

Magnetic fields and UHECR propagation

Katia FERRIÈRE

Institut de Recherche en Astrophysique et Planétologie,
Observatoire Midi-Pyrénées, Toulouse, France

*6th International Symposium
on Ultra High Energy Cosmic Rays
L'Aquila – 3-7 October, 2022*

Outline

- 1 Introduction
- 2 The Galactic magnetic field
 - Dust polarization
 - Synchrotron emission
 - Faraday rotation
 - Faraday tomography
- 3 ExtraGalactic magnetic fields
 - External galaxies
 - Outside galaxies

Outline

- 1 Introduction
- 2 The Galactic magnetic field
 - Dust polarization
 - Synchrotron emission
 - Faraday rotation
 - Faraday tomography
- 3 ExtraGalactic magnetic fields
 - External galaxies
 - Outside galaxies

Magnetic fields and cosmic rays

The existence of interstellar magnetic fields was first predicted as a means to *confine* and *accelerate* Galactic cosmic rays

- **Alfvén (1937)**
Cosmic-ray confinement implies
"the existence of a magnetic field in interstellar space"
- **Fermi (1949)**
"The main process of [cosmic-ray] acceleration is due to [interstellar] magnetic fields ...
The magnetic field in the dilute matter is $\sim 5 \mu\text{G}$, while its intensity is probably greater in the heavier clouds"

☞ The history of interstellar magnetic fields is intimately linked to the history of cosmic rays

Magnetic fields and cosmic rays

The Galactic magnetic field, \vec{B} , deflects the trajectories of cosmic rays

This deflection is measured by the Larmor radius, r_g

- 1 GeV CR

$$r_g \simeq (0.2 \text{ AU}) (B_{\mu\text{G}})^{-1} \simeq (10^{-6} \text{ pc}) (B_{\mu\text{G}})^{-1} \lll L_B, H_{\text{Gal}}$$

- ☞ CR has *helical* motion along field lines (negligible drift)
with some diffusion due to small-scale $\delta\vec{B}$ and collisions

- 1 EeV CR

$$r_g \simeq (1 \text{ kpc}) (B_{\mu\text{G}})^{-1} \sim L_B, H_{\text{Gal}}$$

- ☞ CR motion is *moderately deflected* by \vec{B}
such that trajectory is neither straight nor helical

Outline

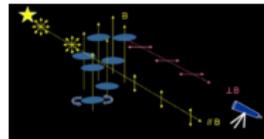
- 1 Introduction
- 2 The Galactic magnetic field
 - Dust polarization
 - Synchrotron emission
 - Faraday rotation
 - Faraday tomography
- 3 ExtraGalactic magnetic fields
 - External galaxies
 - Outside galaxies

Observational tools

● Polarization of starlight & dust thermal emission

Due to *dust grains* → general (dusty) ISM

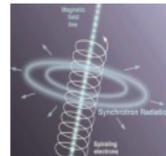
☞ \vec{B}_\perp (orientation only)



● Synchrotron emission

Produced by *CR electrons* → general (CR-filled) ISM

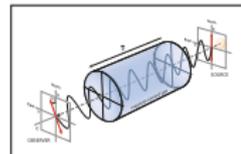
☞ \vec{B}_\perp (strength & orientation)



● Faraday rotation

Caused by *thermal electrons* → ionized regions

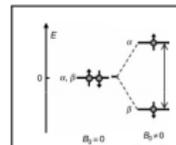
☞ B_\parallel (strength & sign)



● Zeeman splitting

Molecular & atomic *spectral lines* → neutral regions

☞ B_\parallel (strength & sign)



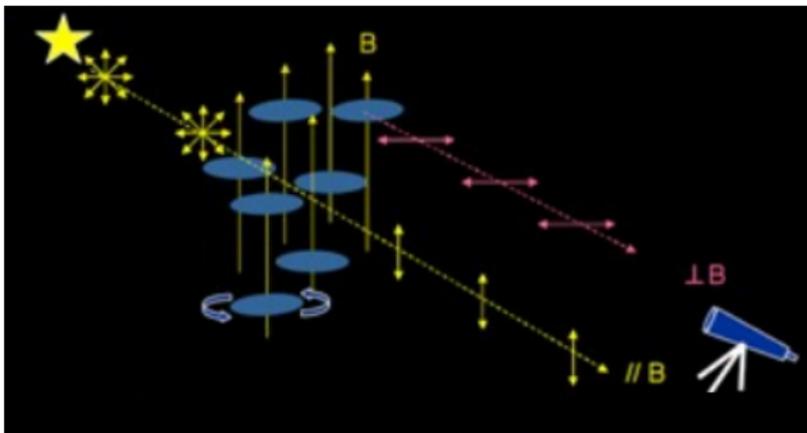
Outline

- 1 Introduction
- 2 The Galactic magnetic field
 - Dust polarization
 - Synchrotron emission
 - Faraday rotation
 - Faraday tomography
- 3 ExtraGalactic magnetic fields
 - External galaxies
 - Outside galaxies

Physical concept

Dust grains tend to **spin** about their short axes
& to **align** their spin axes with \vec{B}

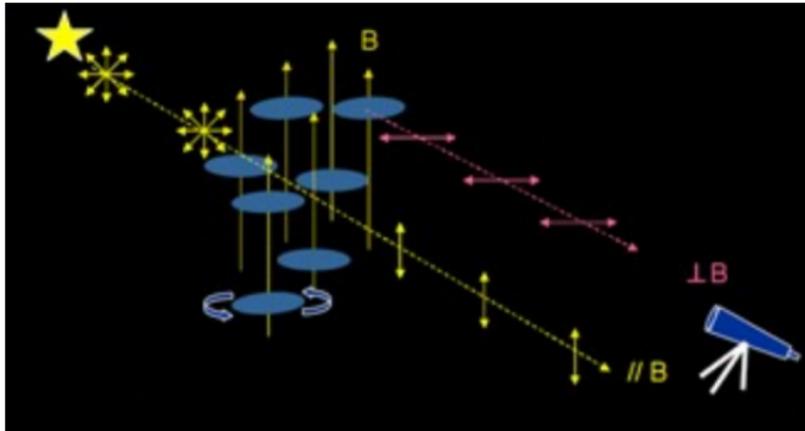
This grain alignment leads to *linear polarization*



Credit: Wen-Ping Chen

Polarization orientation

- Starlight attenuated by dust (*optical*) is polarized $\parallel \vec{B}_\perp$
- Dust thermal emission (*infrared*) is polarized $\perp \vec{B}_\perp$



Credit: Wen-Ping Chen

Polarization fraction

$$p \equiv \frac{P}{I}$$

- Starlight attenuated by dust : $p \simeq \tau p_0 \cos^2 \gamma$

- Dust thermal emission : $p = p_0 \cos^2 \gamma$

$$\hookrightarrow p_0 = p_{\max} F_{\text{align}} F_{\delta B}$$

$$\vec{B} \in \text{PoS}$$
$$(\cos^2 \gamma = 1)$$

$$\Rightarrow p = p_0$$



$$\vec{B} \perp \text{PoS}$$
$$(\cos^2 \gamma = 0)$$

$$\Rightarrow p = 0$$



Credit: Vincent Guillet

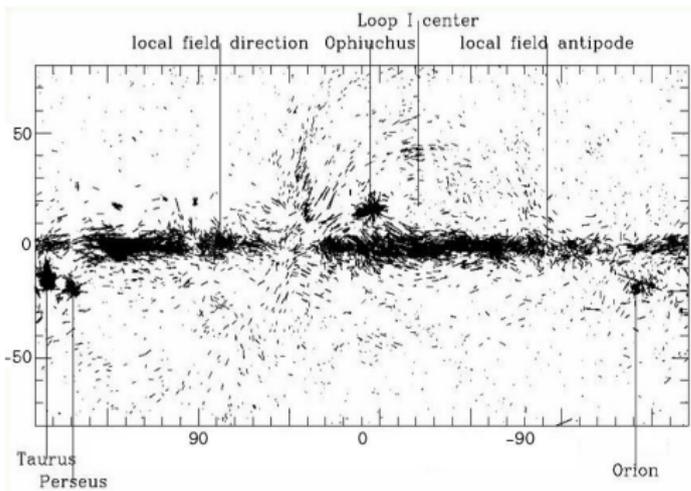
Dust polarization

Altogether

- Polarization *orientation*  *orientation* of \vec{B} in PoS
- Polarization *fraction*  *inclination* of \vec{B} to PoS (for ideal conditions)

Polarization of starlight

\vec{B}_\perp sectors from 8 662 stars



Heiles (2000)

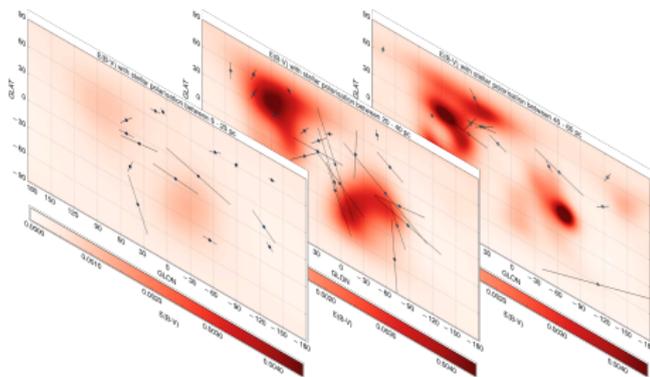
- Near the Sun - In the disk : \vec{B}_{ord} is horizontal
 \vec{B}_{ord} is nearly azimuthal ($\simeq -7^\circ$ from \hat{e}_ϕ)
- In the halo : \vec{B}_{ord} has a vertical component

Polarization of starlight

Stars have accurately measured distances (with Gaia)

➡ Possible to probe \vec{B} in 3D

Stellar polarization cube of nearby ISM



3 layers at

0 – 20 pc

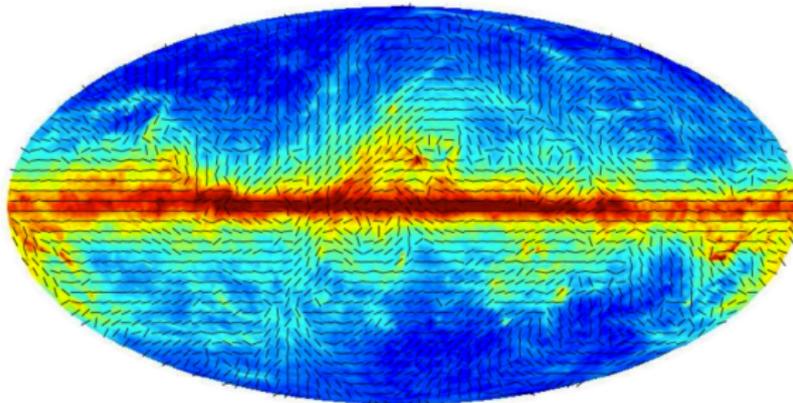
20 – 40 pc

40 – 60 pc

Credit: Marta Alves

Polarization of dust thermal emission

Total intensity & \vec{B}_\perp segments at 353 GHz (Planck)

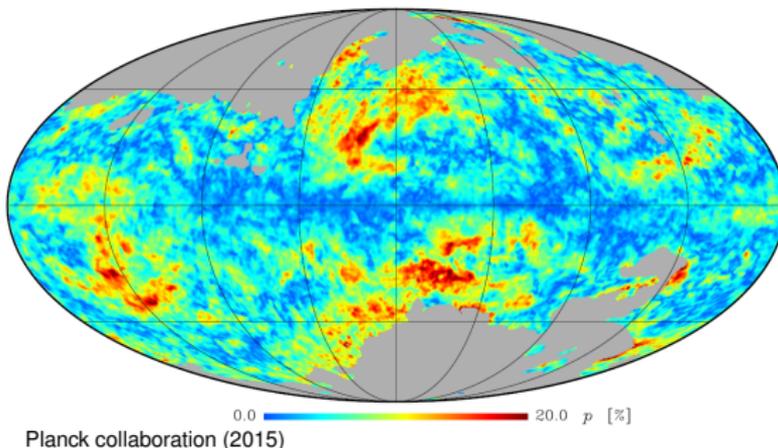


Planck collaboration (2015)

- In the disk : \vec{B}_{ord} is horizontal
- In the halo : \vec{B}_{ord} has a vertical component

Polarization of dust thermal emission

Polarization fraction at 353 GHz (Planck)



Info on - Inclination of \vec{B}_{ord} to PoS : $\cos^2 \gamma$

- Magnetic fluctuations : $\frac{B_{\text{fluct}}}{B_{\text{ord}}}$

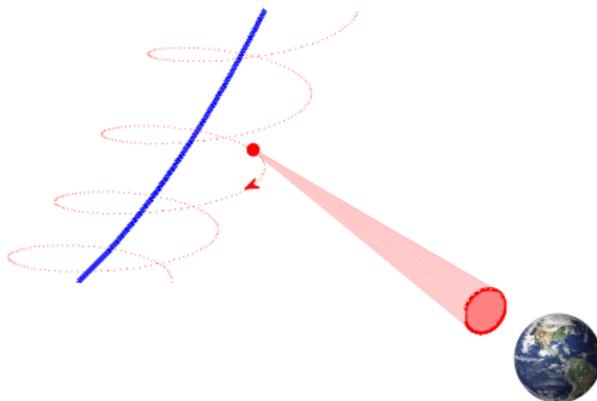
- Grain properties & alignment efficiency : p_{max} & F_{align}

Outline

- 1 Introduction
- 2 **The Galactic magnetic field**
 - Dust polarization
 - **Synchrotron emission**
 - Faraday rotation
 - Faraday tomography
- 3 ExtraGalactic magnetic fields
 - External galaxies
 - Outside galaxies

Physical concept

Relativistic electrons gyrating about magnetic field lines emit *synchrotron radiation*

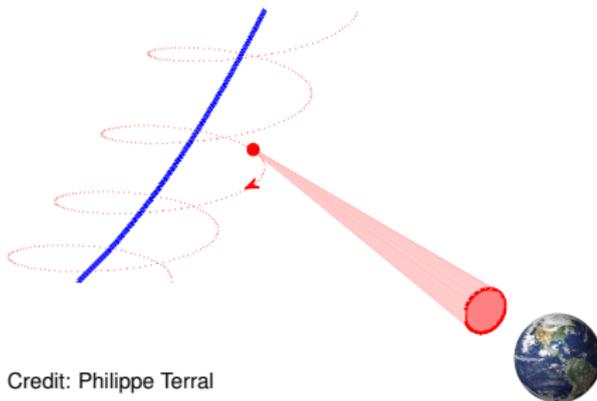


Credit: Philippe Terral

Total & polarized intensities

Emissivity : $\mathcal{E} = f(\alpha) n_{\text{CRE}} B_{\perp}^{\alpha+1} \nu^{-\alpha}$ & $\mathcal{E}_{\text{pol}} = p_{\text{syn}} \mathcal{E}$ & $\vec{\mathcal{E}}_{\text{pol}} \perp \vec{B}_{\perp}$

- Total intensity : $I = \int \mathcal{E} ds$ $\Rightarrow B_{\perp}$
- Polarized intensity : $\vec{P} = \int \vec{\mathcal{E}}_{\text{pol}} ds$ $\Rightarrow (\vec{B}_{\perp})_{\text{ord}}$



Credit: Philippe Terral

Total & polarized intensities

Emissivity : $\mathcal{E} = f(\alpha) n_{\text{CRE}} B_{\perp}^{\alpha+1} \nu^{-\alpha}$ & $\mathcal{E}_{\text{pol}} = p_{\text{syn}} \mathcal{E}$ & $\overleftrightarrow{\mathcal{E}}_{\text{pol}} \perp \overrightarrow{B}_{\perp}$

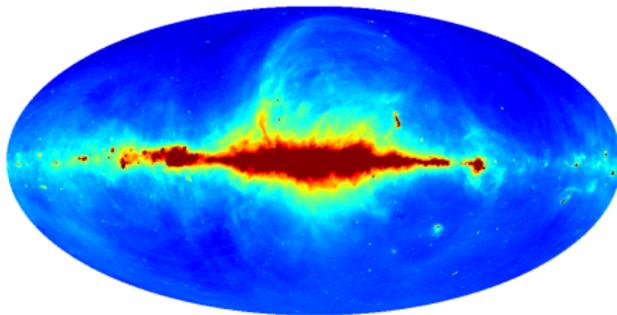
- Total intensity : $I = \int \mathcal{E} ds$ $\Rightarrow B_{\perp}$

- Polarized intensity : $\overleftrightarrow{P} = \int \overleftrightarrow{\mathcal{E}}_{\text{pol}} ds$ $\Rightarrow (\overrightarrow{B}_{\perp})_{\text{ord}}$

$$\hookrightarrow Q + iU = \int \mathcal{E}_{\text{pol}} e^{2i\psi} ds$$

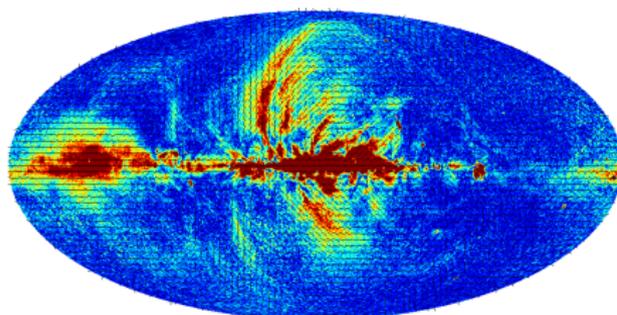
Total & polarized intensities

TI at 1.4 GHz (25m Stockert + 30m Villa Elisa)



Credit: Tess Jaffe

PI & \vec{B}_\perp sectors at 23 GHz (WMAP)



Credit: Tess Jaffe

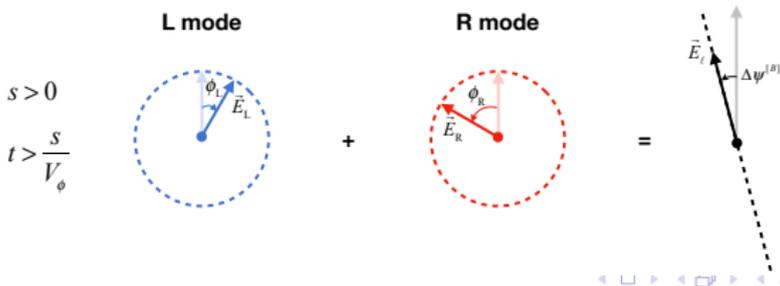
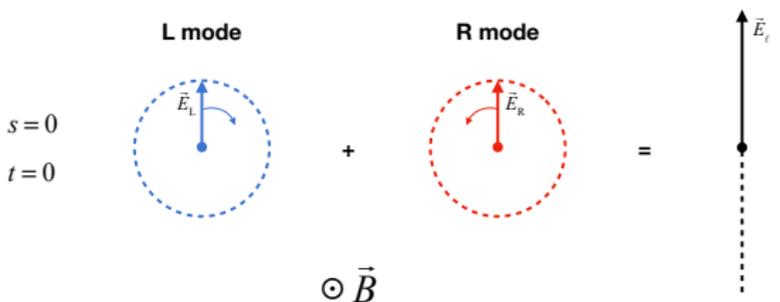
- ☞ - Near the Sun : $B_{\text{ord}} \sim 3 \mu\text{G}$ & $B_{\text{tot}} \sim 5 \mu\text{G}$
- In the disk : \vec{B}_{ord} is horizontal
- In the halo : \vec{B}_{ord} has a vertical component

Outline

- 1 Introduction
- 2 The Galactic magnetic field
 - Dust polarization
 - Synchrotron emission
 - Faraday rotation
 - Faraday tomography
- 3 ExtraGalactic magnetic fields
 - External galaxies
 - Outside galaxies

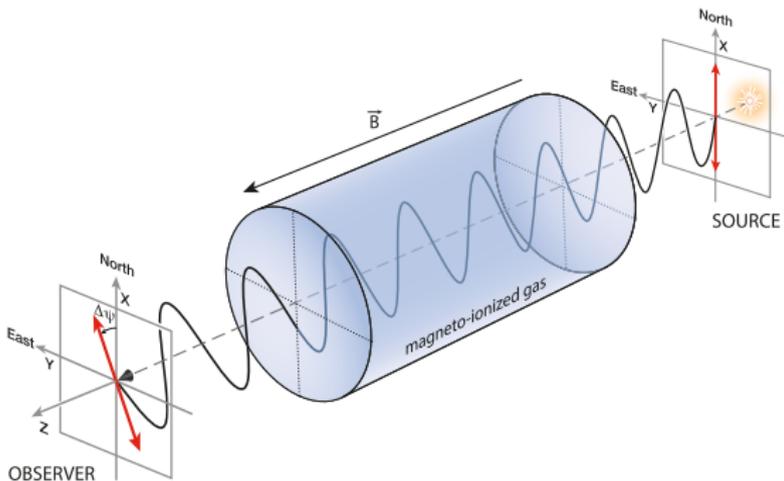
Physical concept

When a linearly polarized radio wave travels through a magneto-ionized medium, the orientation of linear polarization undergoes *Faraday rotation*



Physical concept

When a linearly polarized radio wave travels through a magneto-ionized medium, the orientation of linear polarization undergoes *Faraday rotation*

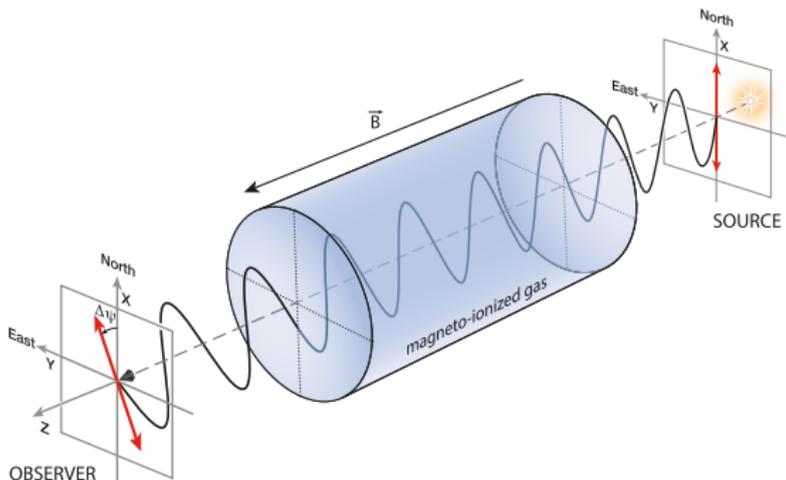


Credit: Theophilus Britt Griswold (NASA Goddard)

Rotation angle & rotation measure

Rotation angle : $\Delta\psi = \text{RM} \lambda^2$

Rotation measure : $\text{RM} = C \int n_e B_{\parallel} ds$ \vec{B}_{\parallel}

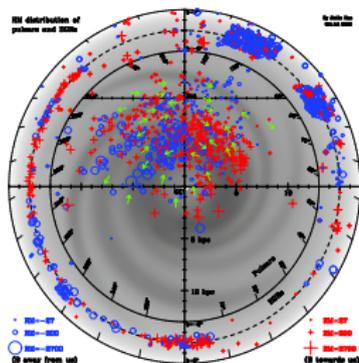


Credit: Theophilus Britt Griswold (NASA Goddard)



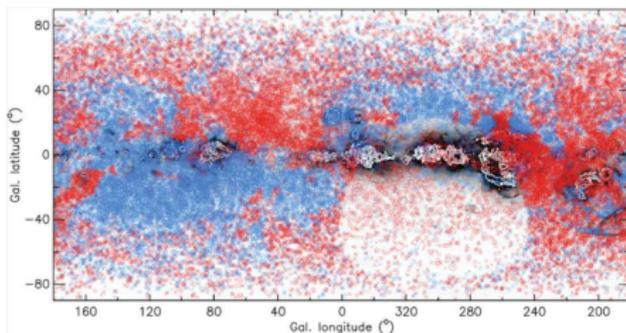
Rotation measures

RMs of pulsars & EGRSs with $|b| < 8^\circ$



Han (2009)

RMs of EGRSs [NVSS ($\delta > -40^\circ$) + S-PASS ($\delta < 0^\circ$)]



Schnitzeler et al. (2019)

☞ - Near the Sun : $B_{\text{reg}} \simeq 1.5 \mu\text{G}$ & $B_{\text{tot}} \sim 5 \mu\text{G}$

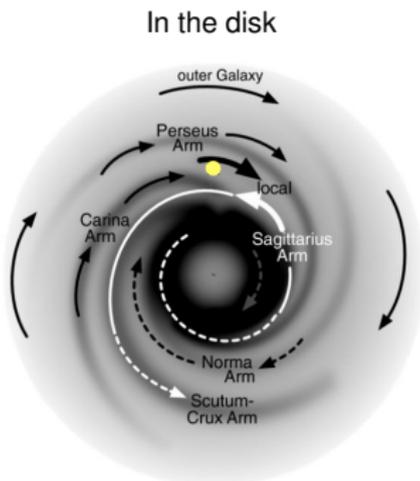
\vec{B}_{reg} is nearly azimuthal ($\simeq -8^\circ$ from \hat{e}_ϕ)

- In the disk : \vec{B}_{reg} is horizontal & mostly azimuthal, with reversals in B_ϕ
 \vec{B}_{reg} probably has a spiral shape

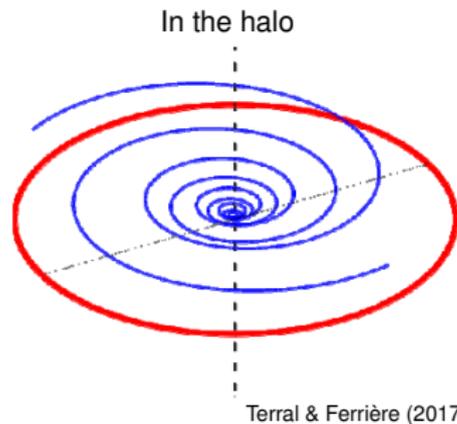
- In the halo : \vec{B}_{reg} is CCW at $z > 0$ & CW at $z < 0$

\vec{B}_{reg} possibly has an upward spiraling shape

Rotation measures



van Eck et al. (2011)



Terral & Ferrière (2017)

☞ - Near the Sun : $B_{\text{reg}} \simeq 1.5 \mu\text{G}$ & $B_{\text{tot}} \sim 5 \mu\text{G}$

\vec{B}_{reg} is nearly azimuthal ($\simeq -8^\circ$ from \hat{e}_ϕ)

- In the disk : \vec{B}_{reg} is horizontal & mostly azimuthal, with reversals in B_ϕ
 \vec{B}_{reg} probably has a spiral shape

- In the halo : \vec{B}_{reg} is CCW at $z > 0$ & CW at $z < 0$

\vec{B}_{reg} possibly has an upward spiraling shape

Outline

- 1 Introduction
- 2 The Galactic magnetic field
 - Dust polarization
 - Synchrotron emission
 - Faraday rotation
 - Faraday tomography
- 3 ExtraGalactic magnetic fields
 - External galaxies
 - Outside galaxies

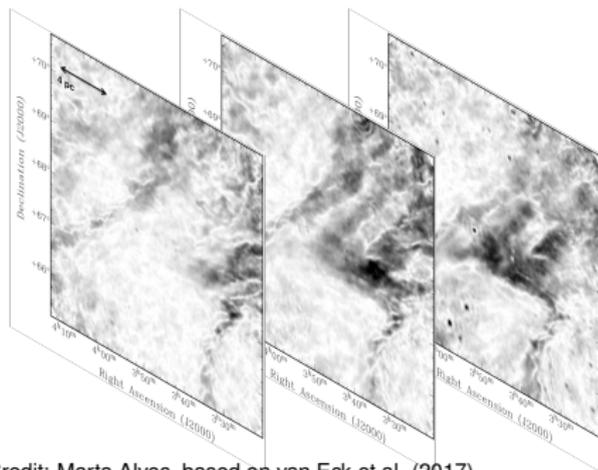
Physical concept

- Underlying processes
 - Galactic **synchrotron emission** : linearly polarized
 - **Faraday rotation** : λ -dependent
- General idea
 - Measure synchrotron polarized intensity at many different λ
 - Convert λ -dependence into s -dependence [Φ -dependence]
- Output

Faraday cube = 3D cube of synchrotron polarized emission as $fc(\alpha, \delta, \Phi)$

Output

Faraday cube toward Fan region (LOFAR)



3 slices at

$$\Phi_1 = -2.0 \text{ rad m}^{-2}$$

$$\Phi_2 = -1.5 \text{ rad m}^{-2}$$

$$\Phi_3 = -1.0 \text{ rad m}^{-2}$$

Credit: Marta Alves, based on van Eck et al. (2017)

What can be learned

- In Faraday cube
 - Uncover **synchrotron-emitting** & **Faraday-rotating** features
 - Identify these features with interstellar matter structures

- For **synchrotron-emitting** regions

$$\int \vec{F}(\Phi) d\Phi \propto \vec{B}_{\perp}$$

- For **Faraday-rotating** regions

$$\Delta\Phi \propto B_{\parallel}$$

Outline

- 1 Introduction
- 2 The Galactic magnetic field
 - Dust polarization
 - Synchrotron emission
 - Faraday rotation
 - Faraday tomography
- 3 **ExtraGalactic magnetic fields**
 - External galaxies
 - Outside galaxies

Outline

- 1 Introduction
- 2 The Galactic magnetic field
 - Dust polarization
 - Synchrotron emission
 - Faraday rotation
 - Faraday tomography
- 3 **ExtraGalactic magnetic fields**
 - **External galaxies**
 - Outside galaxies

Observational tools

- Synchrotron emission
 - ☞ \vec{B}_\perp (strength & orientation)
- Faraday rotation
 - ☞ B_\parallel (strength & sign)
- Polarization of dust thermal emission
 - ☞ \vec{B}_\perp (orientation only)

Main characteristics

Spiral galaxies

- \vec{B} has an **ordered** component
- $B_{\text{tot}} \sim \text{a few } \mu\text{G}$
- * *Face on*
 - Disk : \vec{B}_{ord} follows the **spiral arms**
- * *Edge on*
 - Disk : \vec{B}_{ord} is **horizontal**
 - Halo : \vec{B}_{ord} has an **X shape**

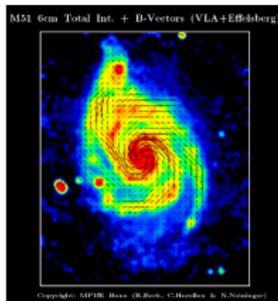
Elliptical galaxies

- \vec{B} has only **fluctuating** component
- $B_{\text{tot}} \sim \text{a few } \mu\text{G}$

M51



HST

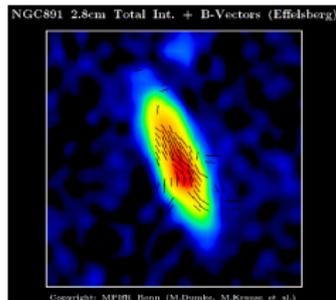


Effelsberg + VLA (6.2 cm)

NGC 891



CFHT



Effelsberg (2.8 cm)

Main characteristics

Spiral galaxies

- \vec{B} has an **ordered** component
- $B_{\text{tot}} \sim \text{a few } \mu\text{G}$
- * *Face on*
 - Disk : \vec{B}_{ord} follows the **spiral arms**
- * *Edge on*
 - Disk : \vec{B}_{ord} is **horizontal**
 - Halo : \vec{B}_{ord} has an **X shape**

Elliptical galaxies

- \vec{B} has only **fluctuating** component
- $B_{\text{tot}} \sim \text{a few } \mu\text{G}$

M51



HST

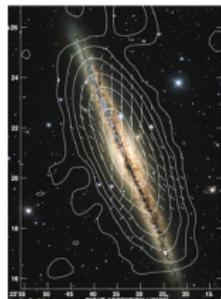


Effelsberg + VLA (6.2 cm)

NGC 891



CFHT



Effelsberg (3.6 cm)

Outline

- 1 Introduction
- 2 The Galactic magnetic field
 - Dust polarization
 - Synchrotron emission
 - Faraday rotation
 - Faraday tomography
- 3 **ExtraGalactic magnetic fields**
 - External galaxies
 - **Outside galaxies**

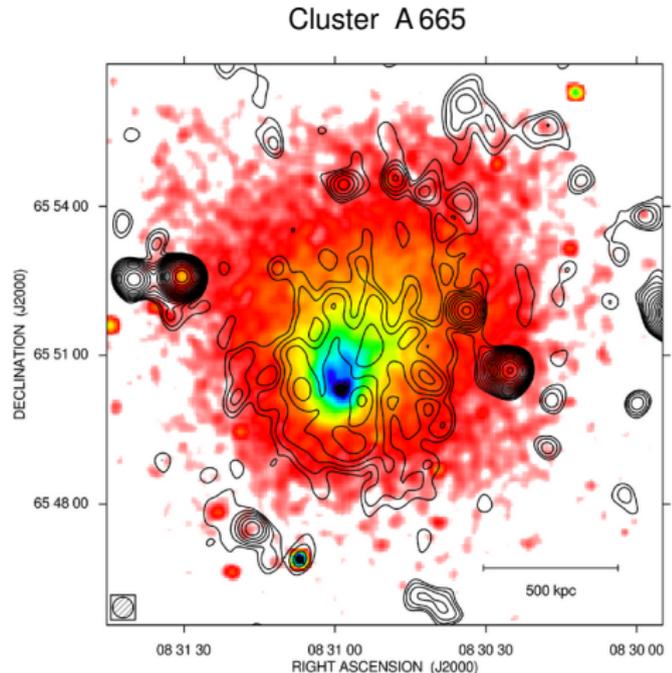
Clusters of galaxies

Observational tools

- Synchrotron emission
- Faraday rotation

Main characteristics

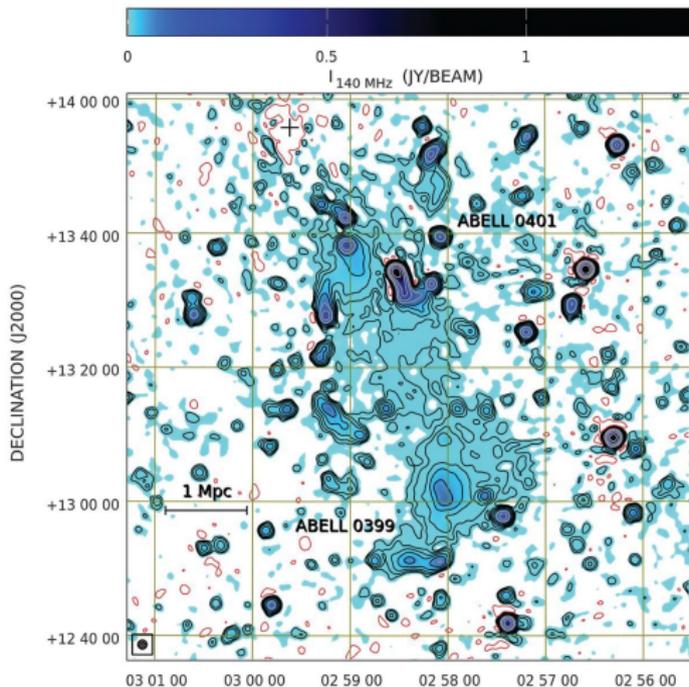
- \vec{B} has only **fluctuating** component
- $B_{\text{tot}} \sim (0.1 - \text{a few}) \mu\text{G}$



Radio 1.4 GHz + X-ray 0.8-4 keV (Vacca et al. 2018)

Cosmic filaments

Radio synchrotron filament between Abell 399 & Abell 401 (LOFAR)



Govoni et al. (2019)

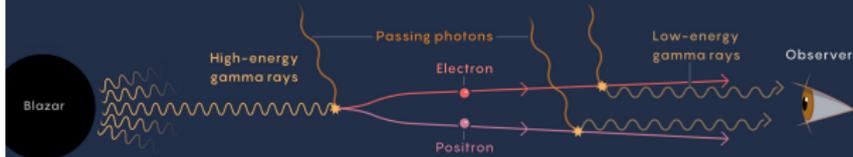
Cosmic voids

The Hidden Blazar Effect

Blazars are bright beams of energy powered by supermassive black holes. These beams can be used to detect the presence of cosmic magnetic fields.

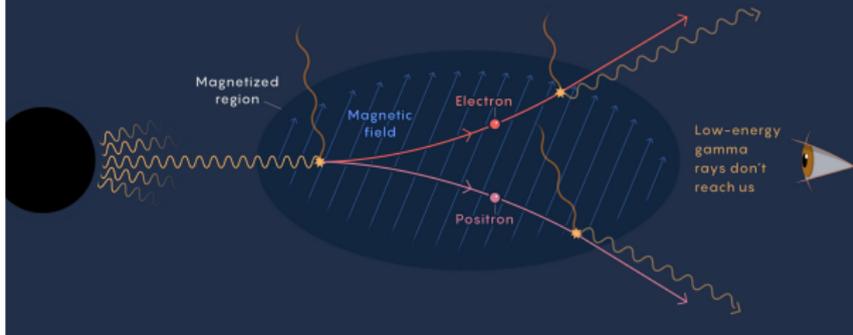
A BEAM'S LONG JOURNEY

A blazar creates high-energy gamma rays that occasionally morph into an electron and a positron. These then create lower-energy gamma rays.



A GAMMA-RAY SWERVE

If the high-energy gamma rays pass through a magnetic field, the field will deflect any electrons and positrons that get created. The resulting low-energy gamma rays will point away from the observer.



- Delay of GeV echo

$$B \gtrsim 10^{-20.5} \text{ G}$$

(Takahashi et al. 2013)

- Size of GeV halo

$$B \sim (10^{-17} - 10^{-15}) \text{ G}$$

(Chen et al. 2015)