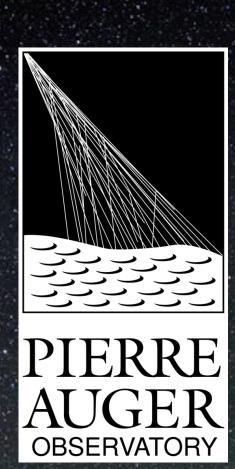


Interpreting the cosmic ray spectrum and composition measurements across the ankle and up to the highest energies with the data of the Pierre Auger Observatory

Eleonora Guido*, on behalf of the Pierre Auger Collaboration

UHECR2022: 6TH INTERNATIONAL SYMPOSIUM ON ULTRA HIGH ENERGY COSMIC RAYS

6th October 2022 GSSI, L'Aquila, Italy





*Universität Siegen, Siegen, Germany



Combined fit of the Pierre Auger Observatory measurements (spectrum and composition) at ultra-high-energy (UHE) • Combined fit above 10^{18.7} eV (above the ankle) already published¹

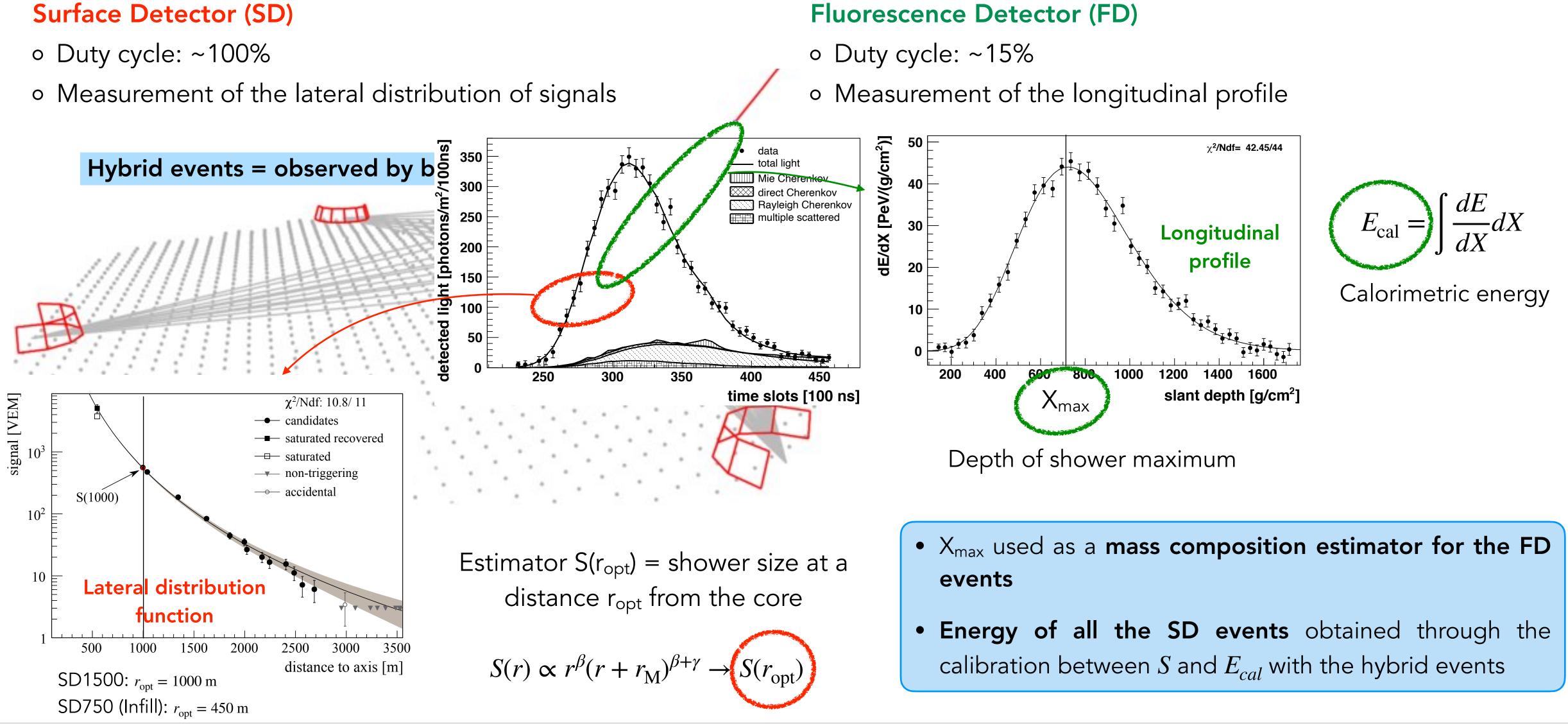
- - \rightarrow extension to low energies to include the ankle feature
- Preliminary results already shown at ICRC2021²
- Paper by the Pierre Auger Collaboration to be soon submitted to a journal







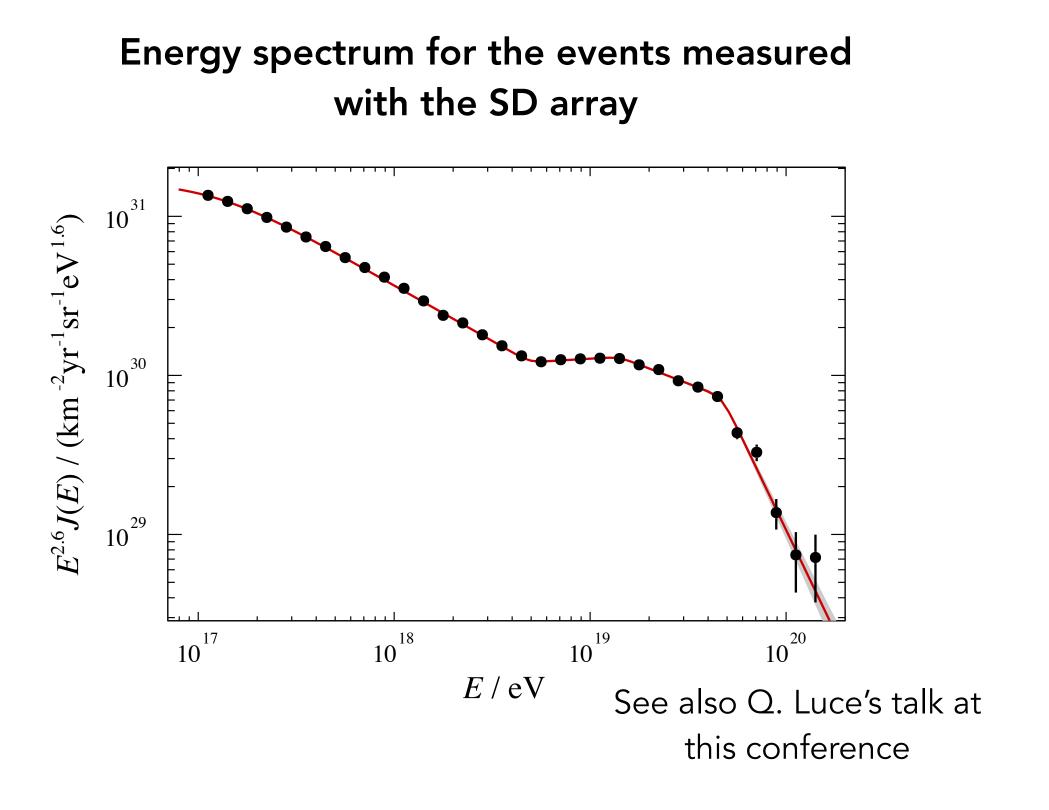
Measurements of the energy and mass composition



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Energy spectrum and mass composition measurements



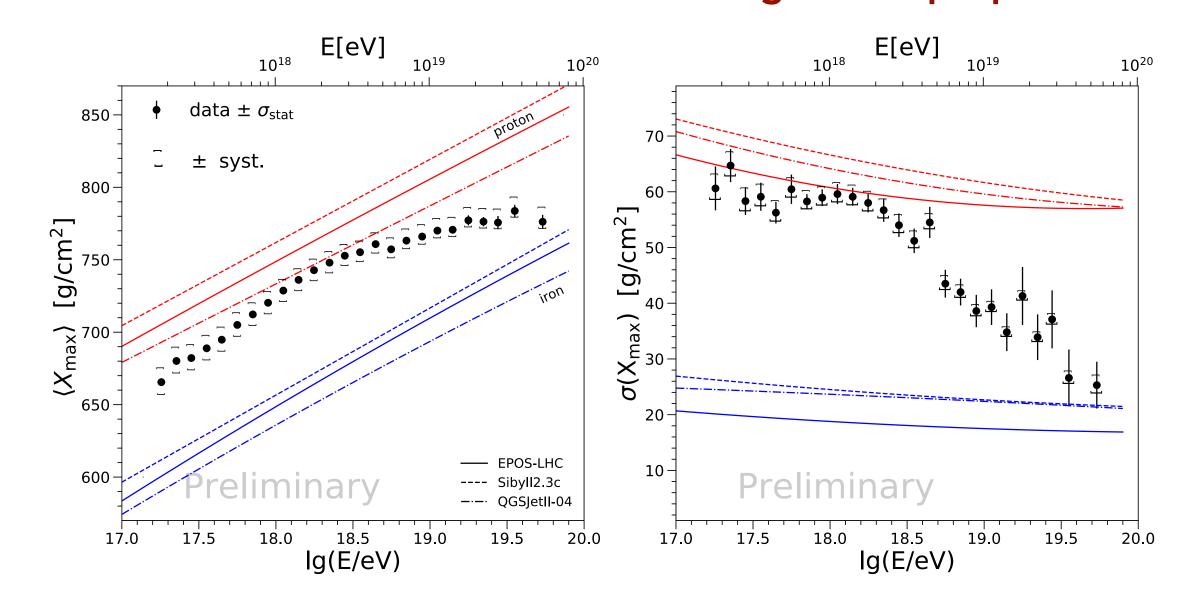
Data in $log_{10}(E/eV)$ bins of 0.1 width:

- Energy spectrum up to 10^{20.2} eV
- * X_{max} distributions: up to 10^{19.7} eV (+ 1 additional bin for events above), binned in intervals of X_{max} of 20 g cm⁻²

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The X_{max} distribution in each energy bin is sensitive to the mass composition \rightarrow first two moments shown for figurative purposes





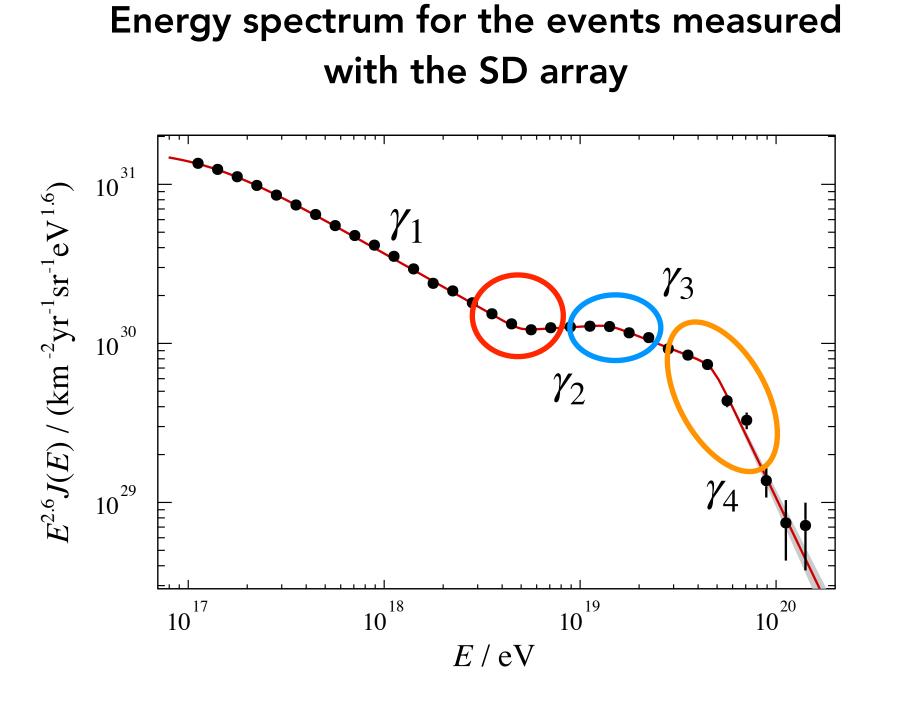


The Pierre Auger Collaboration, Eur. Phys. J. C 81, 966 (2021) A.Yushkov for the Pierre Auger Collaboration PoS(ICRC2019)482





Energy spectrum and mass composition measurements

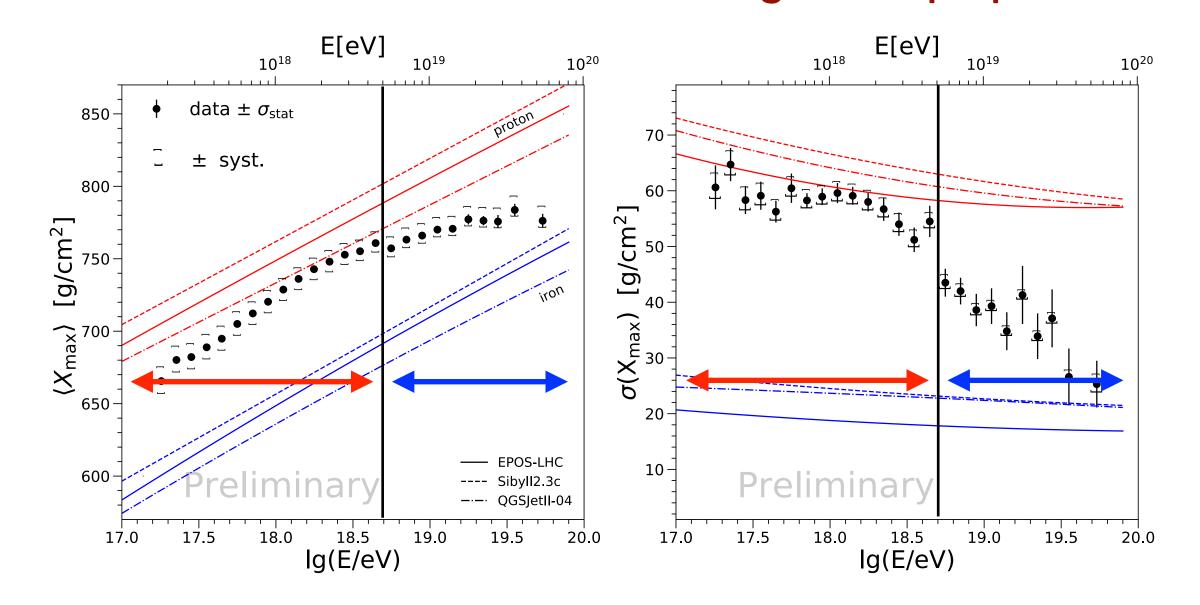


- + Hardening at ~6×10¹⁸ eV (ankle)
- + Recently observed softening at ~1×10¹⁹ eV (instep)
- + Suppression at $\sim 5 \times 10^{19} \text{ eV} \rightarrow \text{energy cut off}$

Propagation effect and/or maximum energy at the acceleration

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The X_{max} distribution in each energy bin is sensitive to the mass composition \rightarrow first two moments shown for figurative purposes



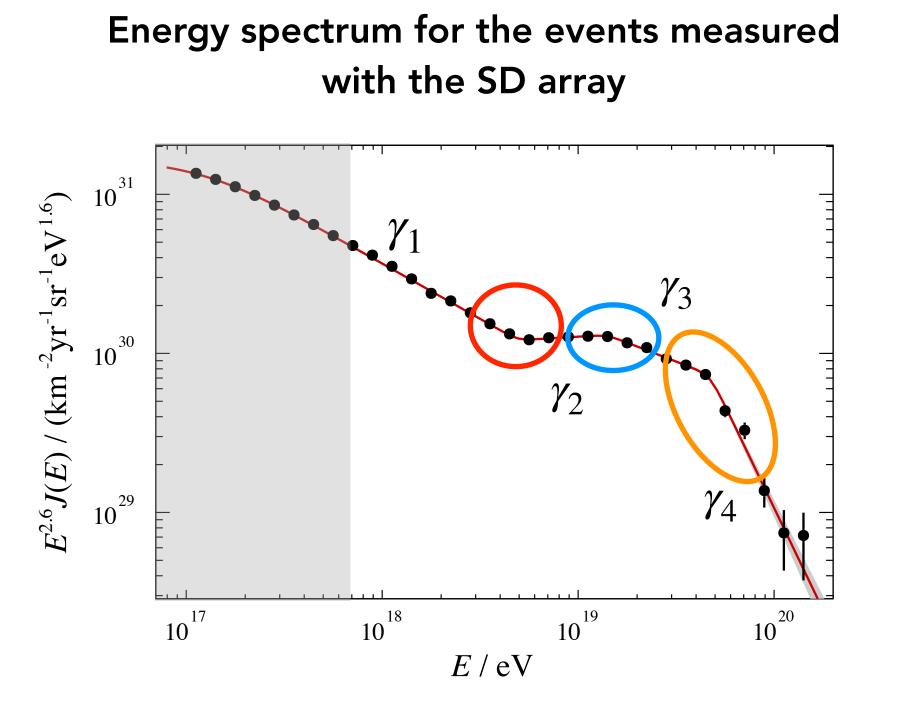
- Below the ankle: mass composition gets increasingly lighter
- At the ankle: mixed composition
- Above the ankle: increasingly heavier and less mixed
 - \rightarrow superposition of alternating and heavier groups of elements
 - \rightarrow increasingly sparse statistics up to ~10^{19.7} eV







Measurements of the energy spectrum and mass composition

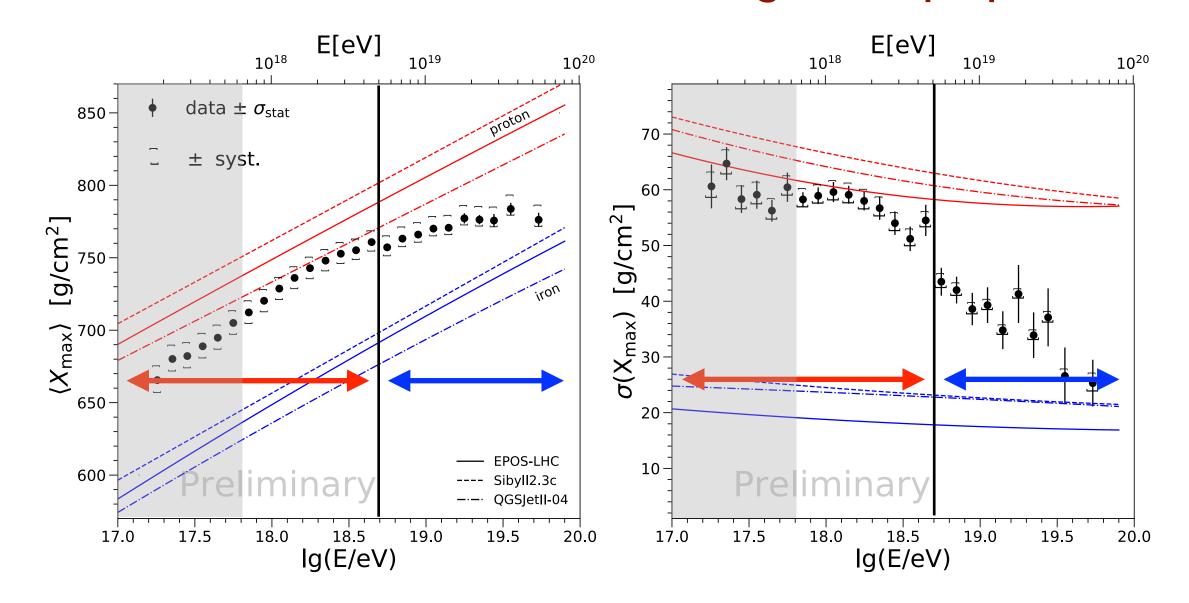


- We aim at including the ankle region
- We want to focus on the energy region the Galactic CRs are not dominant anymore

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The X_{max} distribution in each energy bin is sensitive to the mass composition \rightarrow first two moments shown for figurative purposes



Combining the information from the two data sets is crucial to interpret the features



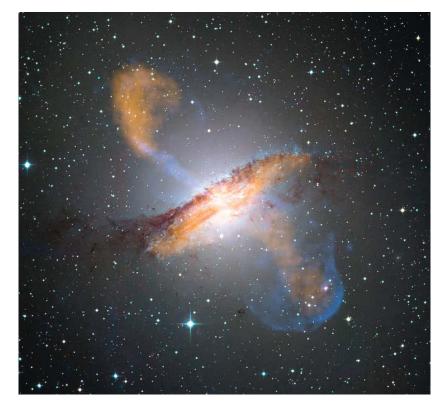
The Pierre Auger Collaboration, Eur. Phys. J. C 81, 966 (2021) A.Yushkov for the Pierre Auger Collaboration PoS(ICRC2019)482

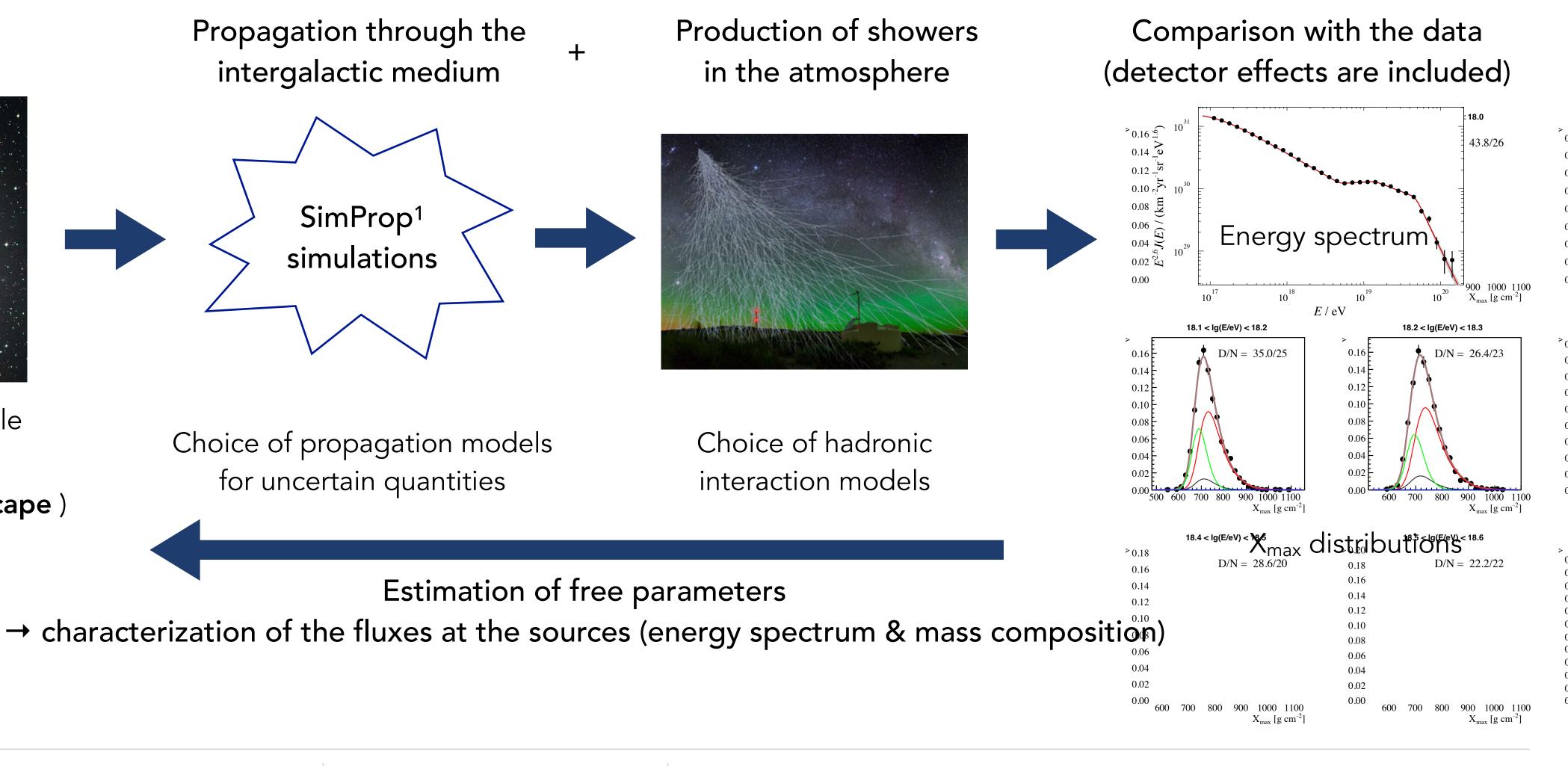




The combined fit

CRs ejected by EG accelerators





Assumptions on a simple astrophysical model (CRs considered at the escape)

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¹ R. Aloisio et al, JCAP 11 (2017) 009

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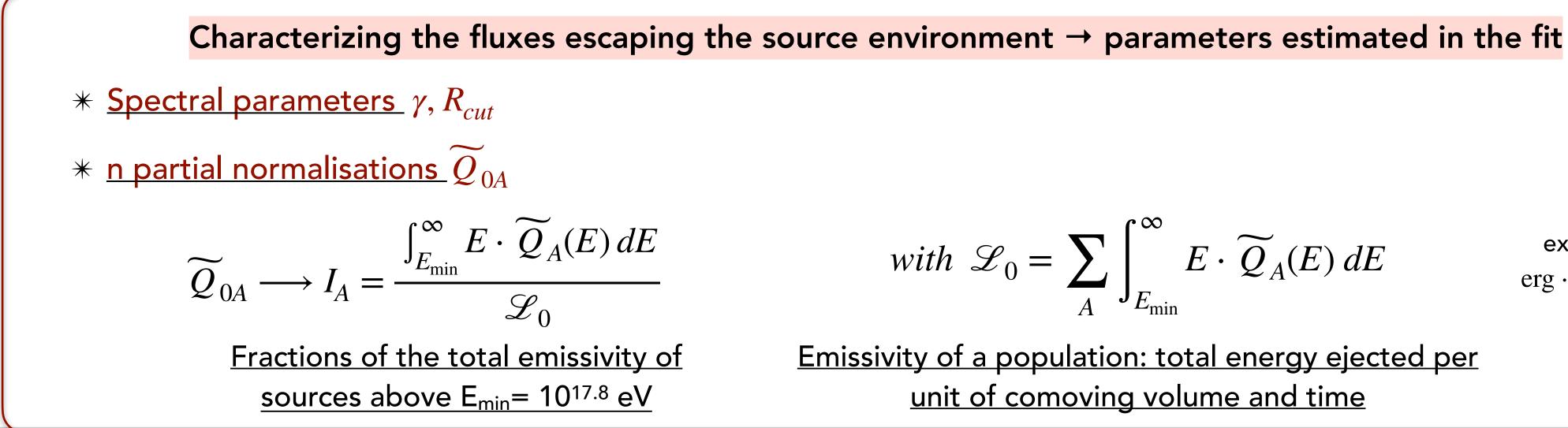
Astrophysical model

Generic population of extragalactic sources

- * population of identical sources
- * uniform distribution (except for a local overdensity for d < 30 Mpc)
- * ejection of n representative nuclear species A, chosen among ¹H, ⁴He, ¹⁴N, ²⁸Si, ⁵⁶Fe

Generation rate at the sources for each mass A (number of nuclei ejected per unit of energy, volume and time) :

$$\widetilde{Q}_{A}(E) = \widetilde{Q}_{0A} \cdot \left(\frac{E}{E_{0}}\right)^{-\gamma} \cdot \begin{cases} 1, & E \leq Z_{A} \cdot R_{\text{cut}}; \\ \exp\left(1 - \frac{E}{Z_{A} \cdot R_{\text{cut}}}\right), & E > Z_{A} \cdot R_{\text{cut}}, \end{cases}$$



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with
$$\mathscr{L}_0 = \sum_A \int_{E_{\min}}^{\infty} E \cdot \widetilde{Q}_A(E) dE$$

expressed in
$$erg \cdot Mpc^{-3} \cdot yr^{-1}$$

Emissivity of a population: total energy ejected per unit of comoving volume and time



Propagation model

Propagation through the IGM and the Earth's atmosphere

- SimProp simulations for the propagation in the IGM
 - Adiabatic energy losses (expansion of the Universe)
 - Interactions of nuclei with background photons (EBL,
 - Photo-pion production
 - Pair production
 - Photo-disintegration
- Hadronic interaction model for the propagation in the atmosphere
- 1D propagation \rightarrow intergalactic magnetic fields are here neglected

Model configuration used for our reference results:

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\rightarrow model for the photo-disintegration cross sections $\sigma_{\rm pd}$ \rightarrow model for the EBL spectrum and evolution

$$-\left(\frac{1}{E}\frac{dE}{dt}\right)_{ad} = H_0\sqrt{(1+z)^3\Omega_m + \Omega_\Lambda}$$

$$N + \gamma \rightarrow N + \pi^0 / N + \pi^{\pm}$$
$$N + \gamma \rightarrow N + e^+ + e^-$$
$$(A, Z) + \gamma \rightarrow (A - n, Z - n') + nN$$

Talys for σ_{pd} , Gilmore model for EBL, EPOS-LHC as hadronic interaction model



J. Koning et al., vol. 769 of American Institute of Physics Conference Series, 2005 R. Gilmore et al., Mon. Not. Roy. Astron. Soc., 422 (2012) 3189 T. Pierog et al., Phys. Rev. C 92 (2015) 034906





Fit procedure

Combined fit of the energy spectrum and X_{max} distributions above $\sim 6 \times 10^{17} eV$

→ compare simulated and measured fluxes at the Earth with the maximum likelihood method

$$D = D(J) + D(X_{\max}) = -2 \ln\left(\frac{\mathscr{L}}{\mathscr{L}^{\text{sat}}}\right) = -2 \ln\left(\frac{\mathscr{L}_{\text{J}}}{\mathscr{L}^{\text{sat}}}\right) - 2 \ln\left(\frac{\mathscr{L}_{X_{\max}}}{\mathscr{L}^{\text{sat}}_{X_{\max}}}\right)$$

Energy spectrum \rightarrow Gaussian distributions

$$L_{\rm J} = \prod_{i} \frac{1}{\sqrt{2\pi\sigma_i^2}} \exp\left(-\frac{(J_i^{\rm ob})}{\sqrt{2\pi\sigma_i^2}}\right)$$

• X_{max} distributions \rightarrow multinomial distributions

$$L_{X_{\max}} = \sum_{i} n_i^{\text{obs}}! \sum_{j} \overline{k_{i,j}^{\text{obs}}}$$

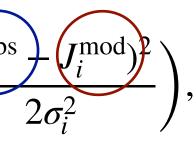
$$i = log_{10}(E) bin, j = X_{max} bin$$

$$D = D(J) + D(X_{\max}) = \sum_{i} \frac{(J_i^{\text{obs}} - J_i^{\text{mod}})^2}{\sigma_i^2} + 2 \cdot \sum_{i} \sum_{k_{i,j}^{\text{obs}}} k_{i,j}^{\text{obs}} \cdot \ln\left(\frac{k_{i,j}^{\text{obs}}}{n_i^{\text{obs}} \cdot G_{i,j}^{\text{mod}}}\right)$$

 \rightarrow The observed and simulated fluxes are compared by minimising the deviance D

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observed unfolded flux (detector effects) expected simulated flux

observed events model probability

(Gumbel distribution + detector effects)



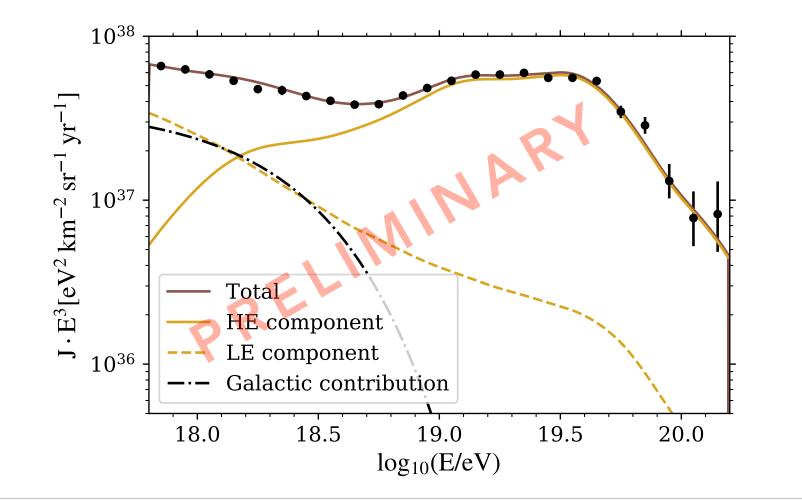
The reference scenarios

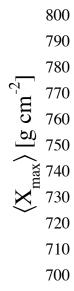
* Superposition of two (or more) populations to describe the ankle feature

* The extragalactic components ejected according to a power law with a rigidity dependent cutoff (with different parameters)

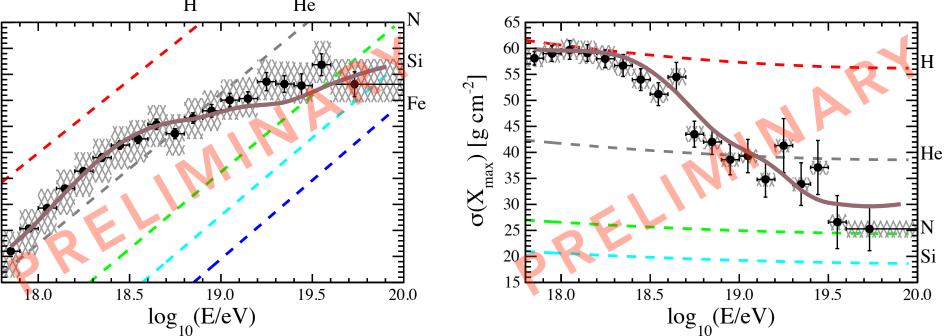
SCENARIO 1 : EXTRAGALACTIC AND GALACTIC POPULATIONS

- **Extragalactic populations with mixed mass composition** dominating at high energy (HE)
- **Extragalactic population of pure protons** dominating at low energy (LE)
 - → Possibly produced by decay of neutrons from photodisintegrations of nuclei in the same source environment
- **Galactic additional contribution** at low energy (considered at the Earth \rightarrow no propagation included)
 - → the best fit is given by a **nitrogen component** extending up to $Z \cdot R_{cut}^{Gal} \approx 2 \cdot 10^{18} \text{ eV}$





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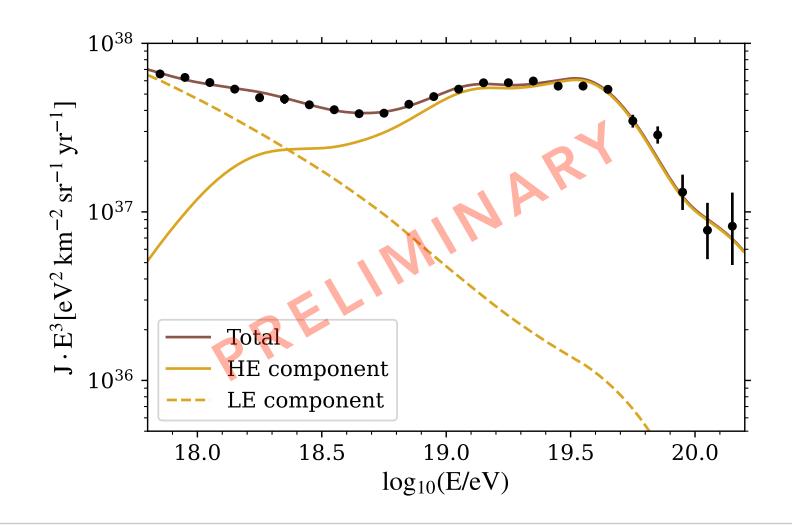
The reference scenarios

* Superposition of two (or more) populations to describe the ankle feature

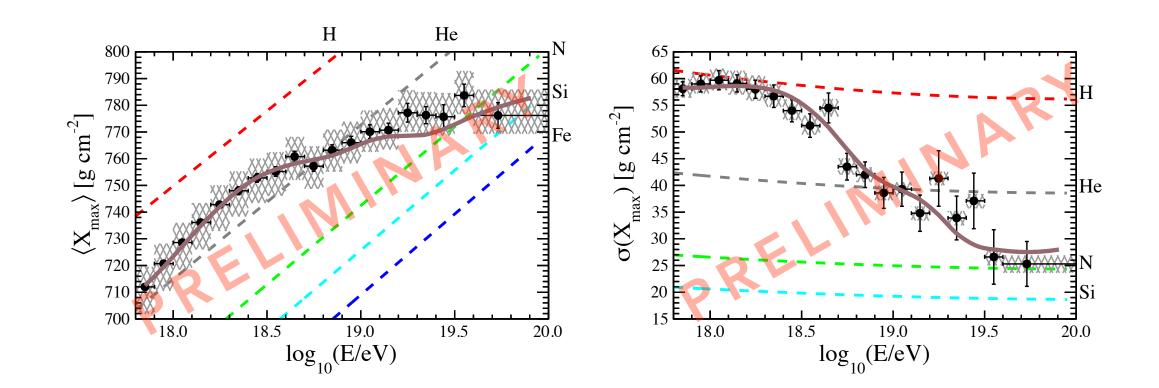
* The extragalactic components ejected according to a power law with a rigidity dependent cutoff (with different parameters)

SCENARIO 2 : TWO MIXED EXTRAGALACTIC POPULATIONS

- **Extragalactic populations with mixed mass composition** dominating at high energy (HE)
- **Extragalactic population with mixed mass composition** dominating at low energy (LE)
 - → produced by two different populations of sources
 - → Galactic contributions are subdominant in this energy range



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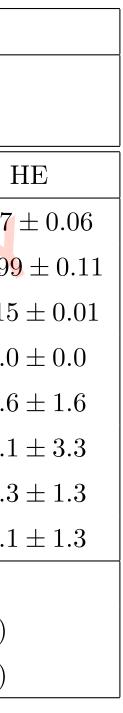


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	Scen	ario 1	Scenario 2	
Galactic contribution (at Earth)	pure N			
$J_0^{\text{Gal}} \left[\text{eV}^{-1} \cdot \text{km}^{-2} \cdot \text{sr}^{-1} \cdot \text{yr}^{-1} \right]$	$(1.06 \pm 0.04) \cdot 10^{-13}$			
$\log_{10}(R_{\rm cut}^{\rm Gal}/{\rm V})$	17.48 ± 0.02		-	
EG components (at the escape)	LE	HE	LE	H
$\mathcal{L}_0 \left[10^{44} \cdot \mathrm{erg} \cdot \mathrm{Mpc}^{-3} \cdot \mathrm{yr}^{-1} \right]^*$	6.54 ± 0.36	5.00 ± 0.35	11.35 ± 0.15	5.07
γ	3.34 ± 0.07	-1.47 ± 0.13	3.52 ± 0.03 (-1.99
$\log_{10}(R_{\rm cut}/{ m V})$	> 19.3	18.19 ± 0.02	> 19.4	18.15
$I_{\rm H} \ (\%)$	100 (fixed)	0.0 ± 0.0	48.7 ± 0.3	0.0
I_{He} (%)		24.5 ± 3.0	7.3 ± 0.4	23.6
$I_{\rm N}$ (%)		68.1 ± 5.0	44.0 ± 0.4	72.1
$I_{\rm Si}$ (%)		4.9 ± 3.9	0.0 ± 0.0	1.3
I_{Fe} (%)	C-K	2.5 ± 0.2	0.0 ± 0.0	3.1
$D_J (N_J)$	48.6 (24)		56.6 (24)	
$D_{X_{\max}} (N_{X_{\max}})$	537.4(329)		516.5(329)	
D(N)	586.0(353)		573.1 (353)	

* from $E_{\min} = 10^{17.8}$ eV.



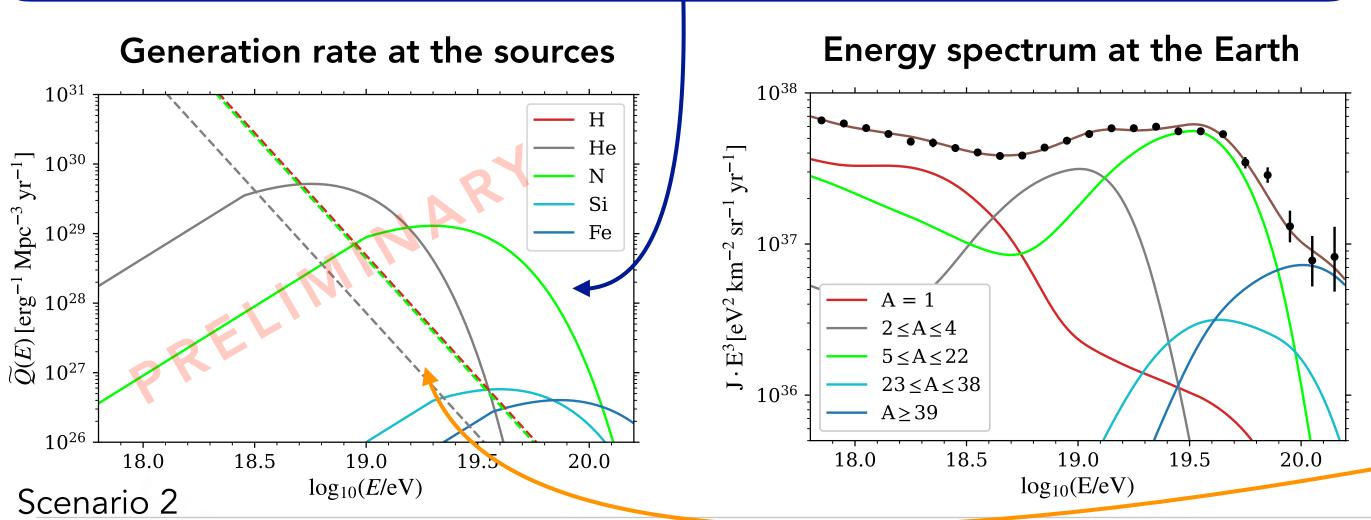


Some common findings between the two scenarios:

Very hard energy spectrum for the HE extragalactic component

- little overlap between different masses
 - → description of very pronounced spectral features and narrow distributions.
- Considering only the extragalactic propagation
 - → energy-dependent effects in the source environment are not incl
- "Magnetic horizon" effect

→ observed harder spectrum because of the suppression of the energy fluxes



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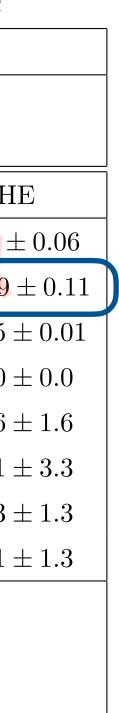
	Scenario 1		Scenario 2		
	Galactic contribution (at Earth)	pure N			
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	$\log_{10}(R_{\rm cut}^{\rm Gal}/{\rm V})$	17.48 ± 0.02			
ow X _{max}	EG components (at the escape)	LE	HE	LE	Η
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	γ	3.34 ± 0.07	-1.47 ± 0.13	3.52 ± 0.03	-1.99
	$\log_{10}(R_{\rm cut}/{ m V})$	> 19.3	18.19 ± 0.02	> 19.4	18.15
cluded	$I_{\rm H} \ (\%)$	100 (fixed)	0.0 ± 0.0	48.7 ± 0.3	0.0 :
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	$I_{ m N}$ (%)	_	68.1 ± 5.0	44.0 ± 0.4	72.1 :
ne low-	I_{Si} (%)	-0	4.9 ± 3.9	0.0 ± 0.0	1.3 :
	I_{Fe} (%)	0	2.5 ± 0.2	0.0 ± 0.0	3.1 :
	$D_J (N_J)$	$48.6 (24) \\537.4 (329)$		56.6(24)	
ne Earth	$D_{X_{\max}} (N_{X_{\max}})$			516.5(329)	
	D(N)	586.0(353)		573.1 (353)	
	* from $E_{\min} = 10^{17.8}$ eV.				

Very soft energy spectrum for the LE extragalactic component Possible explanation:

- Sources with different maximal energies (not identical)
 - \rightarrow the energy spectrum of each source may be less steep











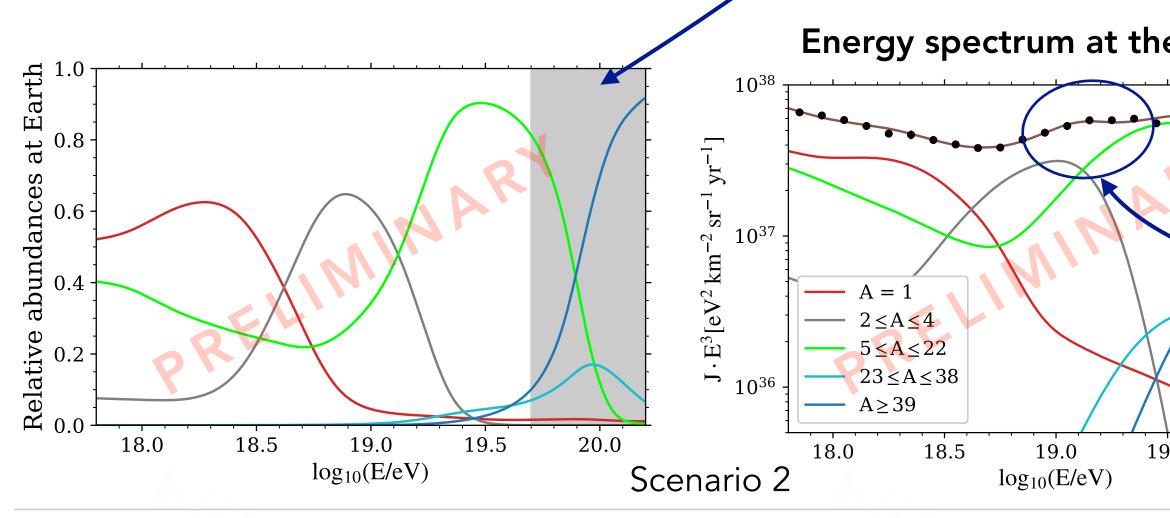
Some common findings between the two scenarios:

Low rigidity cutoff of the HE component

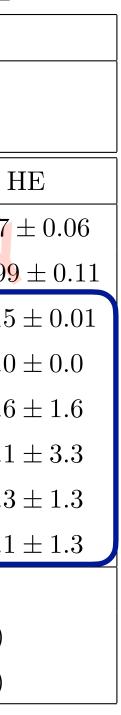
- It affects the observed fluxes ($< 10^{18.5} \text{ eV}$)
 - → <u>but</u> not low enough to make propagation effects negligible

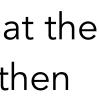
Mixed mass composition of the HE component

- No mass composition information at the highest energies
 - \rightarrow fit based on the shape of the energy spectrum



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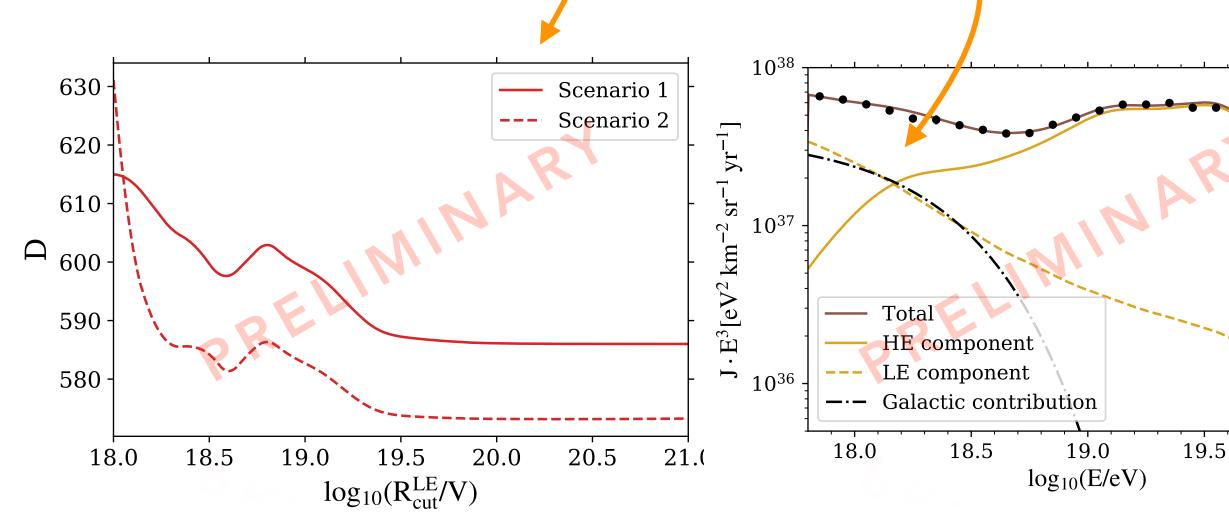




Some common findings between the two scenarios:

Very high rigidity cutoff of the LE component

- Degenerate fit for $R_{cut}^{LE} \gg 10^{19.5} \text{eV}$
 - \rightarrow fixing the parameter to any much higher value does not change
 - → only the lower bound
- The LE component is very steep
 - \rightarrow dominant only in the first energy bins
 - \rightarrow not very sensitive to the energy spectrum shape

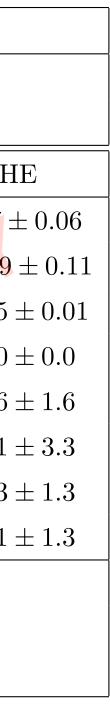


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		Scen	ario 1	Scena	ario 2
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	D(N)	586.0 (353)		573.1 (353)	
	* from $E_{\min} = 10^{17.8}$ eV.				
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20.0





The mass composition in the LE region

Mixture of H+N below the ankle in both scenarios Galactic component in Scenario 1 :

- power law modified by an exponential cutoff with some free param
- Models with Galactic Fe/Si right below the ankle are strongly disfar
- a N-dominated composition is preferred

 \rightarrow contribution from explosions in the winds of Wolf-Rayet-like sta provide N up to ~10¹⁸ eV

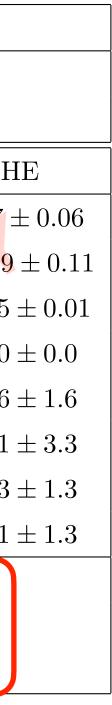
It is not possible to choose a favored scenario

- Scenario 2: better X_{max} distributions and worse spectrum description
- The differences are encompassed within the systematic uncertai
- In Scenario 2, photodisintegration is negligible for the LE component \rightarrow light-to-intermediate masses (similar to the one at the sources)
- Further investigation of the Galactic-to-extragalactic transition reg necessary

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		Scenario 1		Scenario 2	
	Galactic contribution (at Earth)	pure N $(1.06 \pm 0.04) \cdot 10^{-13}$ 17.48 ± 0.02			
	$J_0^{\text{Gal}} [\text{eV}^{-1} \cdot \text{km}^{-2} \cdot \text{sr}^{-1} \cdot \text{yr}^{-1}]$				
	$\log_{10}(R_{\rm cut}^{\rm Gal}/{\rm V})$			_	_
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on	D(N)	586.0	0 (353)	573.1	(353)
on	* from $E_{\min} = 10^{17.8}$ eV.				
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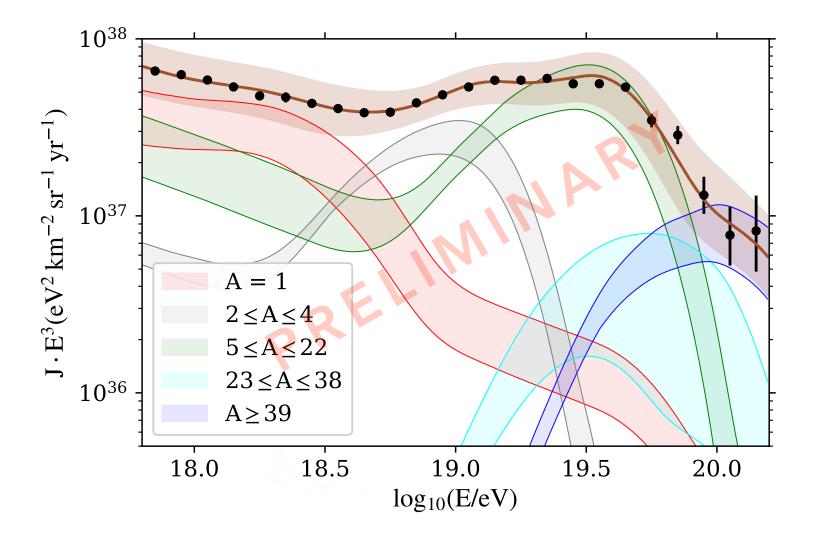




Effect of the systematic uncertainties

Experimental systematic uncertainties:

Effect of the uncertainties on the predicted total fluxes and on the partial contributions from different mass groups



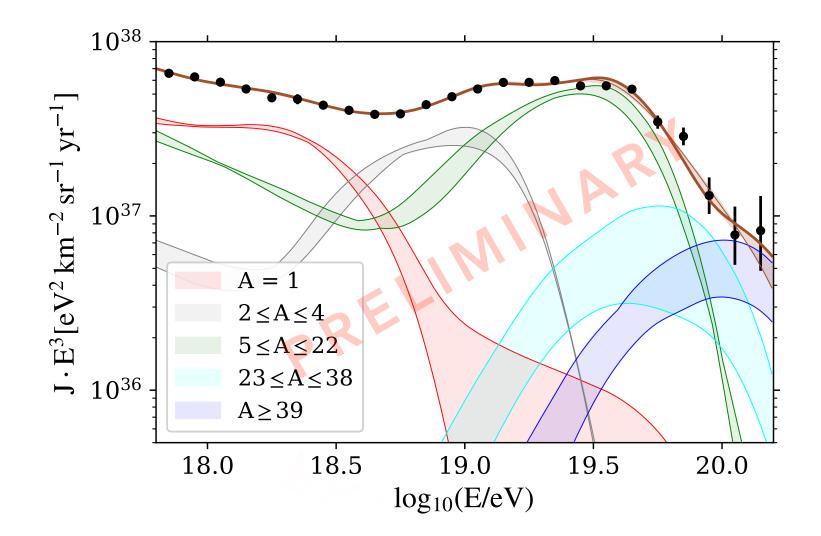
- The **dominant effect** is the one from the **experimental uncertainties** (mainly from the X_{max} scale)
- The systematic uncertainties do not spoil our conclusions in the reference scenarios

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effect is tested in the Scenario 2

Systematic uncertainties from models:



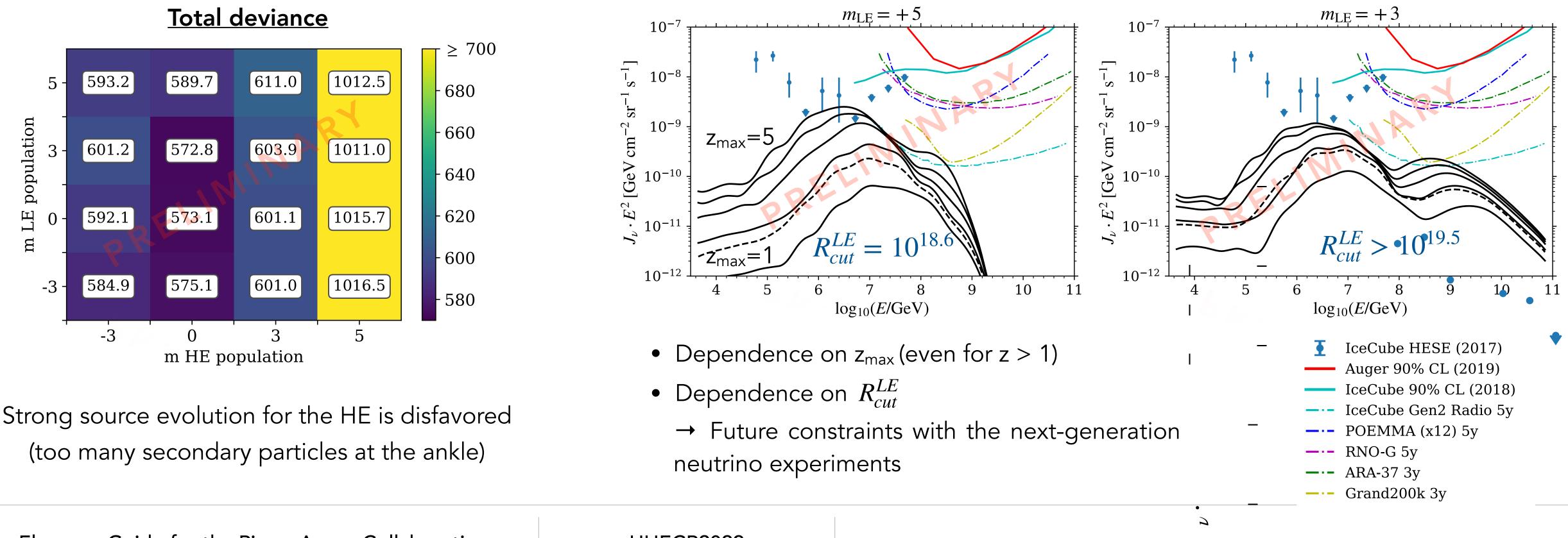
The systematic uncertainty





Cosmological evolution of sources

- Three alternative models for the evolution of the source emissivity, parameterized as $\propto (1+z)^m$ \rightarrow m=-3, m=+3, m=+5 (m=0 was used in the reference scenarios)
- The behavior at z>1 has only a negligible impact on the LE component (no impact on the HE one)
- All the possible combinations have been tested



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<u>Neutrinos fluxes for a strong evolution of the LE component</u>



Conclusions

- - \rightarrow description of the **ankle feature** at $\sim 6 \cdot 10^{18} \, \text{eV}$ as the superposition of different components
 - \rightarrow description of the **instep** at $\sim 10^{19} \, \text{eV}$ and of the **suppression** at the highest energies
 - → similar results in terms of deviance in the two scenarios
- Galactic component at LE (if present) : composition heavier than N strongly disfavored
- The systematic uncertainties do not spoil our conclusions
- \circ Very strong evolution (m=5) for the HE component is excluded
- The cosmogenic neutrino fluxes in some scenarios may reach the sensitivity of next-generation experiments

Collaboration paper about this analysis almost ready to be submitted to a journal

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• Simple astrophysical model with two extragalactic components (with or without a Galactic contribution at LE)











Outlook

Telescopes) in progress

 \rightarrow further insights on the Galactic-to-extragalactic transition region

- Possible additional information **including arrival directions** in the fit \rightarrow preliminary study with a combined fit above the ankle (presented at ICRC2021¹ and in T. Bister's poster at this conference)
- Future mass composition estimates with machine learning techniques on SD data
- Improvement of the mass composition at the highest energies from the detector upgrade (AugerPrime)
 - → same analysis could be performed with *much more statistics*

 \rightarrow mass composition information at the high-energy suppression

THANK YOU FOR YOUR ATTENTION!

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• Update of the X_{max} analysis including also data from **low-energy extension of Auger** (HEAT → High-Elevation Auger





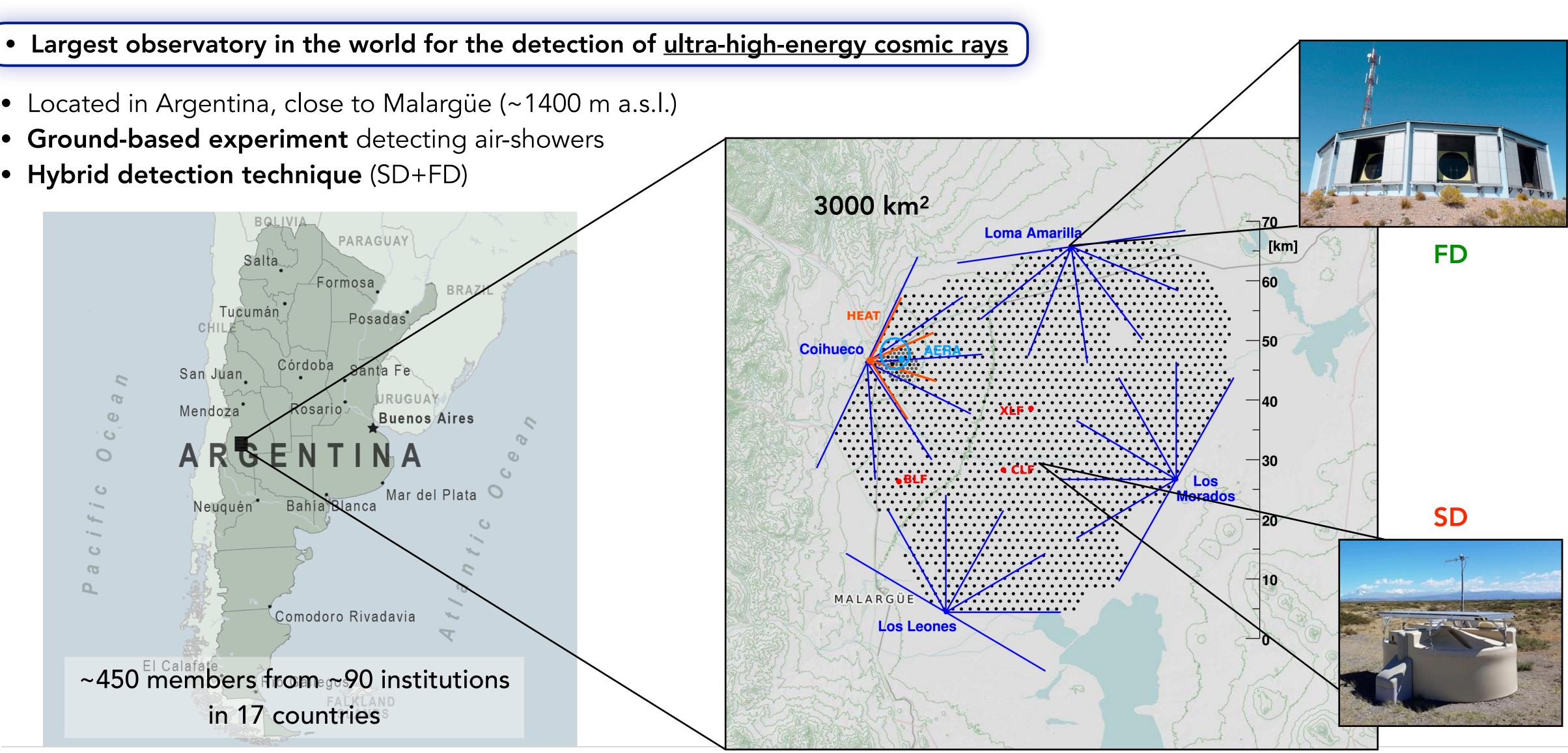




Back-up slides

The Pierre Auger Observatory

- Located in Argentina, close to Malargüe (~1400 m a.s.l.)
- Ground-based experiment detecting air-showers
- Hybrid detection technique (SD+FD)



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Effect of the systematic uncertainties from measurements

Two main sources of experimental systematic uncertainties:

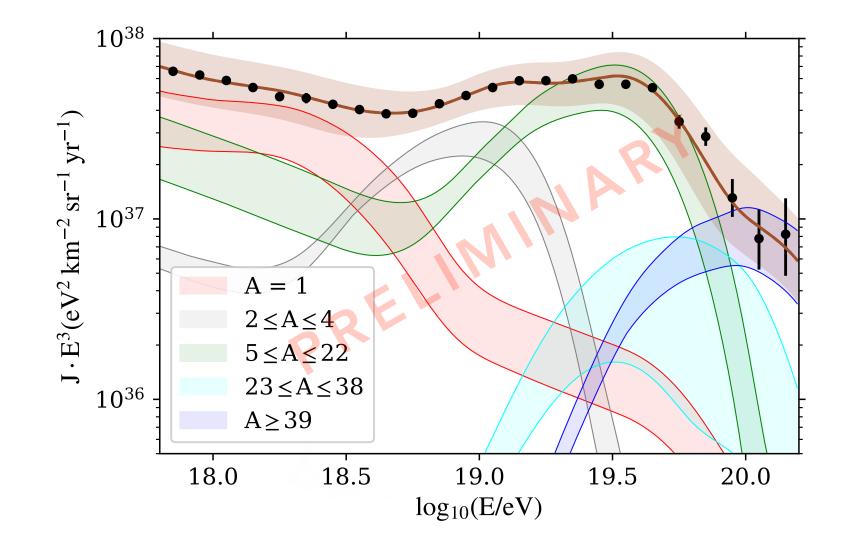
- Energy scale: σ_{sys}(E)/E = 14 %
 X_{max} scale: σ_{sys}(X_{max}) = 6 ÷ 9 g cm⁻²

- Energy scale \rightarrow shift all the energies of $\pm 1\sigma_E$ in each direction
- X_{max} scale \rightarrow the correlations among the energy bins are taken into account allowing for different shifts at different energies
 - * The X_{max} values are shifted by $a \cdot v_1(E) + b \cdot v_2(E)$
 - * a, b are two additional nuisance parameters in the fit
 - * A term $D_{syst} = a^2 + b^2$ has to be added to deviance
- Large band around the total flux due to the energy scale uncertainty → impact mainly on the estimated emissivity of sources
- The strongest impact on the predicted fluxes and on the deviance is due to the X_{max} scale uncertainty

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The systematic uncertainty effect is tested in the Scenario 2





Effect of the systematic uncertainties from models

Models for propagation in the IGM and in the atmosphere

Hadronic interaction model: Sibyll2.3d / EPOS-LHC / intermediate models

- Nuisance parameter $\delta_{
 m HIM}$ to interpolate each Gumbel parameter as $\alpha_{\text{HIM}} = \delta_{\text{HIM}} \cdot \alpha_{\text{EPOS}} + (1 - \delta_{\text{HIM}}) \cdot \alpha_{\text{Sib}}$
- If δ_{HIM} is close to 0 \rightarrow Sibyll2.3d is dominant
- If δ_{HIM} is close to 1 \rightarrow EPOS-LHC is dominant

Propagation model effect:

fit repeated considering different model configurations

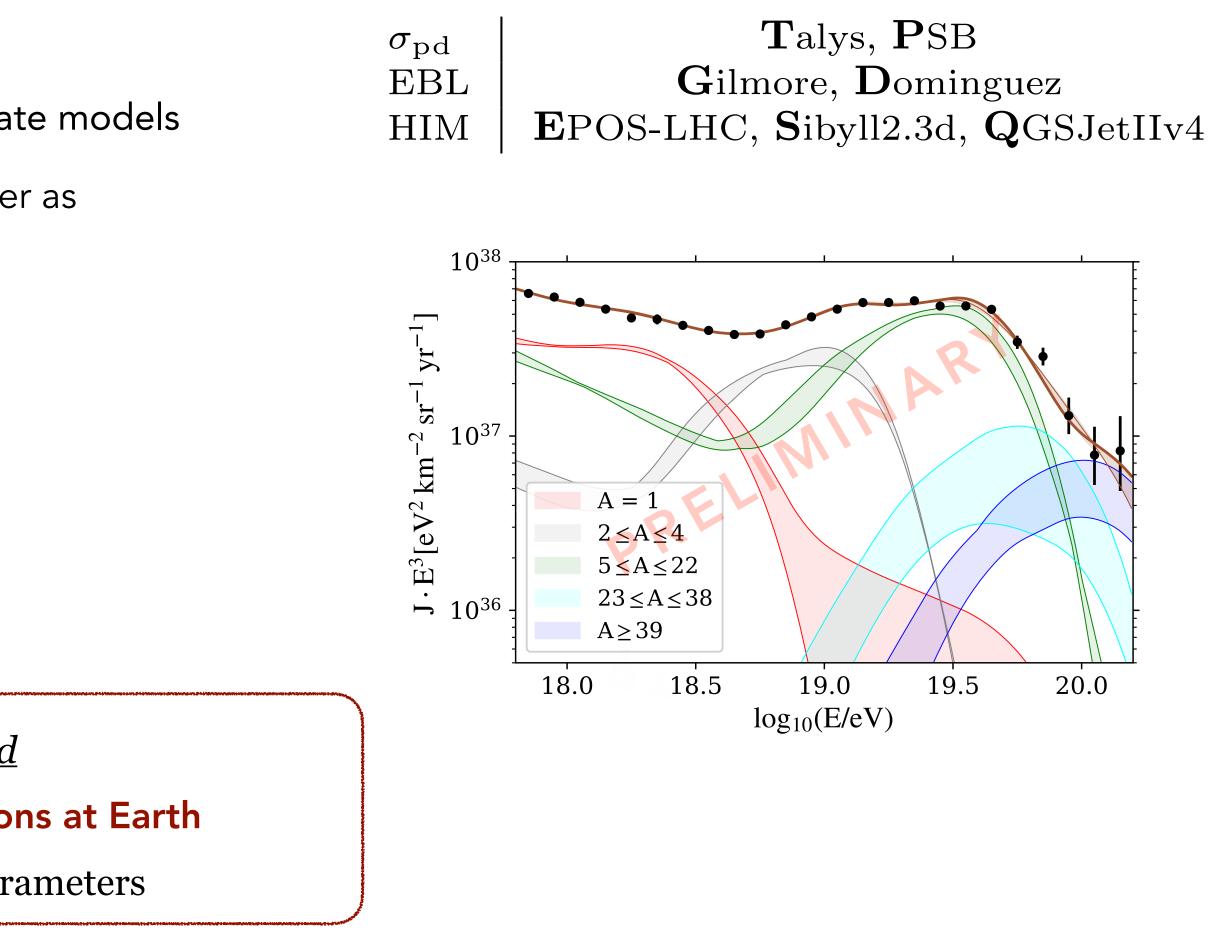
<u>EPOS-LHC or models compatible with it are always preferred</u>

 \rightarrow HIM choice: stronger impact on D and on the predictions at Earth

Propagation models: some expected changes in the best fit parameters

<u>The dominant effect on the the predicted fluxes and on the deviance is the one from the experimental uncertainties</u>

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Intergalactic magnetic fields

 $r_{\rm L} \approx 1.08 \cdot (E/{\rm EeV}) \cdot Z^{-1} \cdot (B_{\perp}/{\rm nG})^{-1} \, {\rm Mpc}$ Larmor radius:

Propagation theorem: the effect of intergalactic magnetic fields is negligible if the distance among sources is much lower than $r_{\rm L}$

- The lowest relevant rigidity $\sim E/Z$ in our model is that of N (Z=7) at $\sim 10^{17.8}$ eV
- Typical distance among sources is ≤ 10 Mpc

 \rightarrow magnetic fields should have $B_{\perp} \ll 10^{-11}$ G to be negligible