

# Magnetic fields and UHECR propagation

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*6<sup>th</sup> International Symposium  
on Ultra High Energy Cosmic Rays  
L'Aquila – 3-7 October, 2022*

# Outline

- 1 Introduction
- 2 The Galactic magnetic field
  - Dust polarization
  - Synchrotron emission
  - Faraday rotation
  - Faraday tomography
- 3 ExtraGalactic magnetic fields
  - External galaxies
  - Outside galaxies

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# Magnetic fields and cosmic rays

The existence of interstellar magnetic fields was first predicted as a means to *confine* and *accelerate* Galactic cosmic rays

- **Alfvén** (1937)

Cosmic-ray confinement implies

"the existence of a magnetic field in interstellar space"

- **Fermi** (1949)

"The main process of [cosmic-ray] acceleration is due to [interstellar] magnetic fields ...

The magnetic field in the dilute matter is  $\sim 5 \mu\text{G}$ , while its intensity is probably greater in the heavier clouds"

☞ The history of interstellar magnetic fields is intimately linked to the history of cosmic rays



# Magnetic fields and cosmic rays

The Galactic magnetic field,  $\vec{B}$ , deflects the trajectories of cosmic rays

This deflection is measured by the Larmor radius,  $r_g$

## ● 1 GeV CR

$$r_g \simeq (0.2 \text{ AU}) (B_{\mu\text{G}})^{-1} \simeq (10^{-6} \text{ pc}) (B_{\mu\text{G}})^{-1} \lll L_B, H_{\text{Gal}}$$

- ☞ CR has *helical* motion along field lines (negligible drift)  
with some diffusion due to small-scale  $\delta\vec{B}$  and collisions

## ● 1 EeV CR

$$r_g \simeq (1 \text{ kpc}) (B_{\mu\text{G}})^{-1} \sim L_B, H_{\text{Gal}}$$

- ☞ CR motion is *moderately deflected* by  $\vec{B}$   
such that trajectory is neither straight nor helical

# Outline

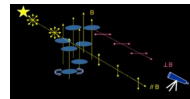
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# Observational tools

## • Polarization of starlight & dust thermal emission

Due to *dust grains* → general (dusty) ISM

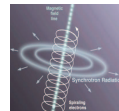
☞  $\vec{B}_\perp$  (orientation only)



## • Synchrotron emission

Produced by *CR electrons* → general (CR-filled) ISM

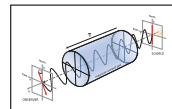
☞  $\vec{B}_\perp$  (strength & orientation)



## • Faraday rotation

Caused by *thermal electrons* → ionized regions

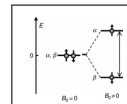
☞  $B_\parallel$  (strength & sign)



## • Zeeman splitting

Molecular & atomic *spectral lines* → neutral regions

☞  $B_\parallel$  (strength & sign)



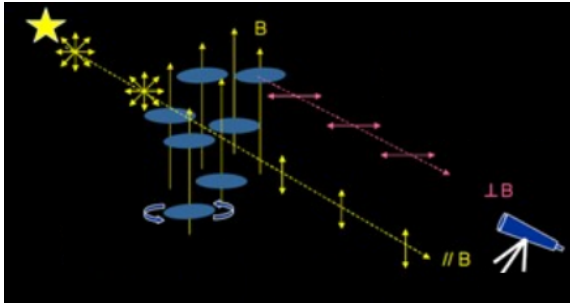
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# Physical concept

Dust grains tend to **spin** about their short axes  
& to **align** their spin axes with  $\vec{B}$

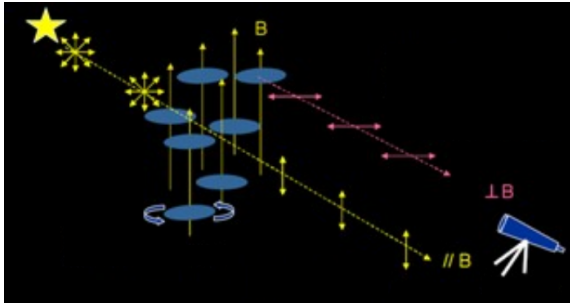
This grain alignment leads to *linear polarization*



Credit: Wen-Ping Chen

# Polarization orientation

- Starlight attenuated by dust (*optical*) is polarized  $\parallel \vec{B}_\perp$
- Dust thermal emission (*infrared*) is polarized  $\perp \vec{B}_\perp$



Credit: Wen-Ping Chen

# Polarization fraction

$$p \equiv \frac{P}{I}$$

- Starlight attenuated by dust :  $p \simeq \tau p_0 \cos^2 \gamma$
- Dust thermal emission :  $p = p_0 \cos^2 \gamma$   
 $\hookrightarrow p_0 = p_{\max} F_{\text{align}} F_{\delta B}$

$$\vec{B} \in \text{PoS}$$

$$(\cos^2 \gamma = 1)$$

$$\Rightarrow p = p_0$$



$$\vec{B} \perp \text{PoS}$$

$$(\cos^2 \gamma = 0)$$



$$\Rightarrow p = 0$$



Credit: Vincent Guillet

# Dust polarization

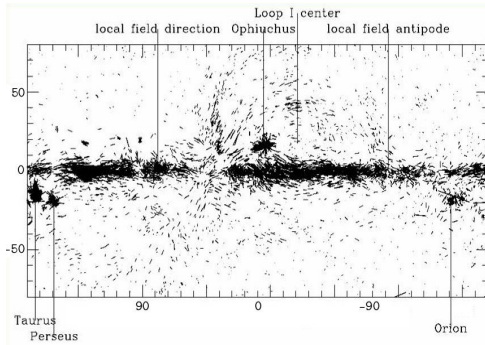
Altogether

- Polarization *orientation*  *orientation* of  $\vec{B}$  in PoS
- Polarization *fraction*  *inclination* of  $\vec{B}$  to PoS (for ideal conditions)



# Polarization of starlight

$\vec{B}_\perp$  sectors from 8 662 stars



Heiles (2000)

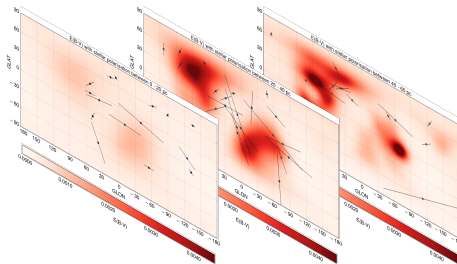
- ☞ Near the Sun
  - In the disk :  $\vec{B}_{\text{ord}}$  is horizontal  
 $\vec{B}_{\text{ord}}$  is nearly azimuthal ( $\simeq -7^\circ$  from  $\hat{e}_\phi$ )
  - In the halo :  $\vec{B}_{\text{ord}}$  has a vertical component

# Polarization of starlight

Stars have accurately measured distances (with Gaia)

☞ Possible to probe  $\vec{B}$  in 3D

Stellar polarization cube of nearby ISM

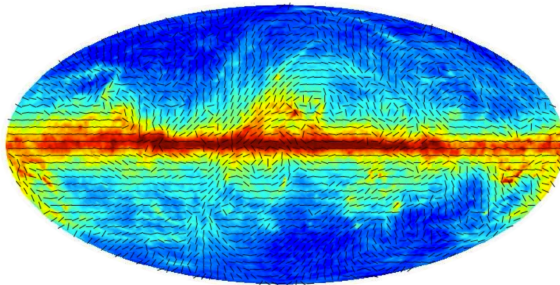


3 layers at  
0 – 20 pc  
20 – 40 pc  
40 – 60 pc

Credit: Marta Alves

# Polarization of dust thermal emission

Total intensity &  $\vec{B}_\perp$  sectors at 353 GHz (Planck)

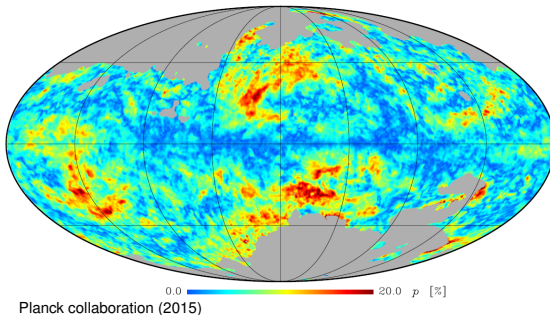


Planck collaboration (2015)

- In the disk :  $\vec{B}_{\text{ord}}$  is horizontal
- In the halo :  $\vec{B}_{\text{ord}}$  has a vertical component

# Polarization of dust thermal emission

Polarization fraction at 353 GHz (Planck)



Info on - Inclination of  $\vec{B}_{\text{ord}}$  to PoS :  $\cos^2 \gamma$

- Magnetic fluctuations :  $\frac{B_{\text{fluct}}}{B_{\text{ord}}}$

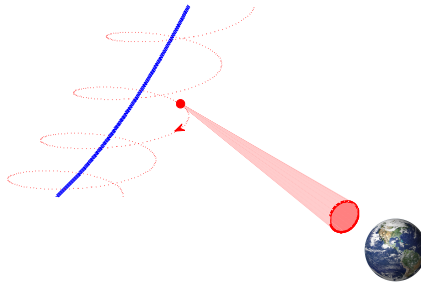
- Grain properties & alignment efficiency :  $p_{\text{max}}$  &  $F_{\text{align}}$

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# Physical concept



Relativistic electrons gyrating about magnetic field lines  
emit *synchrotron radiation*

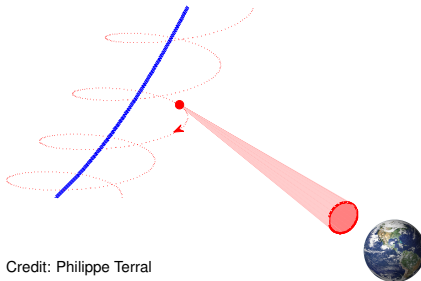


Credit: Philippe Terral

# Total & polarized intensities

Emissivity :  $\mathcal{E} = f(\alpha) n_{\text{CRE}} B_{\perp}^{\alpha+1} \nu^{-\alpha}$  &  $\mathcal{E}_{\text{pol}} = p_{\text{syn}} \mathcal{E}$  &  $\vec{\mathcal{E}}_{\text{pol}} \perp \vec{B}_{\perp}$


- Total intensity :  $I = \int \mathcal{E} ds$    $B_{\perp}$
- Polarized intensity :  $\vec{P} = \int \vec{\mathcal{E}}_{\text{pol}} ds$    $(\vec{B}_{\perp})_{\text{ord}}$




Credit: Philippe Terral

# Total & polarized intensities

Emissivity :  $\mathcal{E} = f(\alpha) n_{\text{CRE}} B_{\perp}^{\alpha+1} \nu^{-\alpha}$  &  $\mathcal{E}_{\text{pol}} = p_{\text{syn}} \mathcal{E}$  &  $\overleftrightarrow{\mathcal{E}}_{\text{pol}} \perp \overrightarrow{B}_{\perp}$

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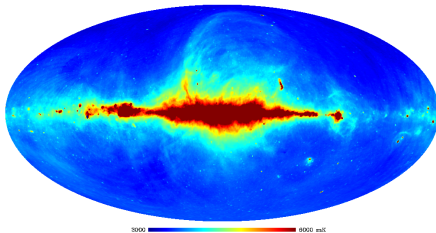
- Polarized intensity :  $\overleftrightarrow{P} = \int \overleftrightarrow{\mathcal{E}}_{\text{pol}} ds$    $(\overrightarrow{B}_{\perp})_{\text{ord}}$

$$\hookrightarrow Q + iU = \int \mathcal{E}_{\text{pol}} e^{2i\psi} ds$$



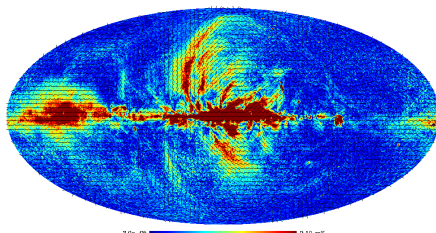
# Total & polarized intensities

TI at 1.4 GHz (25m Stockert + 30m Villa Elisa)



Credit: Tess Jaffe

PI &  $\vec{B}_\perp$  sectors at 23 GHz (WMAP)



Credit: Tess Jaffe

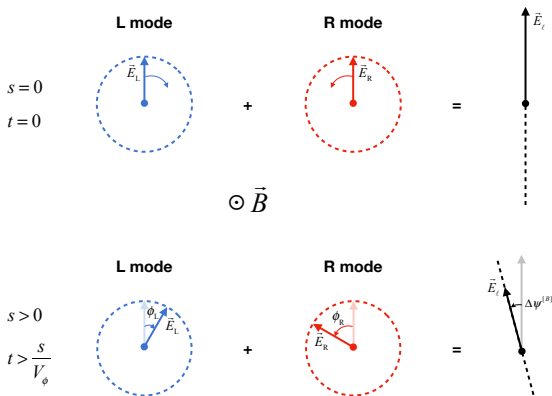
- ☞ - Near the Sun :  $B_{\text{ord}} \sim 3 \mu\text{G}$  &  $B_{\text{tot}} \sim 5 \mu\text{G}$
- In the disk :  $\vec{B}_{\text{ord}}$  is horizontal
- In the halo :  $\vec{B}_{\text{ord}}$  has a vertical component

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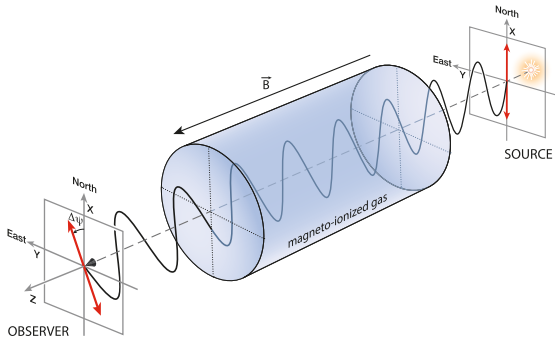
# Physical concept

When a linearly polarized radio wave travels through a magneto-ionized medium, the orientation of linear polarization undergoes *Faraday rotation*



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When a linearly polarized radio wave travels through a magneto-ionized medium, the orientation of linear polarization undergoes *Faraday rotation*

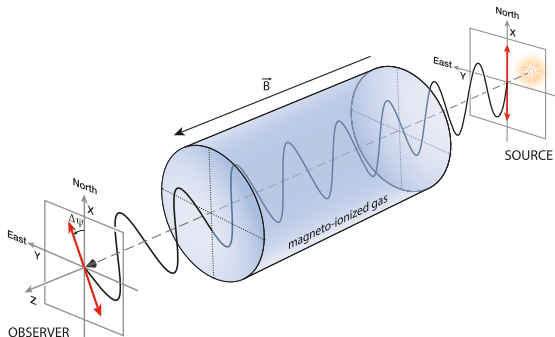


Credit: Theophilus Britt Griswold (NASA Goddard)

# Rotation angle & rotation measure

Rotation angle :  $\Delta\psi = \text{RM} \lambda^2$

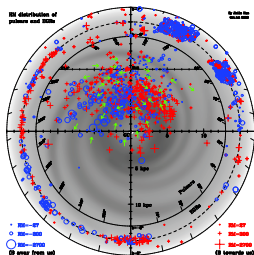
Rotation measure :  $\text{RM} = C \int n_e B_{\parallel} ds$



Credit: Theophilus Britt Griswold (NASA Goddard)

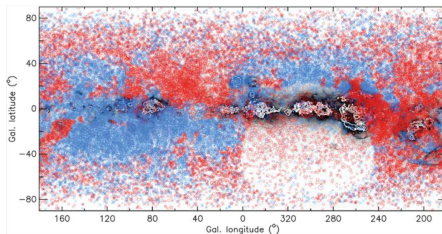
# Rotation measures

RMs of pulsars & EGRSs with  $|b| < 8^\circ$



Han (2009)

RMs of EGRSs [NVSS ( $\delta > -40^\circ$ ) + S-PASS ( $\delta < 0^\circ$ )]



Schnitzeler et al. (2019)

☞ - Near the Sun :  $B_{\text{reg}} \simeq 1.5 \mu\text{G}$  &  $B_{\text{tot}} \sim 5 \mu\text{G}$

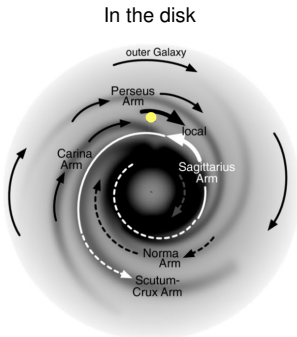
$\vec{B}_{\text{reg}}$  is **nearly azimuthal** ( $\simeq -8^\circ$  from  $\hat{e}_\phi$ )

- In the disk :  $\vec{B}_{\text{reg}}$  is **horizontal** & **mostly azimuthal**, with **reversals** in  $B_\phi$   
 $\vec{B}_{\text{reg}}$  probably has a **spiral** shape

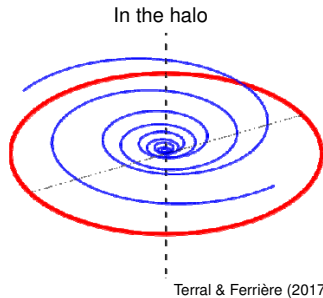
- In the halo :  $\vec{B}_{\text{reg}}$  is **CCW** at  $z > 0$  & **CW** at  $z < 0$

$\vec{B}_{\text{reg}}$  possibly has an **upward spiraling** shape

# Rotation measures



van Eck et al. (2011)



Terral & Ferrière (2017)

- Near the Sun :  $B_{\text{reg}} \simeq 1.5 \mu\text{G}$  &  $B_{\text{tot}} \sim 5 \mu\text{G}$

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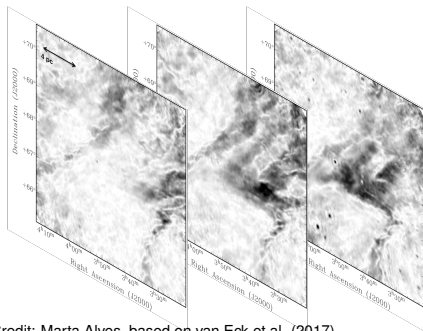
# Physical concept

- Underlying processes
  - Galactic **synchrotron emission** : linearly polarized
  - **Faraday rotation** :  $\lambda$ -dependent
- General idea
  - Measure synchrotron polarized intensity at many different  $\lambda$
  - Convert  $\lambda$ -dependence into  $s$ -dependence [ $\Phi$ -dependence]
- Output

**Faraday cube** = 3D cube of synchrotron polarized emission as  $fc(\alpha, \delta, \Phi)$

# Output

Faraday cube toward Fan region (LOFAR)



Credit: Marta Alves, based on van Eck et al. (2017)

3 slices at

$$\Phi_1 = -2.0 \text{ rad m}^{-2}$$

$$\Phi_2 = -1.5 \text{ rad m}^{-2}$$

$$\Phi_3 = -1.0 \text{ rad m}^{-2}$$

# What can be learned

- In Faraday cube
  - Uncover **synchrotron-emitting** & **Faraday-rotating** features
  - Identify these features with interstellar matter structures

- For **synchrotron-emitting** regions

$$\int \vec{F}(\Phi) d\Phi \quad \Leftrightarrow \quad \vec{B}_\perp$$

- For **Faraday-rotating** regions

$$\Delta\Phi \quad \Leftrightarrow \quad B_\parallel$$

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# Observational tools

- Synchrotron emission

☞  $\vec{B}_\perp$  (strength & orientation)

- Faraday rotation

☞  $B_\parallel$  (strength & sign)

- Polarization of dust thermal emission

☞  $\vec{B}_\perp$  (orientation only)

# Main characteristics

## Spiral galaxies

- $\vec{B}$  has an **ordered** component
- $B_{\text{tot}} \sim \text{a few } \mu\text{G}$
- \* *Face on*
  - Disk :  $\vec{B}_{\text{ord}}$  follows the **spiral arms**
- \* *Edge on*
  - Disk :  $\vec{B}_{\text{ord}}$  is **horizontal**
  - Halo :  $\vec{B}_{\text{ord}}$  has an **X shape**

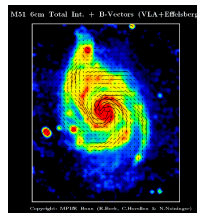
## Elliptical galaxies

- $\vec{B}$  has only **fluctuating** component
- $B_{\text{tot}} \sim \text{a few } \mu\text{G}$

M51



HST

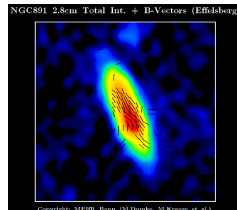


Effelsberg + VLA (6.2 cm)

NGC 891



CFHT



Effelsberg (2.8 cm)

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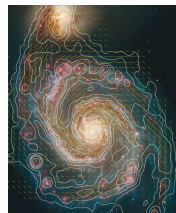
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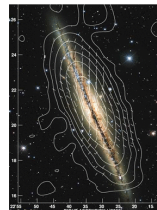


Effelsberg + VLA (6.2 cm)

NGC 891



CFHT



Effelsberg (3.6 cm)



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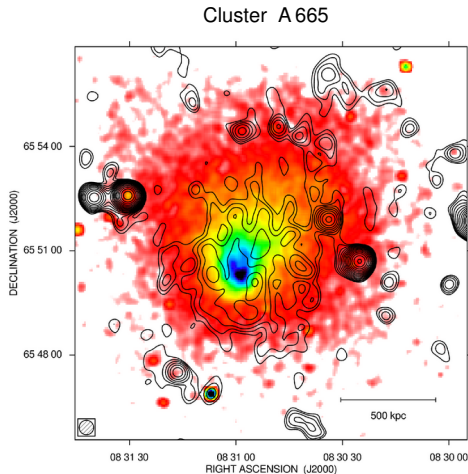
# Clusters of galaxies

## Observational tools

- Synchrotron emission
- Faraday rotation

## Main characteristics

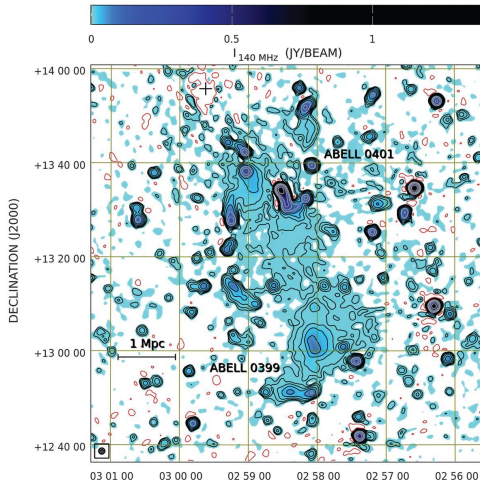
- $\vec{B}$  has only **fluctuating** component
- $B_{\text{tot}} \sim (0.1 - \text{a few}) \mu\text{G}$



Radio 1.4 GHz + X-ray 0.8-4 keV (Vacca et al. 2018)

# Cosmic filaments

Radio synchrotron filament between Abell 399 & Abell 401 (LOFAR)



Govoni et al. (2019)

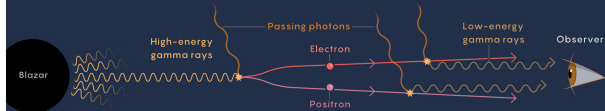
# Cosmic voids

## The Hidden Blazar Effect

Blazars are bright beams of energy powered by supermassive black holes. These beams can be used to detect the presence of cosmic magnetic fields.

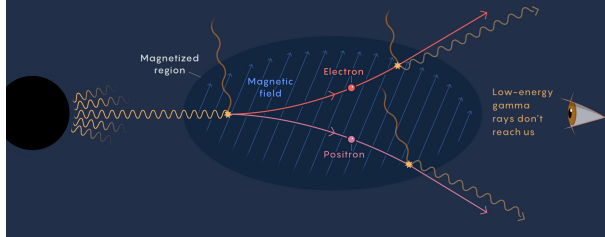
### A BEAM'S LONG JOURNEY

A blazar creates high-energy gamma rays that occasionally morph into an electron and a positron. These then create lower-energy gamma rays.



### A GAMMA-RAY SWERVE

If the high-energy gamma rays pass through a magnetic field, the field will deflect any electrons and positrons that get created. The resulting low-energy gamma rays will point away from the observer.



- Delay of GeV echo

$B \gtrsim 10^{-20.5} \text{ G}$

(Takahashi et al. 2013)

- Size of GeV halo

$B \sim (10^{-17} - 10^{-15}) \text{ G}$

(Chen et al. 2015)