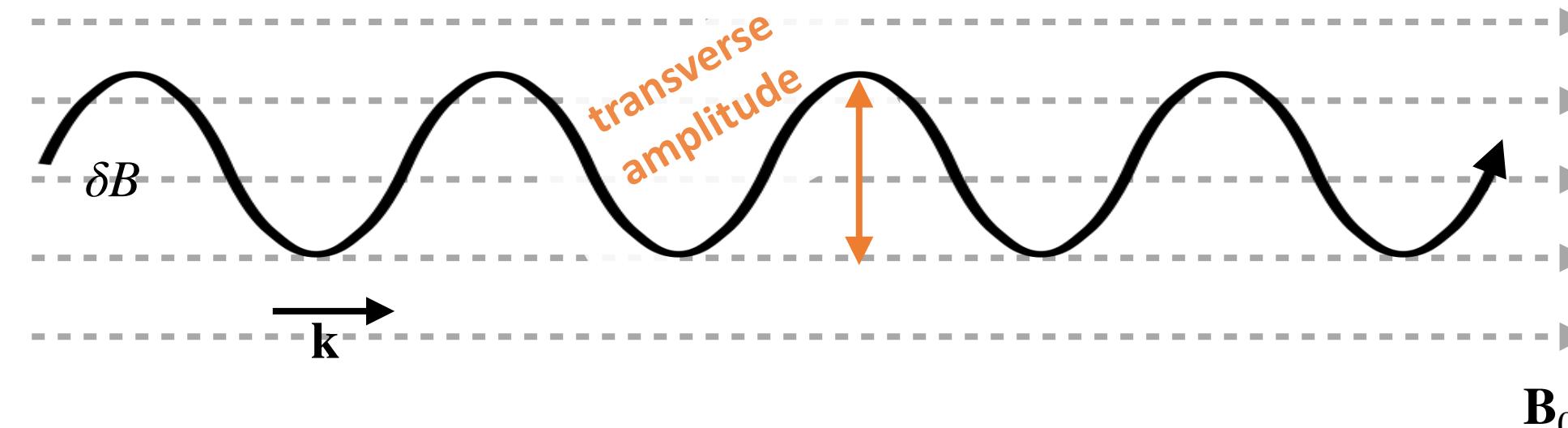


MHD equations



$$\begin{pmatrix} \omega^2 - k^2 v_A^2 - k_\perp^2 c_s^2 & 0 & -k_\perp k_\parallel c_s^2 \\ 0 & \omega^2 - k_\parallel^2 v_A^2 & 0 \\ -k_\perp k_\parallel c_s^2 & 0 & \omega^2 - k_\parallel^2 c_s^2 \end{pmatrix} \begin{pmatrix} \delta u_x \\ \delta u_y \\ \delta u_z \end{pmatrix} = 0$$

- Alfvén modes

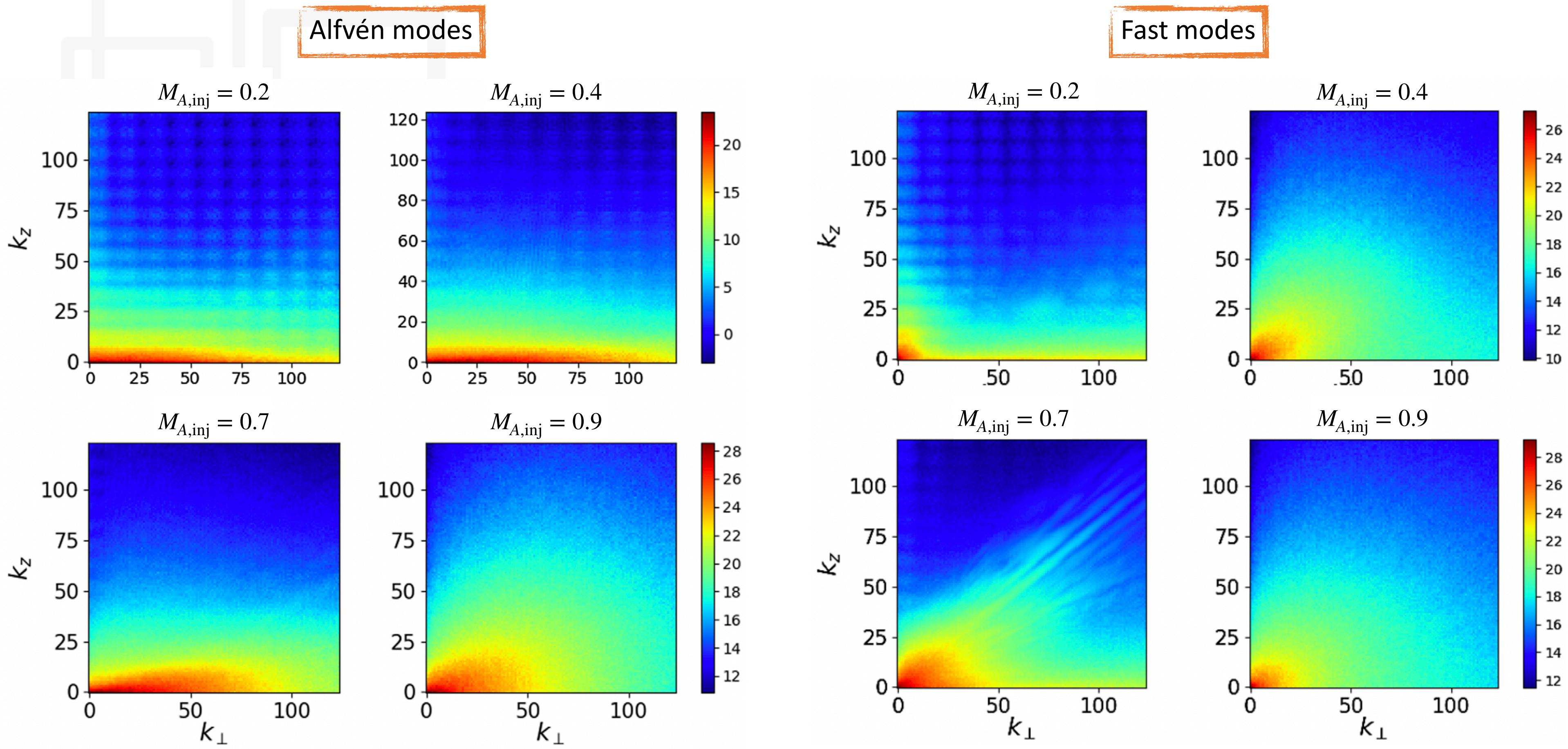


- Magnetosonic modes



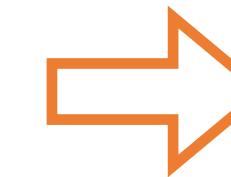
G S
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Distribution of the turbulent power

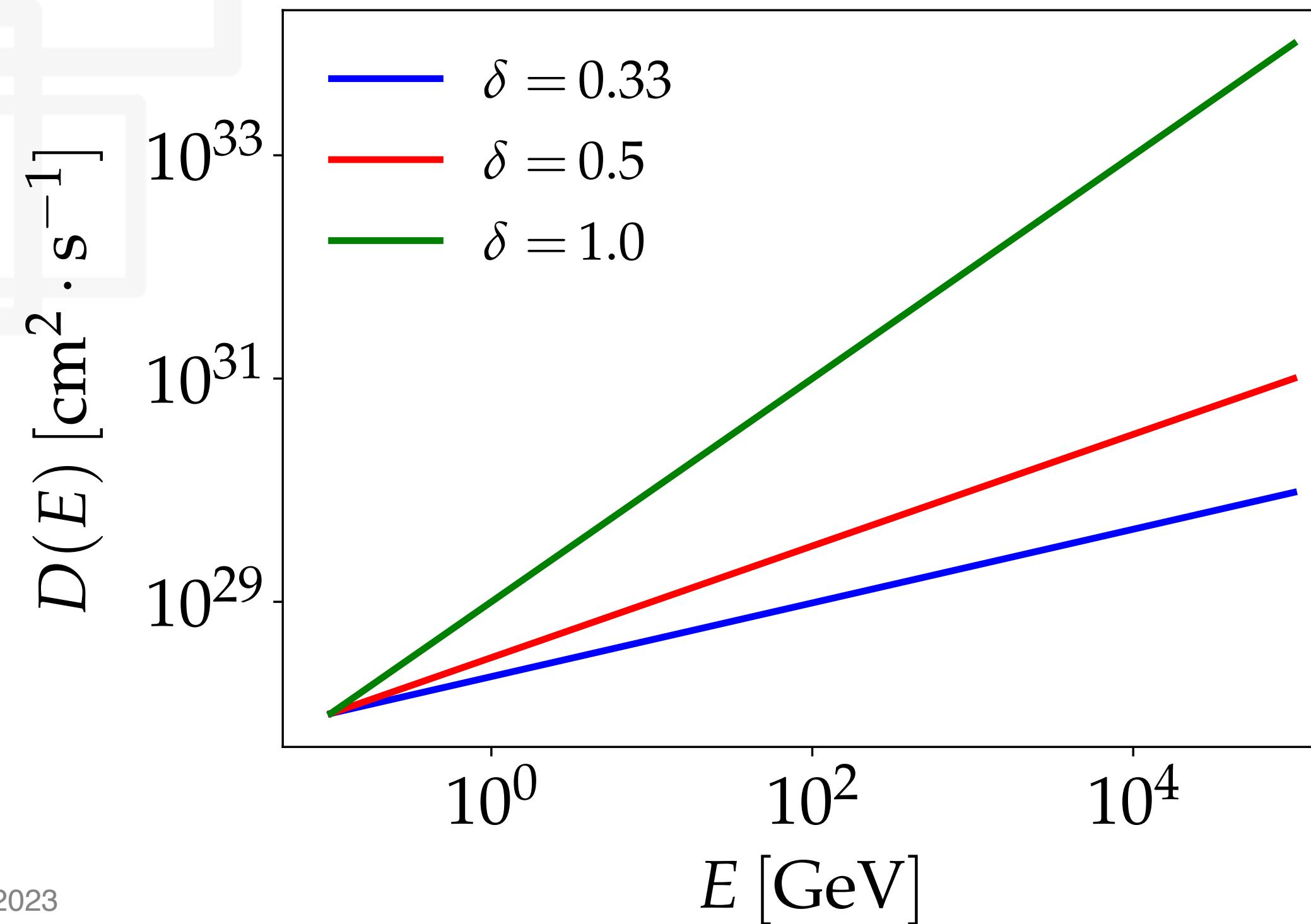


From turbulence to CR diffusion

$$D(E) = \frac{1}{3} \cdot \frac{c r_L}{k_{\text{res}} \cdot E(k_{\text{res}})} \Rightarrow D(E) \sim \frac{r_L}{r_L^{-1} \cdot r_L^{\alpha}}$$

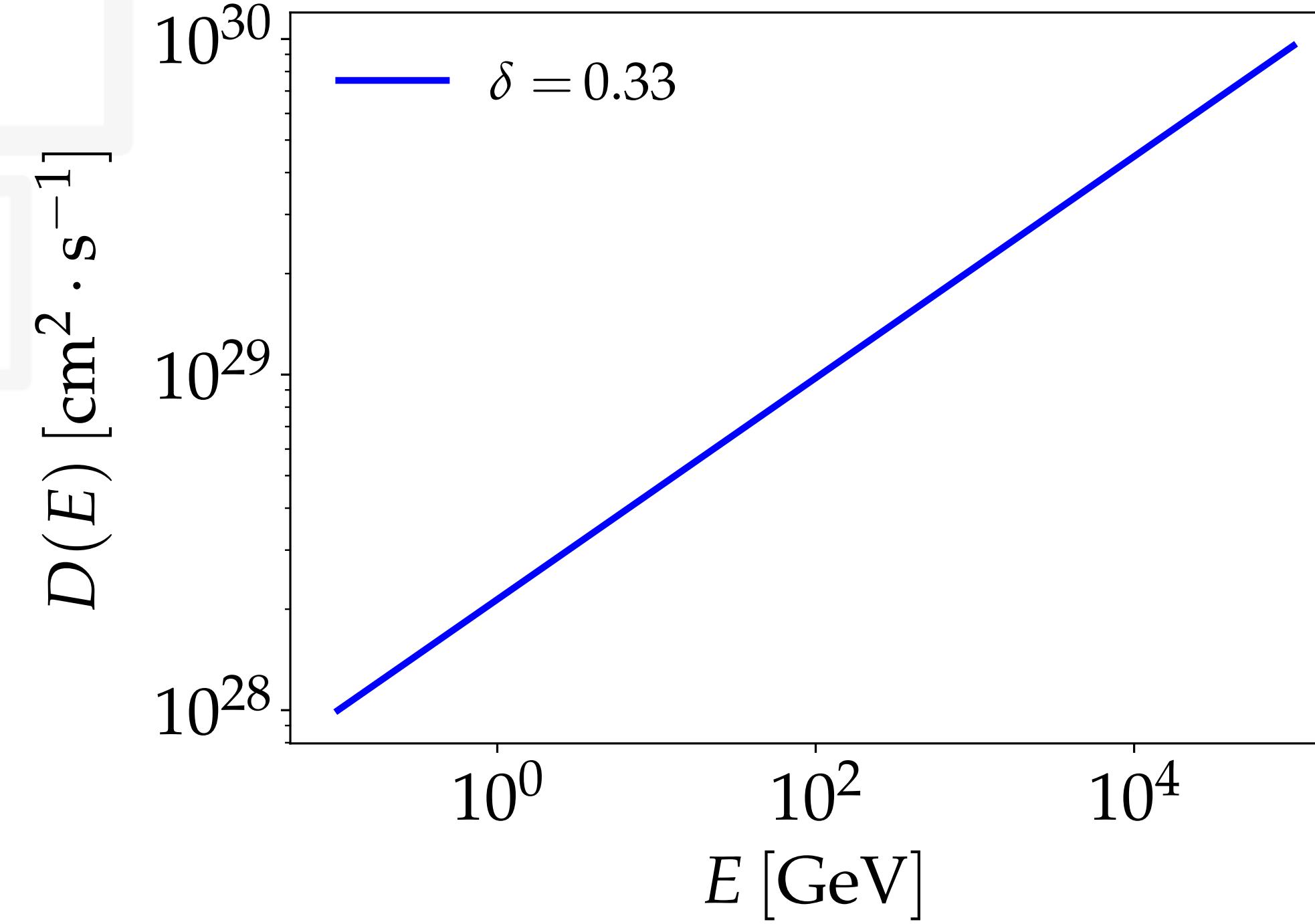


$$D(E) \sim E^{2-\alpha} \equiv E^\delta$$



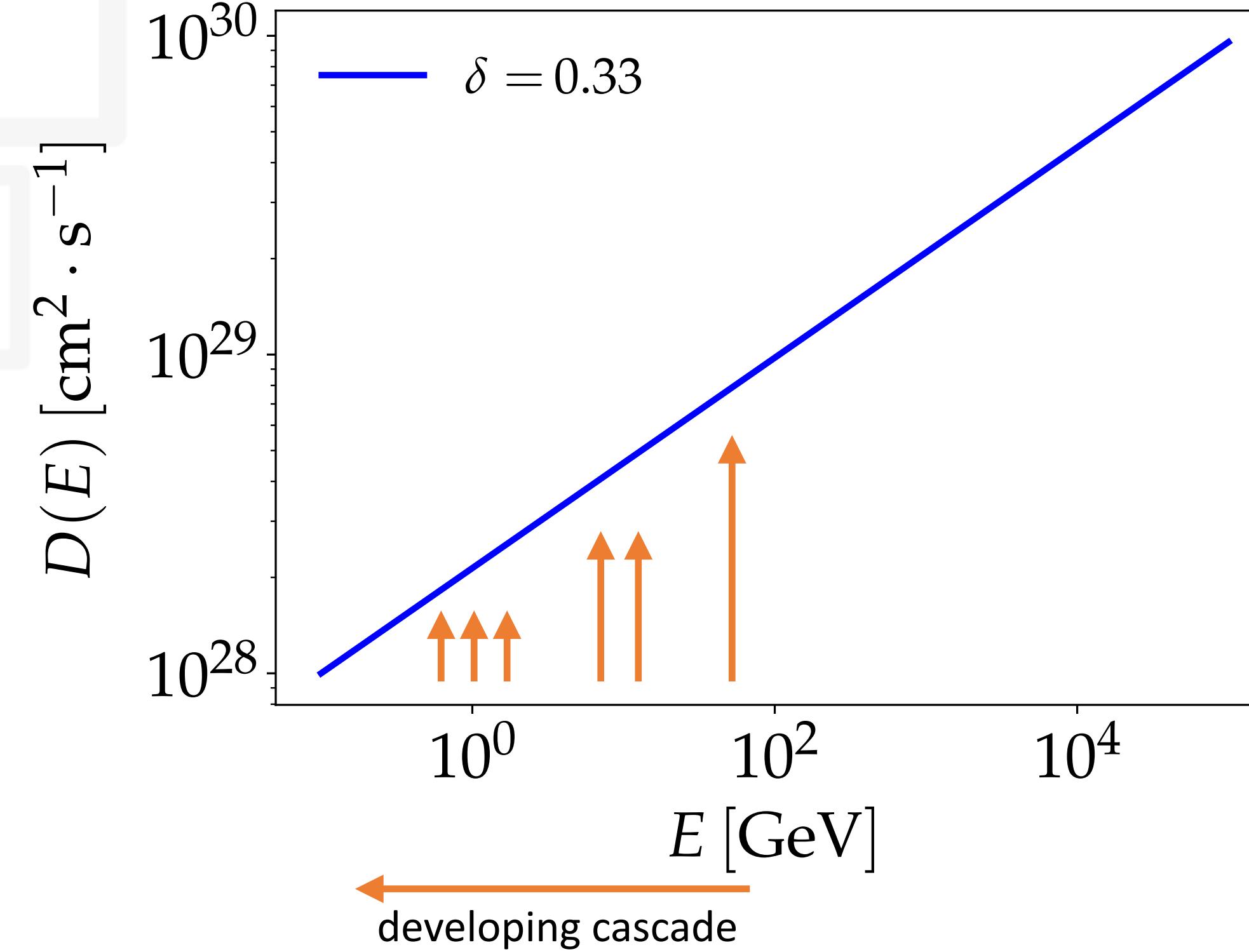
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Damping of turbulence



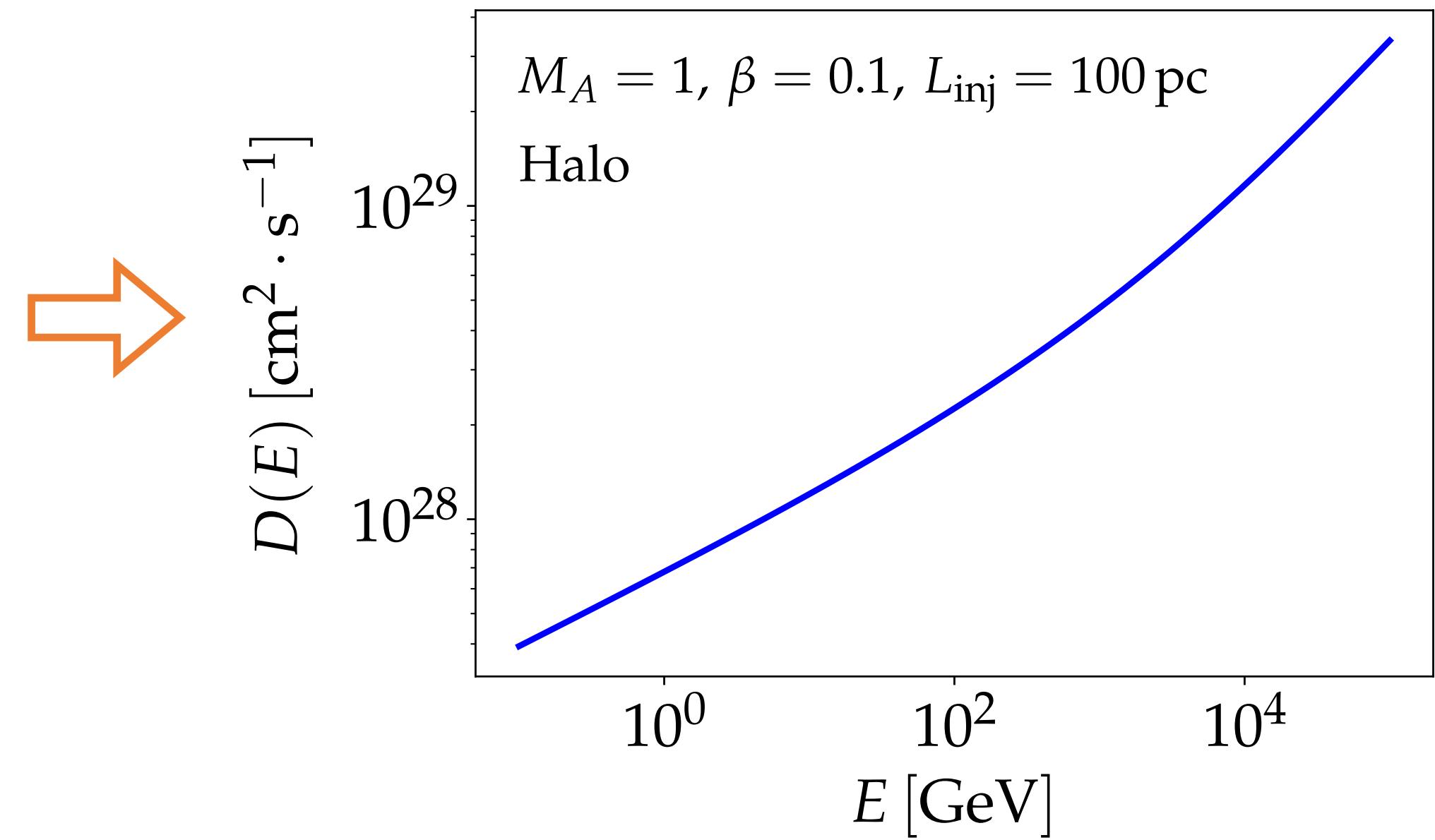
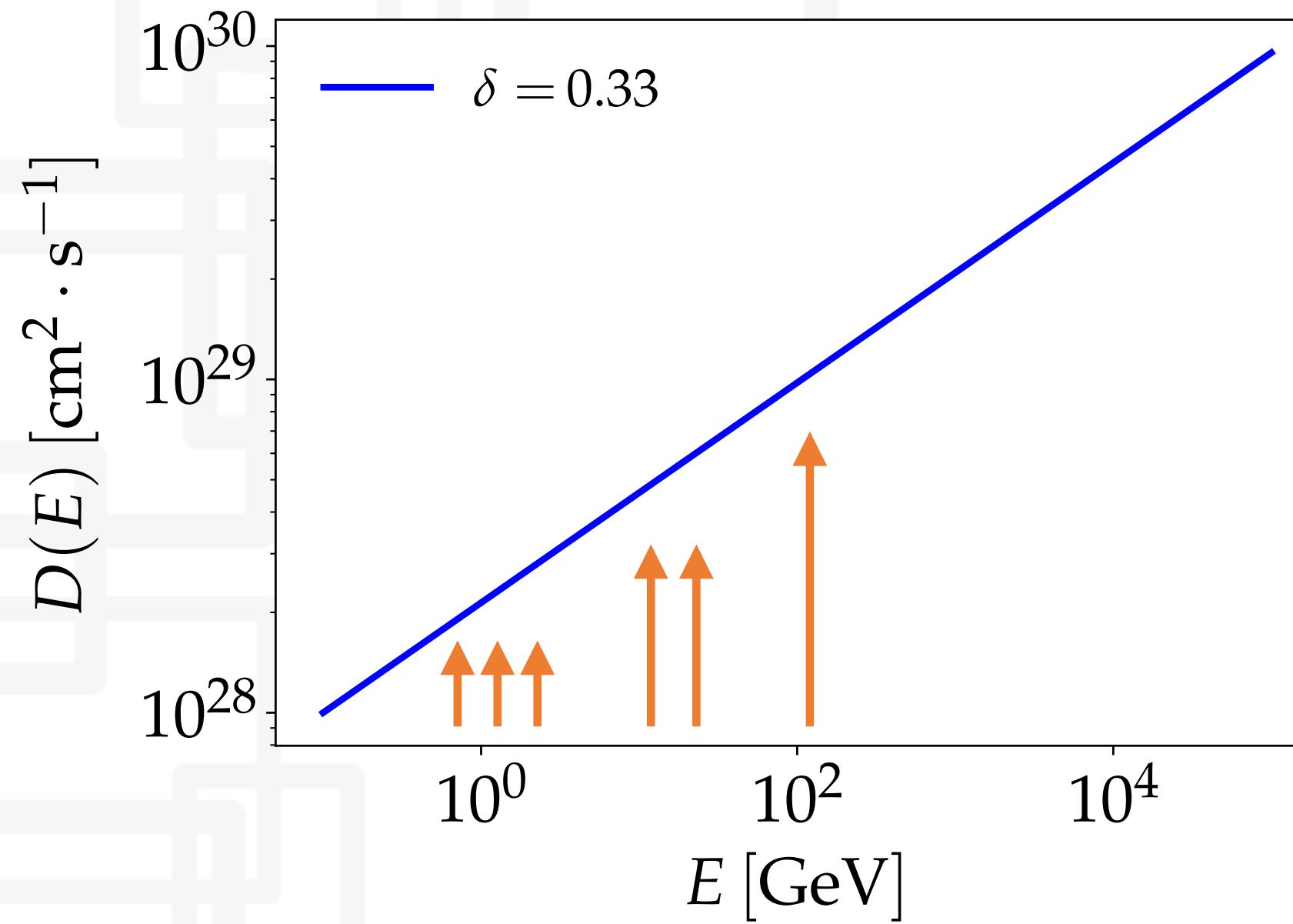
G S
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Damping of turbulence



G S
S I

Damping of turbulence



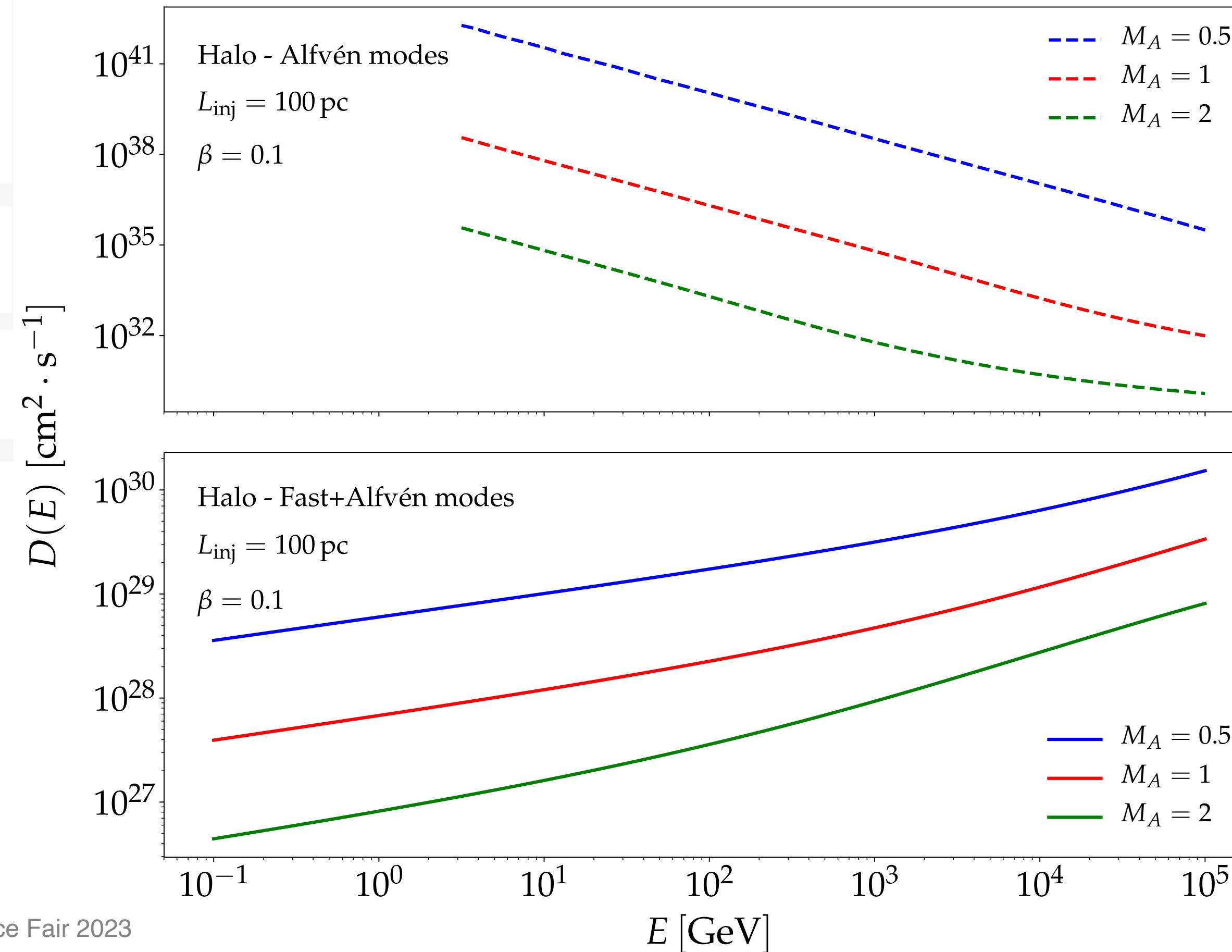
G S
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Conclusion...

Cosmic-ray phenomenology



Resulting CR diffusivity



G S
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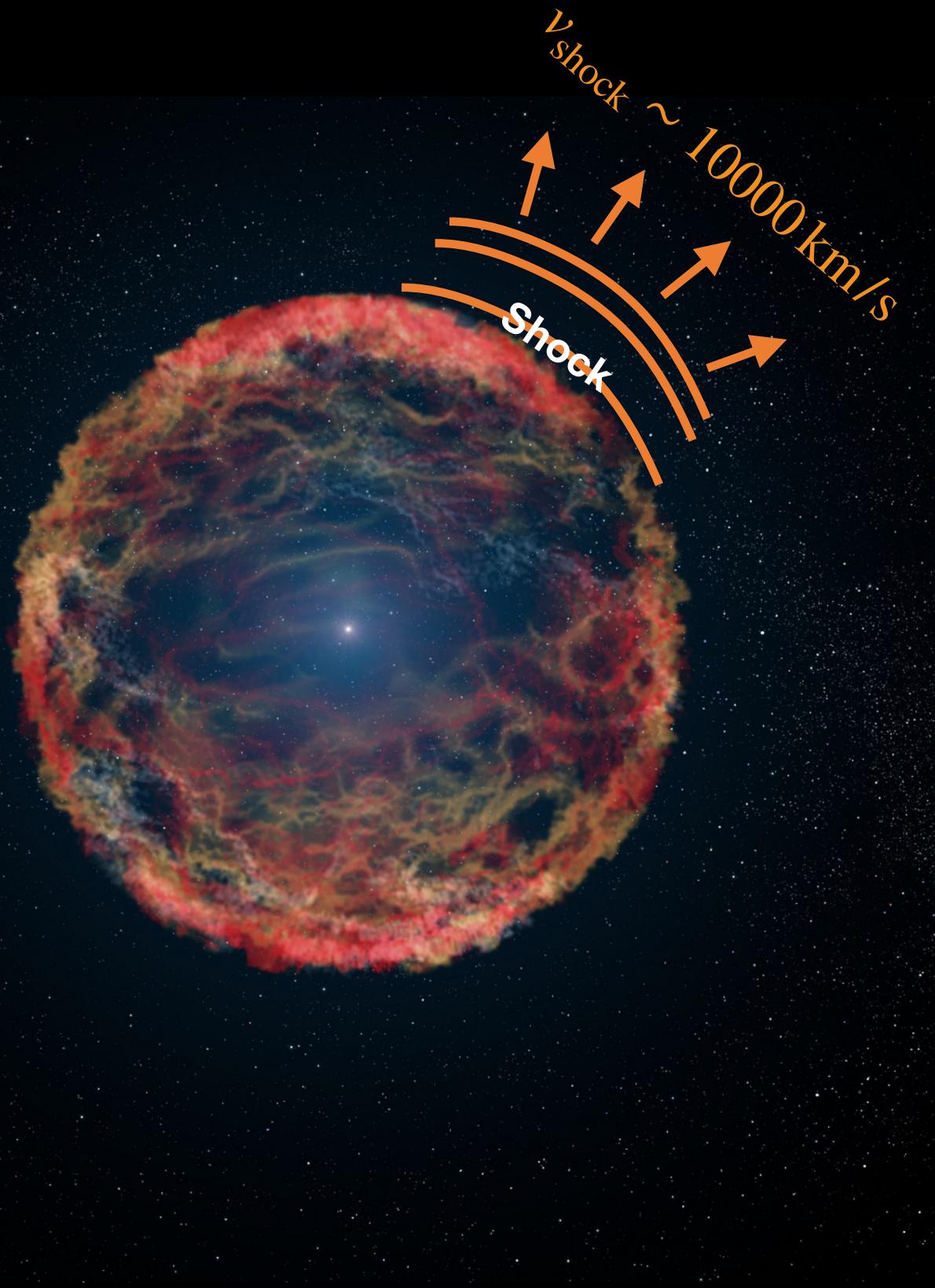
G S
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Take-home message

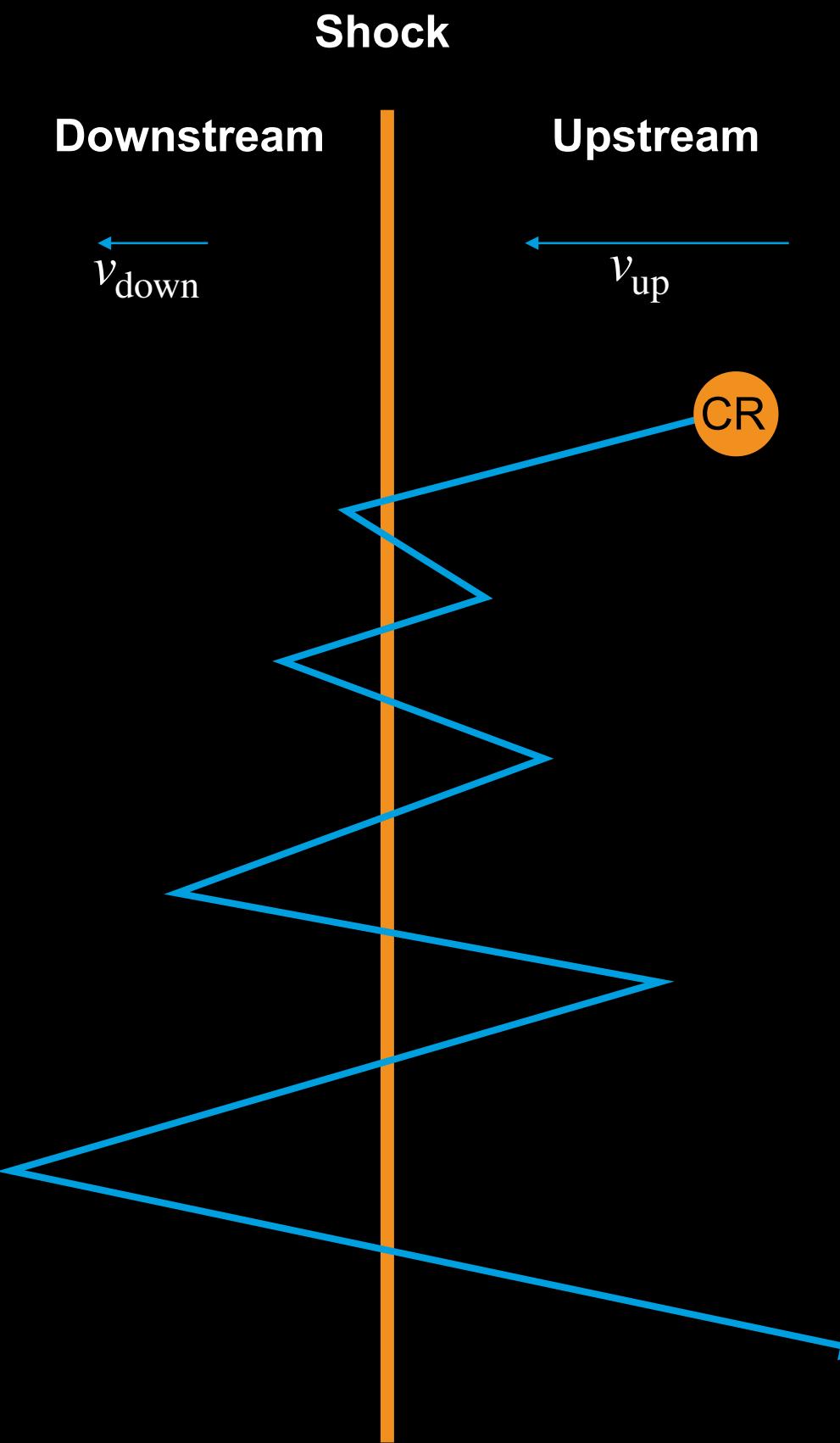
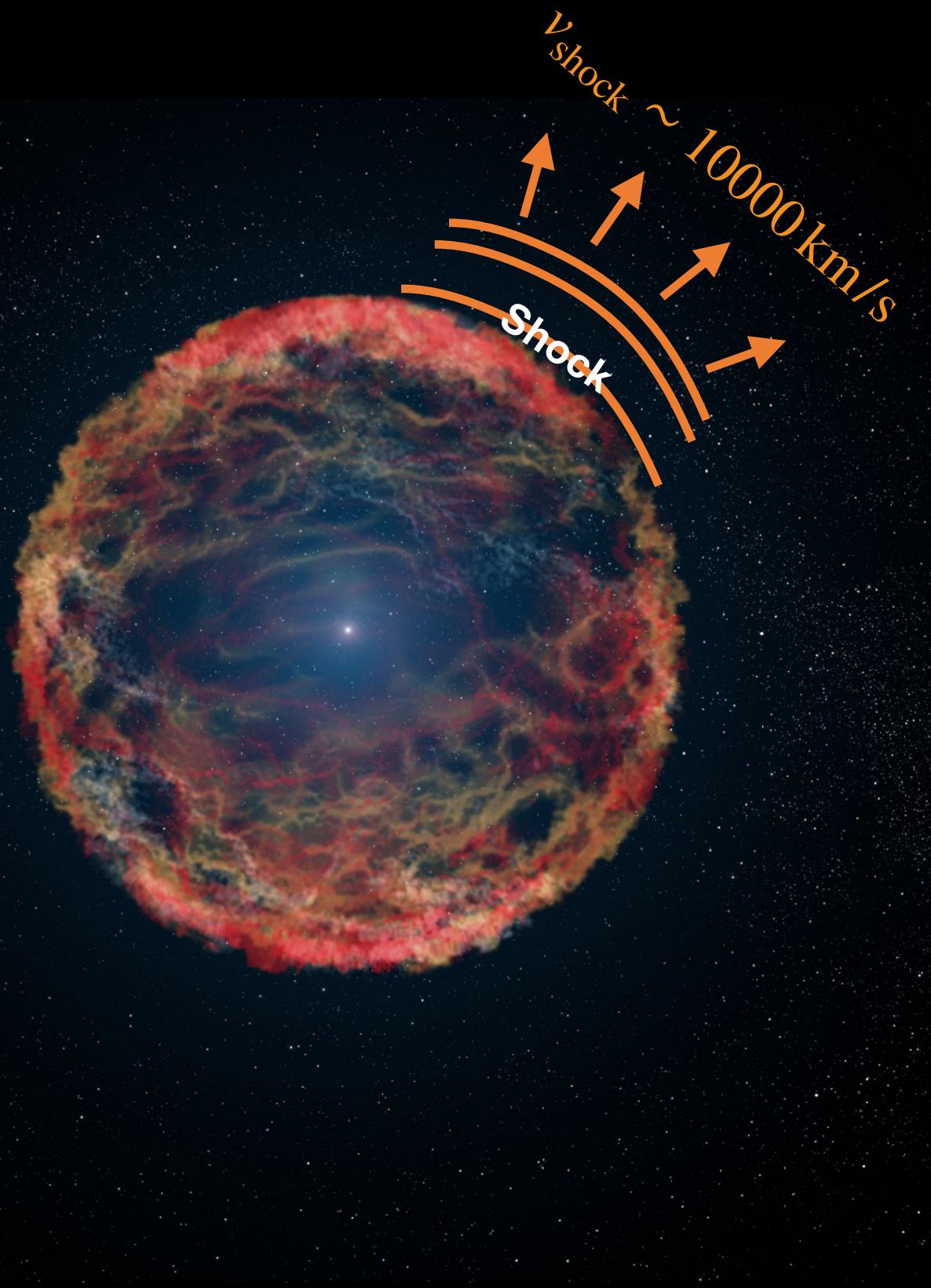
*Accurate measurements
require detailed knowledge of
the microphysics of CR
transport in our Galaxy.*

Backup slides

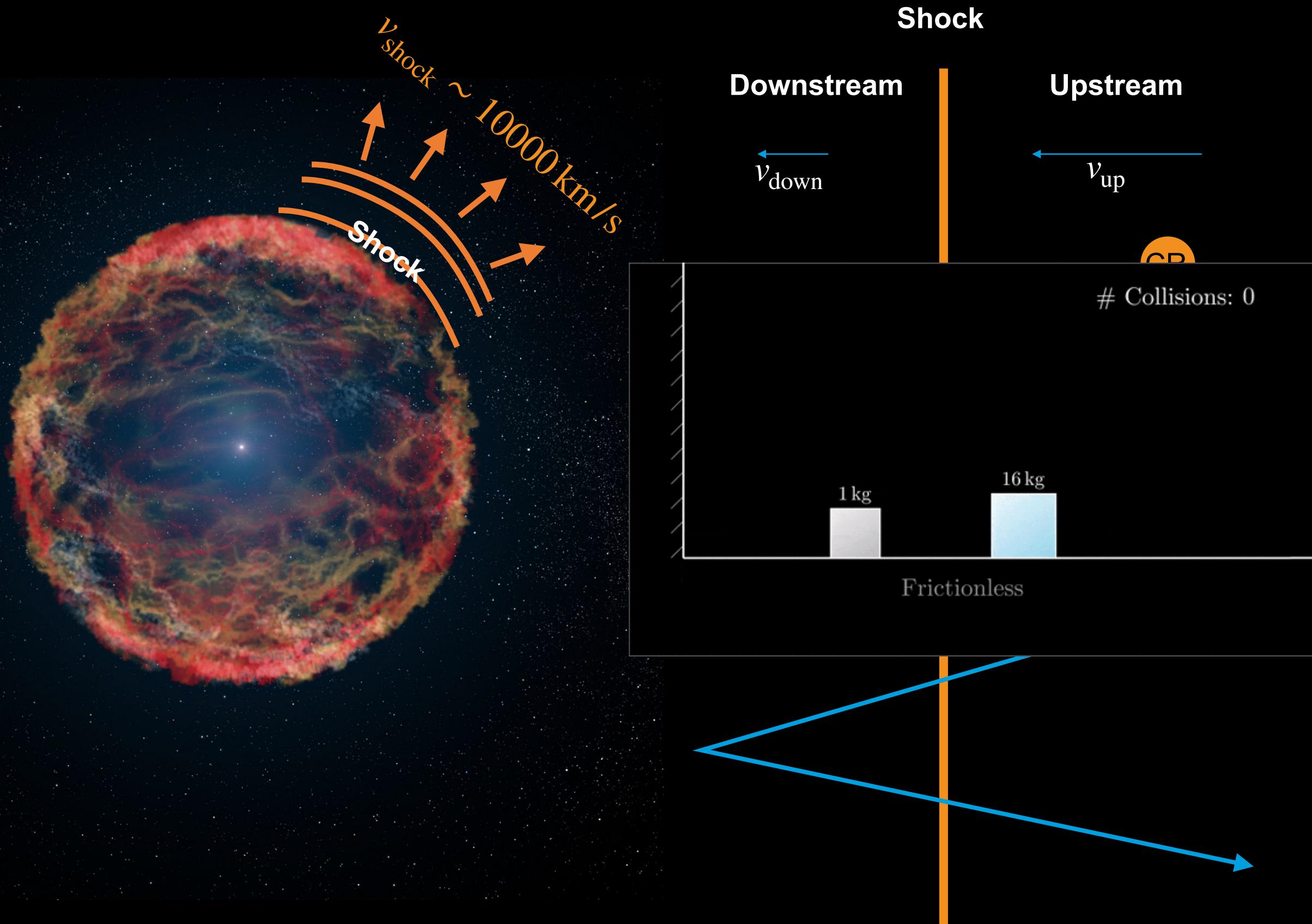
CR production up to high energies



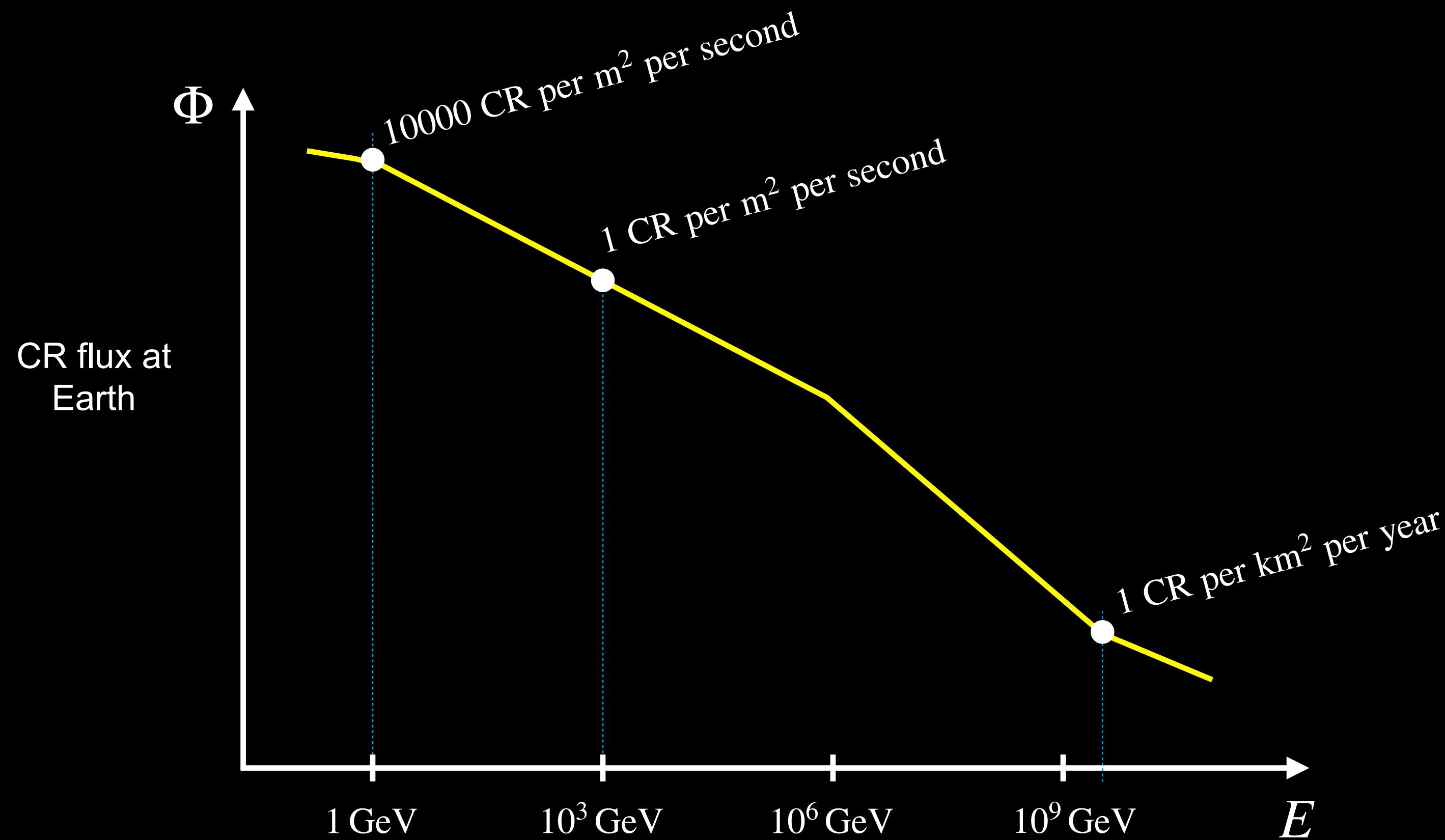
CR production up to high energies



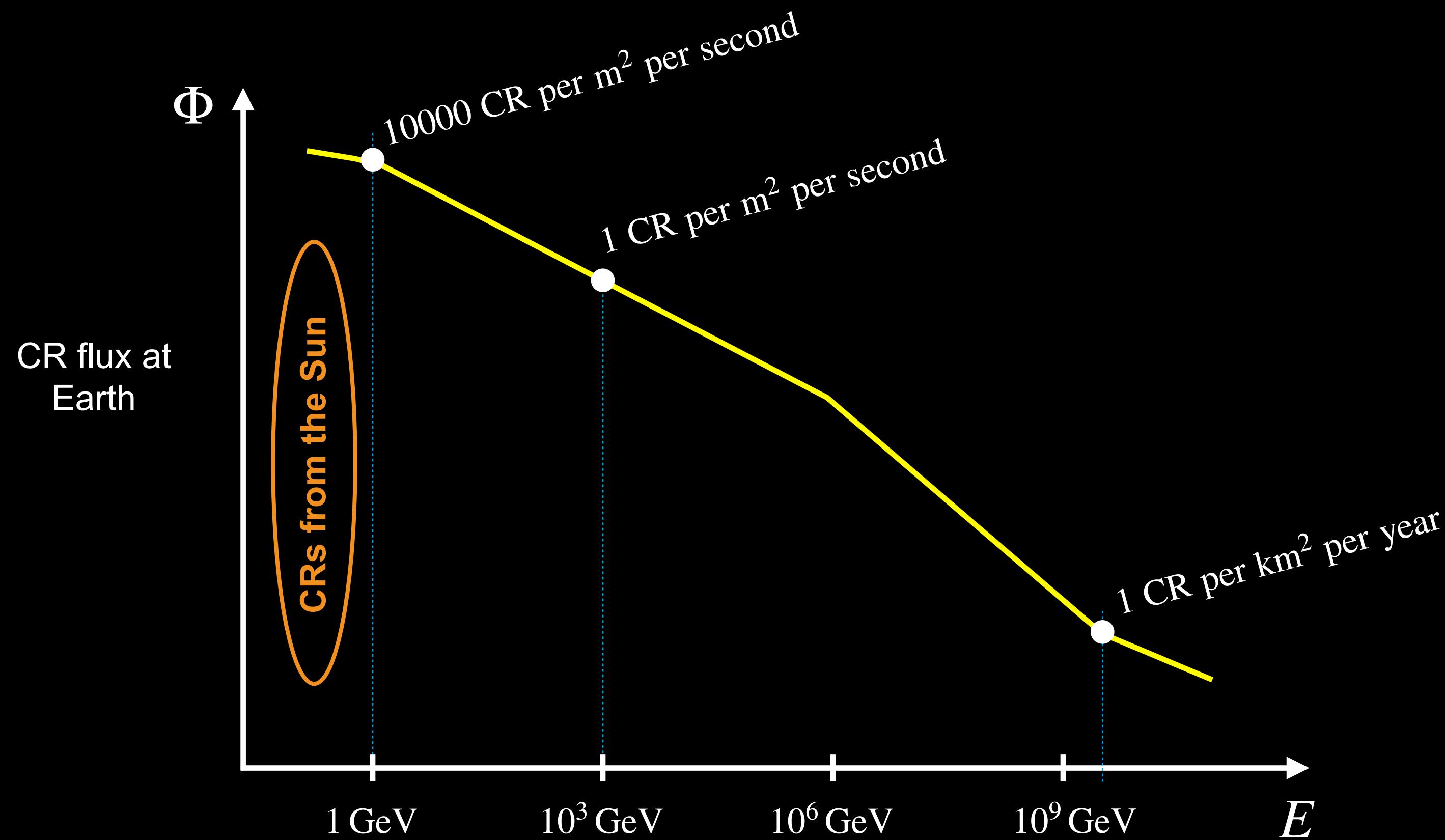
CR production up to high energies



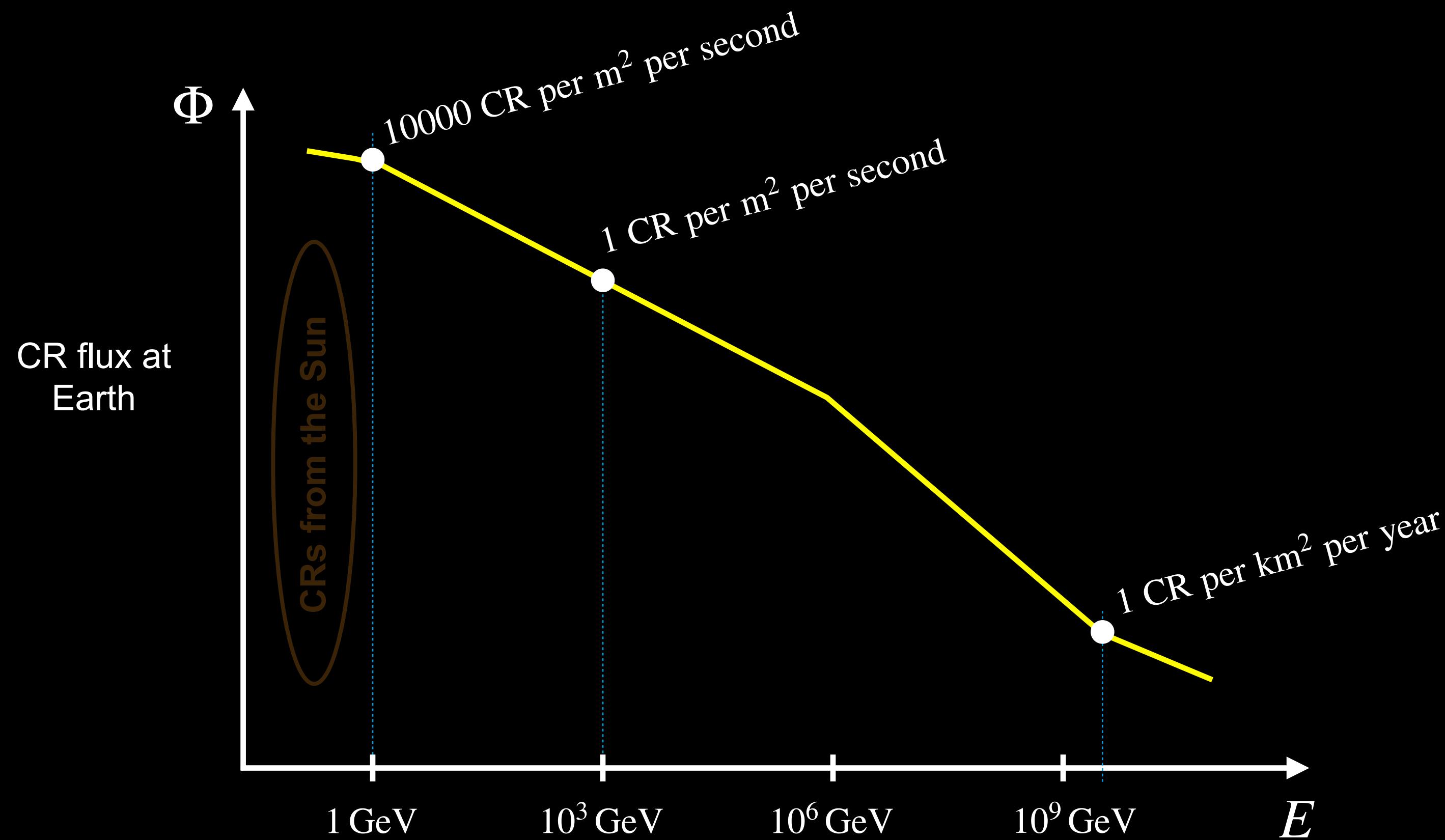
Resulting CR spectrum

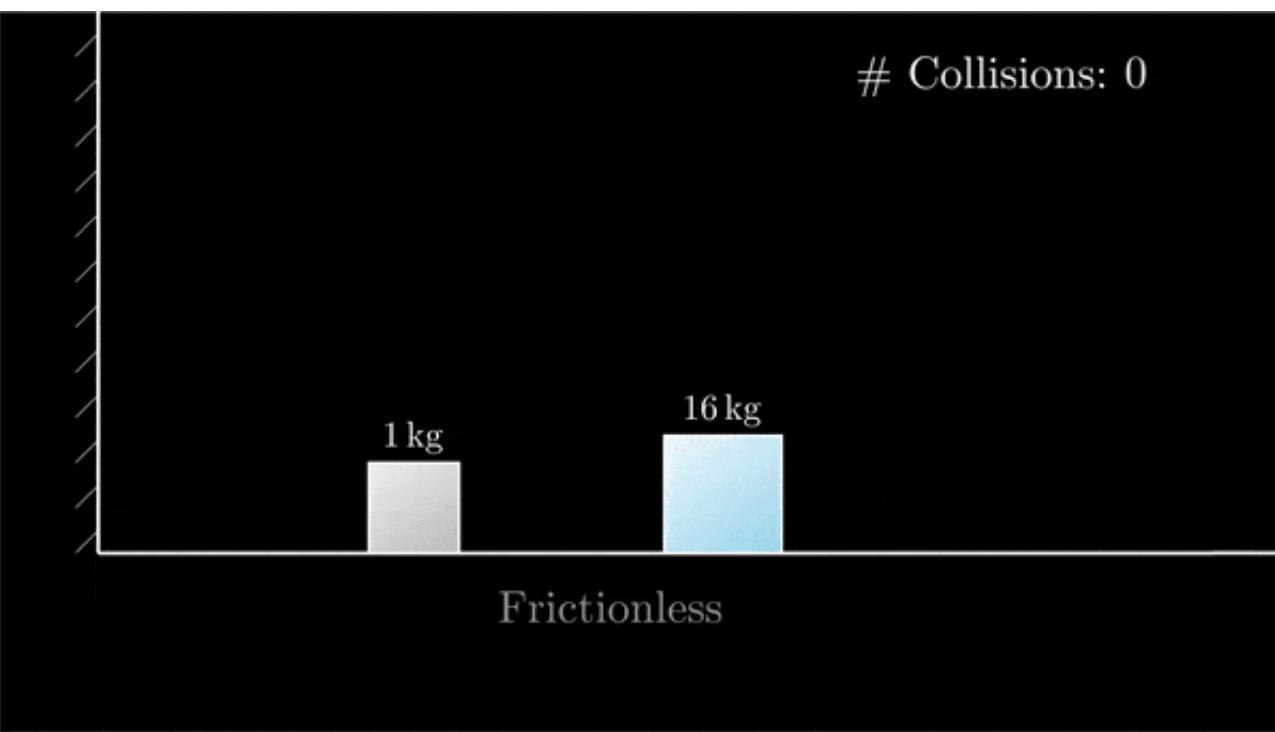


Resulting CR spectrum

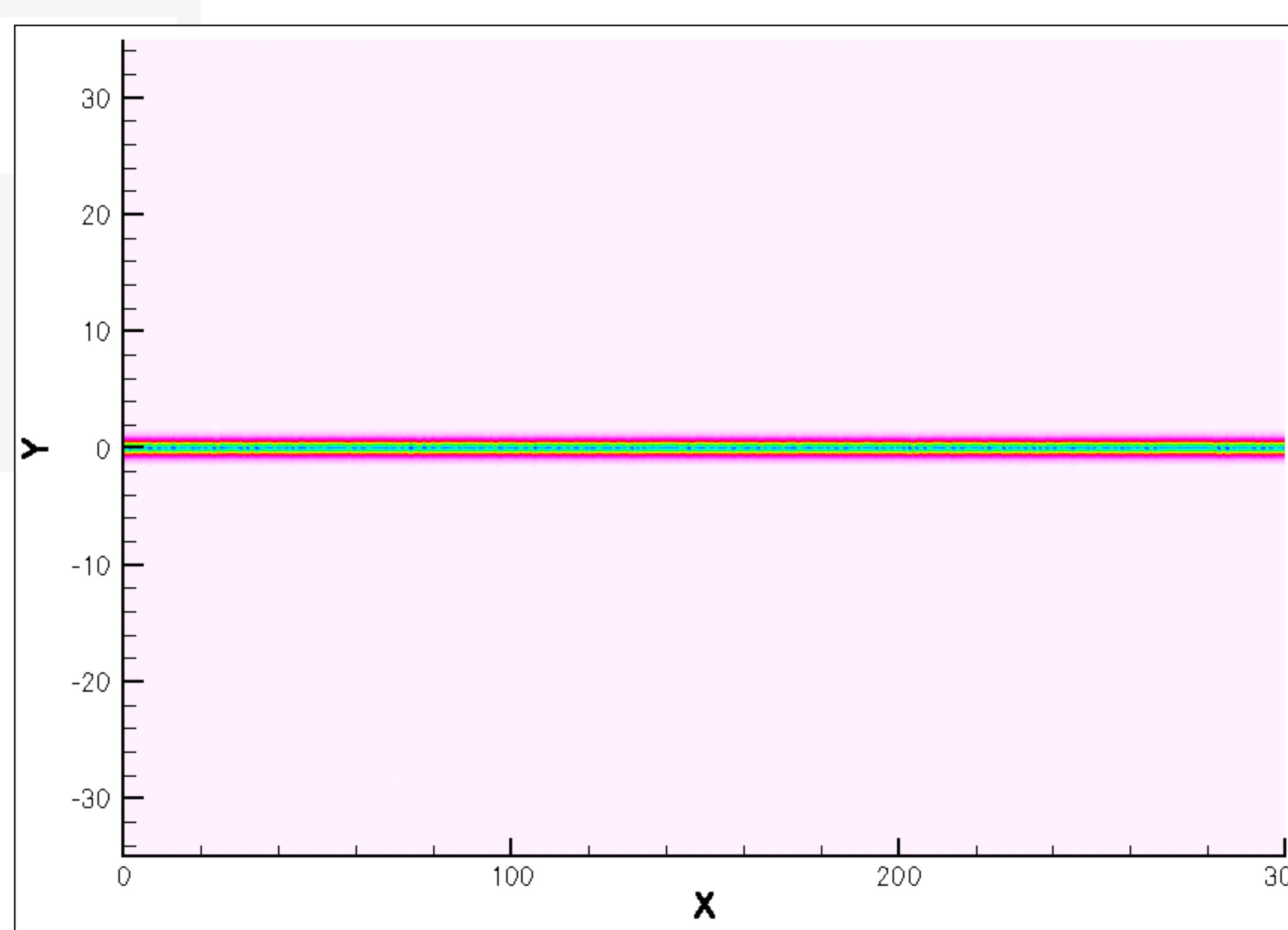


Resulting CR spectrum

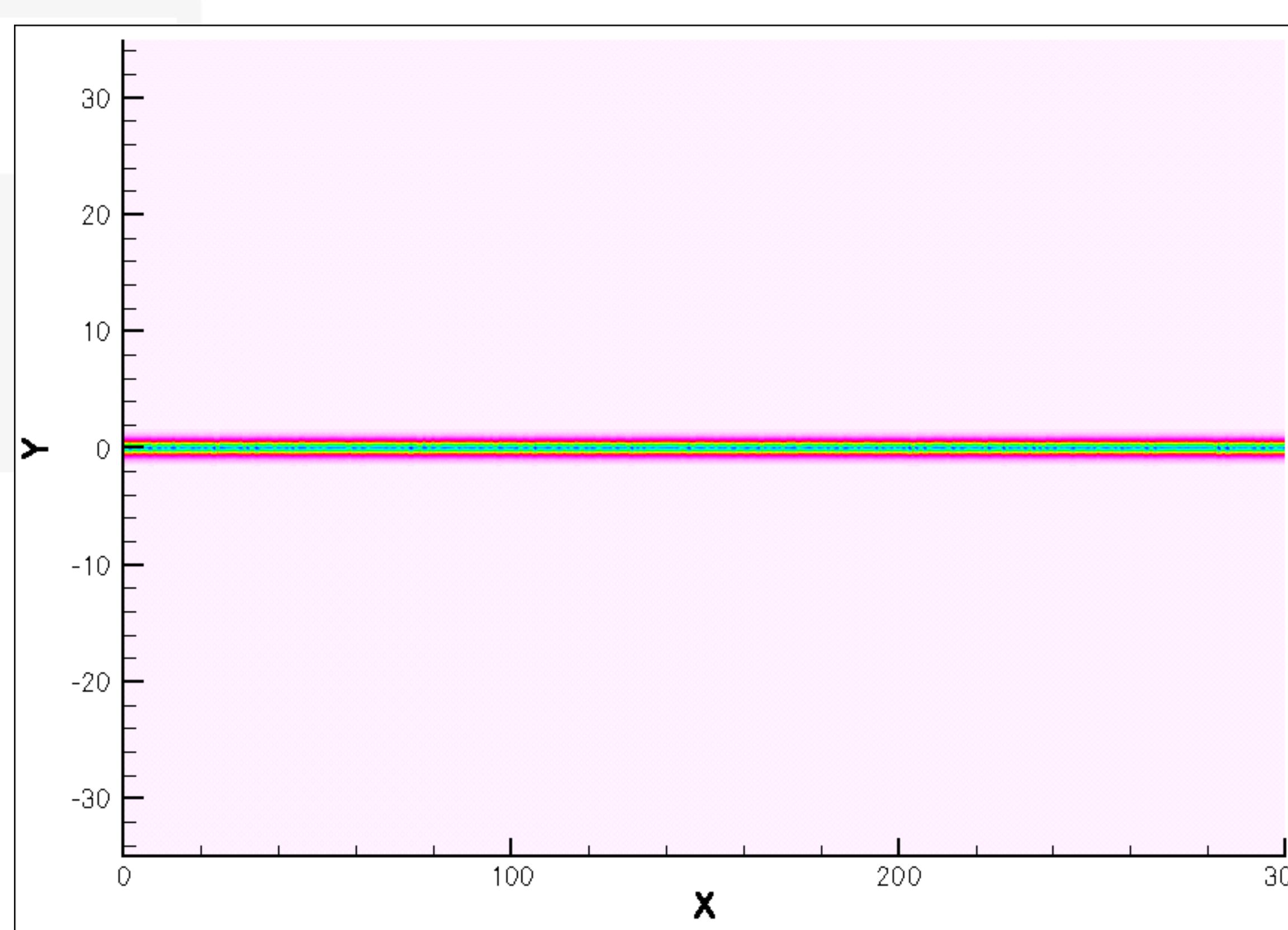




Turbulent cascade in the inertial range



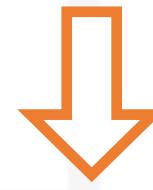
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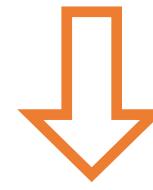
Turbulent cascade in the inertial range

Kolmogorov's approach

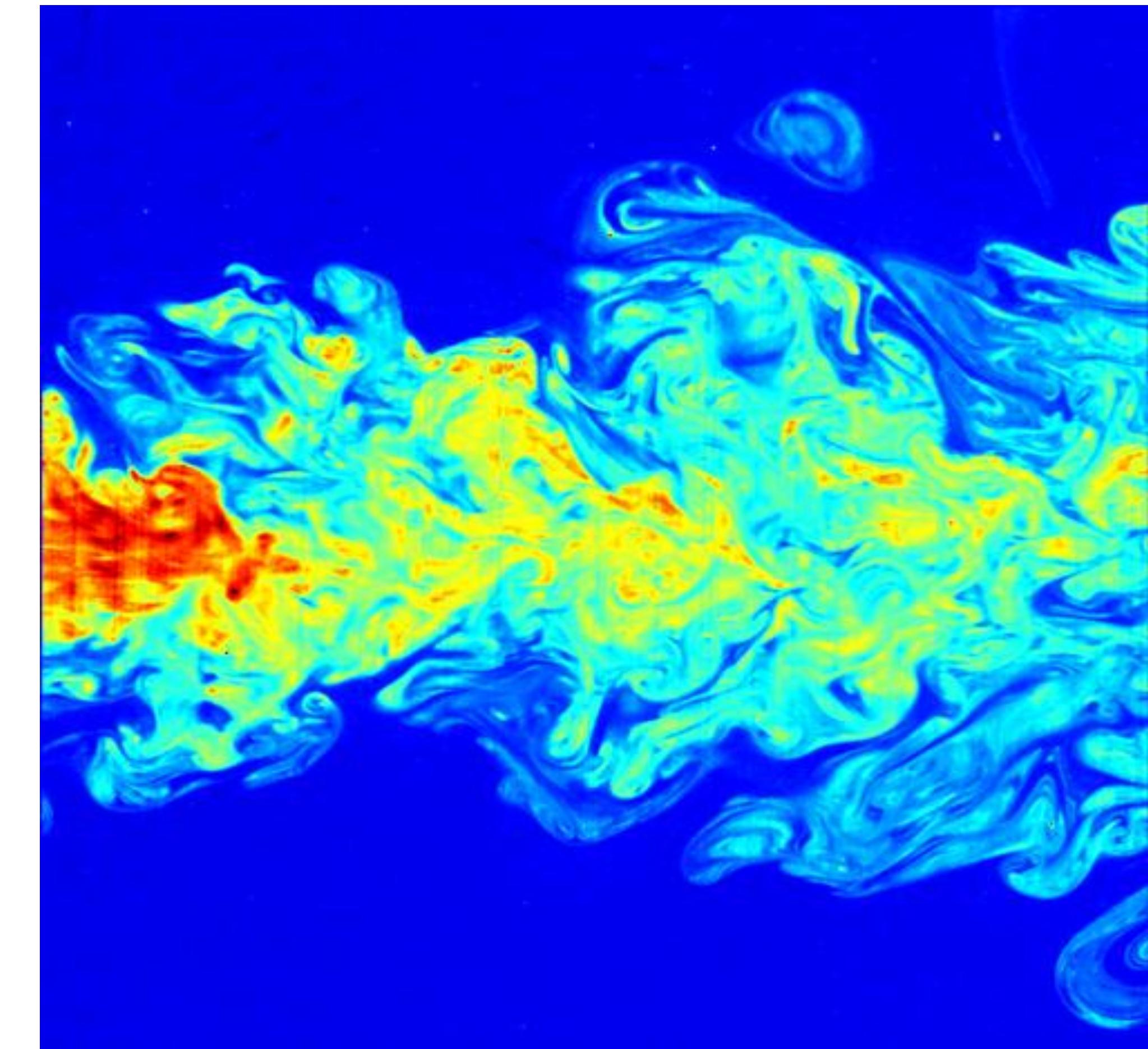
$$\frac{E_K/V}{\tau_{\text{turn}}} \sim \frac{\rho v_\ell^2}{\ell/v_\ell} = \text{const}$$



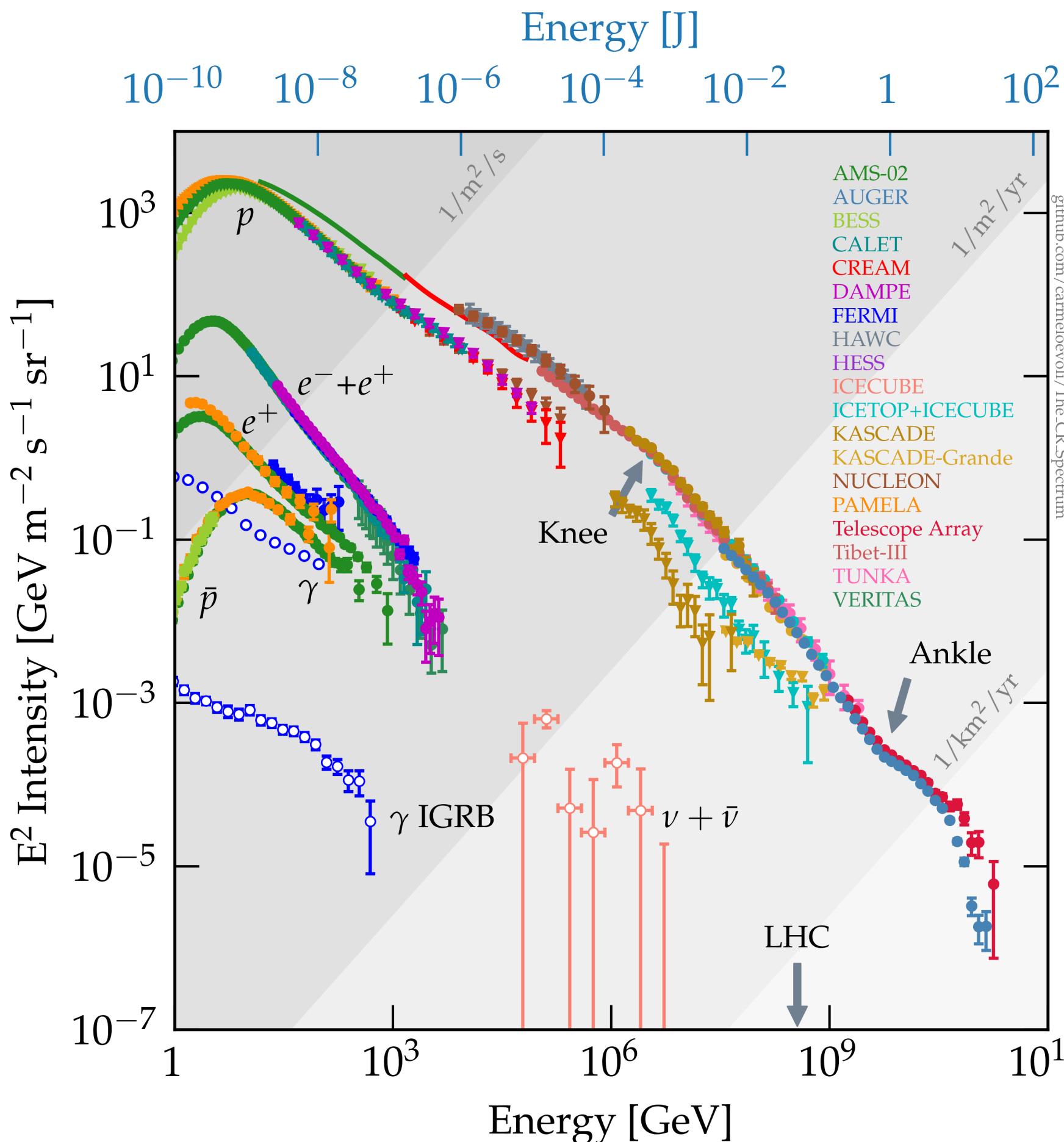
$$v_\ell^3 \sim \ell \quad \Rightarrow \quad v_\ell \sim \ell^{1/3} \quad \Rightarrow \quad v_k \sim k^{-1/3}$$



$$k \cdot E(k) \sim \rho v_k^2 \quad \Rightarrow \quad E(k) \sim k^{-5/3}$$

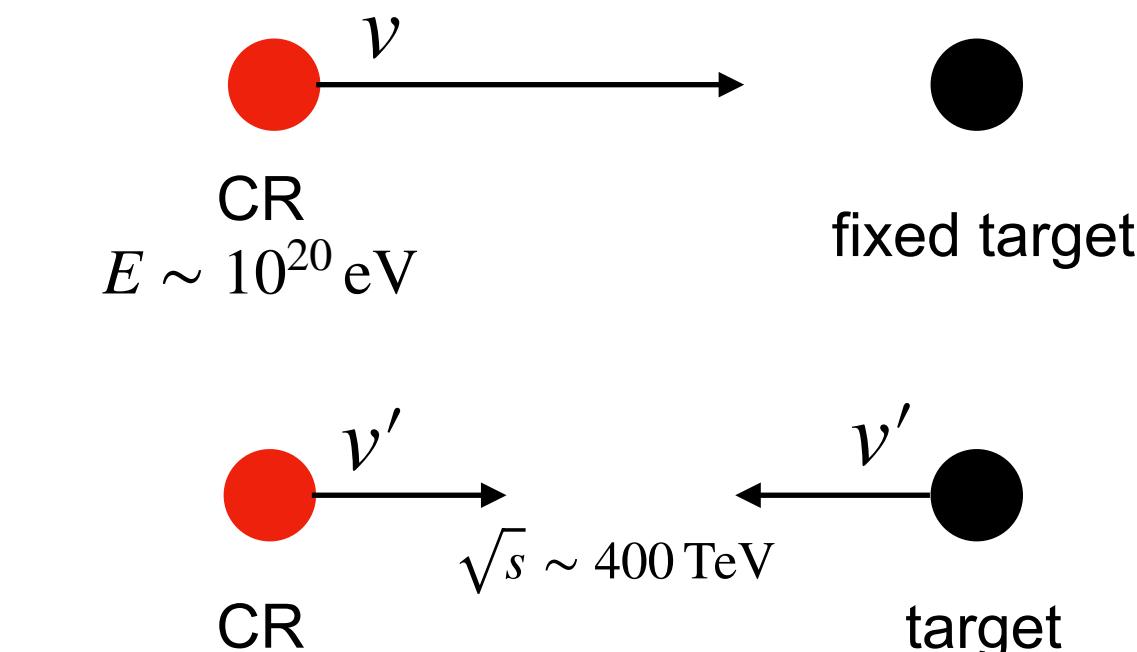


Why studying CR physics?



- $\frac{dN}{dE} \propto E^{-\gamma} \quad 2.7 \lesssim \gamma \lesssim 3.1$

- Very energetic particles



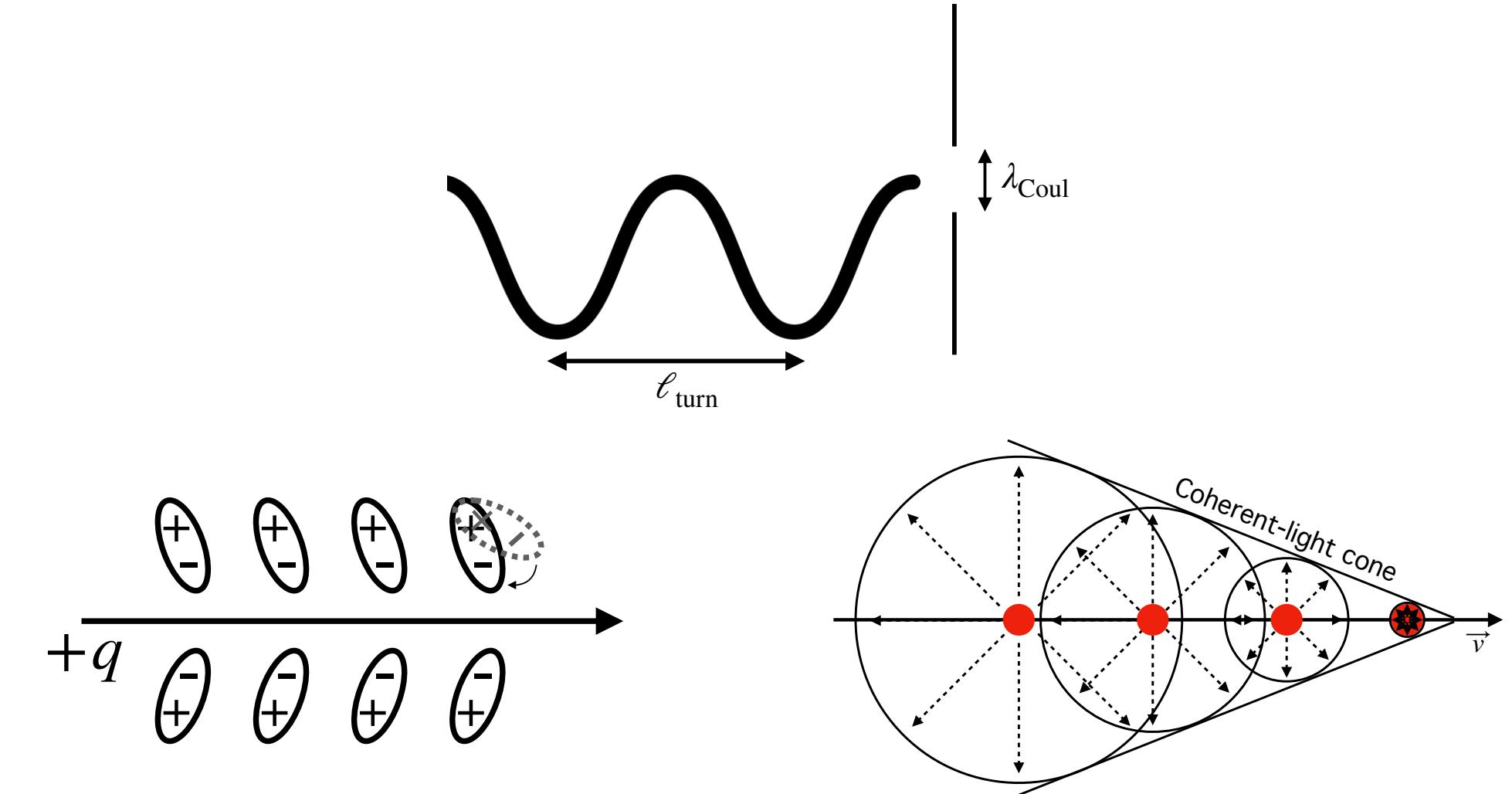
- Unique probe of extreme astrophysical phenomena

G S
S I

Damping of the fast modes

- Collisional damping
- Collisionless damping

$$\lambda_{\text{Coul}} \approx 1.3 \cdot 10^{-5} \left(\frac{\text{cm}^{-3}}{n_{\text{ISM}}} \right) \cdot \left(\frac{T}{10^4 \text{ K}} \right)^2 \text{ pc}$$

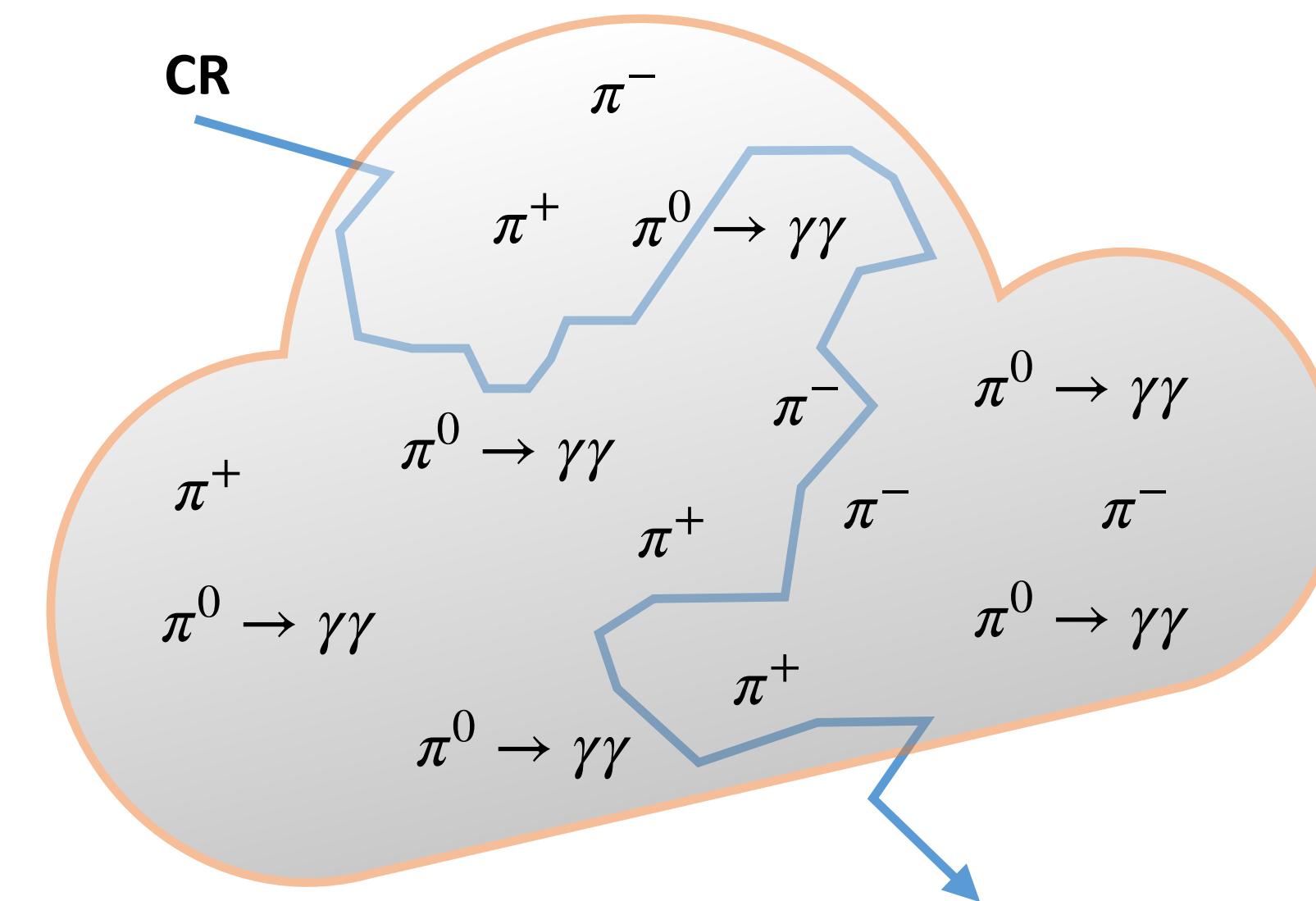


$$\lambda_{\text{Coul}}^{\text{disk}} \approx 1.3 \cdot 10^{-5} \text{ pc}, \quad n_{\text{disk}} = 1 \text{ cm}^{-3}, \quad T = 10^4 \text{ K}$$

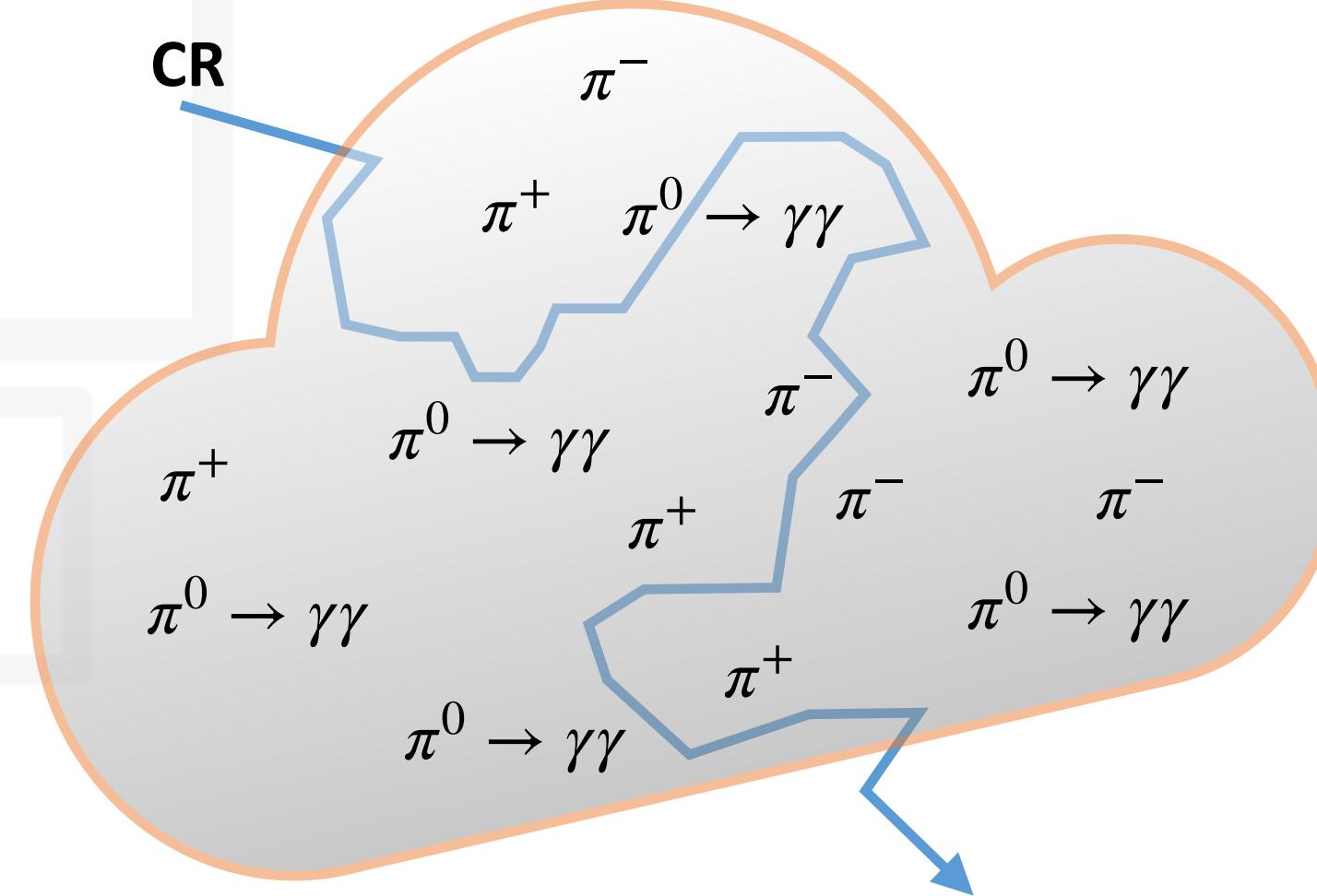
$$\lambda_{\text{Coul}}^{\text{halo}} \approx 1.3 \cdot 10^2 \text{ pc} \simeq L_{\text{inj}}, \quad n_{\text{Halo}} = 10^{-3} \text{ cm}^{-3}, \quad T = 10^6 \text{ K}$$

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Inferred CR density from pion decay



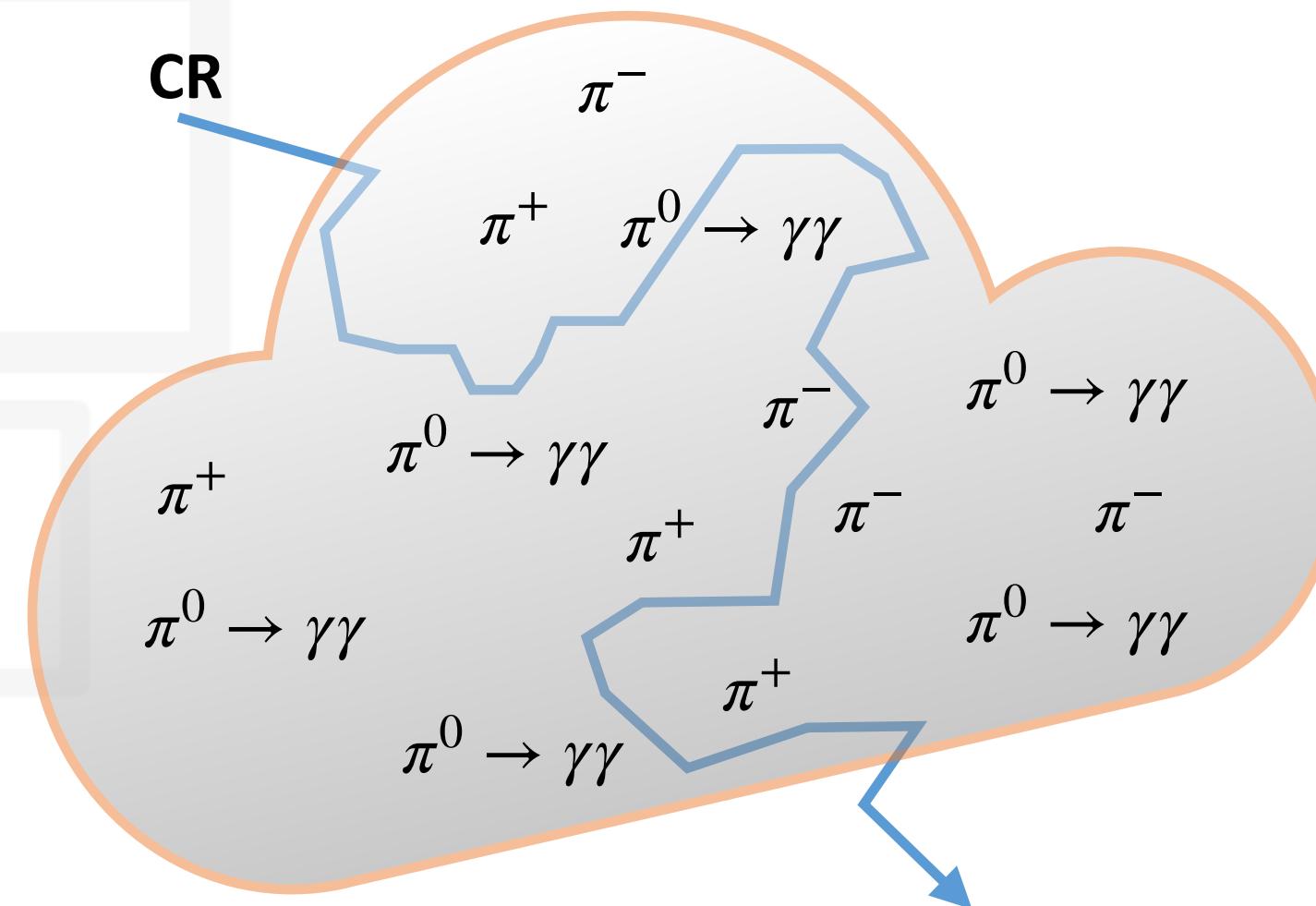
Inferred CR density from pion decay



$$\langle E_\gamma \rangle \simeq 0.1 E_{\text{CR}}$$

$$E_{\text{flux}} = \int_{E_{\text{min}}}^{E_{\text{max}}} dE E \cdot \left(\frac{\Phi_\gamma}{\frac{dN_\gamma}{dE} \frac{c}{4\pi}} \right) \left[\frac{E}{L^2 \cdot T} \right]$$
$$\Rightarrow L_\gamma = E_{\text{flux}} \cdot 4\pi d^2 \left[\frac{E}{T} \right]$$

Inferred CR density from pion decay



$$L_\gamma(\geq E_\gamma) \approx \eta_N \cdot \frac{W_p(\geq 10E_\gamma)}{\tau_{pp \rightarrow \pi^0}} = \eta_N \cdot \frac{W_p(\geq E_\gamma)}{1.6 \cdot 10^8 \text{ yr} \left(\frac{\text{cm}^{-3}}{n_H} \right)}$$

$$\omega_{\text{CR}} \equiv \frac{W_p(\geq 10E_\gamma)}{V_{\text{crossed}}} = \frac{W_p(\geq 10E_\gamma)}{M_{\text{tot}}} \cdot n_H \approx 1.8 \cdot 10^{-2} \left(\frac{\eta_N}{1.5} \right)^{-1} \left(\frac{L_\gamma(\geq E_\gamma)}{10^{34} \text{ erg} \cdot \text{s}^{-1}} \right) \left(\frac{M_{\text{tot}}}{10^6 M_\odot} \right)^{-1} \text{ erg} \cdot \text{cm}^{-3}$$

$$\langle E_\gamma \rangle \simeq 0.1 E_{\text{CR}}$$

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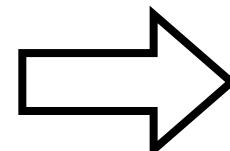
Inefficiency of Alfvén modes

$$E^{\text{GS}}(k_{\perp}) \sim k_{\perp}^{-5/3}$$

$$k_{\parallel} \sim k_{\perp}^{2/3} \Rightarrow k_{\parallel}^{3/2} \sim k_{\perp}$$

$$\rightarrow \frac{dk_{\perp}}{dk_{\parallel}} = \frac{3}{2} k_{\parallel}^{3/2 - 1} = \frac{3}{2} k_{\parallel}^{1/2}$$

$$\int dk_{\parallel} E(k_{\parallel}) = \int dk_{\perp} E(k_{\perp})$$



$$\begin{aligned}\int dk_{\perp} E(k_{\perp}) &= \int \frac{3}{2} k_{\parallel}^{1/2} dk_{\parallel} E(k_{\perp}) = \\ &= \frac{3}{2} \int k_{\parallel}^{1/2} dk_{\parallel} k_{\perp}^{-5/3} = \frac{3}{2} \int dk_{\parallel} k_{\parallel}^{1/2} \left(k_{\parallel}^{3/2} \right)^{-5/3}\end{aligned}$$

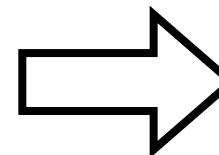
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$$\begin{aligned} \int dk_{\perp} E(k_{\perp}) &= \int \frac{3}{2} k_{\parallel}^{1/2} dk_{\parallel} E(k_{\perp}) = \\ &= \frac{3}{2} \int k_{\parallel}^{1/2} dk_{\parallel} k_{\perp}^{-5/3} = \frac{3}{2} \int dk_{\parallel} k_{\parallel}^{1/2} \left(k_{\parallel}^{3/2} \right)^{-5/3} \end{aligned}$$



$$E^{\text{GS}}(k_{\parallel}) \sim k_{\parallel}^{-2}$$

G S
S I

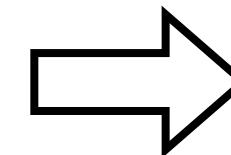
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$$\rightarrow \frac{dk_{\perp}}{dk_{\parallel}} = \frac{3}{2} k_{\parallel}^{3/2 - 1} = \frac{3}{2} k_{\parallel}^{1/2}$$

$$\int dk_{\parallel} E(k_{\parallel}) = \int dk_{\perp} E(k_{\perp})$$



$$\begin{aligned} \int dk_{\perp} E(k_{\perp}) &= \int \frac{3}{2} k_{\parallel}^{1/2} dk_{\parallel} E(k_{\perp}) = \\ &= \frac{3}{2} \int k_{\parallel}^{1/2} dk_{\parallel} k_{\perp}^{-5/3} = \frac{3}{2} \int dk_{\parallel} k_{\parallel}^{1/2} \left(k_{\parallel}^{3/2} \right)^{-5/3} \end{aligned}$$

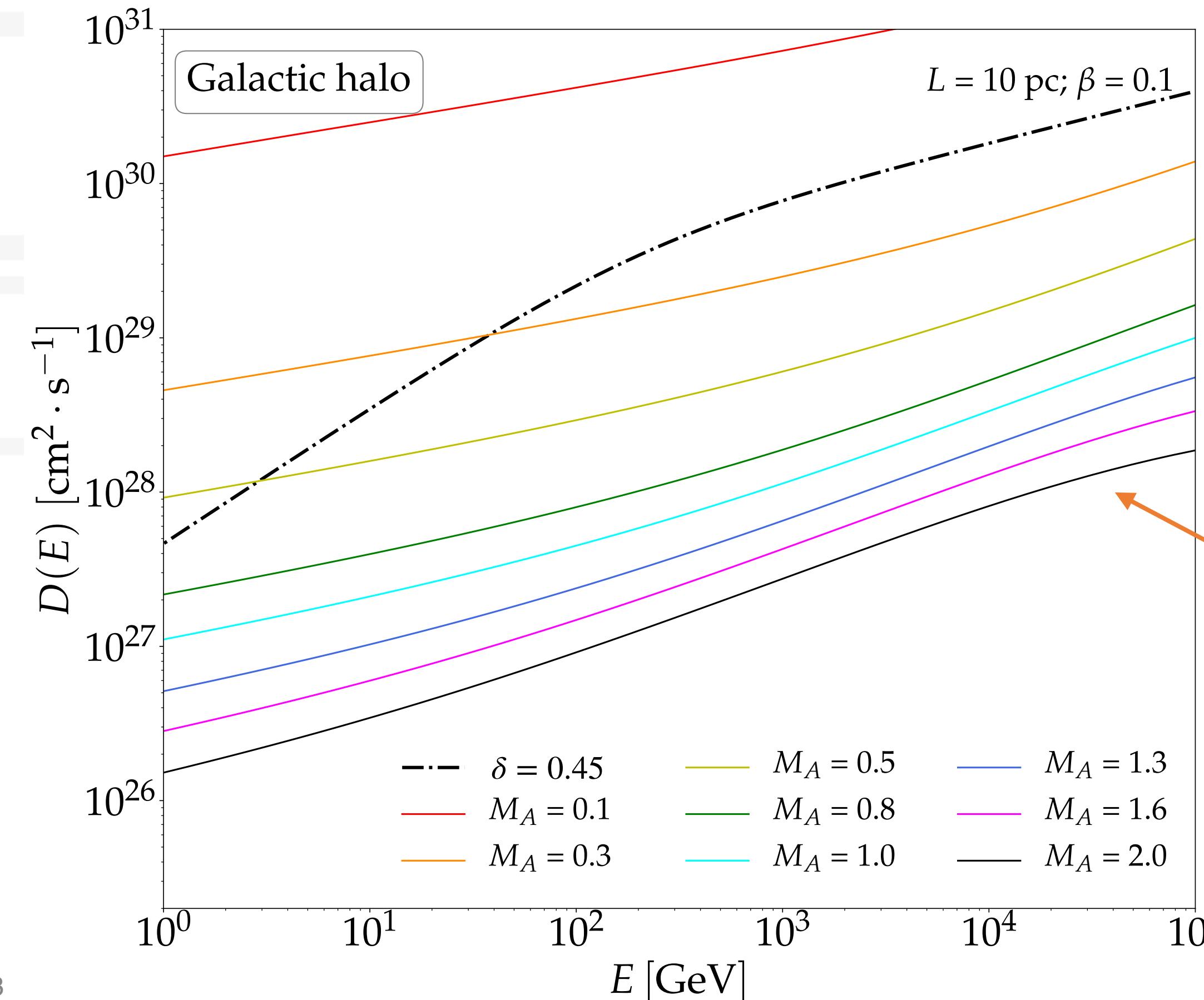


$$E^{\text{GS}}(k_{\parallel}) \sim k_{\parallel}^{-2}$$

But $D_{\mu\mu} \propto R_n(\textcolor{orange}{k}_{\parallel} v_{\parallel} - \omega + n\Omega) \dots$

G S
S I

Resulting CR diffusivity



G S
S I