

Cosmic Ray propagation in the Galaxy: the role of “*self-confinement*”

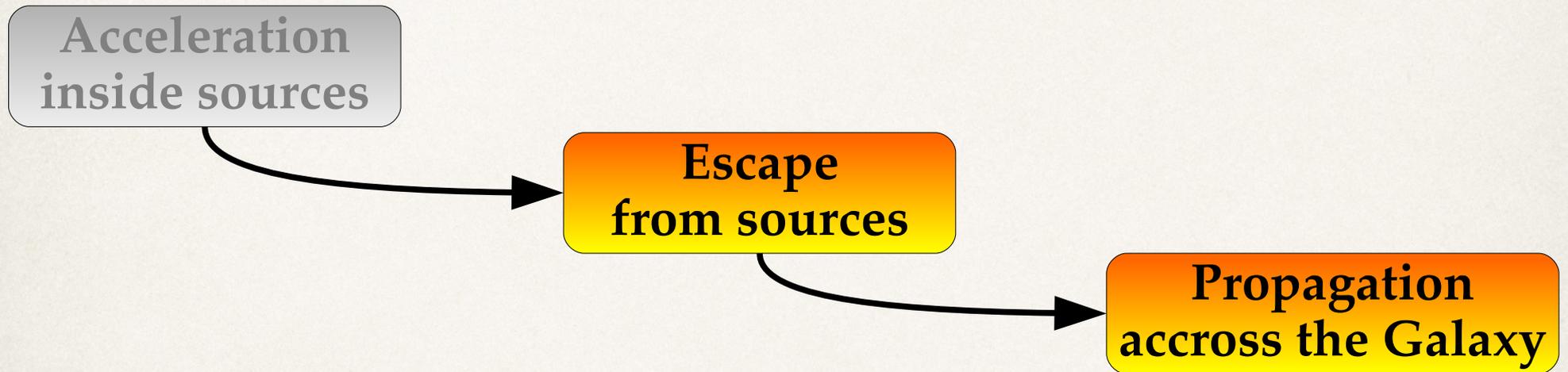


Giovanni Morlino,

INAF/Osservatorio Astrofisico di Arcetri, Firenze, ITALY

Role of self-generated turbulence

In the current picture of CR physics, the magnetic turbulence self-generated by the streaming of CRs plays a role in all basic steps of their life:



- ▶ This talk will cover only the turbulence generated through *Resonant Streaming Instability*
- ▶ *Non-Resonant (Bell) Instability* could also play a relevant role but has not been fully explored yet

Basic *diffusive halo model*

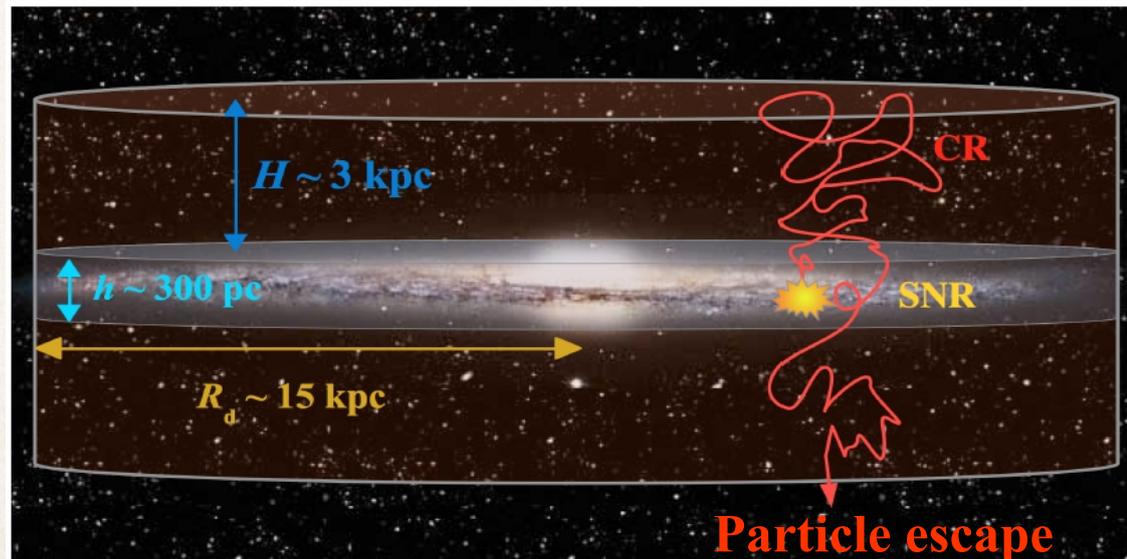
In the basic picture of CR propagation model:

- CR sources are located inside the Galactic disc
- CRs diffuse in a magnetic halo larger than the Galactic disc
- The CR distribution vanishes at $z = H$ ($H \sim 3-4$ kpc from diffuse synchrotron emission)
- The diffusion coefficient $D(E)$ is assumed constant everywhere in the halo
(Diffusion in the disc play a minor role)

Escaping time from the halo

$$\tau_{esc}(E) = \frac{H^2}{2D(E)}$$

Cylindrical symmetry



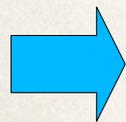
The Secondary/Primary ratio

Secondary/primary ratio:

$$\begin{cases} N_p(E) = Q_{inj}(E) \tau_{esc}(E) \propto E^{-\gamma-\delta} \\ N_s(E) = N_p(E) (n_H \sigma_{sp} c) \tau_{esc}(E) \propto E^{-\gamma-2\delta} \end{cases}$$



$$\frac{N_s}{N_p} \propto \tau_{esc}(E) \sim E^{-\delta}$$

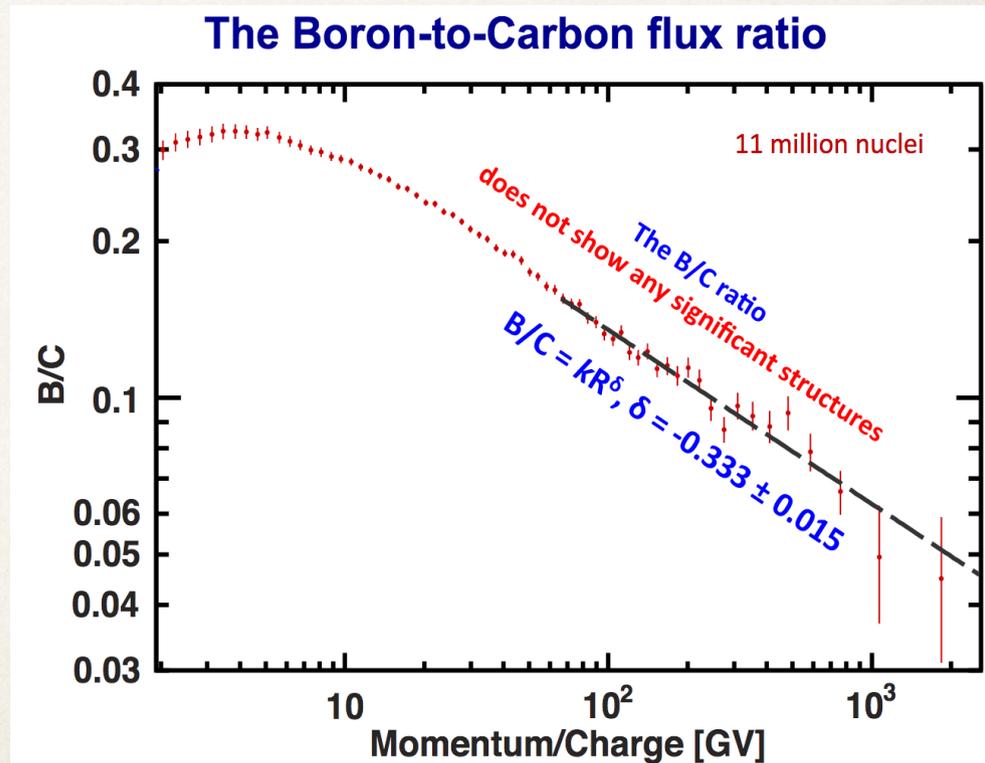
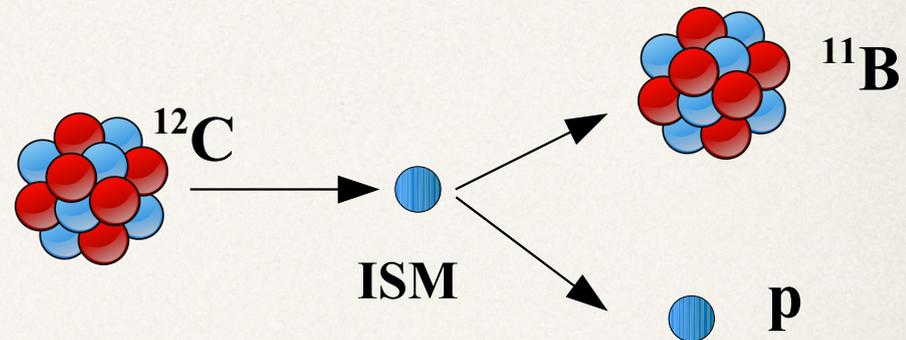


$$D(E) \propto E^{1/3}$$

**Suggesting
Kolmogorov
turbulence**

This picture is unsatisfactory for at least two reasons:

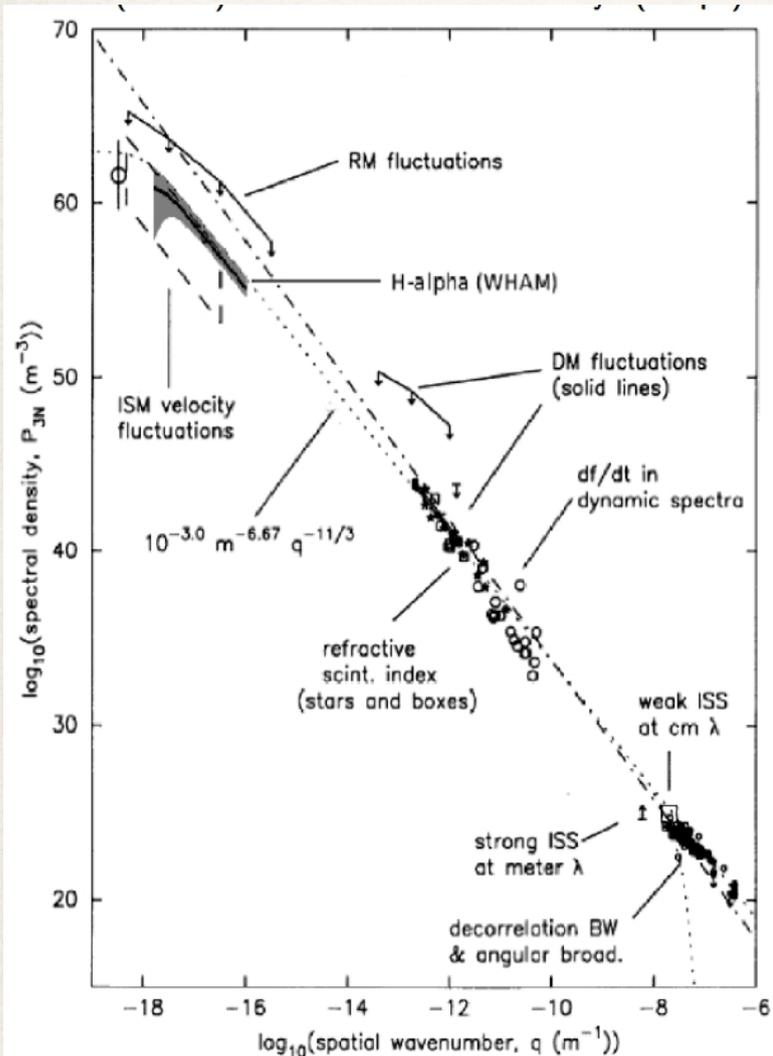
- Which is the physical meaning of H ?
- What generates the diffusion?



The interstellar turbulence

Electron density fluctuation in the ISM

[Armstrong et al. 1995, Lazarian 1995]



BASIC ASSUMPTIONS:

- ▶ Turbulence is stirred by SNe at a typical scale $L \sim 10-100$ pc
- ▶ The spectrum is Kolmogorov like $\sim k^{-5/3}$ (density is a passive tracer so it has the same spectrum: $\delta n \propto \delta B^2$):

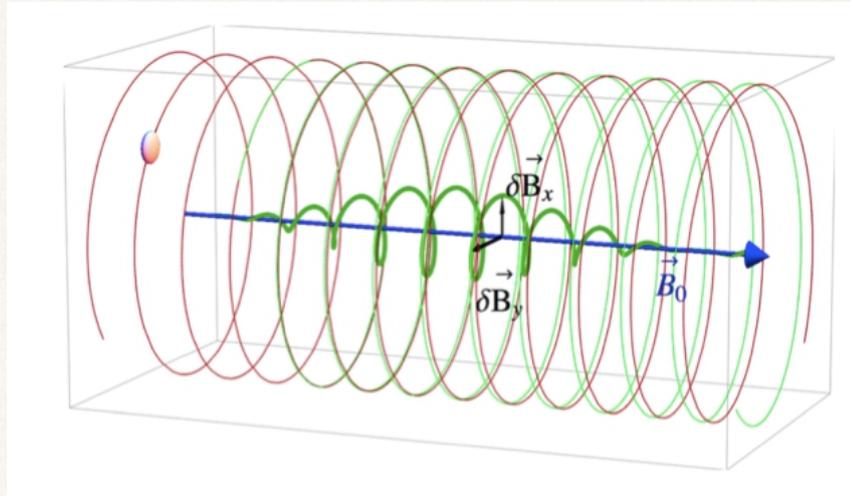
$$W(k) dk = \frac{\langle \delta B(k) \rangle^2}{B_0^2} = \frac{2}{3} \frac{\eta_B}{k_0} \left(\frac{k}{k_0} \right)^{-5/3}$$

where $k_0 = L^{-1}$ and the level of turbulence is

$$\eta_B = \int_{k_0}^{\infty} W(k) dk \sim 0.01 - 0.1$$

- ▶ **Magnetic field turbulence is Alfvénic**

Charged particles in a turbulent field: quasi-linear theory

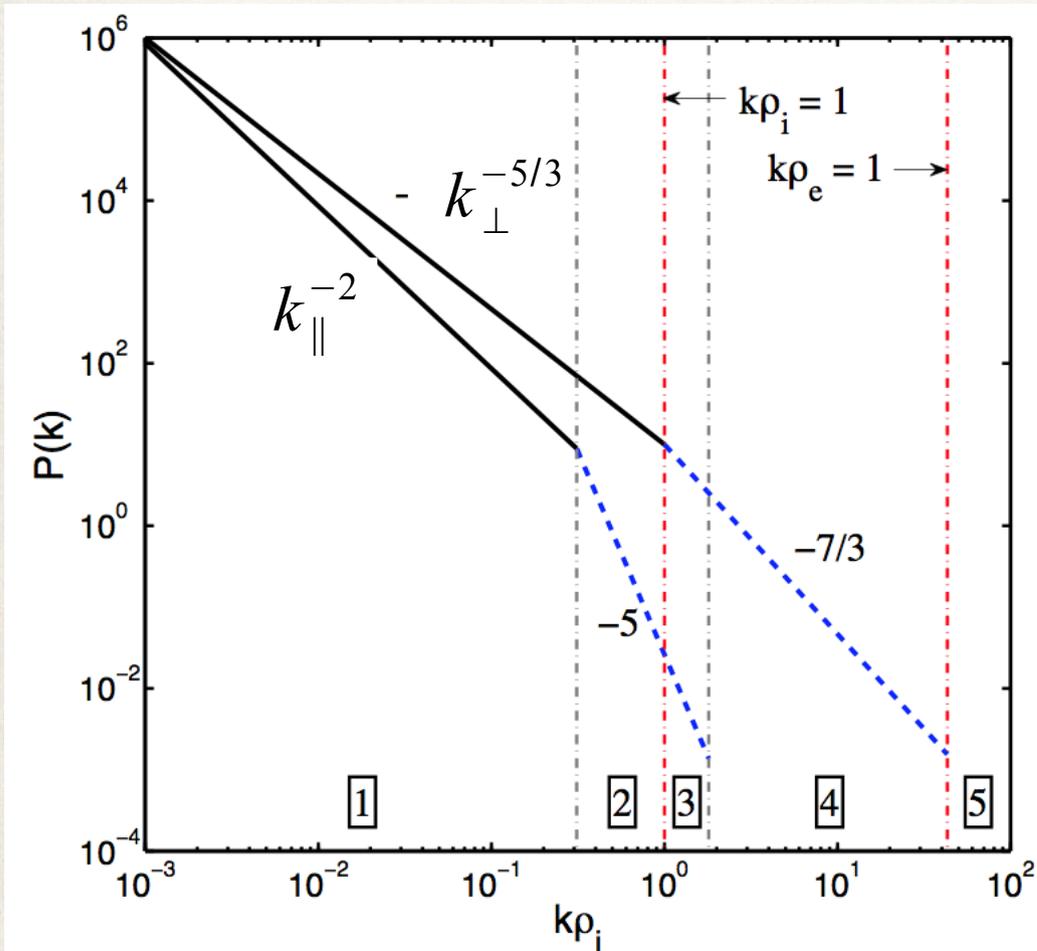


- ▶ The turbulent field is a small perturbation with respect to the regular component
- ▶ Particles interact with parallel waves resonantly: $k_{res} = 1/r_L(p) \quad k \parallel B_0$
- ▶ A diffusion behavior follows with typical diffusion coefficient

$$D_{zz}(p) = \frac{v r_L}{3} \frac{1}{k_{res} W(k_{res})} \sim 3 \times 10^{28} \left(\frac{p}{GeV/c} \right)^{1/3} cm^2 s^{-1}$$

1st problem: anisotropic turbulent cascade

Energy spectra of Alfvénic turbulence from critical balance



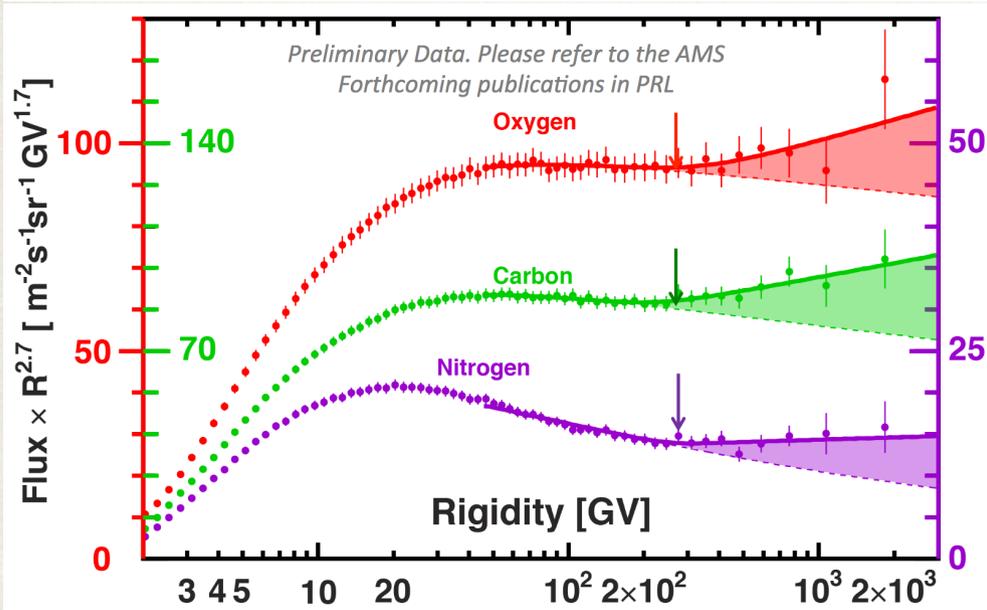
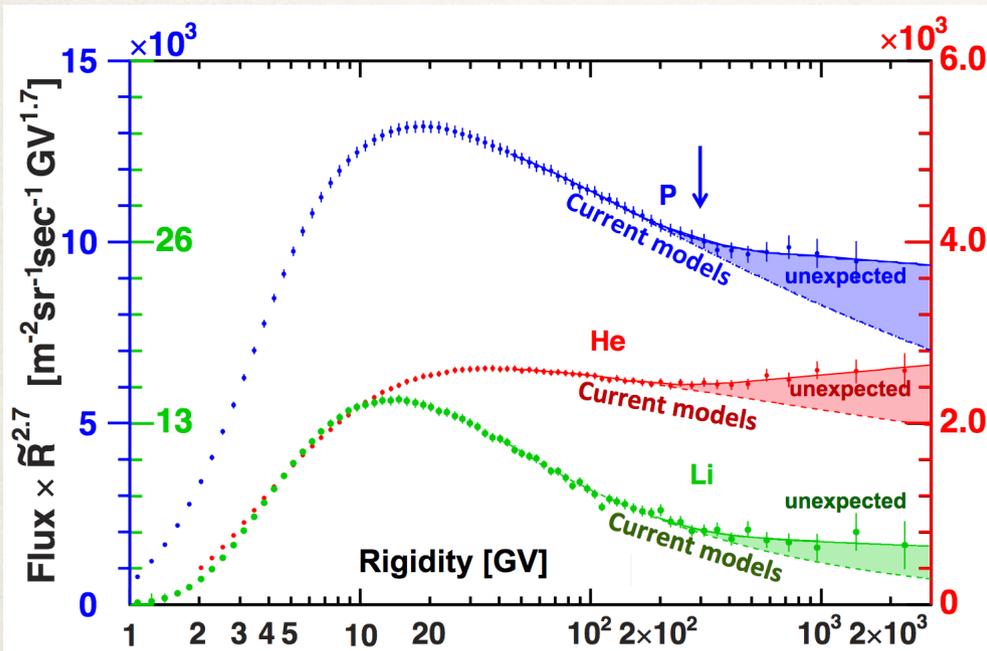
[Chen et al., 2010]

- ▶ The cascade of Alfvénic turbulence becomes anisotropic (Goldreich & Sridhar 1994, 1995)

$$P(k_{\parallel}) \propto k_{\parallel}^{-2}; \quad P(k_{\perp}) \propto k_{\perp}^{-5/3}$$

- ▶ The power is mainly in modes perpendicular to the local magnetic field, which are inefficient to scatter particles
- ▶ At small scale there is not enough power to scatter particles

2nd problem: spectral hardening



Recent measurements by PAMELA and AMS-02 revealed the existence of a fine structure:

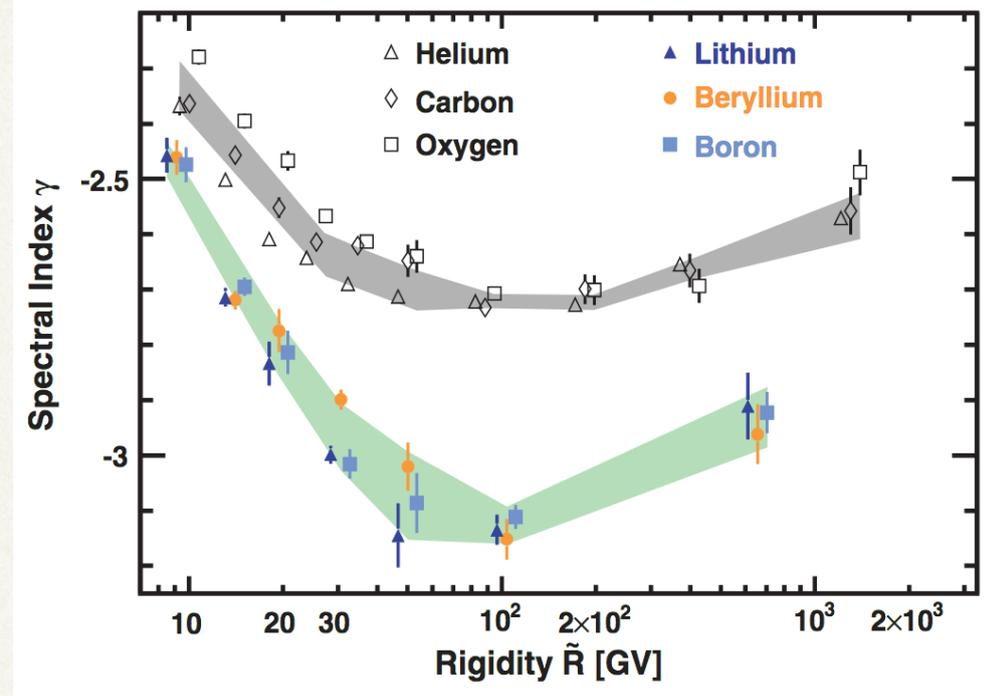
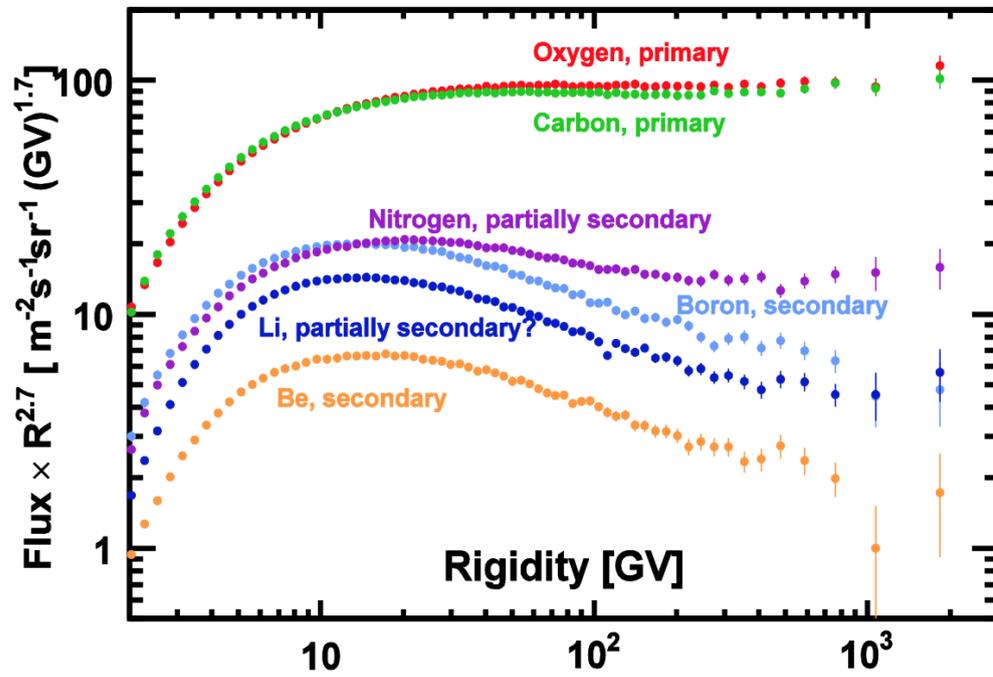
At rigidity of ~ 300 GV all spectra show a spectral hardening

NO MORE A SIMPLE POWER-LAW

$$f_0(E) = \frac{Q_{SN}(E)}{2\pi R_{disc}^2} \frac{H}{D(E)} \propto E^{-\gamma-\delta}$$

Either the injected spectrum or the diffusion present a break at ~ 300 GV

Spectral hardening for secondary CRs



[AMS collaboration, PRL 120, 021101, 2019]

$$f_{sec}(E) = f_{pri} \times \tau_{esc} \propto E^{-\gamma-2\delta}$$

The spectral hardening of secondary species is larger than primaries

^ supports the origin of break due to propagation rather than primary acceleration

Waves from resonant streaming instability: a simple exercise

(Blasi, Amato & Serpico, 2016)

- ▶ Apply the resonant streaming instability to CR escaping from the halo

$$\Gamma_{CR} = \frac{v_A}{B_0^2/8\pi} \frac{1}{F(k)} \frac{\partial P_{CR}(>p)}{\partial z}$$

- ▶ The magnetic turbulence is damped by non-linear processes:

$$\Gamma_{NLD} = v_s k F(k)$$

- ▶ Equating the rate of amplification and damping we get $F(k)$ and the diffusion coefficient

$$\Gamma_{CR} = \Gamma_{NLD}$$

$$D(E) \simeq \frac{r_L v}{3 F(k_{res})} = 6 \times 10^{27} \left(\frac{E}{10 \text{ GeV}} \right)^{0.85} \left(\frac{H}{4 \text{ kpc}} \right)^{1/2} \text{ cm}^2/\text{s}$$

Normalization close to the value
inferred from B/C

Different energy dependence from
the Kolmogorov turbulence

Coupling CR transport with magnetic turbulence evolution

(Blasi, Amato & Serpico, 2016)

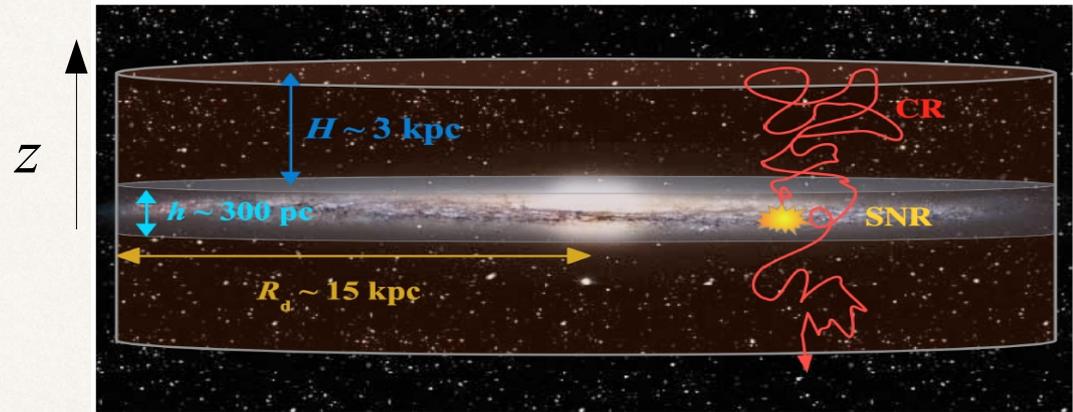
CR transport equation in 1-D

$$\frac{\partial f}{\partial t} + v_A \frac{\partial f}{\partial z} = \frac{\partial}{\partial z} \left[D \frac{\partial f}{\partial z} \right]$$

Self-generated diffusion coefficient

$$D(p, z, t) = \frac{r_L v}{3} \frac{1}{\mathcal{F}(k, z, t)} \Big|_{k=1/r_L}$$

$$\frac{\delta B^2}{B_0^2} = \int \mathcal{F}(k) \frac{dk}{k} \quad \text{Turbulence spectrum}$$



Transport equation for magnetic turbulence

$$\frac{\partial \mathcal{F}}{\partial t} + v_A \frac{\partial \mathcal{F}}{\partial z} = (\Gamma_{CR} - \Gamma_D) \mathcal{F} + Q_w$$

Damping
Injection

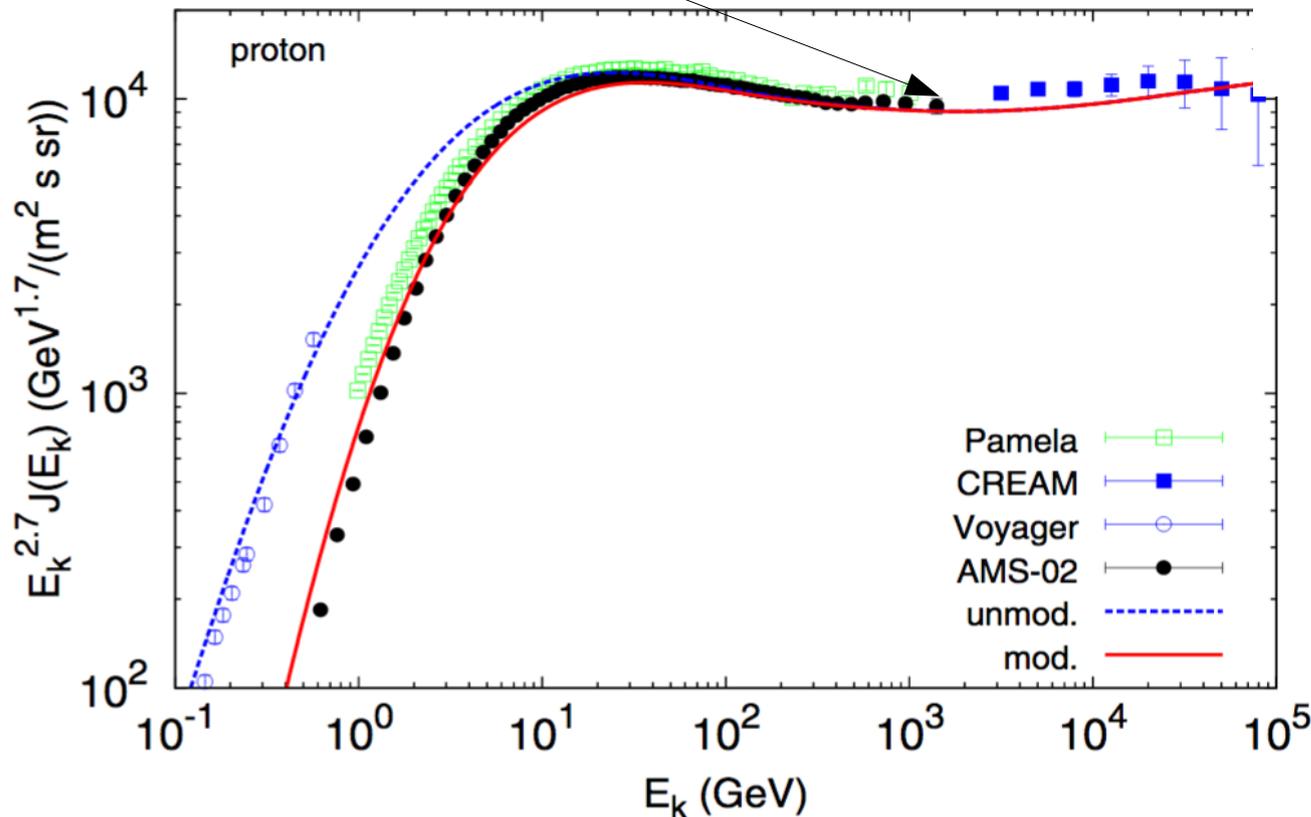
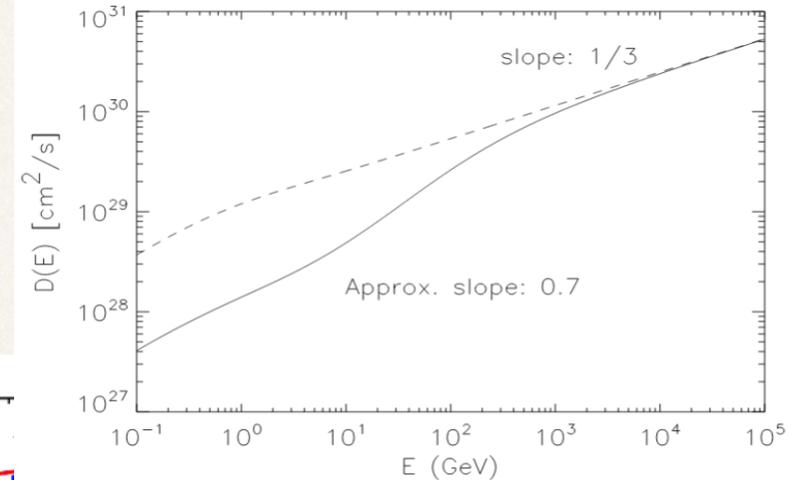
Resonant amplification:

$$\Gamma_{CR} = \frac{16\pi}{3} \frac{v_A}{\mathcal{F} B_0^2} [p^4 v \nabla f]_{p=p_{res}}$$

Position of the break

(Blasi, Amato & Serpico, 2012)

$$E_{\text{tr}} = 228 \text{ GeV} \left(\frac{R_{d,10}^2 H_3^{-1/3}}{\xi_{0.1} E_{51} \mathcal{R}_{30}} \right)^{\frac{3}{2(\gamma_p - 4)}} B_{0,\mu}^{\frac{2\gamma_p - 5}{2(\gamma_p - 4)}}$$



- ▶ Pre-existing waves (Kolmogorov) dominates above the break
- ▶ Self-generated turbulence dominates below ~ 100 GeV

Equation for the turbulence evolution

[Jones 1993, ApJ 413, 619]

- ▶ CR transport equation is now coupled with evolution of turbulence

$$\frac{\partial W}{\partial t} = \frac{\partial}{\partial k} \left[D_{kk} \frac{\partial W}{\partial k} \right] + \frac{\partial}{\partial z} (v_A W) + \Gamma_{CR} W + Q_k$$

- ▶ Diffusion in k -space (non-linear): $D_{kk} = c_k |v_A| k^{7/2} W^{1/2}$
- ▶ Advection of waves at the Alfvén speed
- ▶ Waves growth due to CR streaming: $\Gamma_{CR} \propto \partial f / \partial z$

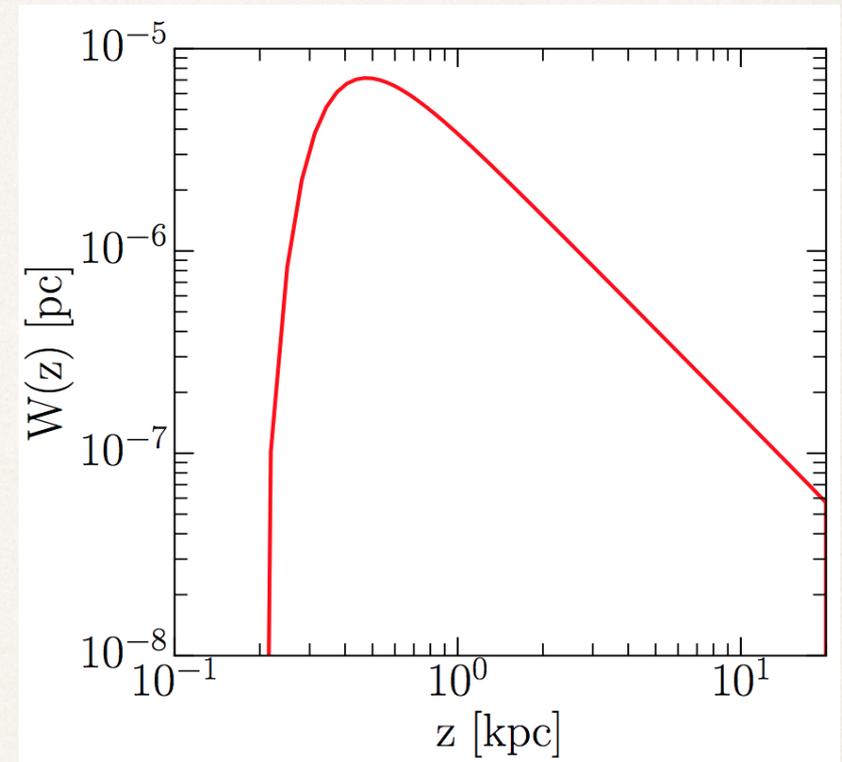
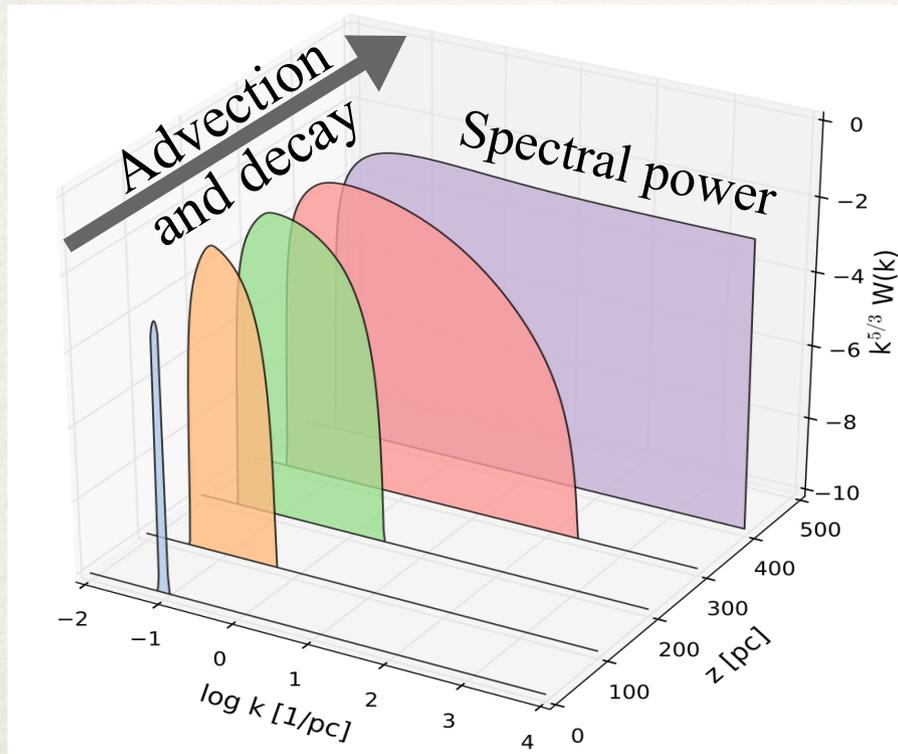
$$\Gamma_{CR} = \frac{16\pi}{3} \frac{v_A}{k W(k) B_0^2} \left[v p^4 \frac{\partial f}{\partial z} \right]$$

- ▶ External (e.g. SNe) source term in the disc: $Q \sim \delta(z) \delta(k - k_0)$
- ▶ In the absence of CR it returns the Kolmogorov spectrum: $W(k) \sim k^{-5/3}$

Developing the turbulent halo

[Evoli, Blasi, GM, Aloisio, 2018, PRL]

Large scale turbulence generated inside the Galactic disc by SN explosion advected and decaying through Kolmogorov cascade.

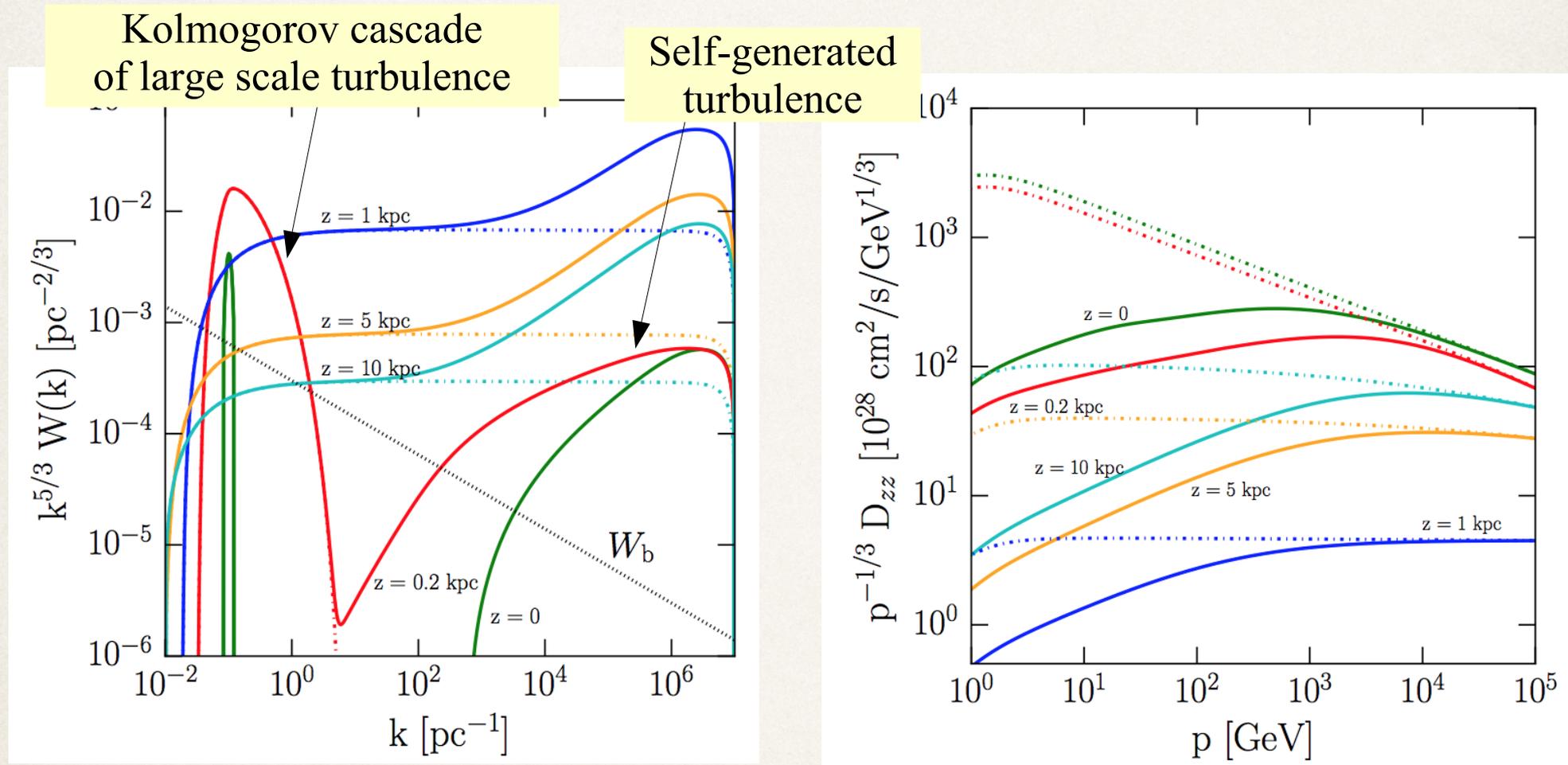


$$\tau_{\text{cascade}} = \tau_{\text{adv}} \rightarrow \frac{k_0^2}{D_{kk}} = \frac{z_c}{v_A} \rightarrow z_c \sim \text{kpc}$$

- ▶ z_c sets the scale where the turbulent cascade develops
- ▶ The boundary H does not have any physical meaning

Non-linear cosmic ray transport: a global picture

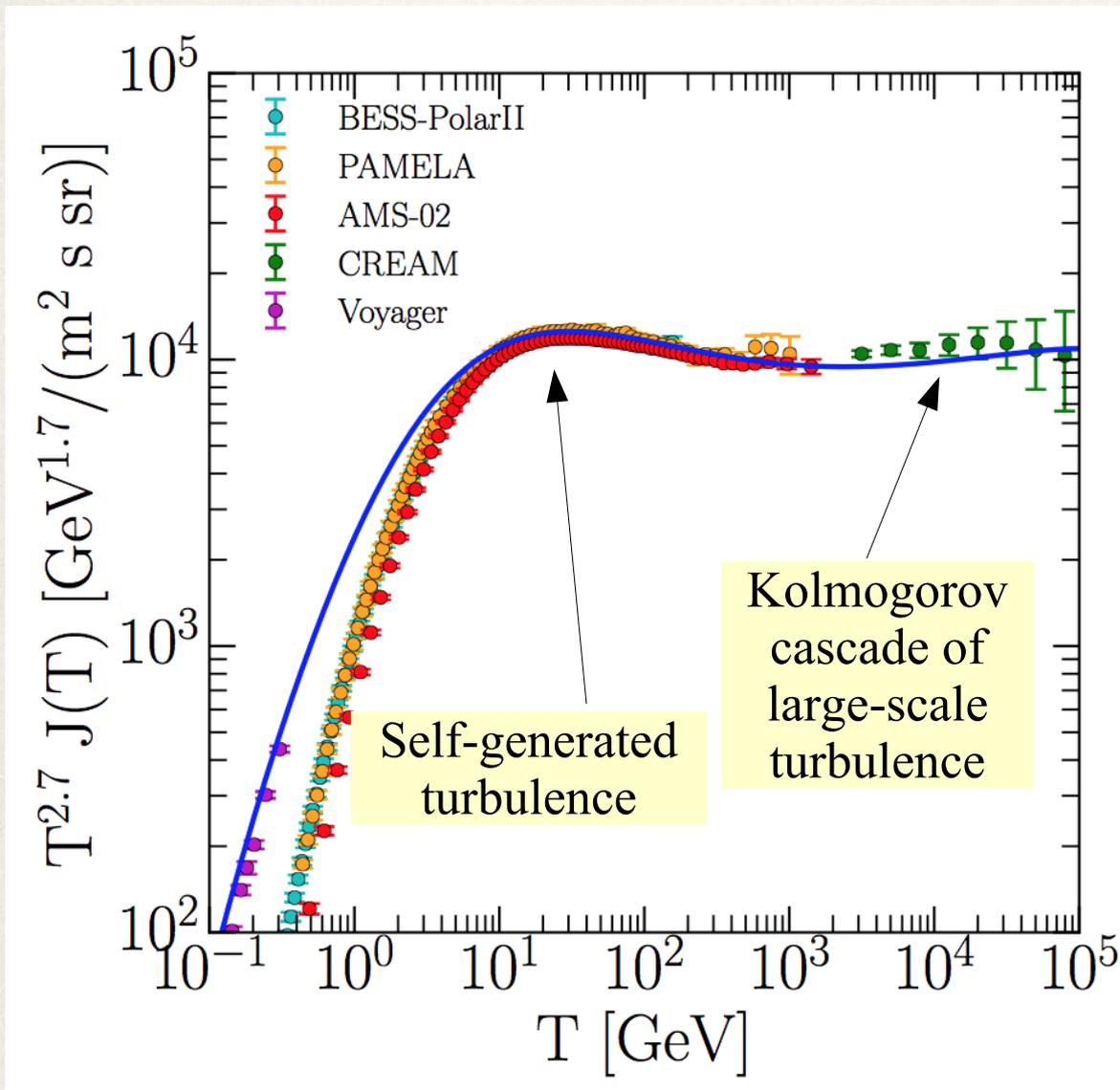
(Evoli, Blasi, GM, Aloisio, 2018, PRL)



Turbulence spectrum without (dotted) and with (solid) CR self-generated waves at different distance from the galactic plane.

Non-linear cosmic ray transport: a global picture

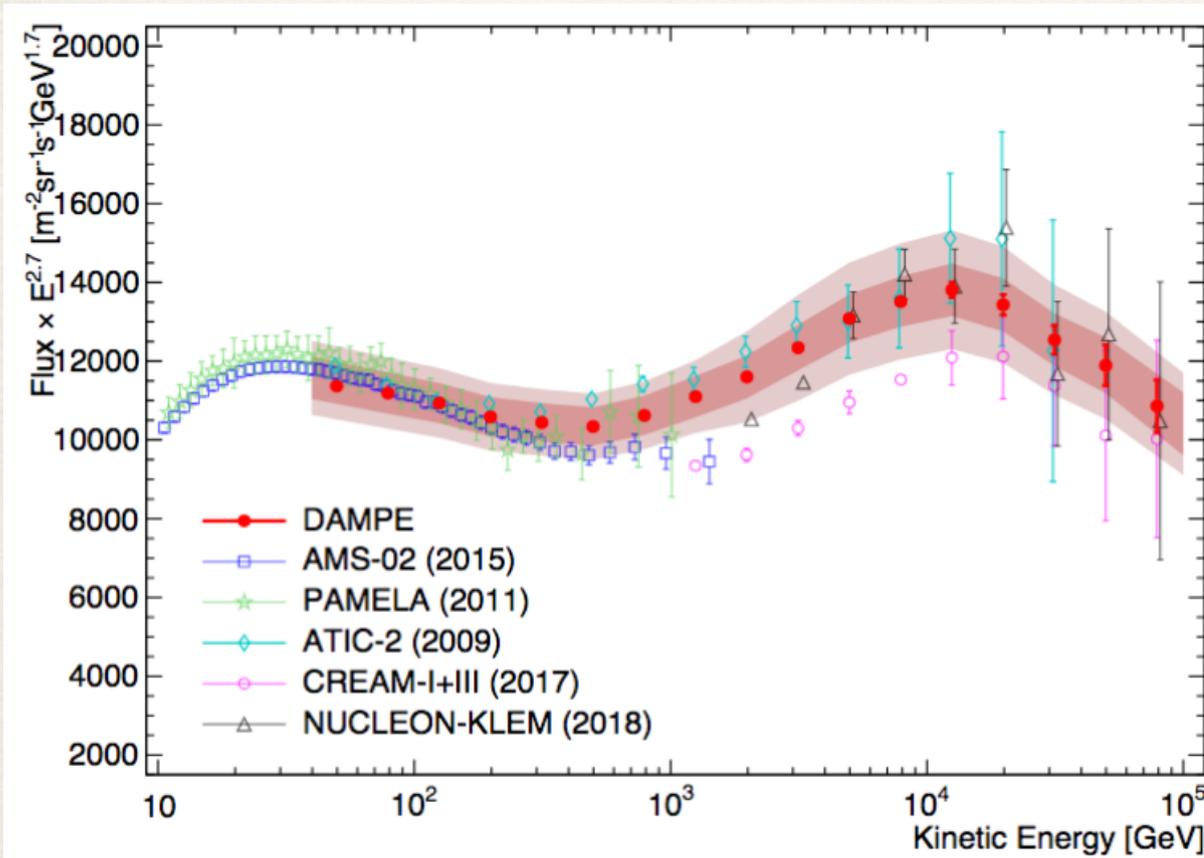
[Evoli, Blasi, GM, Aloisio, 2018, PRL]



- ▶ Pre-existing waves (Kolmogorov) dominates above the break
- ▶ Self-generated turbulence dominates below ~ 100 GeV
- ▶ Voyager data are reproduced with no additional breaks, but due to advection with self-generated waves
- ▶ The boundary ($H = 100$ kpc) has no impact on the result
- ▶ Low energy spectrum is well accounted by advection without introducing *ad hoc* breaks in the primary spectra.

Not the end of the story: more transport regimes?

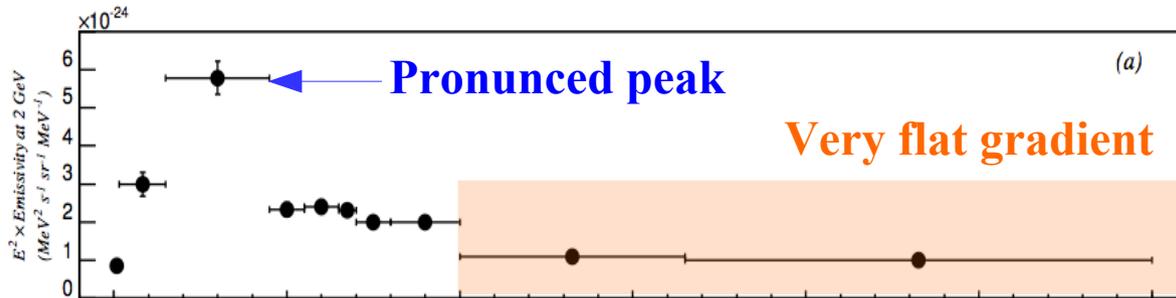
Proton spectrum above 1 TeV



- ▶ Above ~10 TeV the proton spectrum seems to steepen again
- ▶ Possible indication of a further change in transport regime?
- ▶ The answer could come from the measurement of secondaries spectra

[An+, DAMPE colabration 2019, Science]

3rd anomaly: the cosmic ray gradient in the Galactic plane

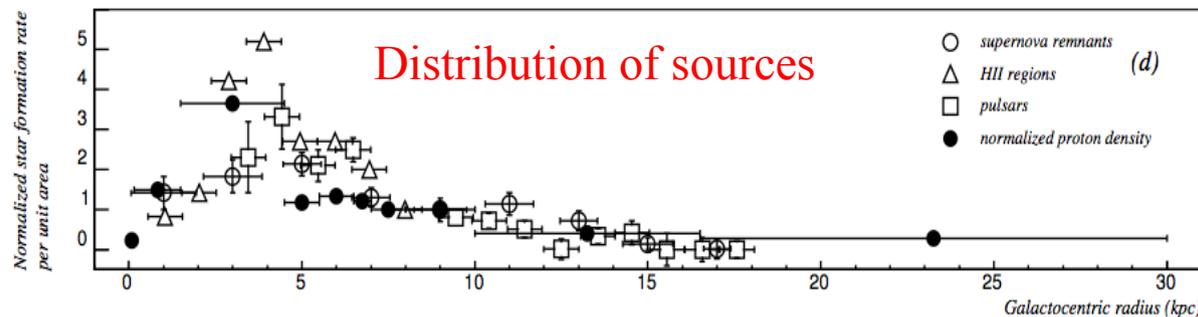
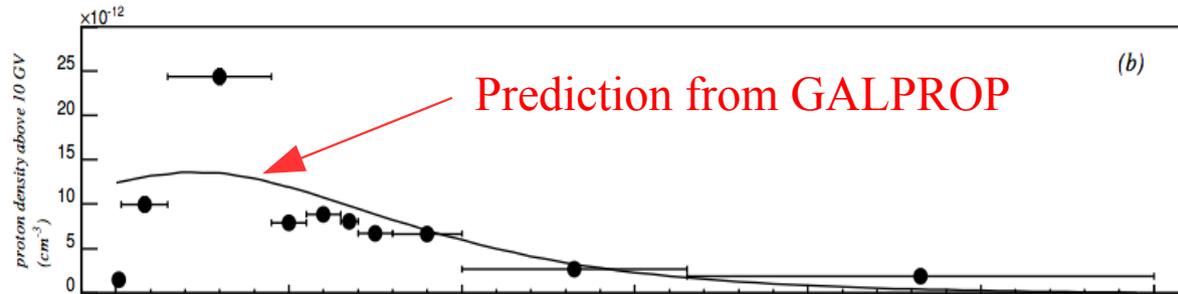
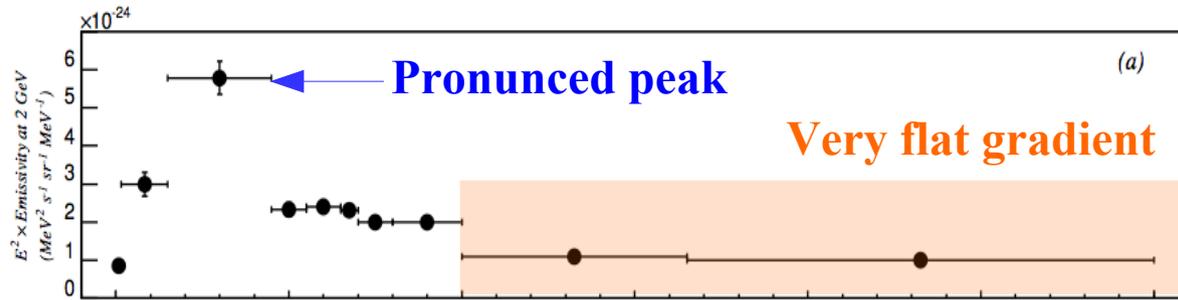


Recent results from FermiLAT collaboration on the CR distribution in the Galactic plane

[Acero et al. arXiv:1602.07246]

- In the outer region ($R > 8\text{kpc}$) the CR density at $\sim 20\text{ GeV}$ is flat (i.e. decreases much slower than the source distribution)
- In the inner region the CR density has a peak at $\sim 3\text{ kpc}$

3rd anomaly: the cosmic ray gradient in the Galactic plane



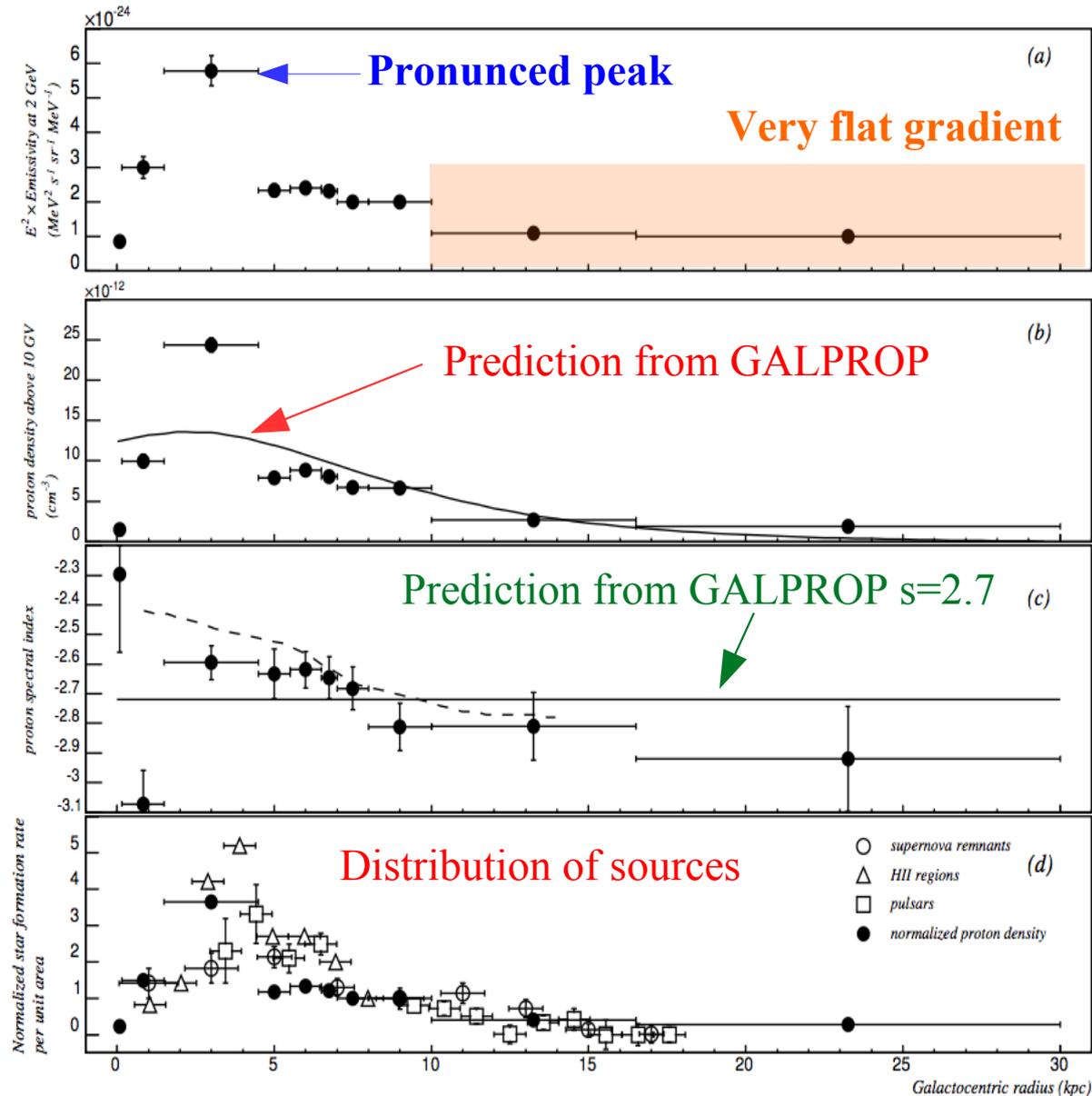
Recent results from FermiLAT collaboration on the CR distribution in the Galactic plane

[Acero et al. arXiv:1602.07246]

- In the outer region ($R > 8\text{kpc}$) the CR density at $\sim 20\text{ GeV}$ is flat (i.e. decreases much slower than the source distribution)

- In the inner region the CR density has a peak at $\sim 3\text{ kpc}$

3rd anomaly: the cosmic ray gradient in the Galactic plane



Recent results from FermiLAT collaboration on the CR distribution in the Galactic plane

[Acero et al. arXiv:1602.07246]

- In the outer region ($R > 8\text{kpc}$) the CR density at ~ 20 GeV is flat (i.e. decreases much slower than the source distribution)

- In the inner region the CR density has a peak at ~ 3 kpc

- The slope @ 20 GeV is not constant

This scenario is difficult to accommodate in a standard diffusion model

Possible solutions

In the context of leaky-box model several solutions have been proposed:

- Extended halo, $H > 4$ kpc
(Dogiel, Uryson, 1988; Strong et al., 1988; Bloemen, 1993, Ackerman et al., 2011)
 - ^ predicts a flat spectrum (but not flat enough)
 - ^ cannot explain the density bump in the inner Galaxy
- Flatter distribution of SNR in the outer Galaxy
(Ackerman et al., 2011)
- Enhancement of CO/H₂ density ratio (X_{CO}) in the outer Galaxy (Strong et al., 2004)
- Injection dependence on the ISM temperature
(Erlykin et al., 2015)
- Advection effects due to the Galactic wind
(Bloemen, 1993; Breitschwerdt, Dogiel, Voelk, 2002)

None of these ideas can simultaneously account for all signatures

- flatness $R > 8$ kpc,
- peak at $R \sim 3-4$ kpc,
- variation in the slope

CAN SELF-GENERATED DIFFUSION EXPLAIN THE OBSERVATIONS?

1-D slab model with self-generated turbulence

[Recchia, Blasi, GM, MNRAS 462, 2016]

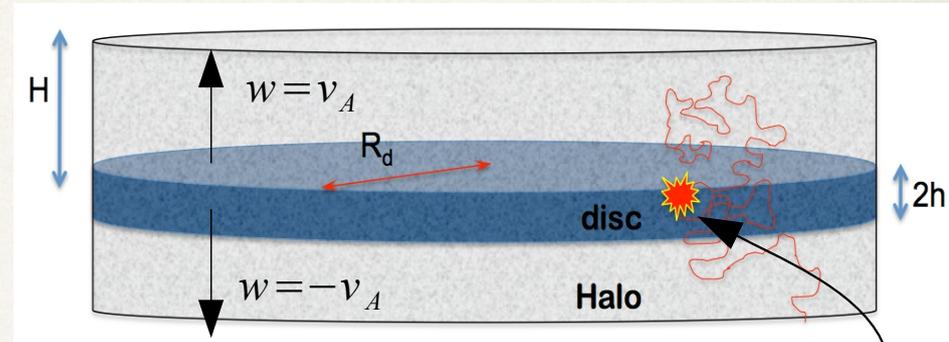
- CR escaping from the Galactic plane produce magnetic turbulence through resonant streaming instability

$$\Gamma_{CR} = \frac{v_A}{B_0^2/8\pi} \frac{1}{F(k)} \frac{\partial P_{CR}(>p)}{\partial z}$$

- Non-linear damping $\Gamma_{NLD} = v_s k F(k)$

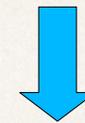
$$\Gamma_{NLD} = \Gamma_{CR}$$

- CRs are also advected by the global motion of the waves at the Alfvén speed



Spectrum injected at the disk

$$Q_0(p, R) \propto N_{SNR}(R) p^{-\gamma}$$



Propagated spectrum in the disk

$$f_{disk}(p) \propto p^7 \left(\frac{Q_0(p, R)}{B_0(R)} \right)^{s(p)} ;$$

$s = 1 \div 3$



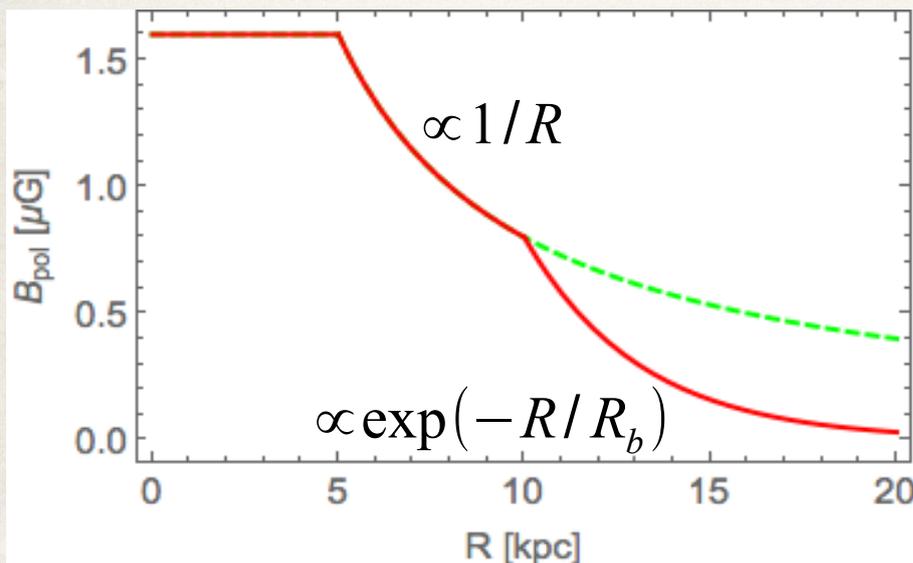
Self-generated turbulence and the gradient problem

[Recchia, Blasi, GM, MNRAS 462, 2016]

Self-generated turbulence could explain the gradient and the spectral index changes because it is more effective where B is smaller

- ^ less effective in the inner Galaxy
- ^ more effective in the outer Galaxy

Strength of large scale magnetic field in the Galaxy – poloidal component



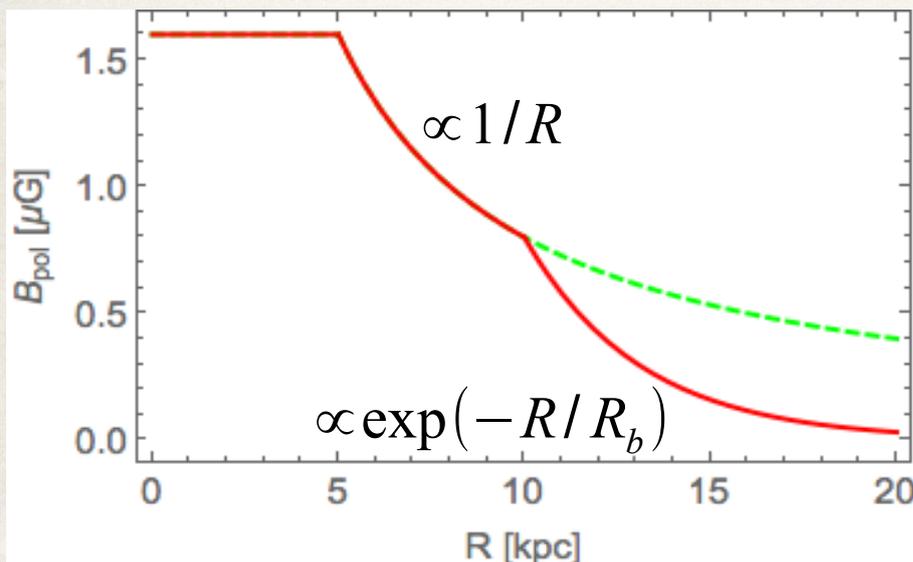
Self-generated turbulence and the gradient problem

[Recchia, Blasi, GM, MNRAS 462, 2016]

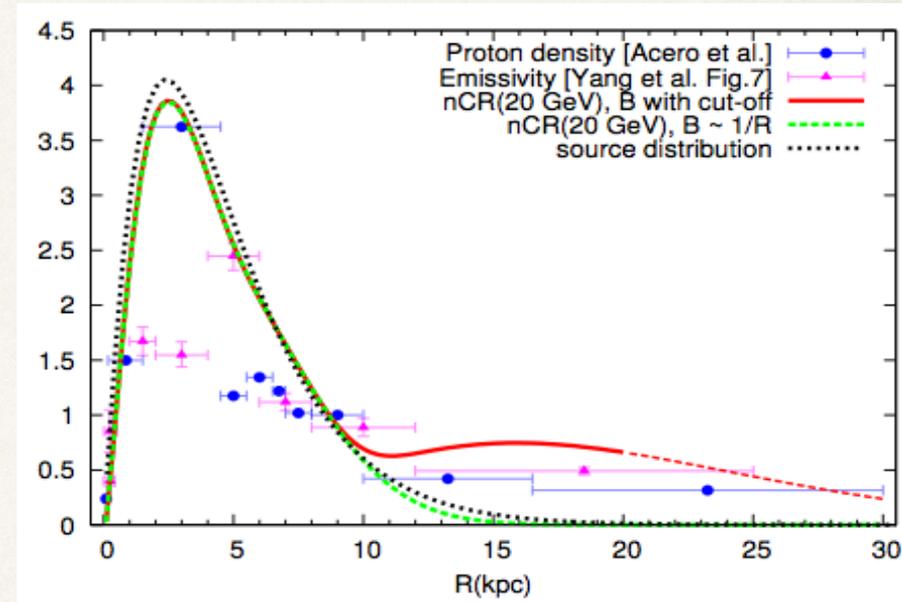
Self-generated turbulence could explain the gradient and the spectral index changes because it is more effective where B is smaller

- ^ less effective in the inner Galaxy
- ^ more effective in the outer Galaxy

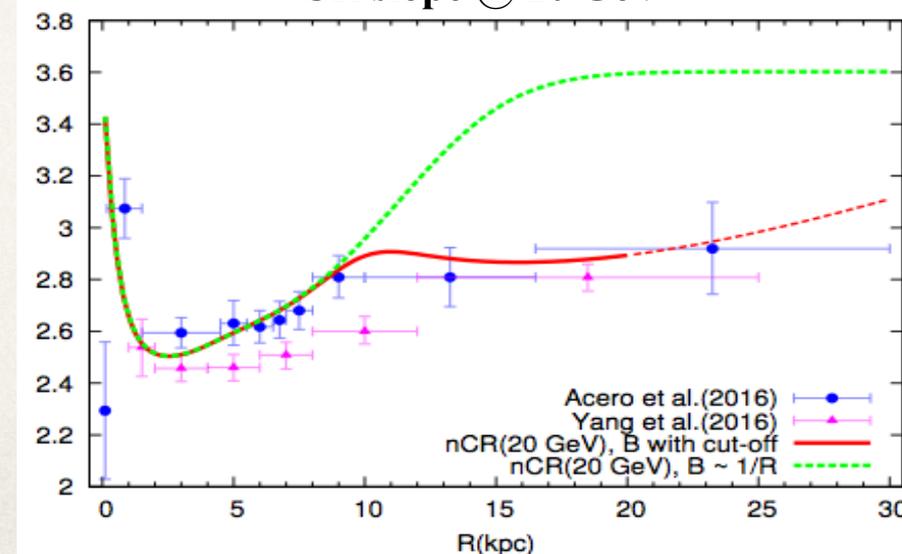
Strength of large scale magnetic field in the Galaxy – popoidal component



CR spectrum density @ 20 GeV



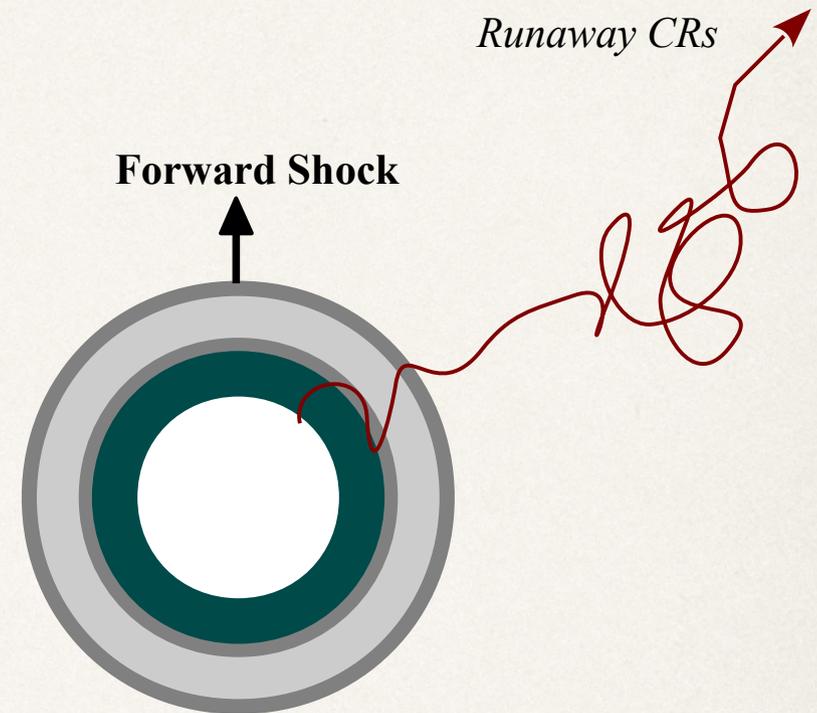
CR slope @ 20 GeV



Effect of self-amplification near the CR sources

During the process of escaping, CR can excite magnetic turbulence (via **streaming instability**) that keep the CR close to the SNR for a long time

[Malkom et al. (2013)
Nava et al. (2015)
D'Angelo et al.(2016)]



Effect of self-amplification near the CR

sources

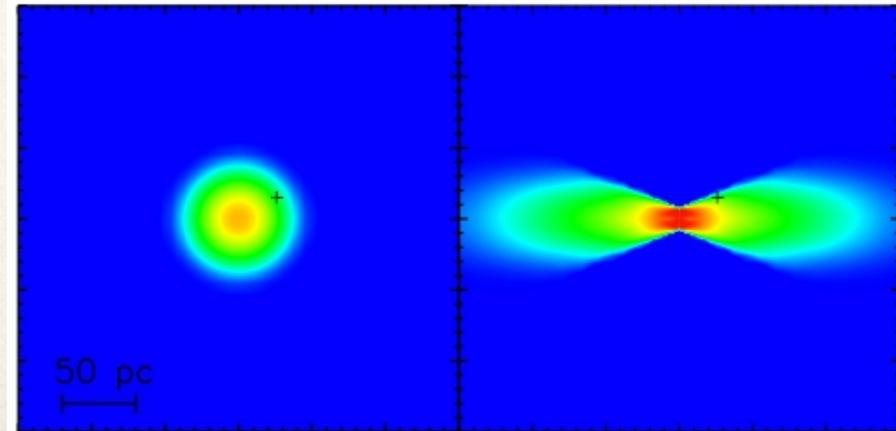
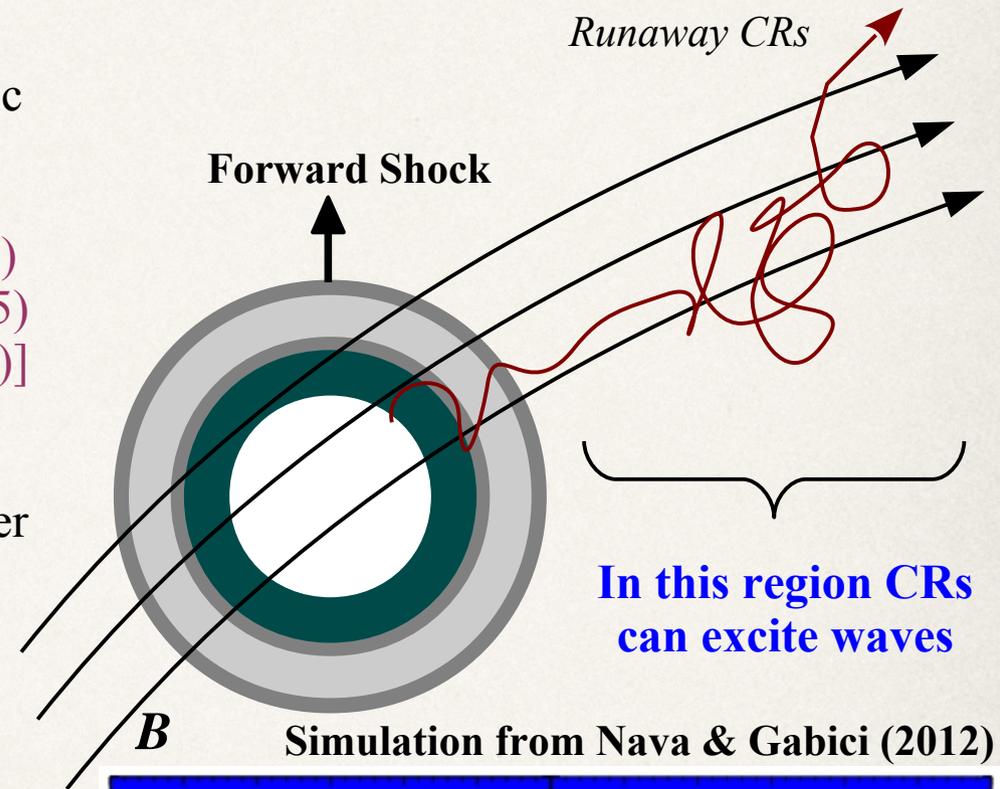
During the process of escaping, CR can excite magnetic turbulence (via **streaming instability**) that keep the CR close to the SNR for a long time

[Malkom et al. (2013)
Nava et al. (2015)
D'Angelo et al.(2016)]

The region where this can happen is at most of the order of the coherence-length of the magnetic field (after this distance the diffusion becomes 3D and the CR density drops rapidly below the average Galactic value)

During the time CRs spend in the vicinity of sources they can:

- 1) accumulate grammage
- 2) produce diffuse emission via $\pi^0 \rightarrow \gamma \gamma$



Effect of self-amplification near the CR sources: basic equations

[D'Angelo, GM, Amato, Blasi, 2018]

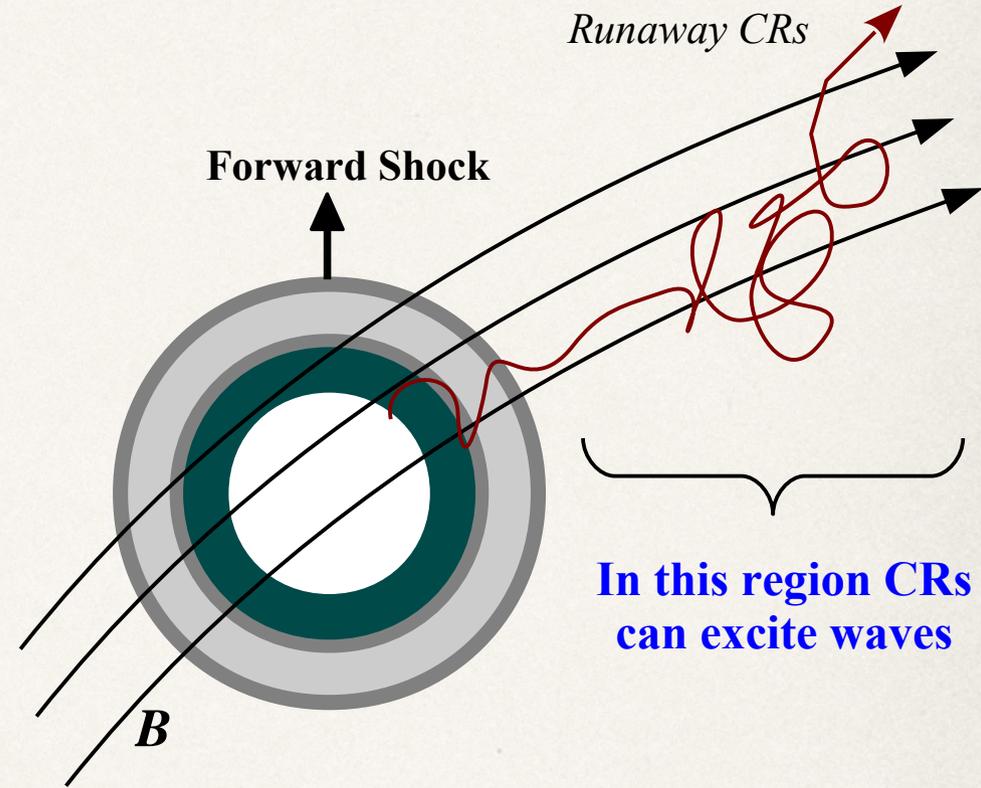
CR transport equation in 1-D

$$\frac{\partial f}{\partial t} + v_A \frac{\partial f}{\partial z} = \frac{\partial}{\partial z} \left[D \frac{\partial f}{\partial z} \right]$$

Self-generated diffusion coefficient

$$D(p, z, t) = \frac{r_L v}{3} \frac{1}{\mathcal{F}(k, z, t)} \Big|_{k=1/r_L}$$

$$\frac{\delta B^2}{B_0^2} = \int \mathcal{F}(k) \frac{dk}{k} \quad \text{Turbulence spectrum}$$



Transport equation for magnetic turbulence

$$\frac{\partial \mathcal{F}}{\partial t} + v_A \frac{\partial \mathcal{F}}{\partial z} = (\Gamma_{CR} - \Gamma_D) \mathcal{F} + Q_w$$

Resonant straming

Injection

Damping mechanisms

- Non linear damping
- Ion-neutral damping

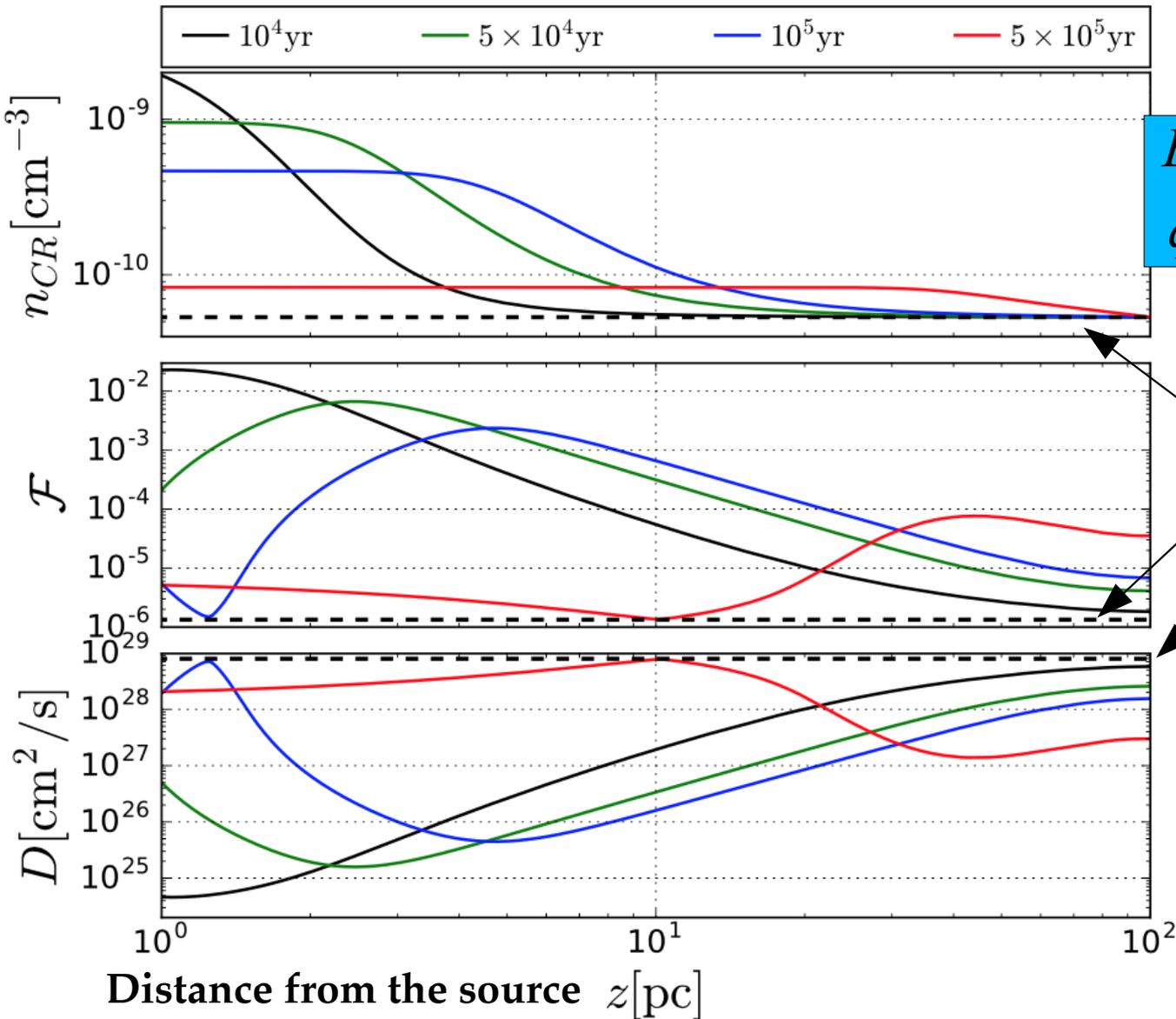
Evolution of CR density close to the source

[D'Angelo, GM, Amato, Blasi, 2018]

CR distribution function @ 10 GeV for several ages

Distribution function of turbulence

Diffusion coefficient



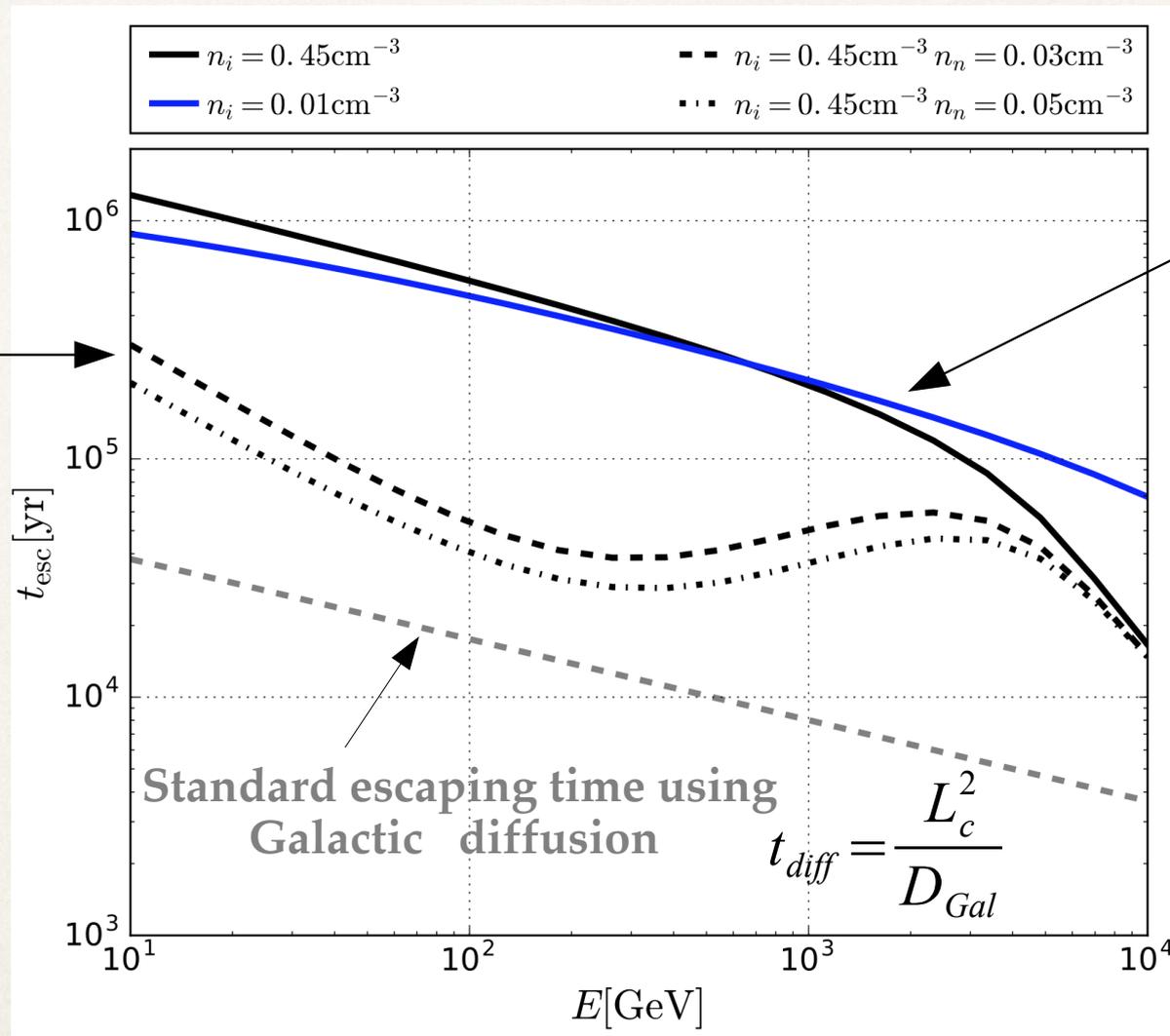
$E_{CR} = 0.2 E_{SN}$
 $q_{inj}(p) \propto p^{-4}$

Average Galactic level

Escaping time from the near-source region

[D'Angelo, GM, Amato, Blasi, 2018]

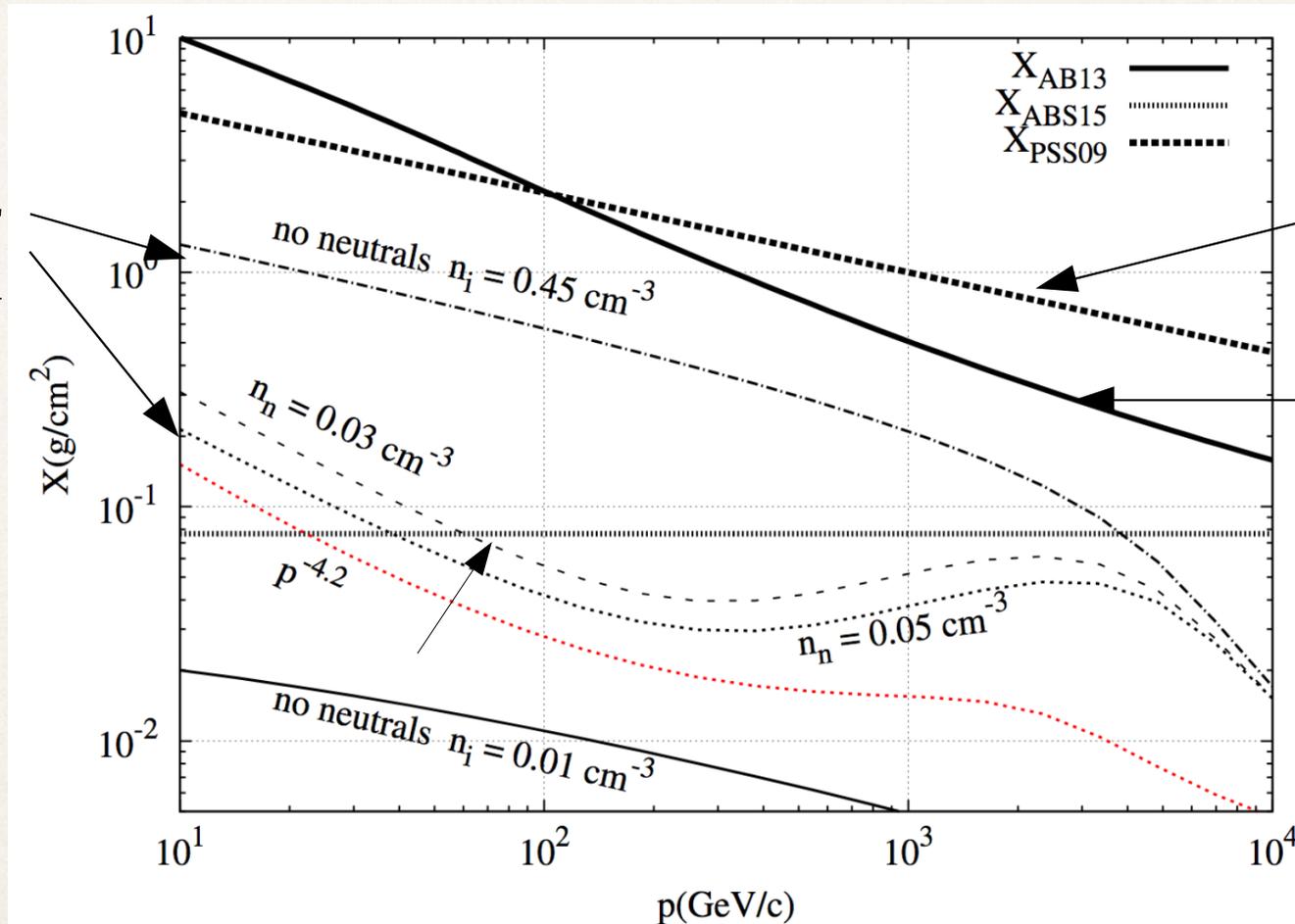
Assuming injected spectrum p^{-4}
 CR energy $\sim 20\% E_{\text{SN}}$



Grammage in the near-source region

[D'Angelo, Blasi, Amato 2016, *PRD*]

Assuming injected spectrum p^{-4}
CR energy $\sim 20\% E_{\text{SN}}$



Grammage from the near source region

Grammage inferred from B/C:

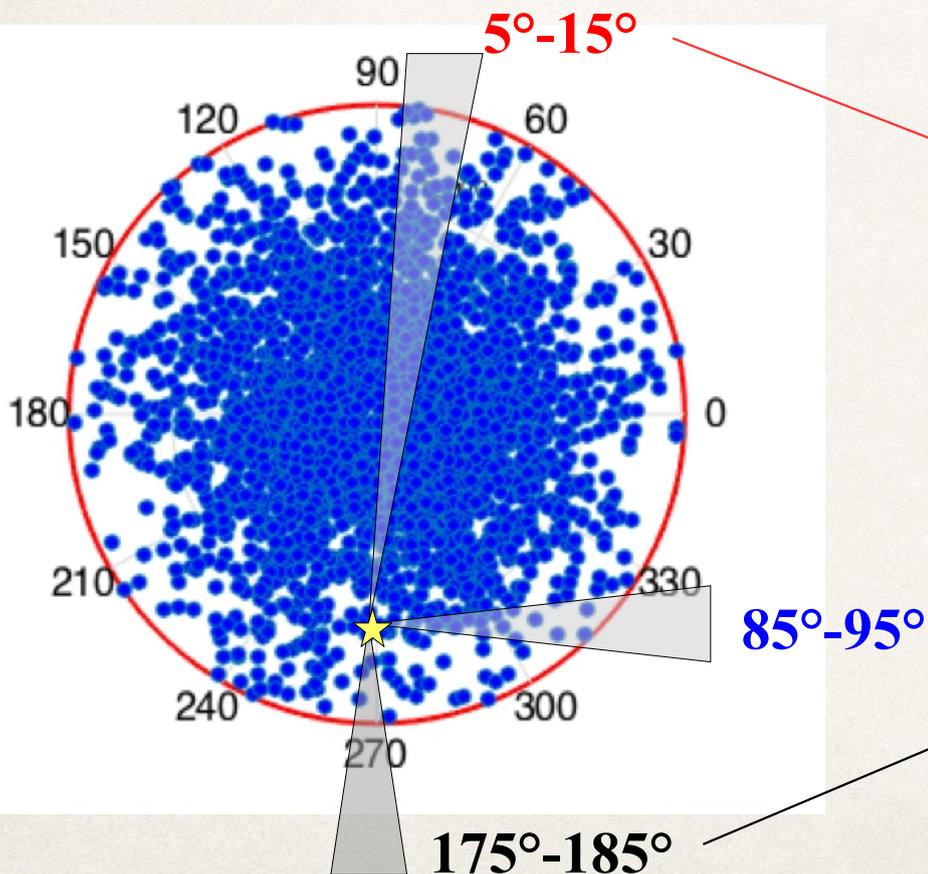
1) Standard halo model

2) Non-linear propagation

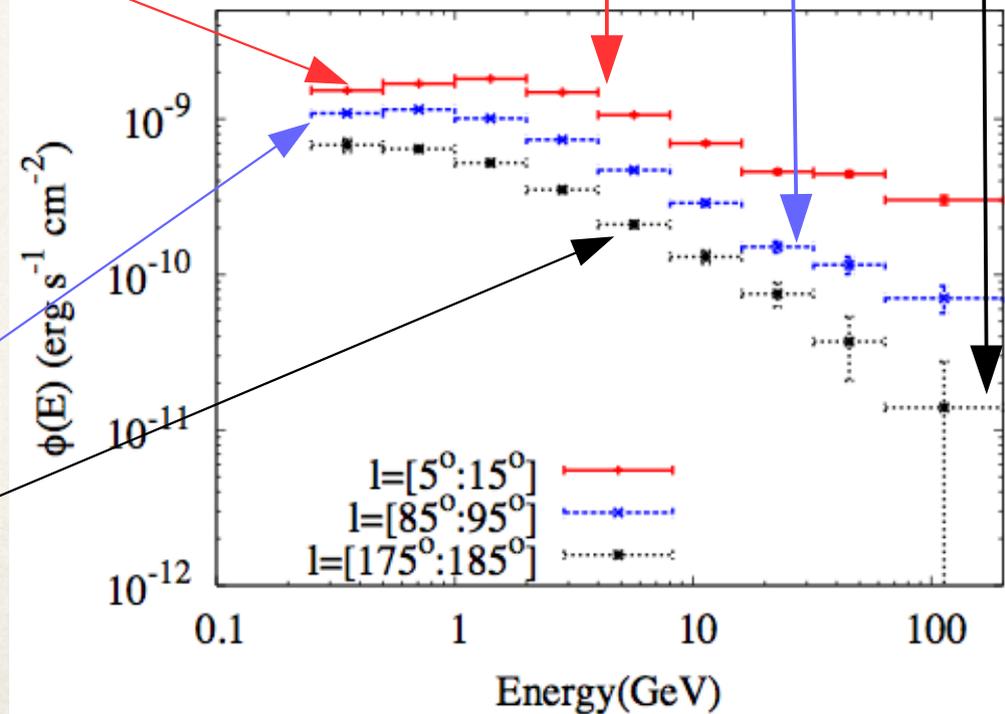
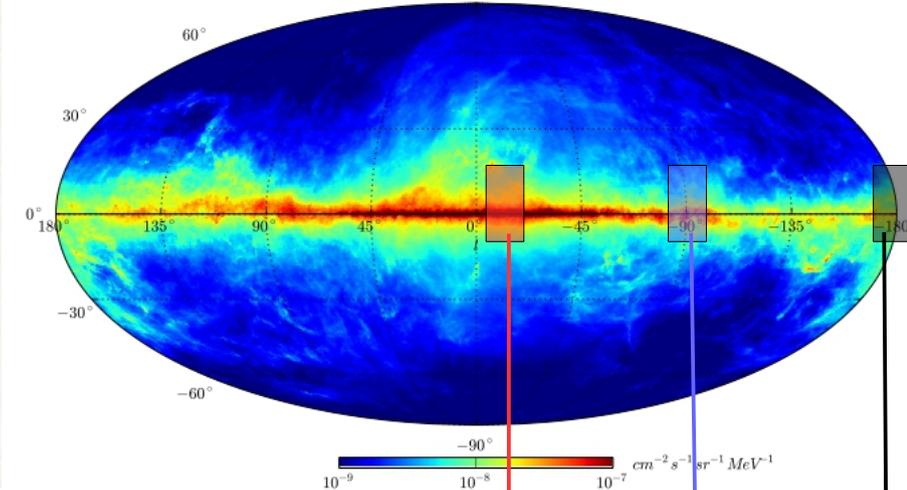
Diffuse Galactic emission from *Fermi*-LAT

[D'Angelo, GM, Amato, Blasi, 2018]

Diffuse Galactic γ -ray flux for three different angular sectors extracted from the *Fermi*-LAT data
[Yang, Aharonian & Evoli, 2016]



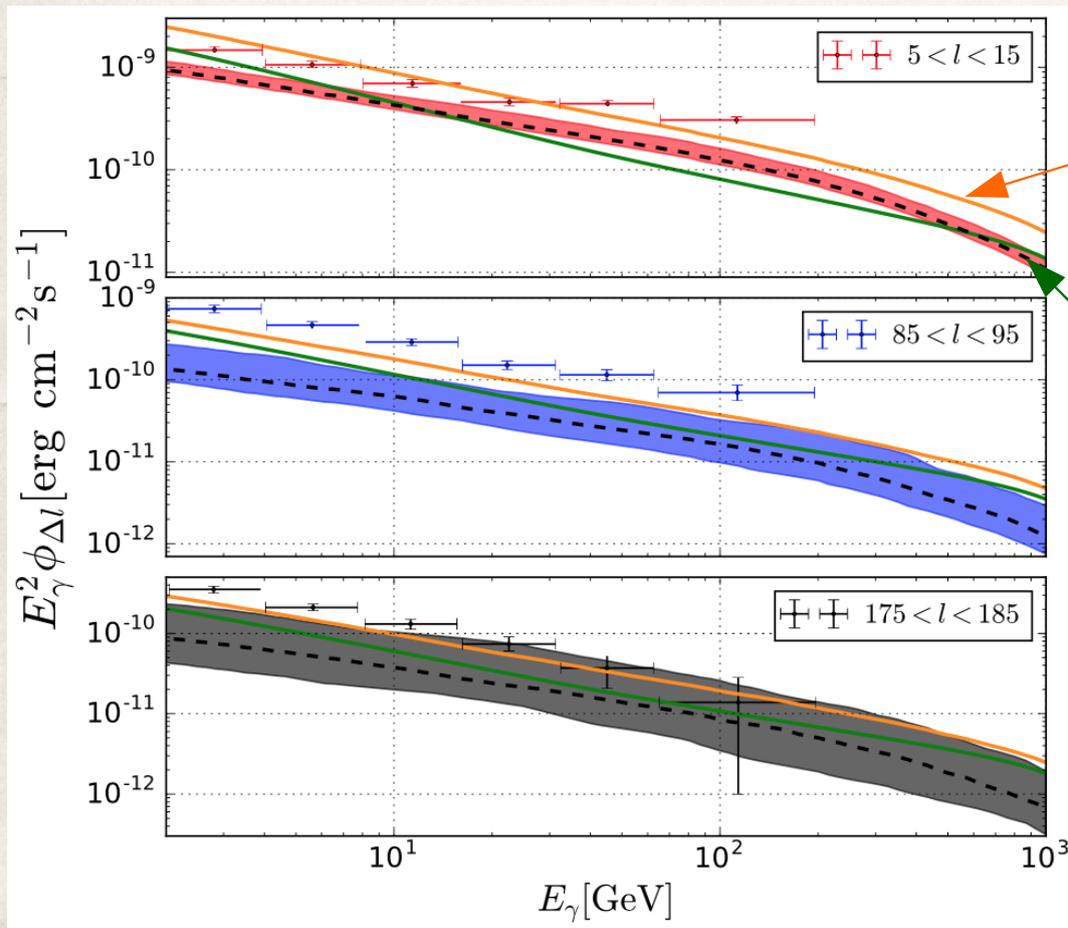
FermiLAT diffuse emission



Contribution of the escaping CRs to the diffuse Galactic emission

[D'Angelo, GM, Amato, Blasi, 2018]

$$n_i = 0.45 \text{ cm}^{-3}; \quad n_H = 0.0 \text{ cm}^{-3}$$



Sum of diffuse emission plus contribution from all the source cocoons

“Real” diffuse contribution assuming AMS spectrum in the whole Galaxy

Contribution of the escaping CRs to the diffuse Galactic emission

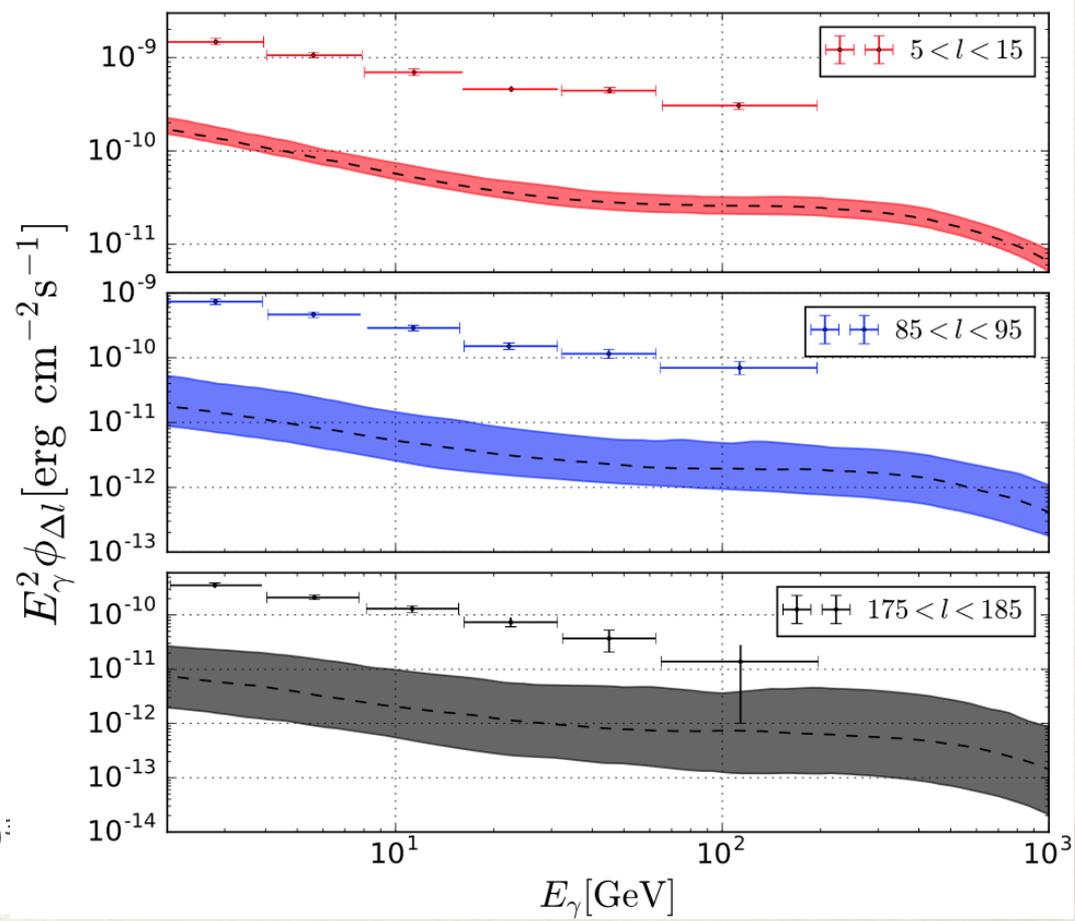
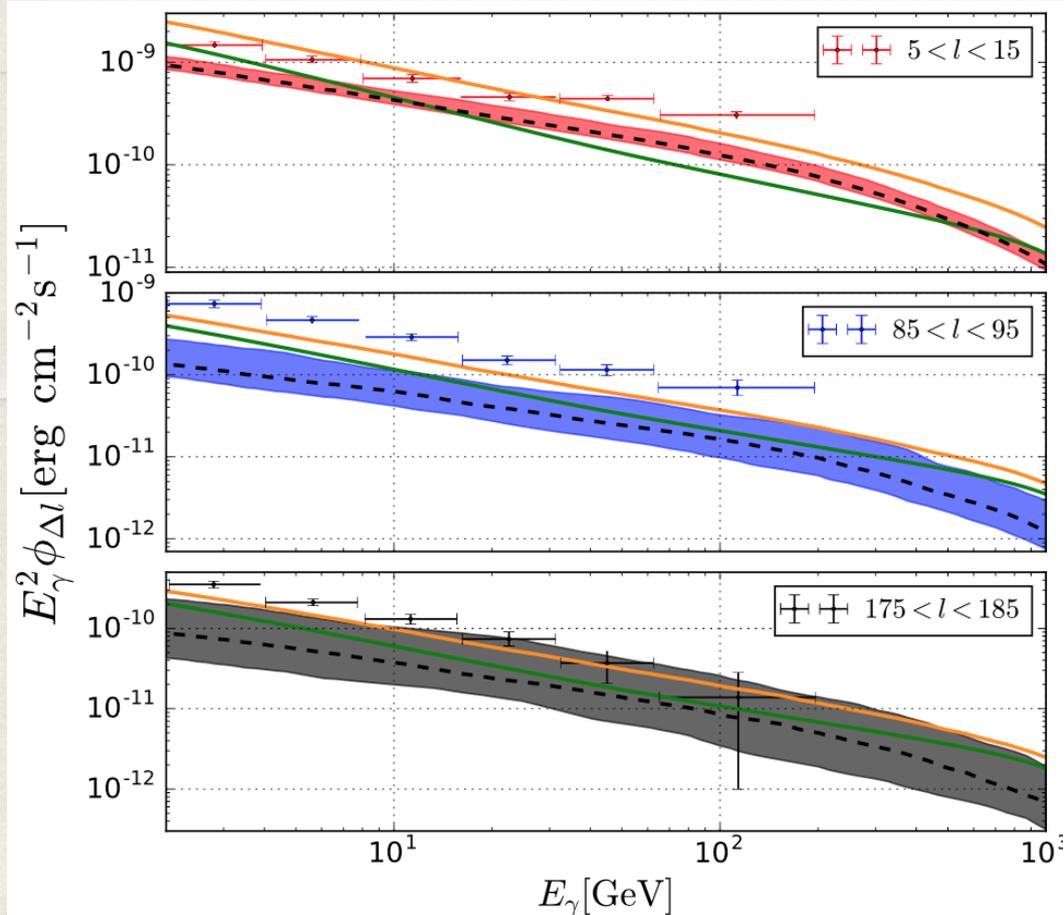
[D'Angelo, GM, Amato, Blasi, 2018]

Number of sources contributing to the emission

Angular sector	Fully ionized	$n_H=0.05$
$5^\circ-15^\circ$	4500	740
$85^\circ-95^\circ$	350	57
$175^\circ-185^\circ$	77	13

$n_i=0.45 \text{ cm}^{-3}$; $n_H=0.0 \text{ cm}^{-3}$

$n_i=0.45 \text{ cm}^{-3}$; $n_H=0.05 \text{ cm}^{-3}$



Presence of Molecular Clouds near the source

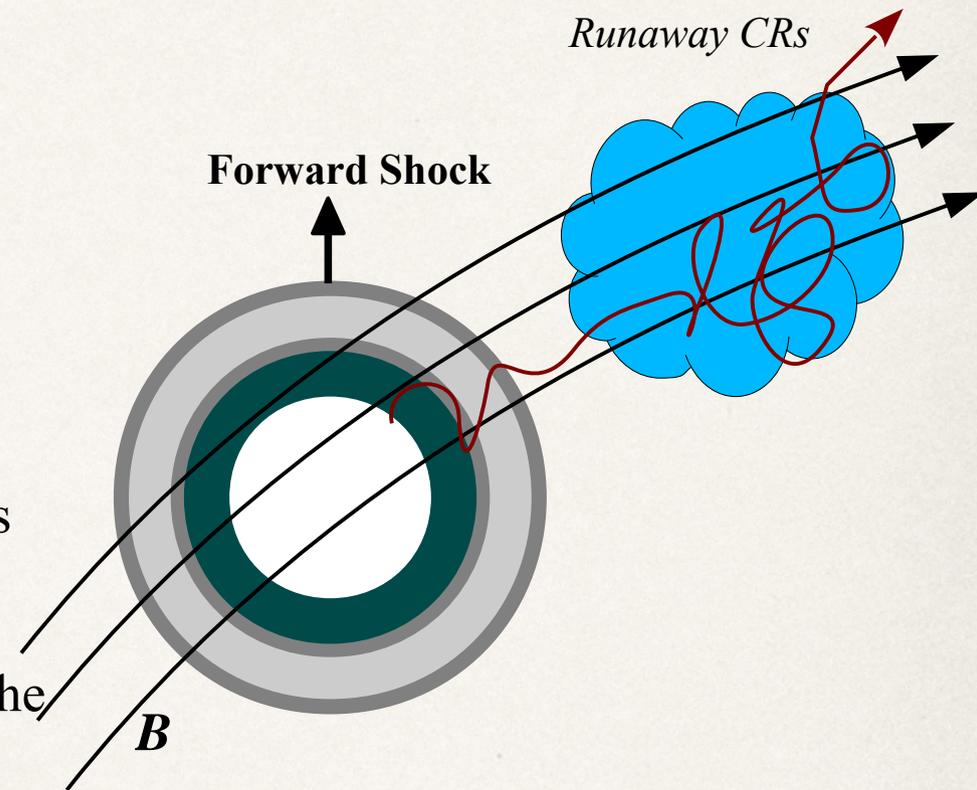
If a molecular cloud is close to the SNR both the **gamma-ray emission and the grammage accumulated can be enhanced**

- Several SNR-MC association are known from gamma-ray observation
- CTA will probably discover tens of such objects

Is it possible to accumulate all the grammage in the near source region?

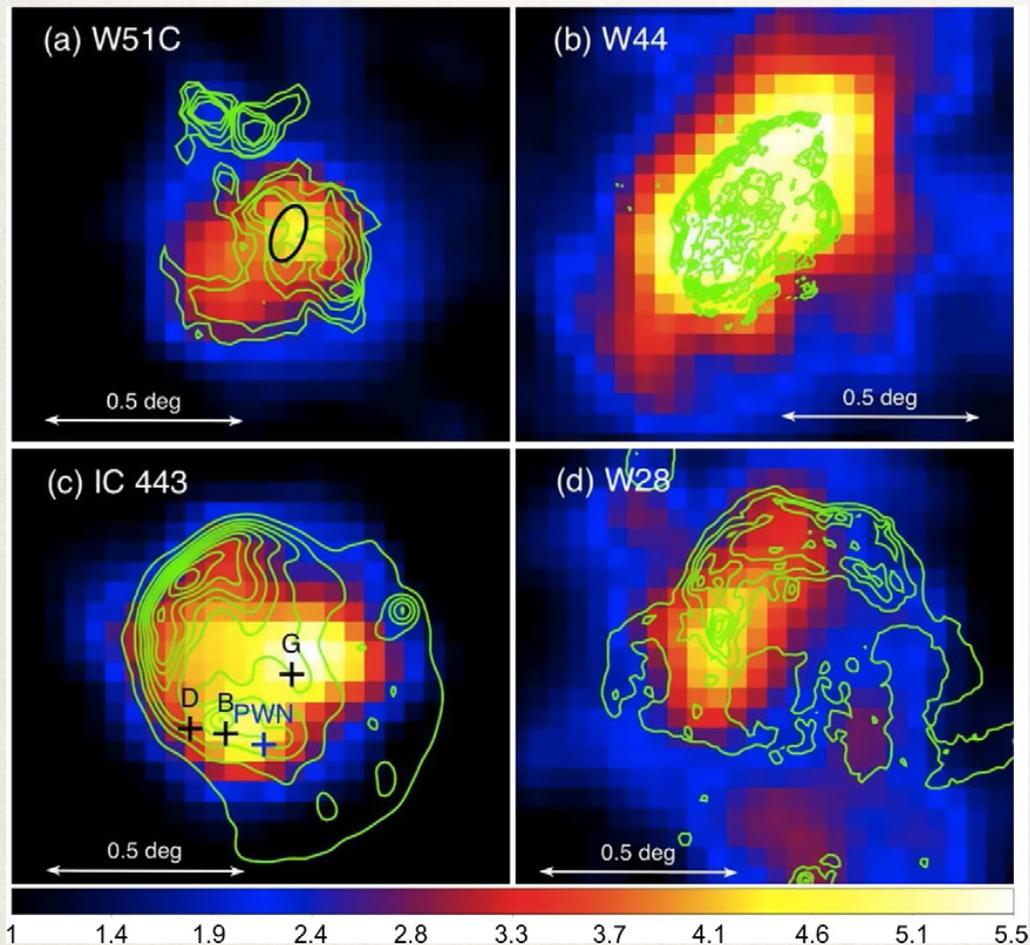
This possibility would support unorthodox models for the CR propagation which require fast escape from the Galaxy

- **Cowsik et al.(1975, 2016) “*Nested Leaky-box*”**
- **Lipari (2016)**



Using molecular clouds as CR barometer

Examples of γ -ray emission from clouds close or interacting with SNRs - [*Fermi*-LAT]



OBSERVATIONS of MCs in γ -RAYS:

- CRs interact inside MCs
$$pp \rightarrow \pi^0 \rightarrow \gamma\gamma$$
- strong emission in GeV range
- γ -emission sensible to CR energy $E > 280$ MeV

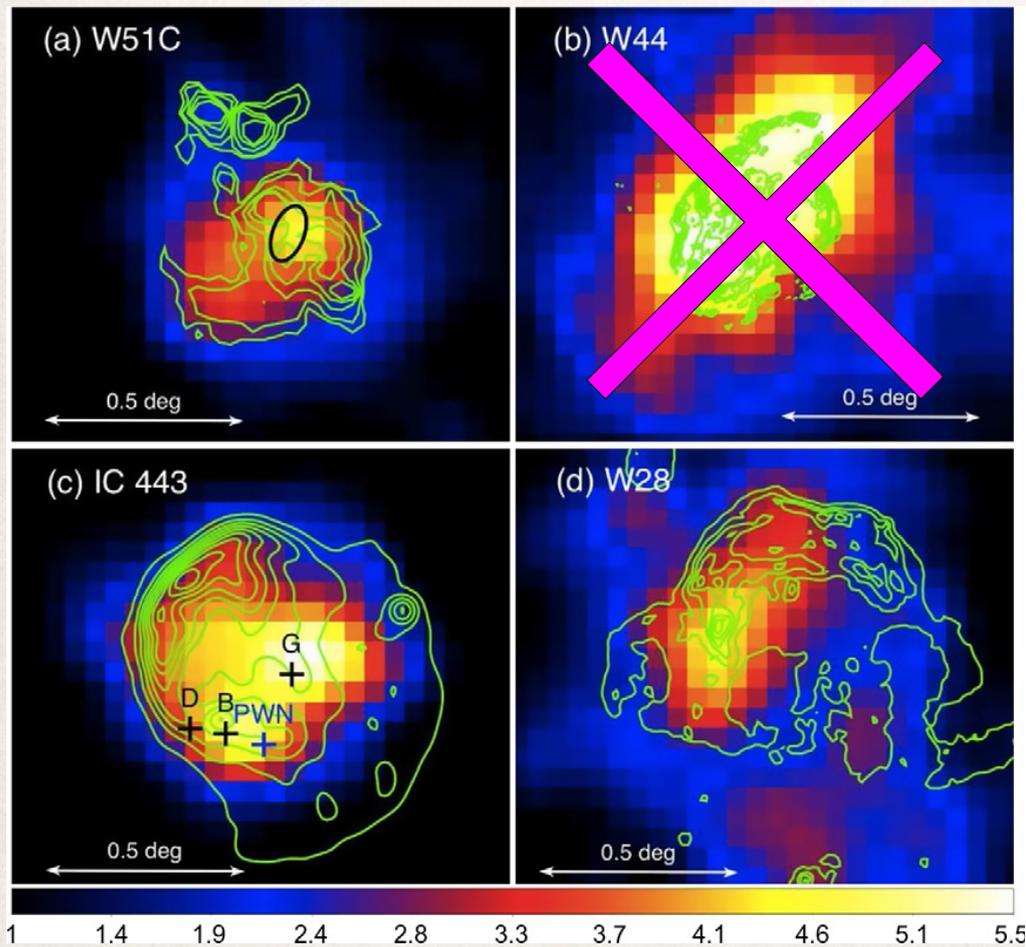
DETECTION OF IONIZATION

- The ionization rate of several molecules depends on the CR flux (H_2 , H_3^+ , CH, OH, C_2 , DCO^+ , HCO^+ ,.....)
- Ionization sensible to CR energy $E > 0.1$ MeV

Is it possible to use combined information from ionization and γ -ray emission to infer the CR spectrum from \sim MeV up to \sim TeV

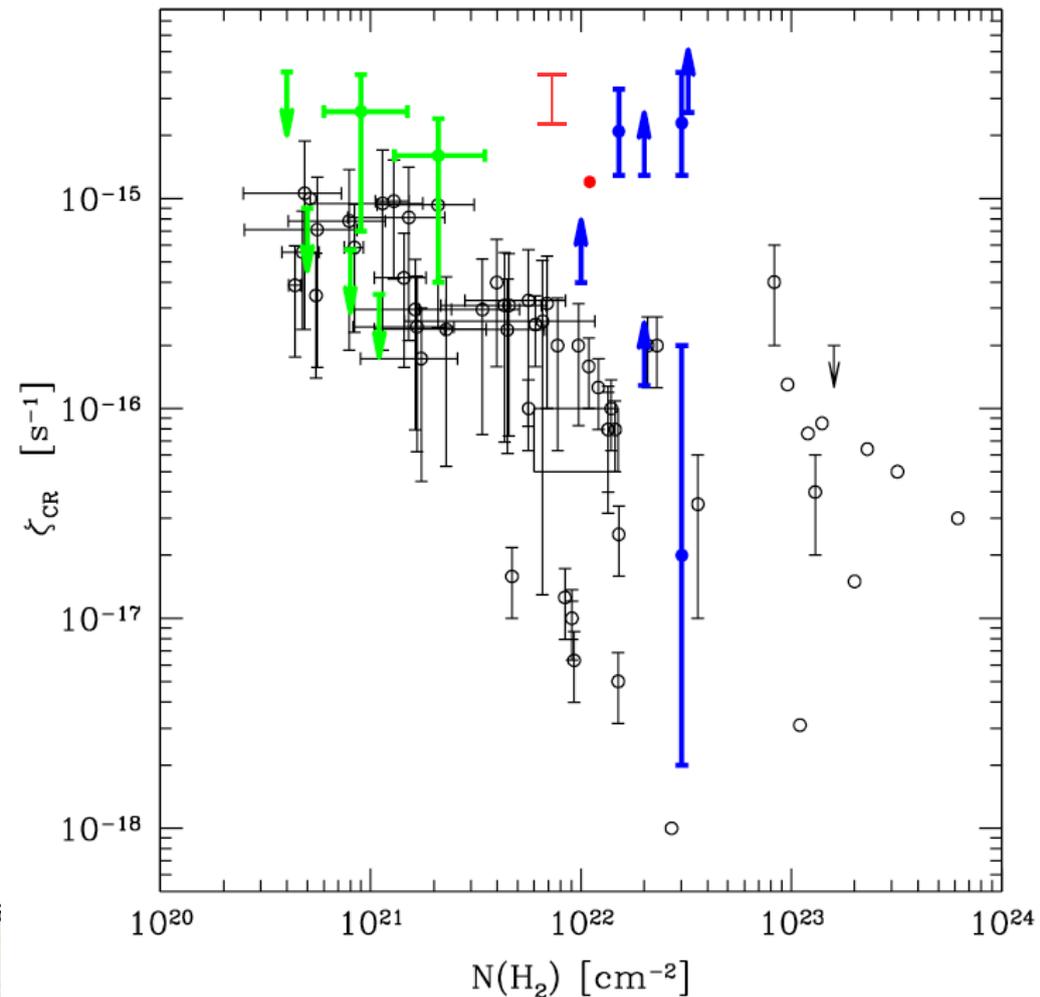
Using molecular clouds as CR barometer

Examples of γ -ray emission from clouds close or interacting with SNRs - [*Fermi*-LAT]



Average level of ionization is larger for MCs close to SNRs

IC 443 W 51C W28



Conclusions

CR self-generated turbulence seems to be an ubiquitous phenomenon

- ▶ **Acceleration at shocks** (not covered here)
- ▶ **Propagation:**
 - Can explain the spectral breaks (@300 GeV)
 - Can account for the gradient problem in the galactic plane
- ▶ **Escaping from sources** → SI increases the escaping time:
 - Increase the accumulated grammage (but not the dominant contribution)
 - Enhance the gamma-ray emission from the cocoon (possible contribution to the diffuse gamma-ray emission)
 - Effect on anisotropy? (see talk by Kachelriess on Friday)