

# Search for double beta decay of $^{130}\text{Te}$ to the $0^+$ states of $^{130}\text{Xe}$ with the CUORE experiment



PhD Candidate  
**G. Fantini<sup>1,2</sup>**

Advisors  
**C. Bucci<sup>2</sup>**  
**C. Tomei<sup>3</sup>**

<sup>1</sup> Gran Sasso Science Institute, L'Aquila I-67100, Italy

<sup>2</sup> INFN - Laboratori Nazionali del Gran Sasso, Assergi (L'Aquila), I-67100, Italy

<sup>3</sup> INFN - Sezione di Roma, Roma, I-00185, Italy



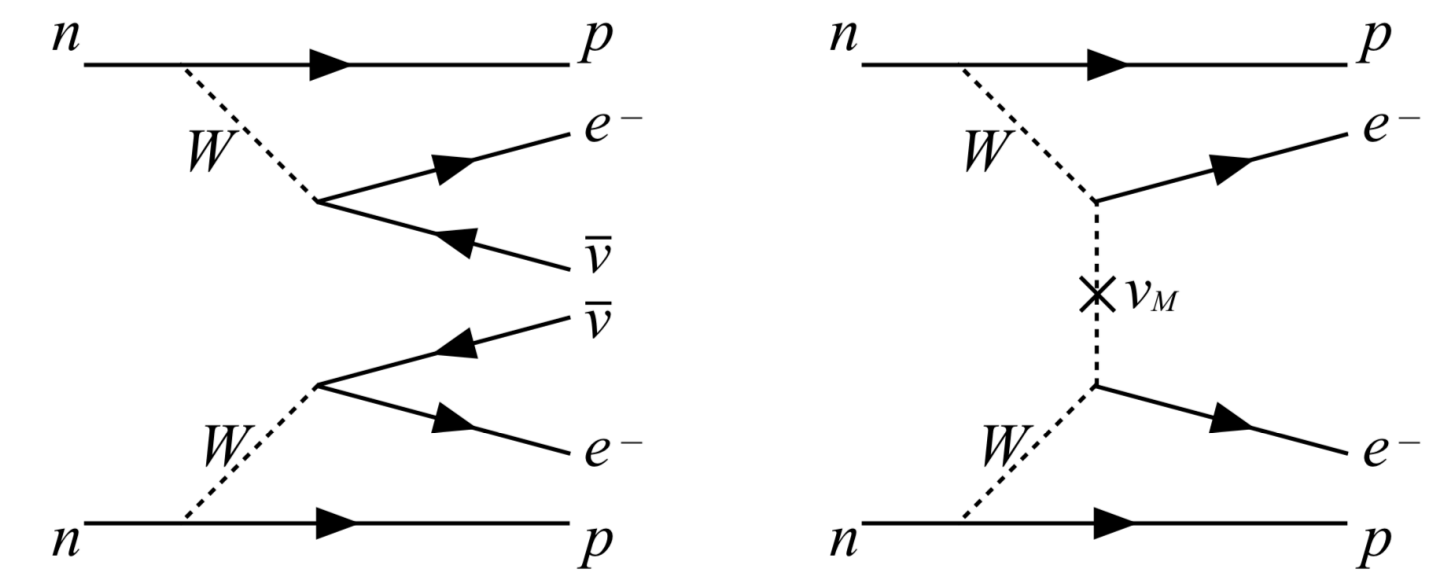
---

# Outline

- Introduction
- The CUORE experiment
  - Operating principles
  - Data processing
  - Detector characterization
- Ground state  $0\nu\beta\beta$  decay search
- Excited state searches

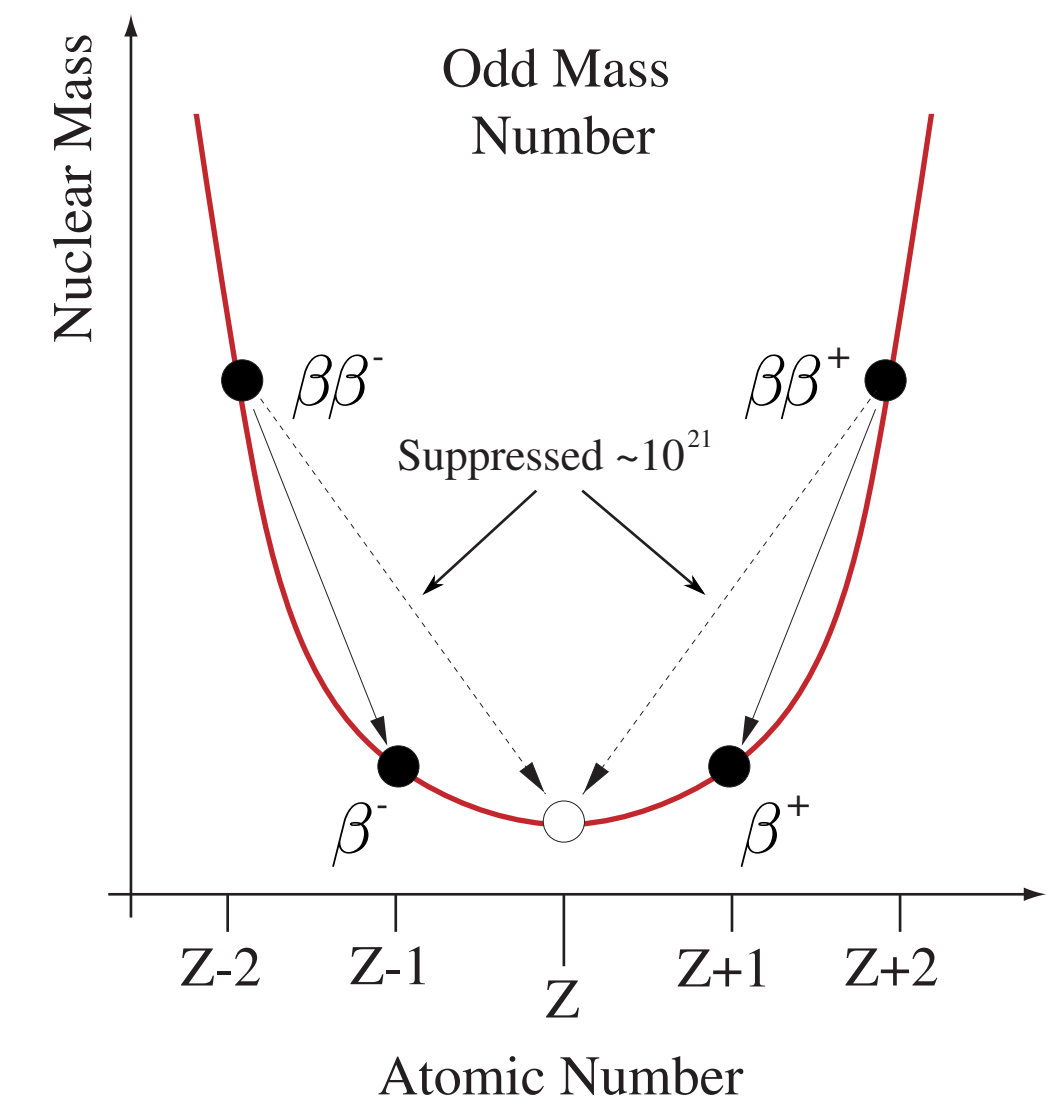
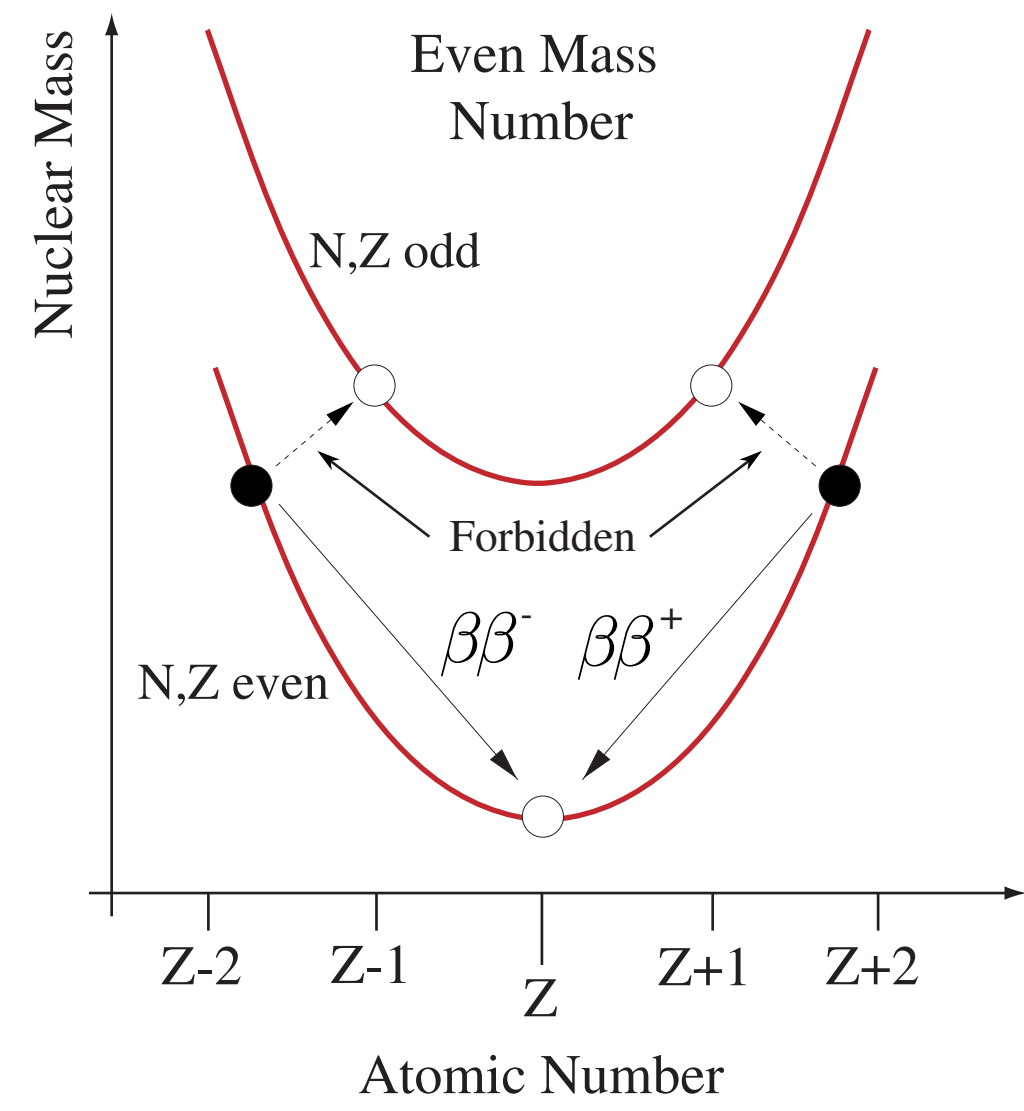
# Introduction

## Double Beta Decay



$$X = 2\bar{\nu}_e$$

- Allowed within the Standard Model
- 2<sup>nd</sup> order process
- observed in 12 isotopes

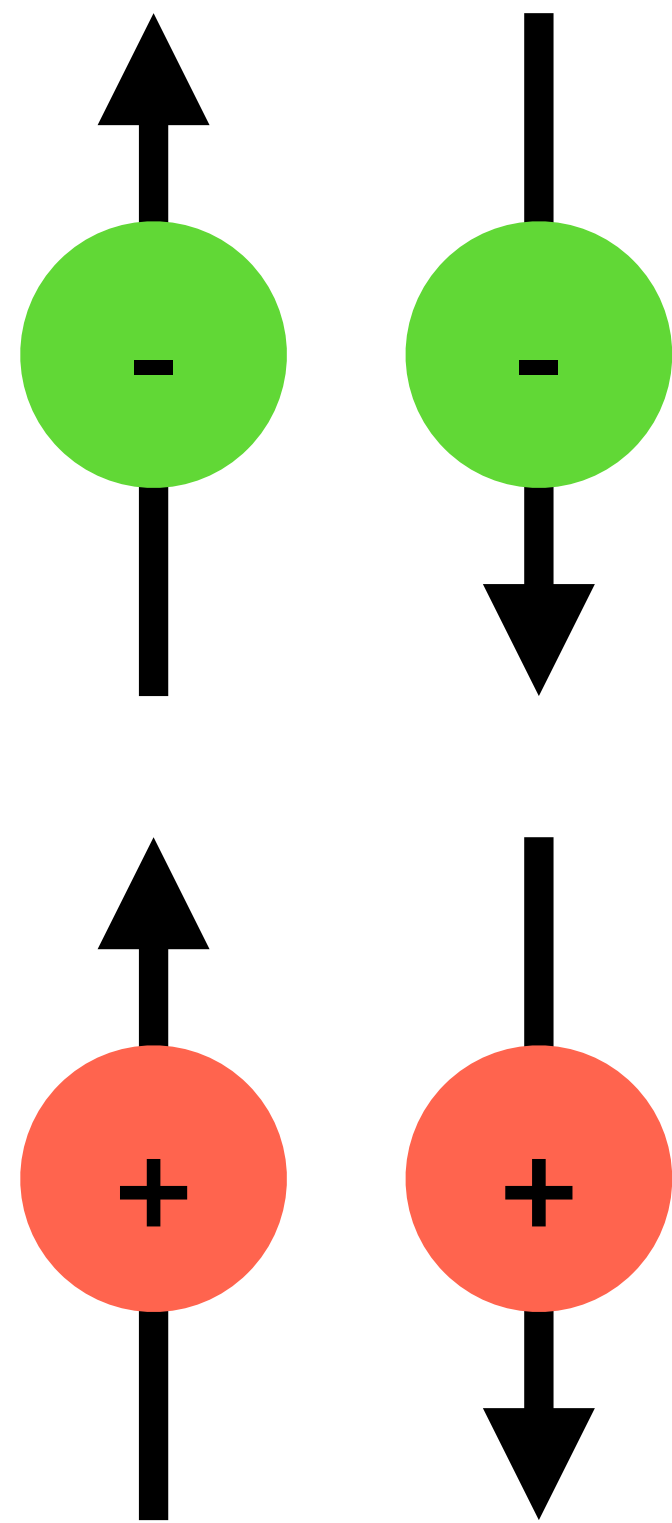


$$X = \emptyset$$

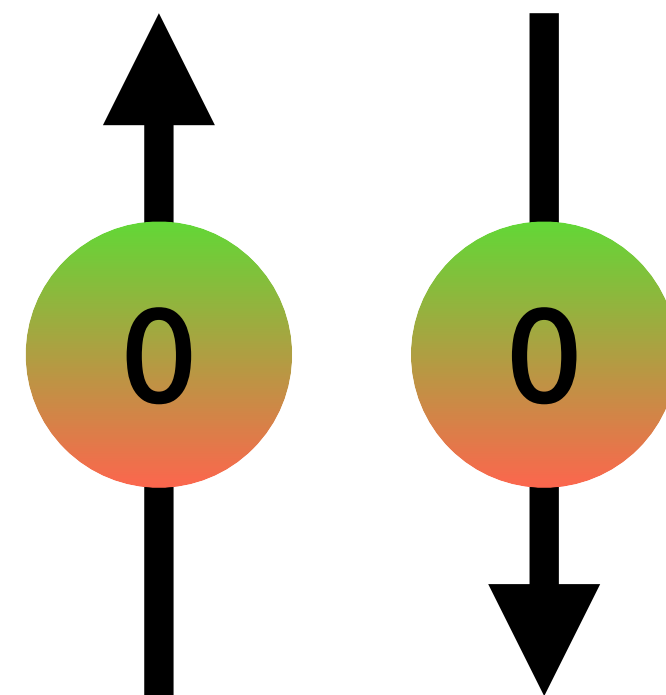
- Beyond Standard Model
- Majorana nature of neutrinos
- Lepton number violation

# Introduction

## Majorana neutrinos



Dirac



Majorana

- Vanishing internal quantum numbers (lepton number)
- Their own anti-particles
- Neutrino mass generation mechanism
- Mediate neutrino-less double beta decay
- Leptogenesis

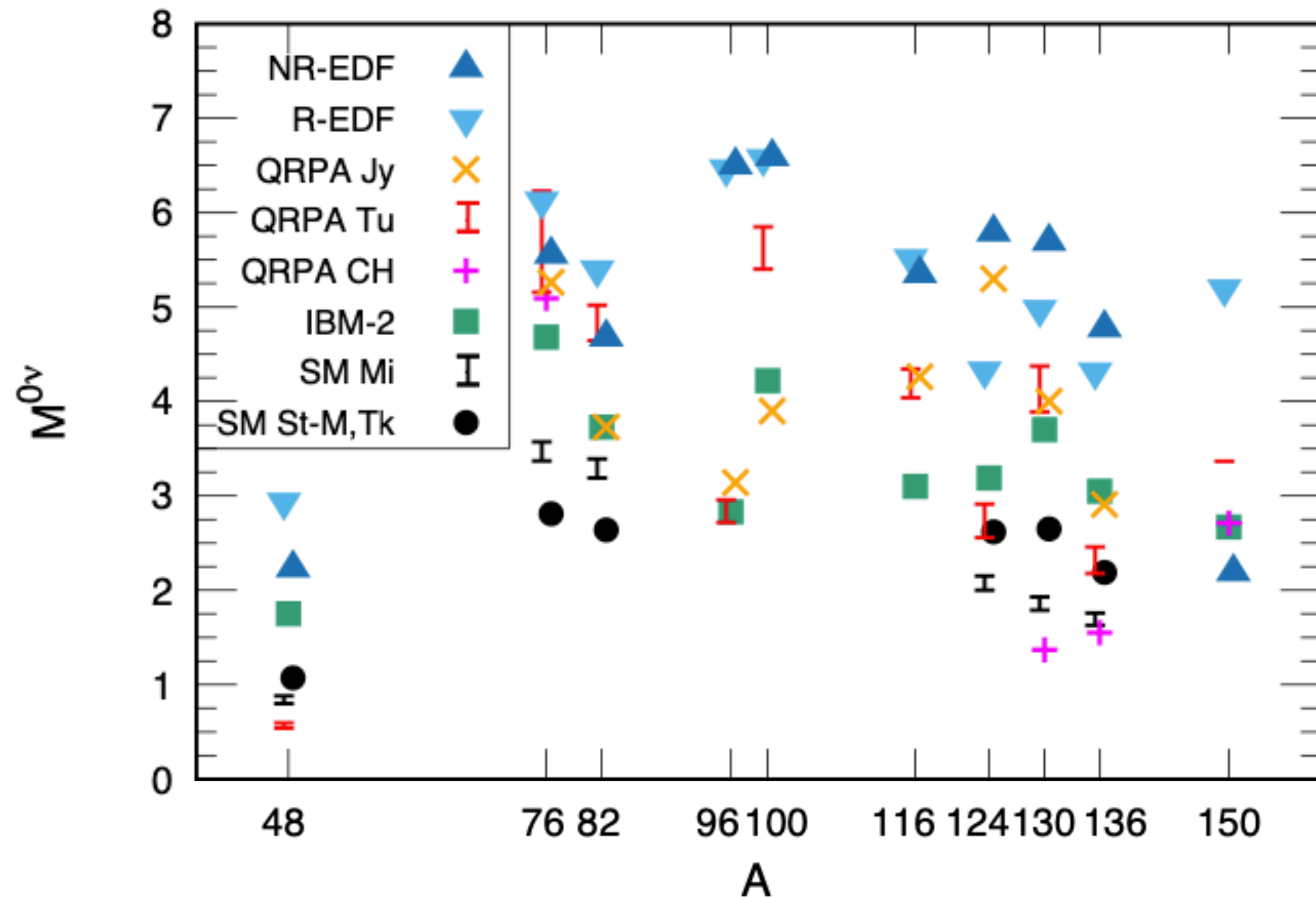


# Introduction

## Neutrinoless Double Beta Decay

$$\langle m_{\beta\beta} \rangle = \left| \sum_i U_{ei}^2 m_i \right|$$

$$\Gamma_{0\nu} = G_{0\nu} g_A^4 |M_{0\nu}|^2 \frac{\langle m_{\beta\beta} \rangle^2}{m_e^2}$$

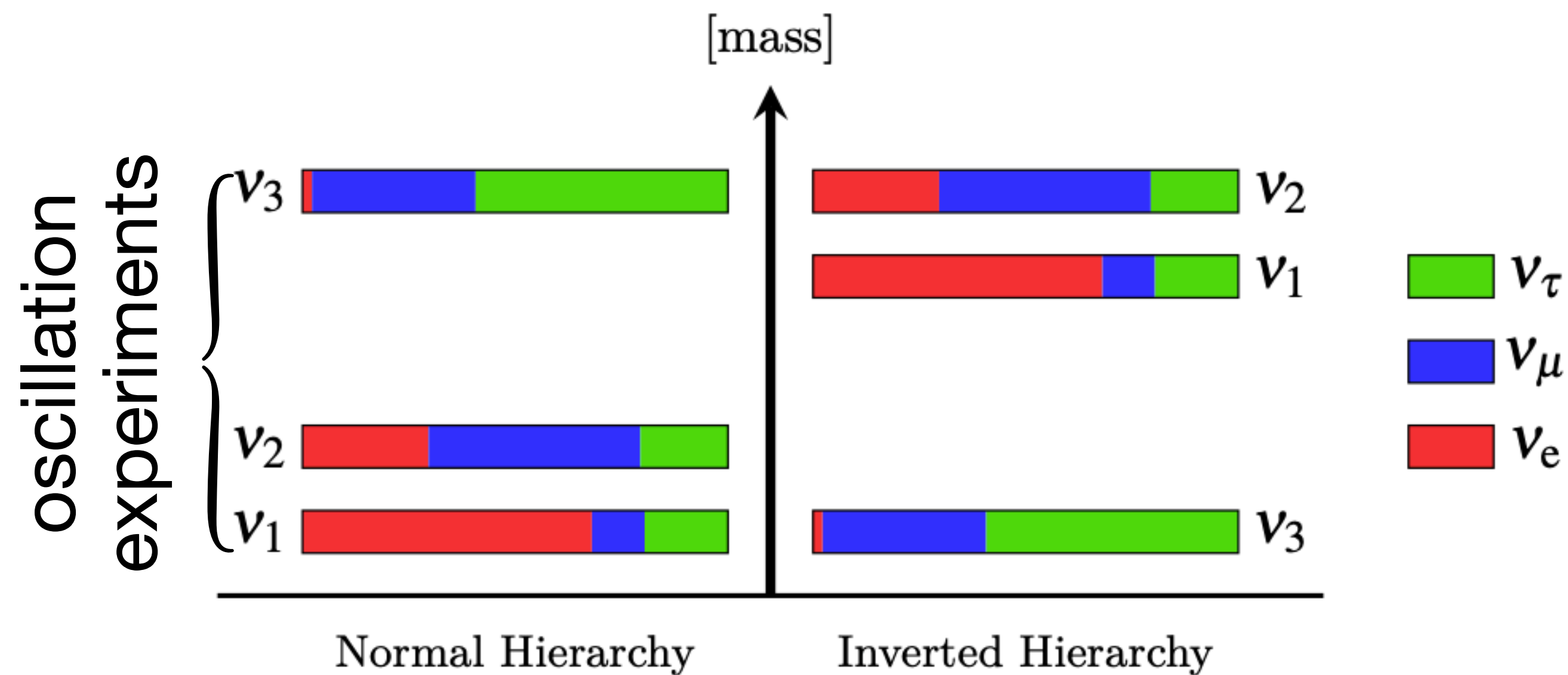


- $G_{0\nu}$  phase space factor  
(well known, small uncertainty)
- Nuclear Matrix Element (NME)  
(model dependence, large uncertainty)
- Effective Majorana mass  
(unknown)

[J.Engel, J.Menéndez, Rep. Prog. Phys. 80, 4 (2017)]



# Introduction



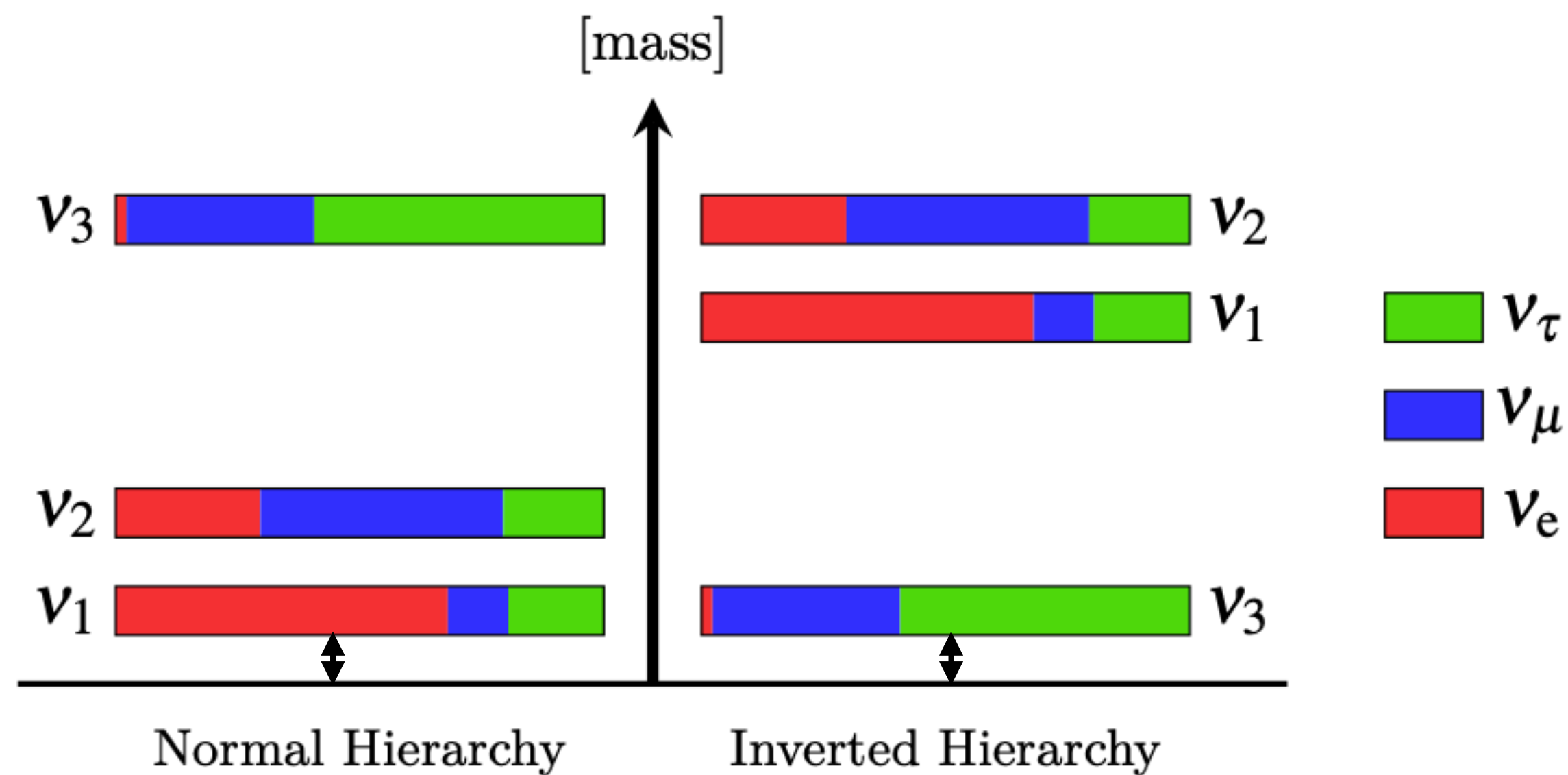
[G. Fantini, A. Gallo Rosso, F. Vissani, V. Zema, *Adv. Ser. Direct. High Energy Phys.* 28 (2018) 37-119]

$$\nu_{f,L} = \sum_{j=1}^3 U_{l,j} \nu_{j,L}$$

- Pontecorvo-Maki-Nakagawa-Sakata mixing matrix
- 3 mixing angles
- 1 (Dirac) / 3 (Majorana) CP violating phases



# Introduction



cosmology  
β decay  
 $0\nu\beta\beta$

[G. Fantini, A. Gallo Rosso, F. Vissani, V. Zema, *Adv. Ser. Direct. High Energy Phys.* 28 (2018) 37-119]

$$\nu_{f,L} = \sum_{j=1}^3 U_{l,j} \nu_{j,L}$$

- Pontecorvo-Maki-Nakagawa-Sakata mixing matrix
- 3 mixing angles
- 1 (Dirac) / 3 (Majorana) CP violating phases

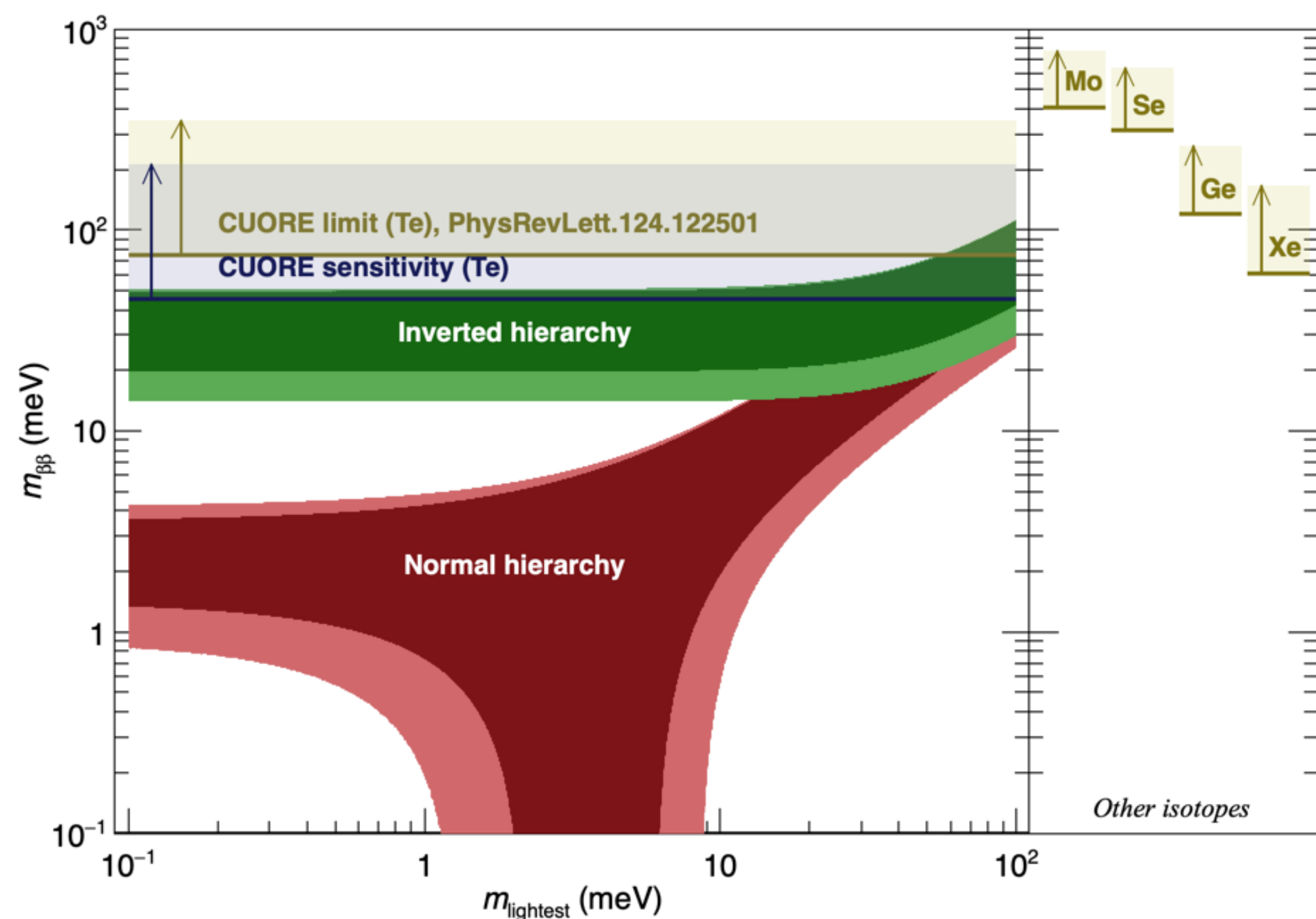


# Introduction

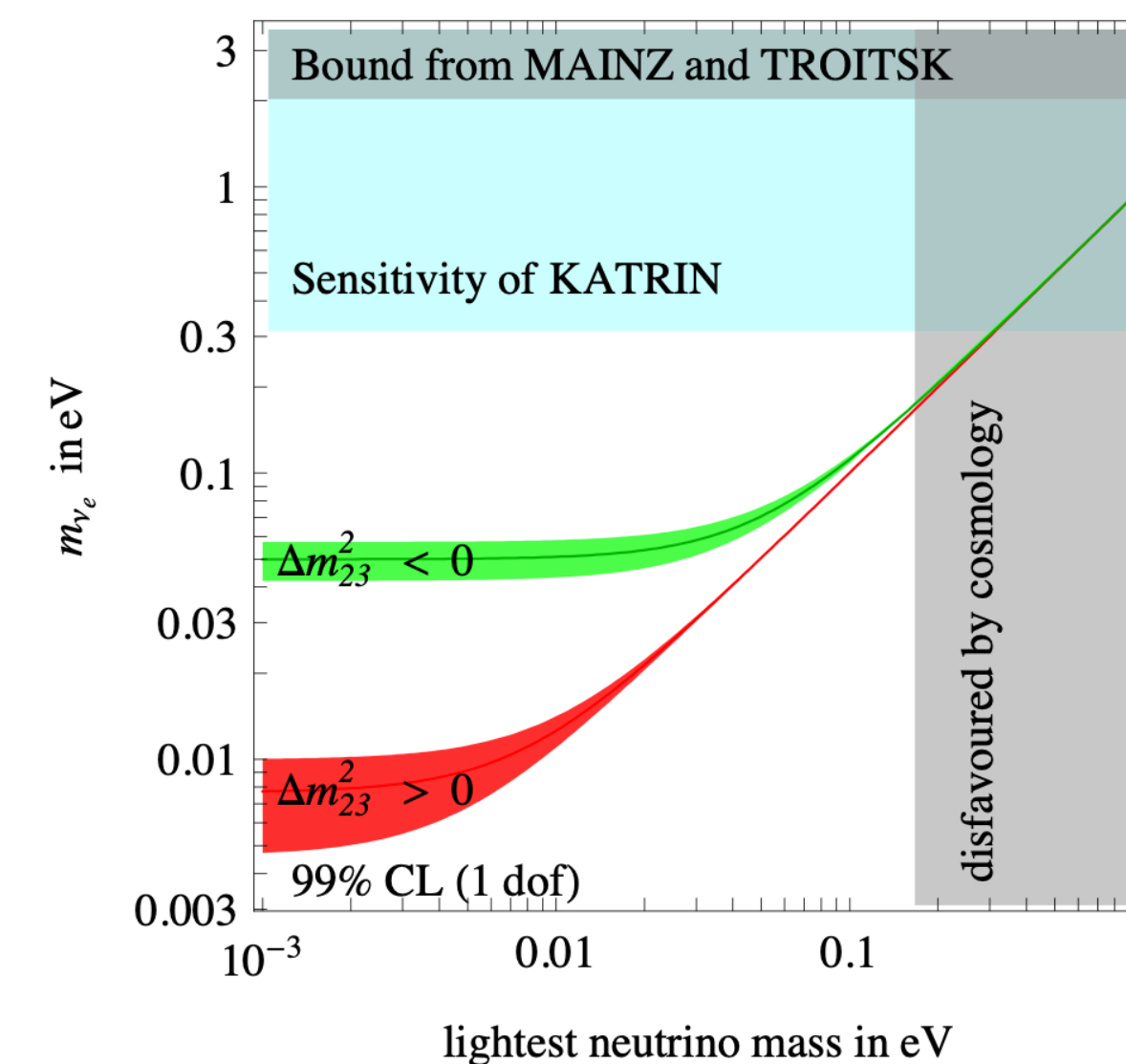
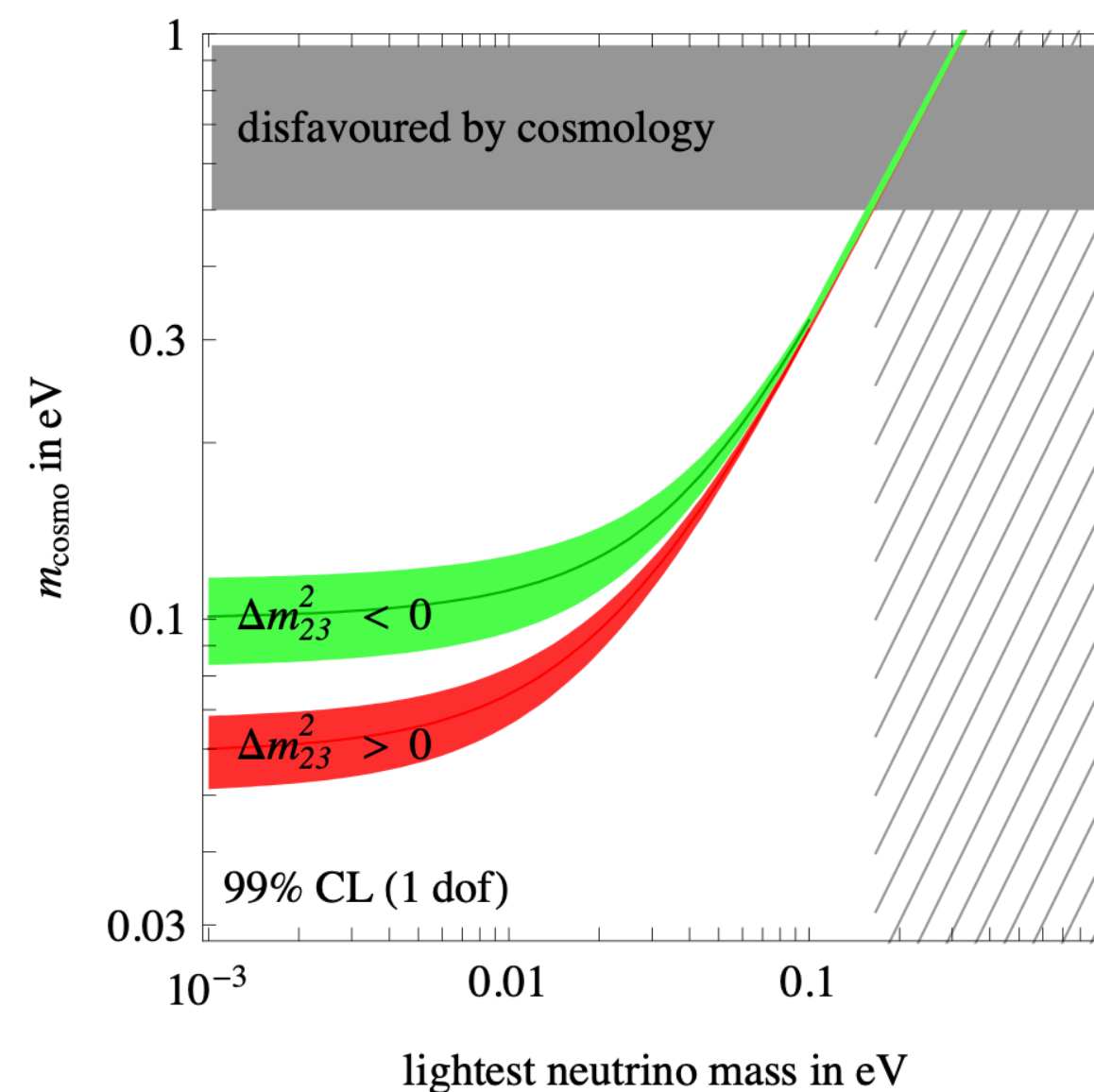
$$\langle m_{\beta\beta} \rangle = \left| \sum_i U_{ei}^2 m_i \right|$$

$$\Gamma_{0\nu} = G_{0\nu} g_A^4 |M_{0\nu}|^2 \frac{\langle m_{\beta\beta} \rangle^2}{m_e^2}$$

- Neutrino-less double beta decay
- Cosmology
- Beta decay endpoint



$$\Sigma = m_1 + m_2 + m_3$$



[A. Strumia and F. Vissani, IFUP-TH2004-1; arXiv:0606054]



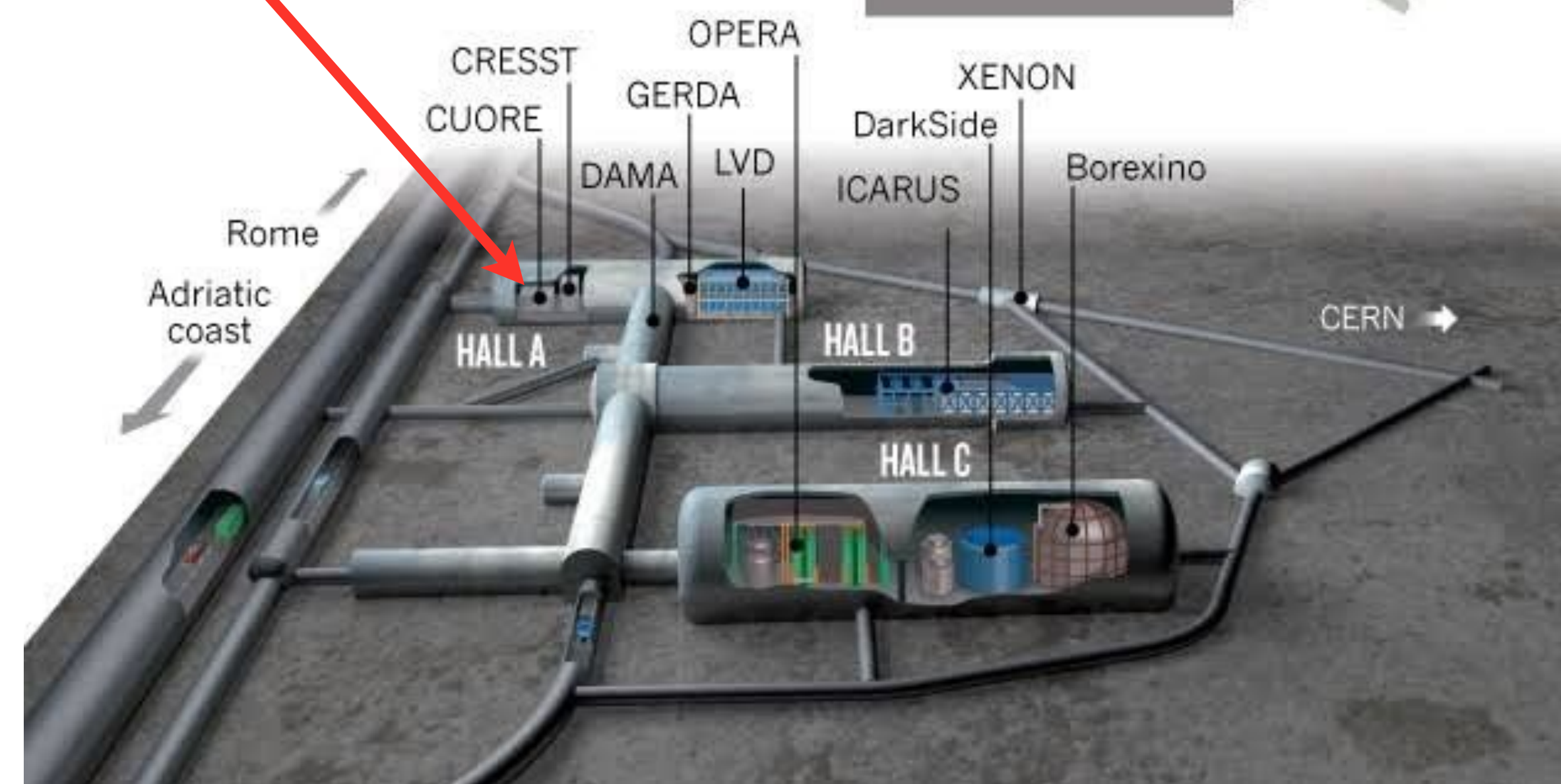
# The CUORE experiment





# The CUORE experiment

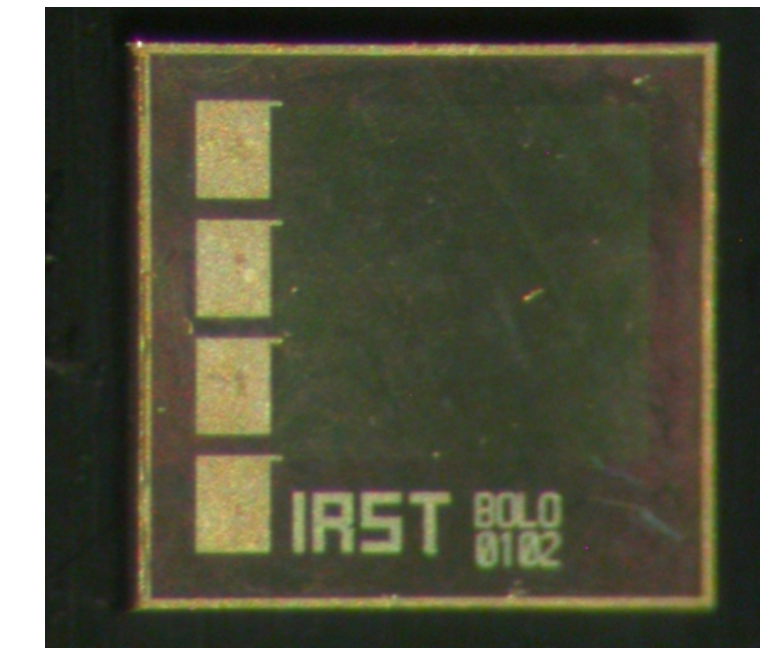
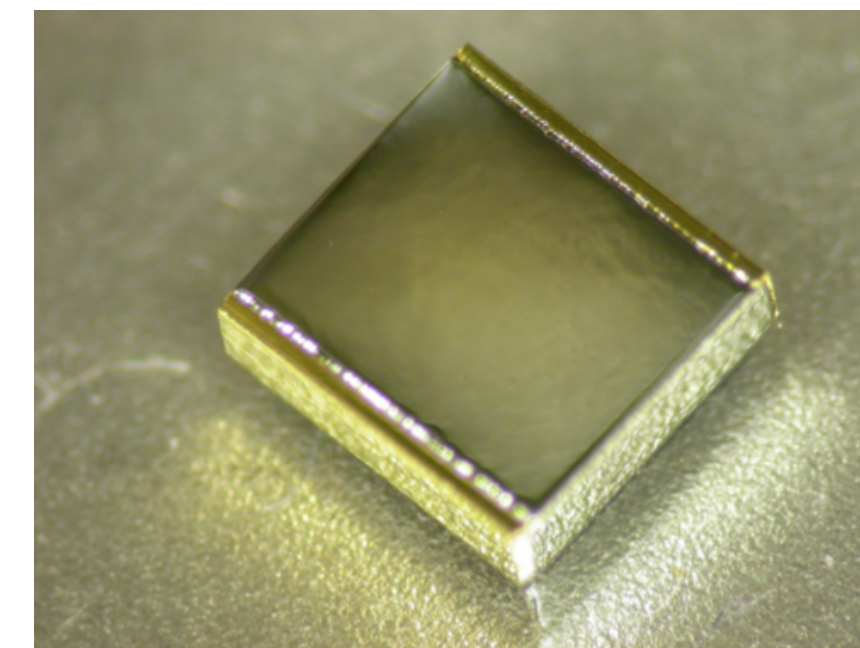
- Cryogenic Underground Observatory for Rare Events
- Hosted underground in the Hall A of the INFN Laboratori Nazionali del Gran Sasso
- 3600 m.w.e. rock overburden
- Main physics goal  $0\nu\beta\beta$
- High efficiency (source = detector)
- Ton scale bolometer array
- $\sim 10$  mK operating temperature



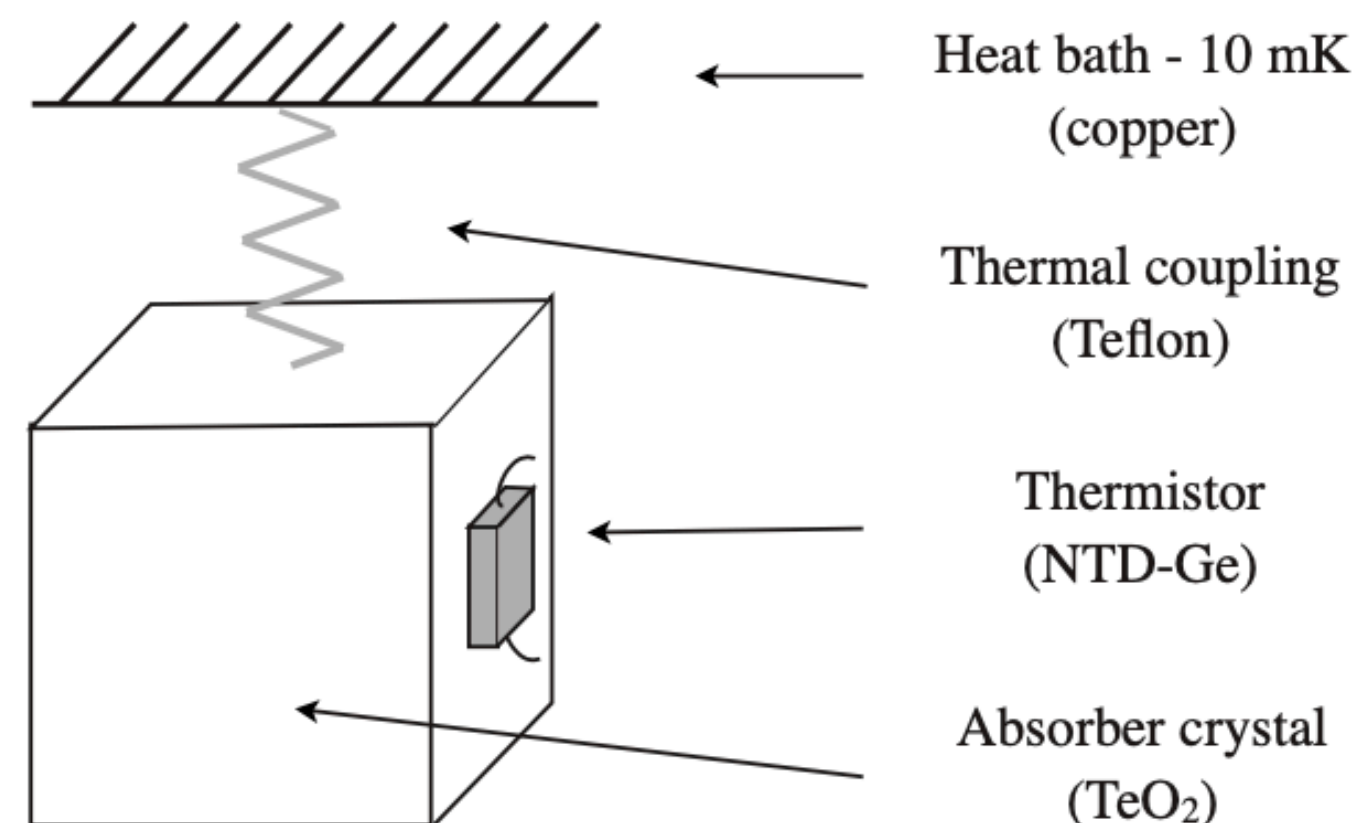


# The CUORE experiment

## Operating principles



[C. Alduino et al., JINST 11 P07009 (2016)]



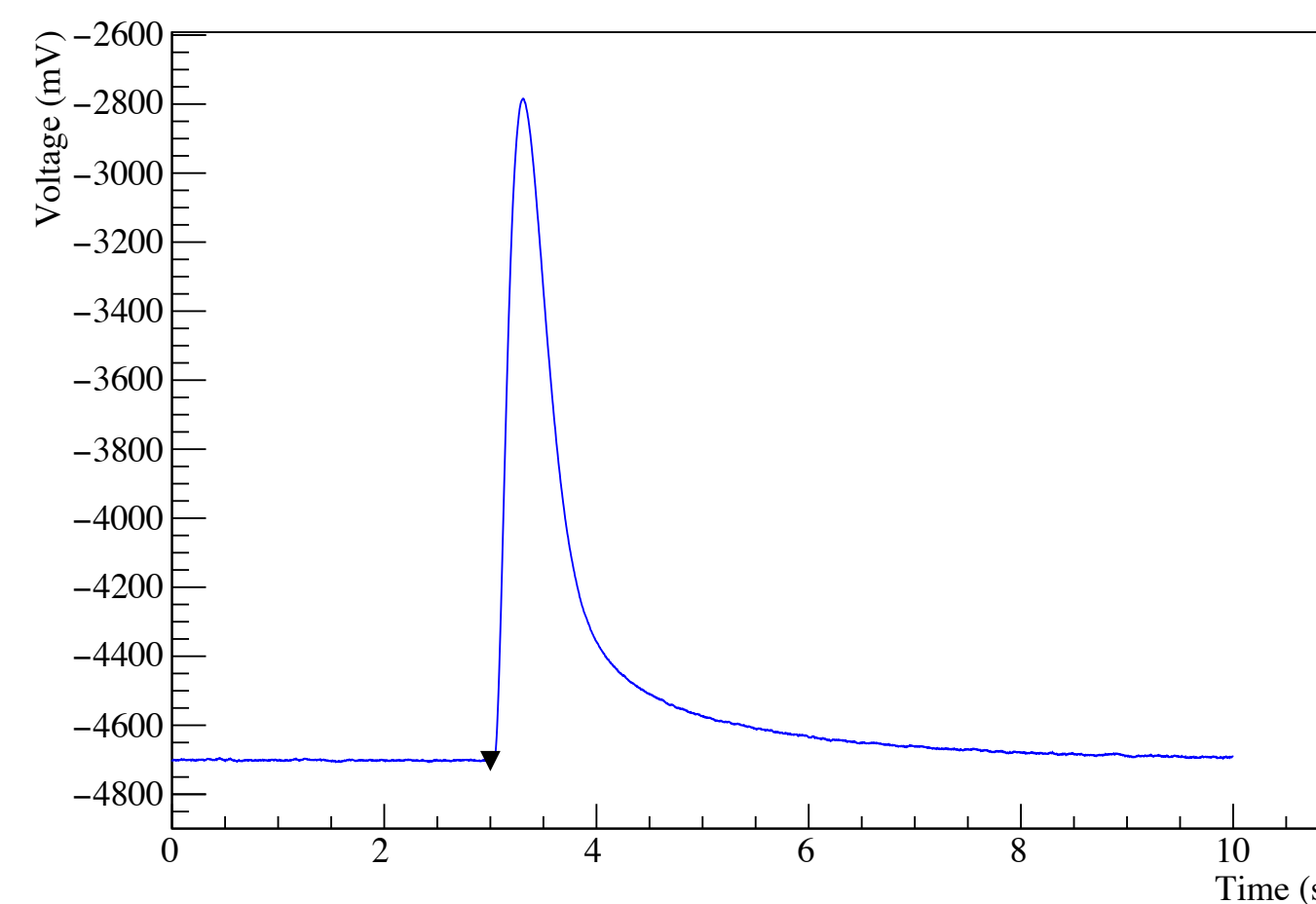
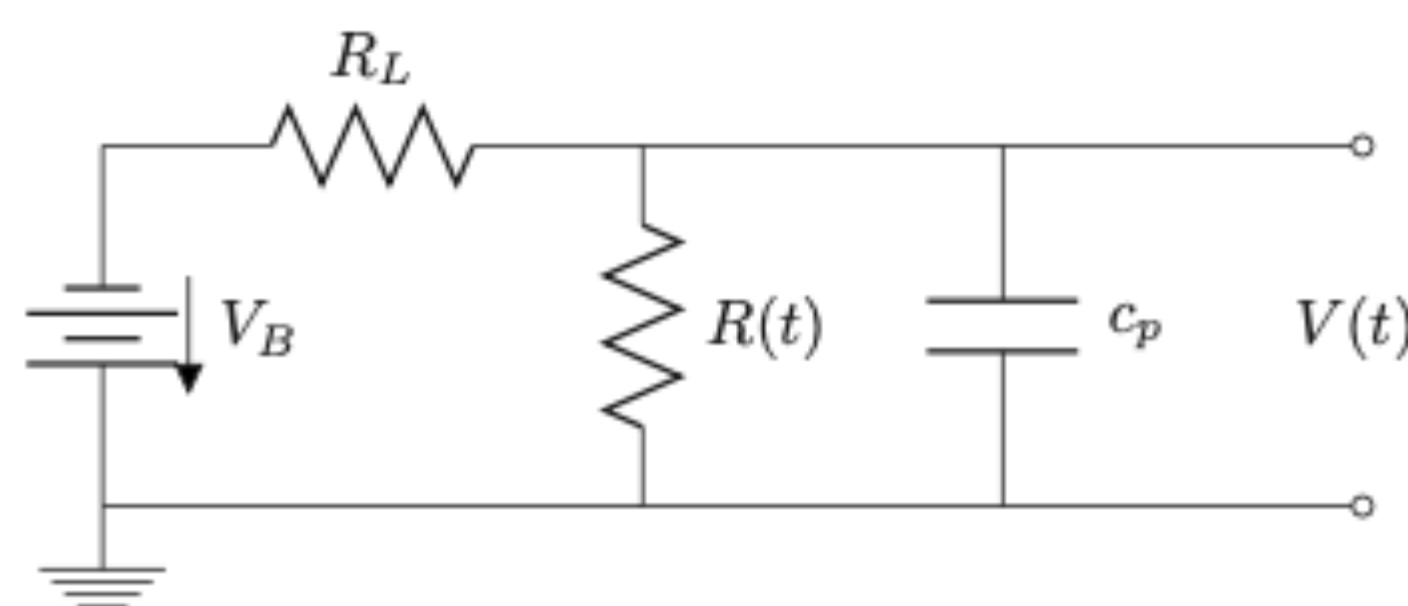
Crystals instrumented with

- NTD Ge thermistor
- Si heater

Weak thermal link to heat bath

Particle interactions in the crystals heat them up

Constant NTD current bias produces voltage pulses from temperature variations

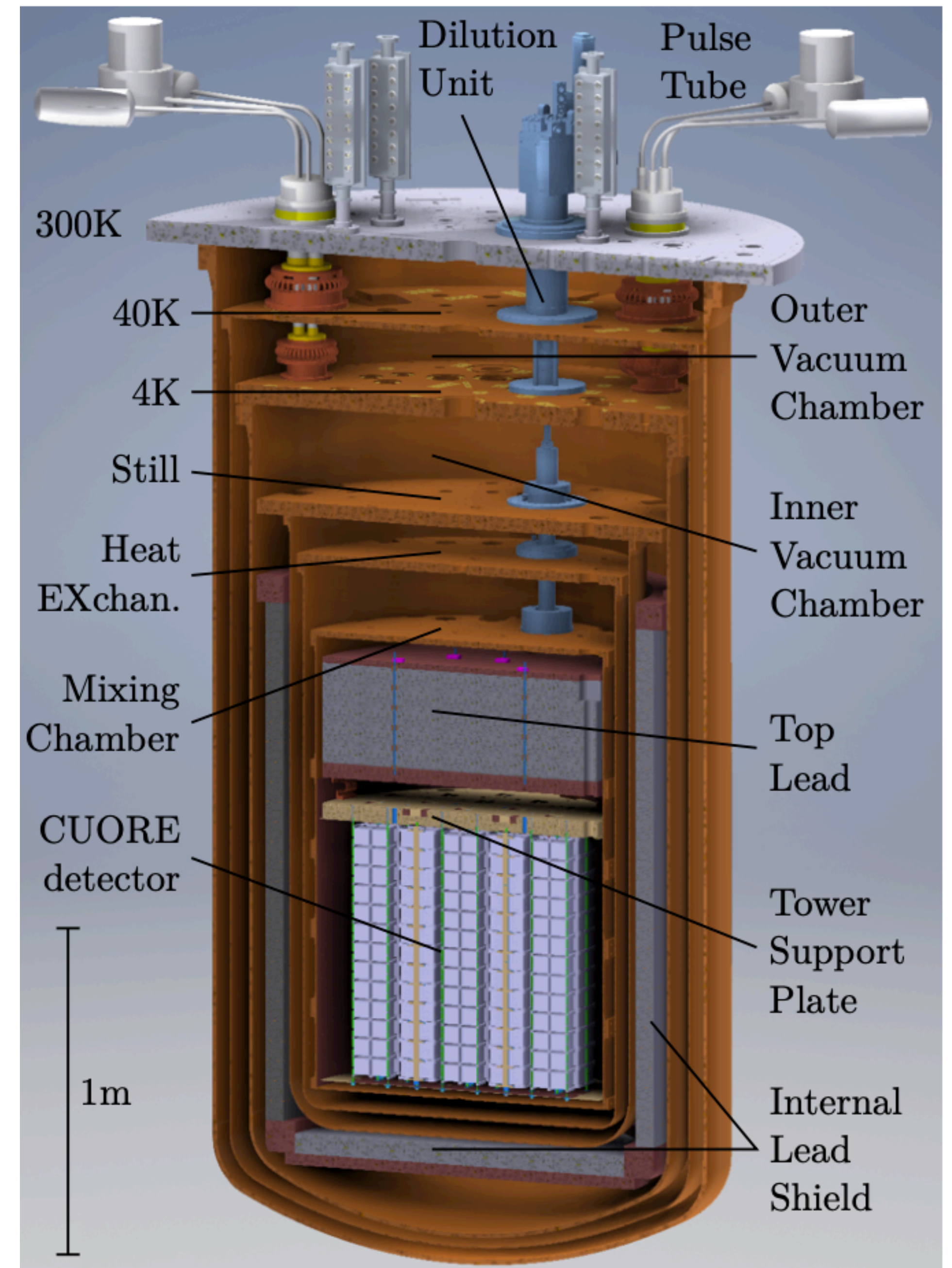


[M. Vignati, Journal of Applied Physics 108, 084903-1]



# The CUORE experiment

- 988  $5 \times 5 \times 5 \text{ cm}^3$  crystals arranged in 19 towers of 52 channels each
- Total active mass 742 kg
- High natural isotopic abundance ( $\sim 206 \text{ kg } ^{130}\text{Te}$ )
- Multi-stage cryogen-free  $^3\text{He}/^4\text{He}$  dilution refrigerator
- Pre-cooling supplied by Pulse Tube cryocoolers
- $\sim 10 \text{ mK}$  operating temperature
- Good energy resolution  $\sim 0.2\% \text{ FWHM/E}$

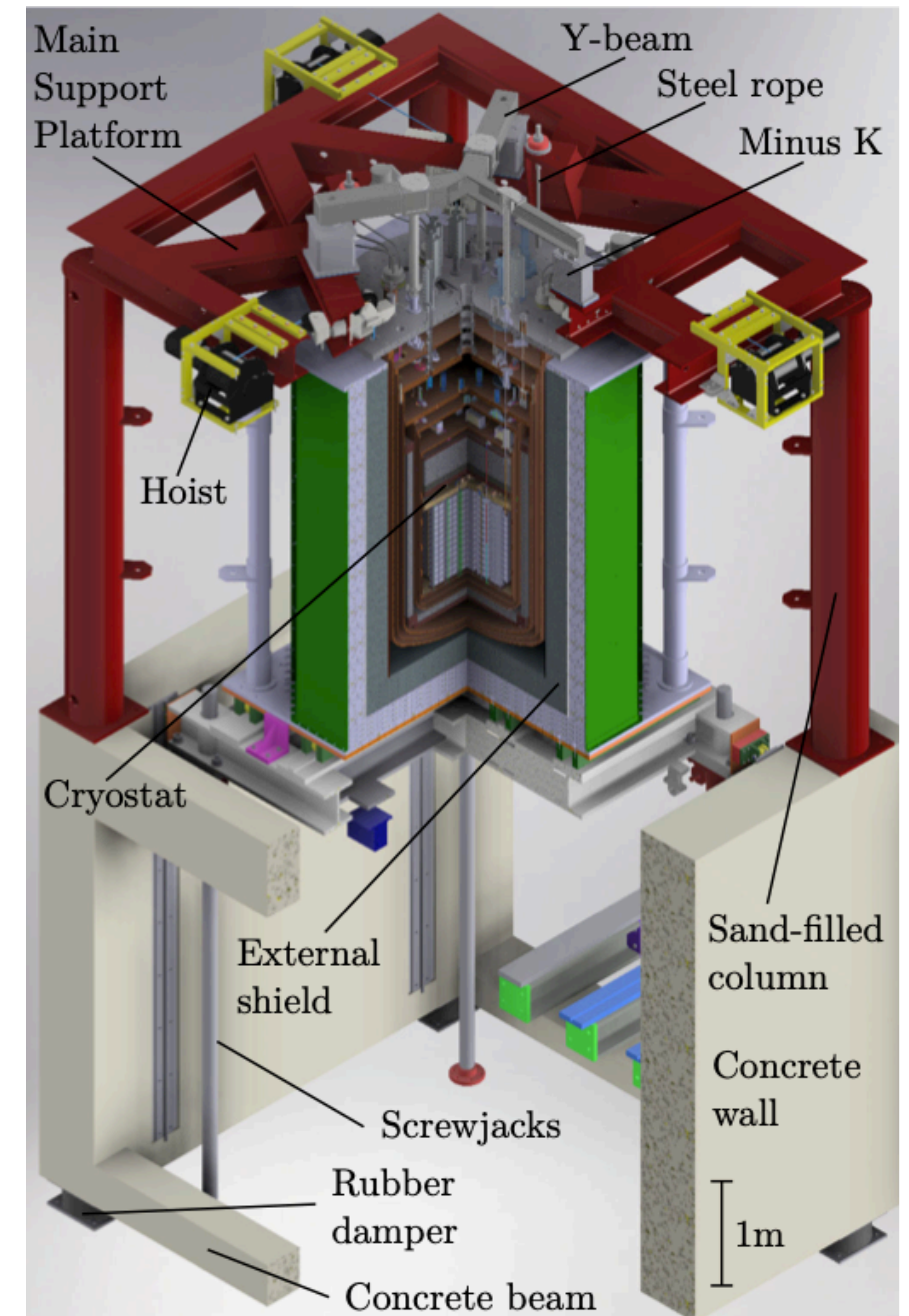


[C. Alduino et al., *Cryogenics* 102, 9 (2019)]



# The CUORE experiment

- Low background goal of 0.01 cts/(keV kg yr) at  $Q_{\beta\beta}$ 
  - 18 cm polyethylene + 2 cm borated material
  - 30 cm lead
  - Inner  $^{210}\text{Pb}$  depleted Roman lead shielding
- Thorough campaign of material assay
- Strict construction, transportation, assembly protocols
- Designed to reduce vibrations



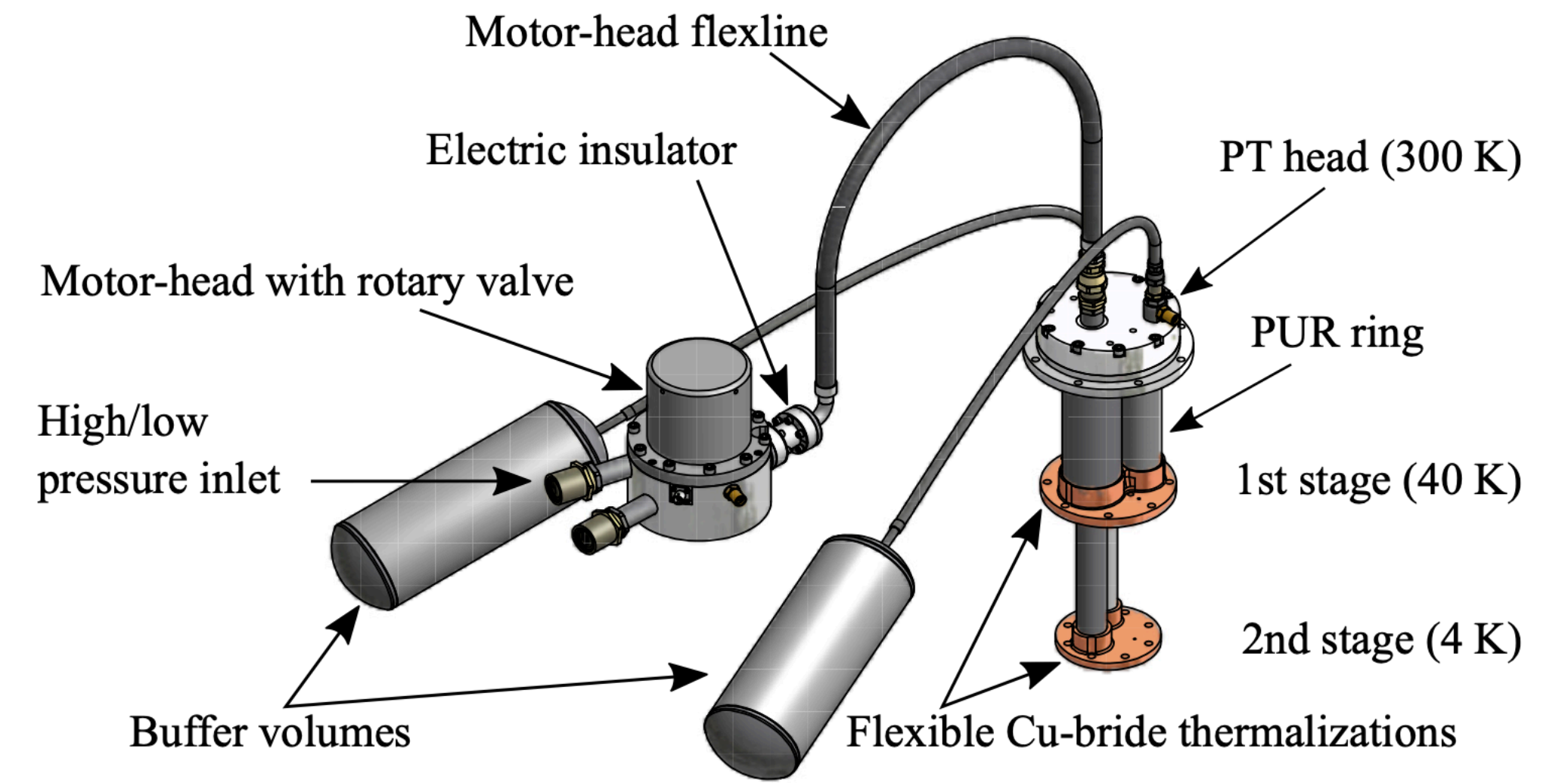
[C. Alduino et al., *Cryogenics* 102, 9 (2019)]



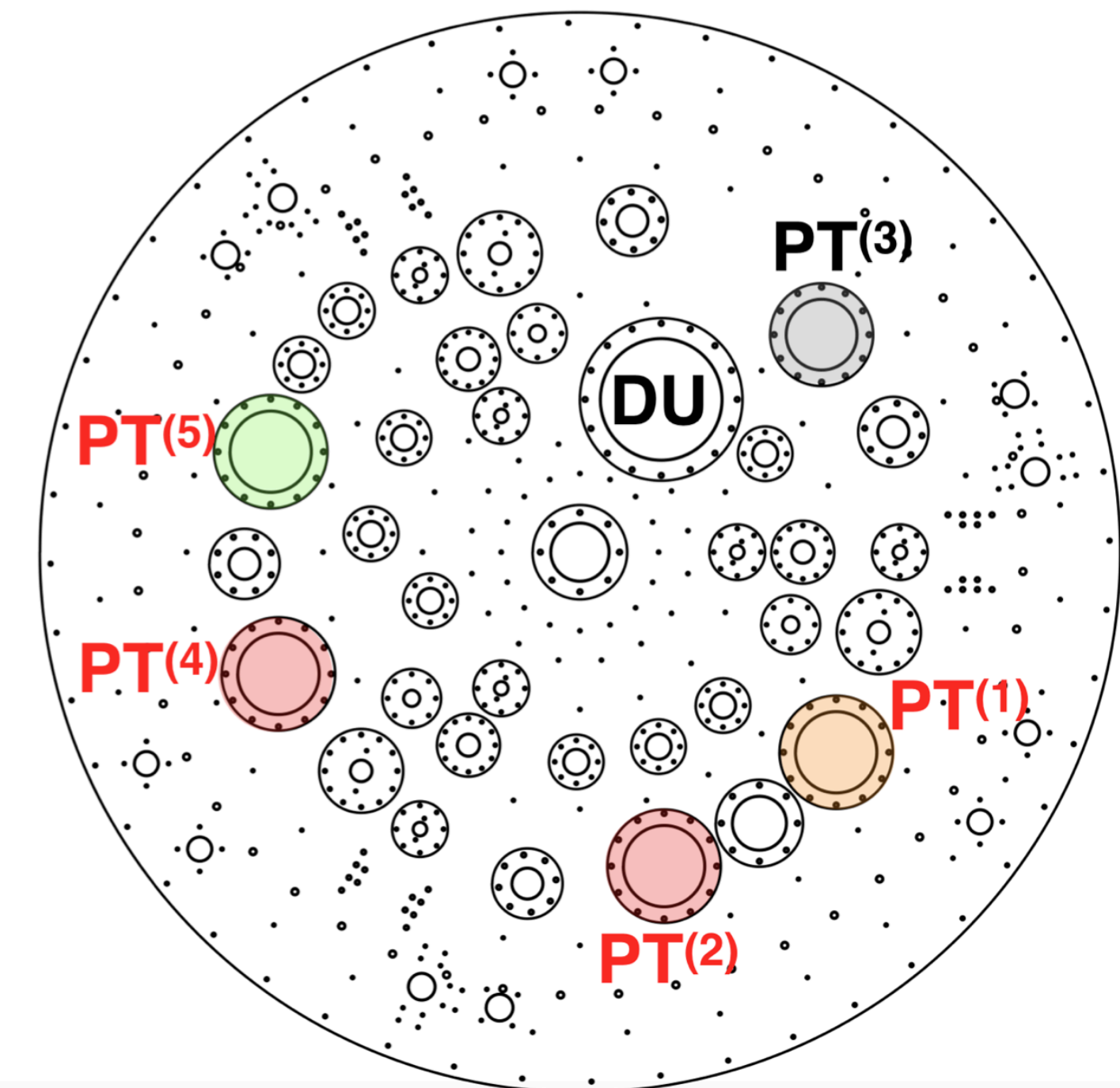
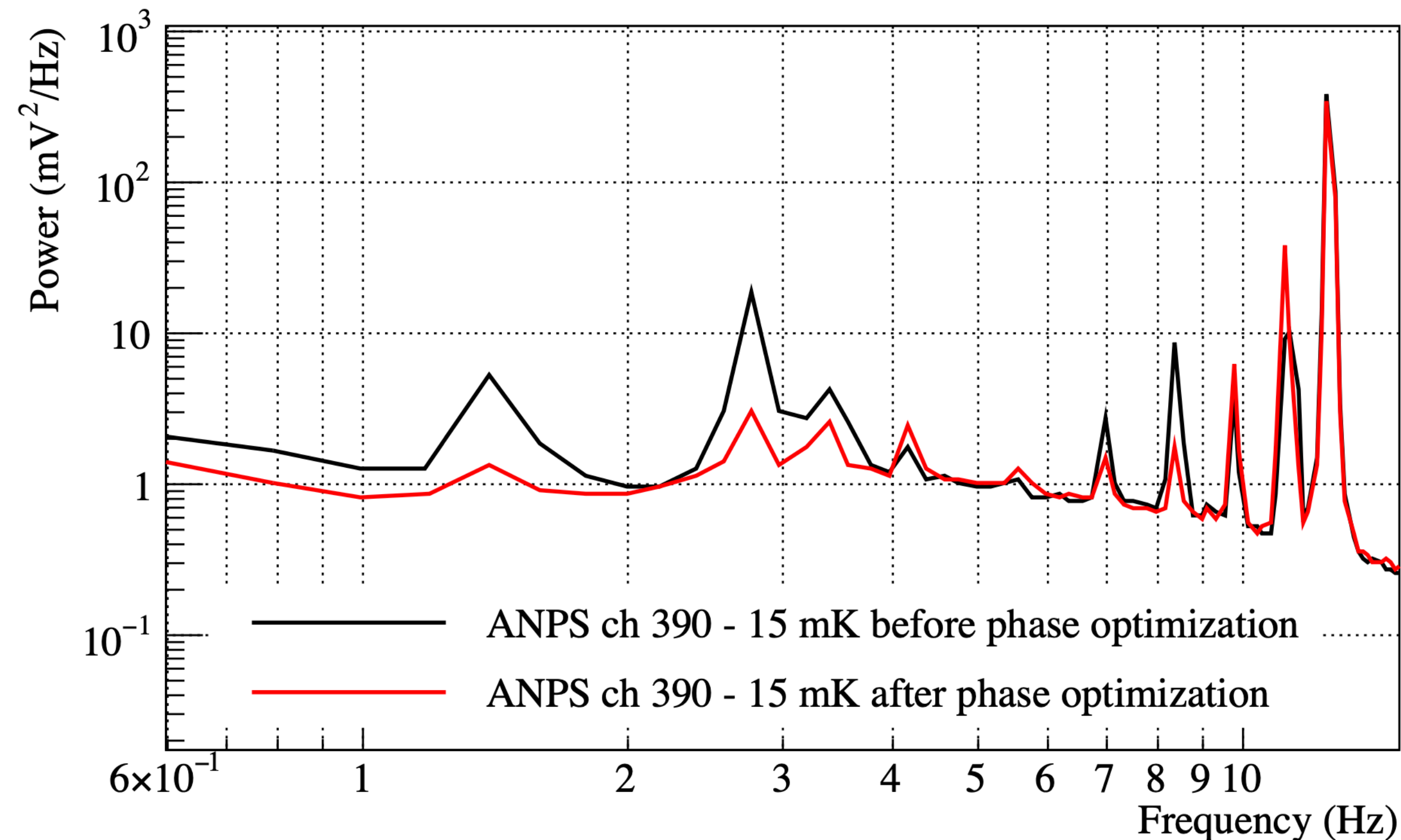
# The CUORE experiment

## Pulse Tubes noise reduction

Tuning the relative phase of the PT rotary valves allows 1.4 Hz (and higher harmonics) coherent noise suppression

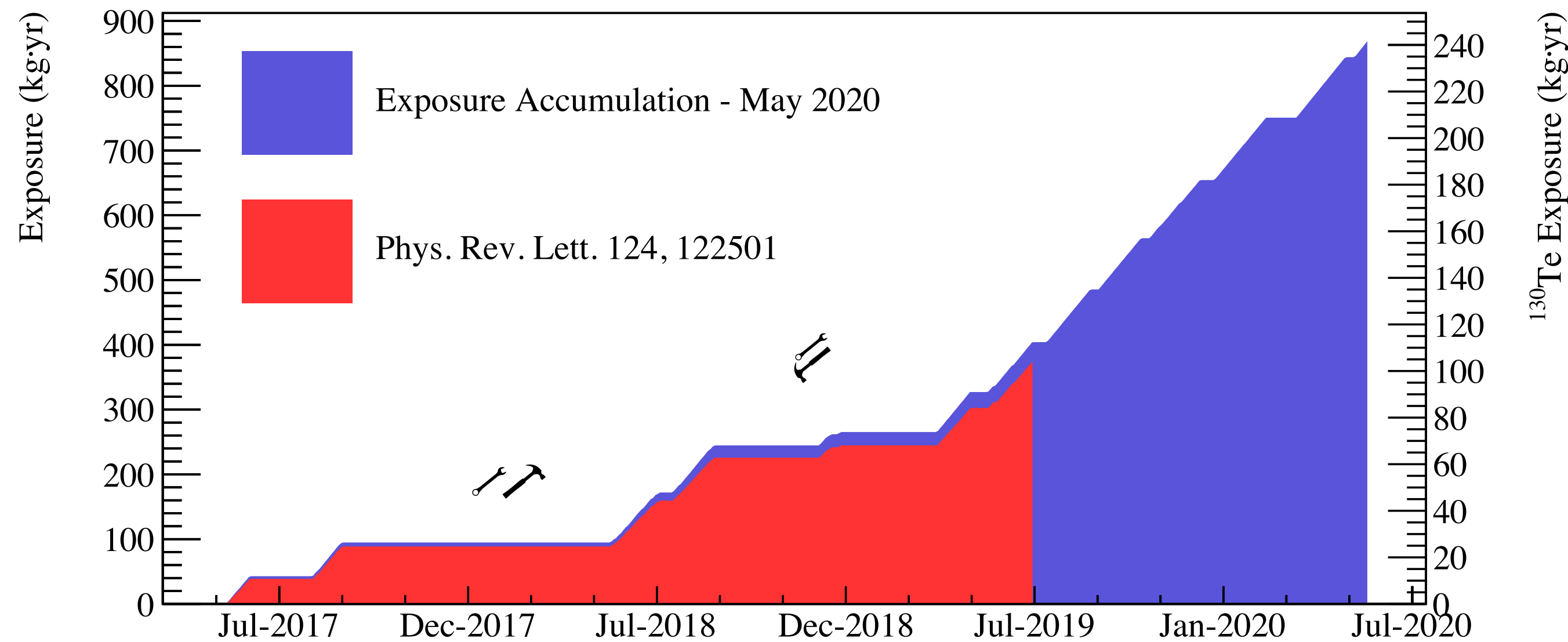


[A. D'Addabbo, Cryogenics 93, 56-65 (2018)]

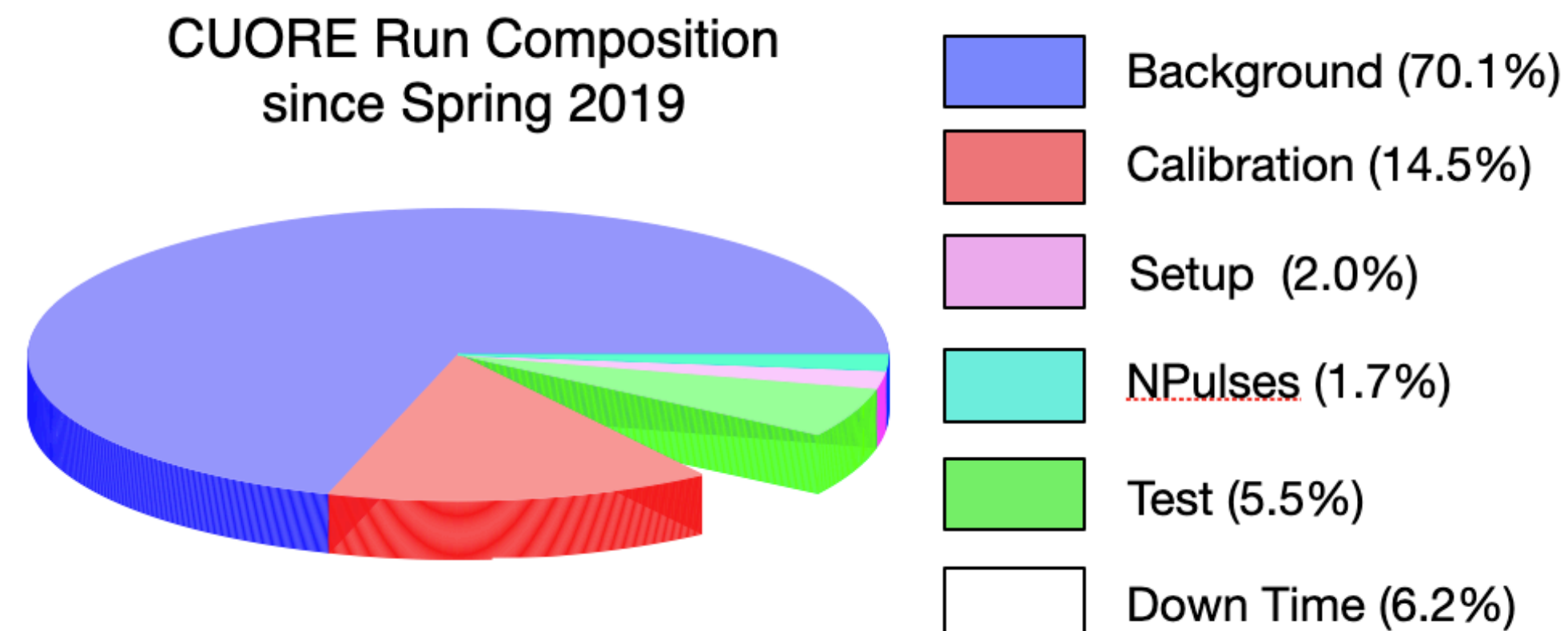


# The CUORE experiment

## Detector operation



[D. Q. Adams et al., Phys. Rev. Lett. 124, 122501 (2020)]



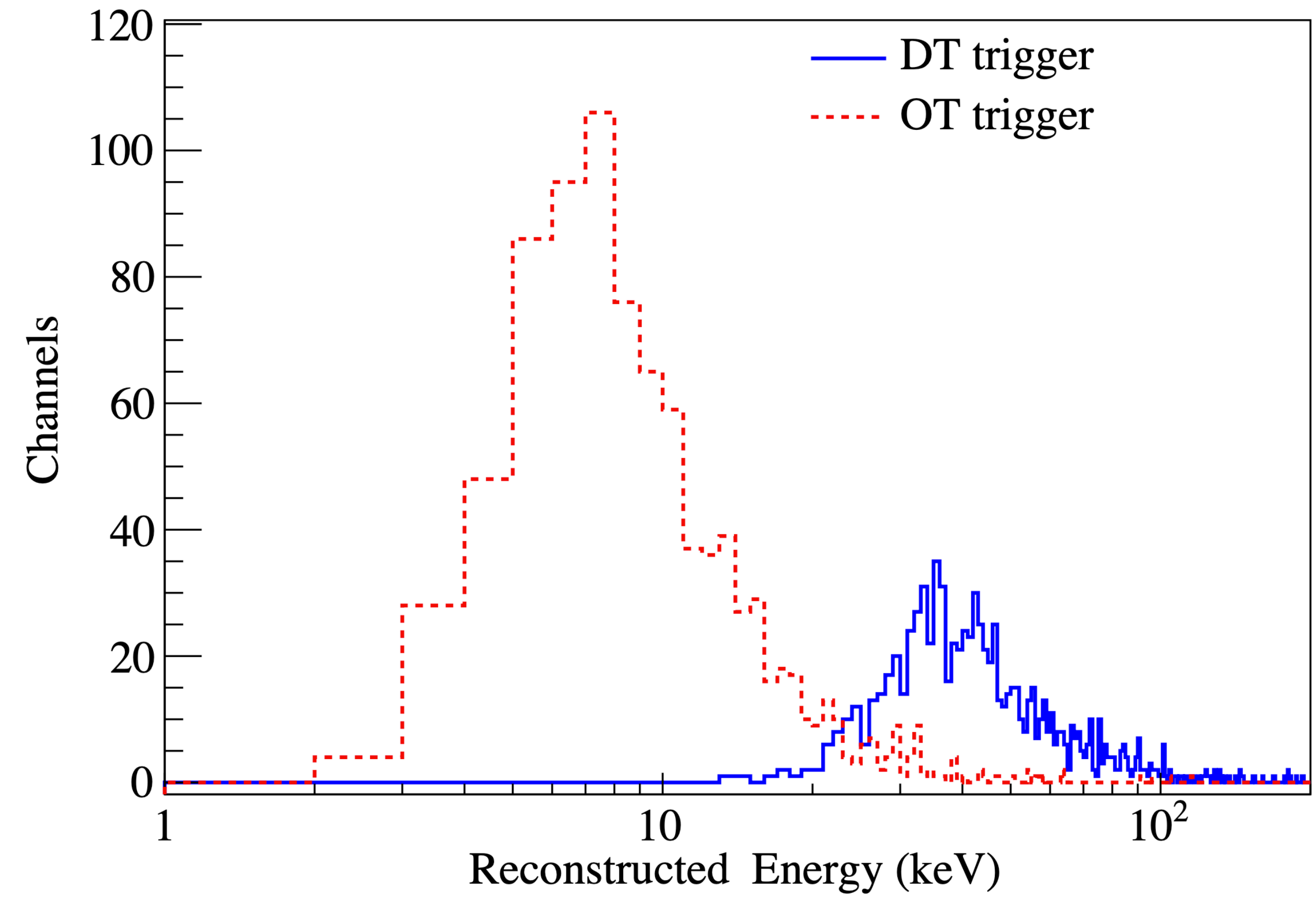
- Data taking started in Spring 2017
- CUORE “data set”: ~1 month of background data taking with a few days of calibration at the start and end
- Unblinded accumulated exposure  $372.5 \text{ kg} \cdot \text{yr}$
- Cryogenic facility for ton-scale bolometer arrays



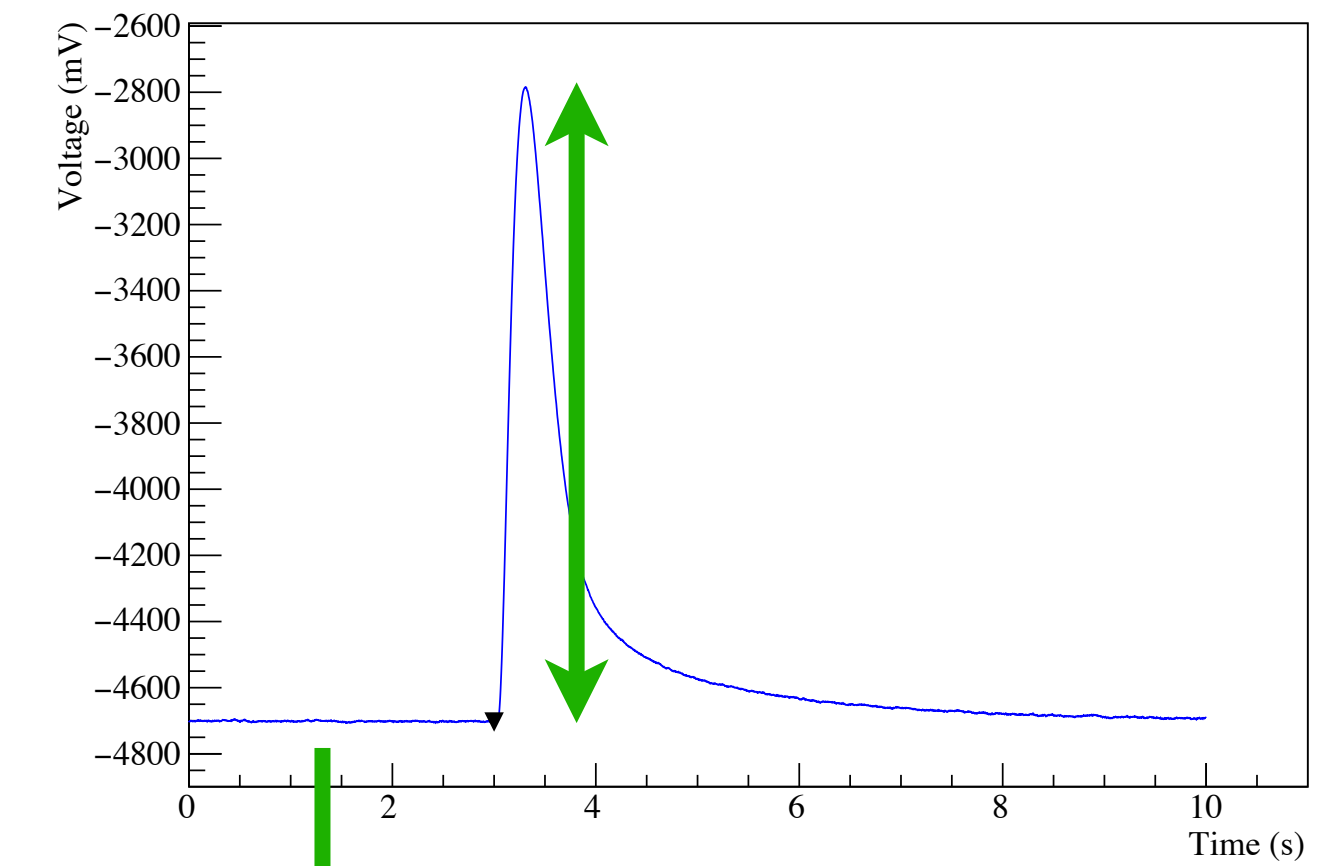
# The CUORE experiment

## Data processing

90% efficiency trigger thresholds



### PREPROCESS



#### Baseline basic parameters

- value (temperature)
- slope (pile-up)
- RMS (noise)

#### Waveform basic parameters

- Amplitude (max - min)
- Number of pulses (pile-up)

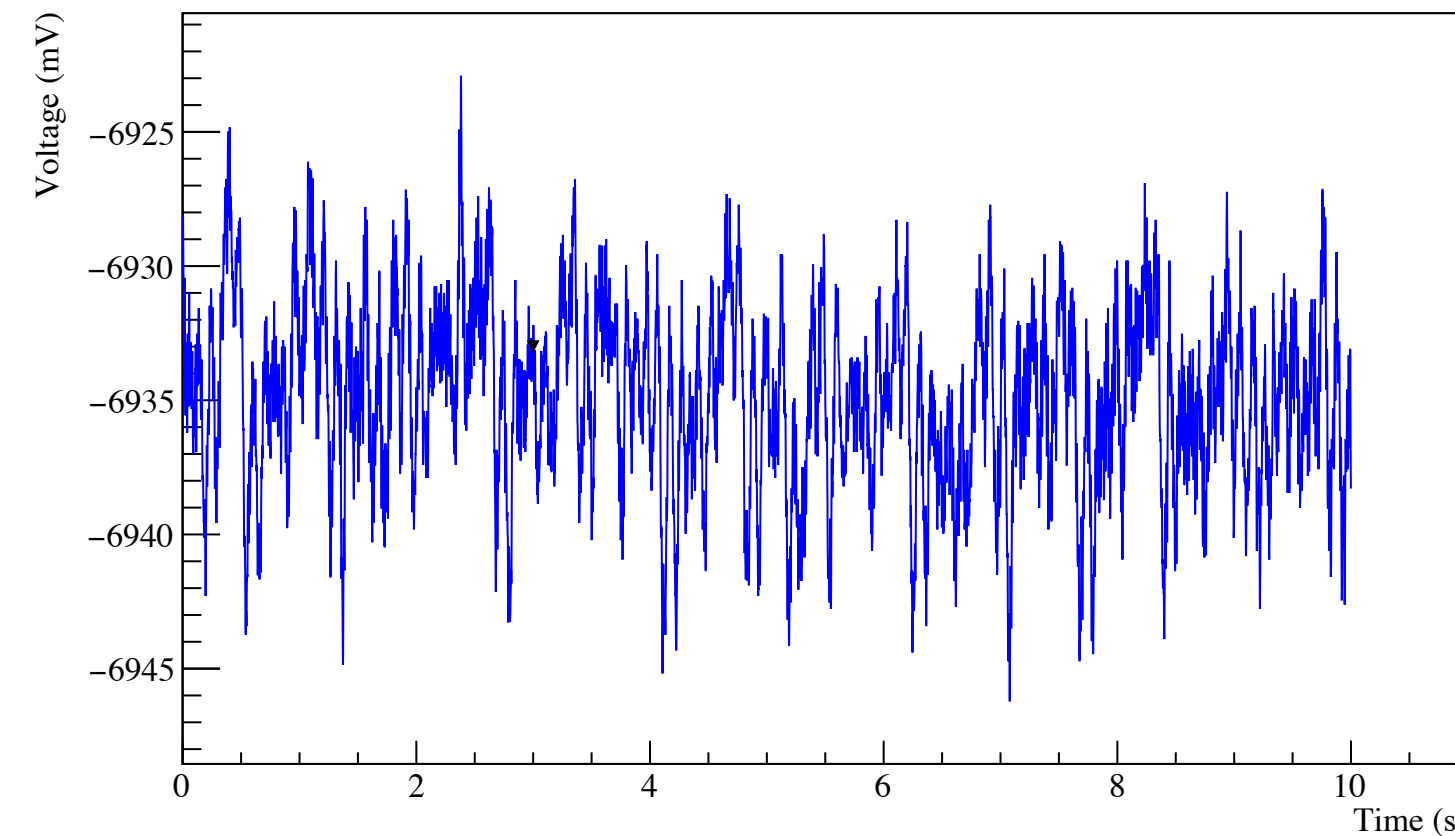
# The CUORE experiment

## Data processing

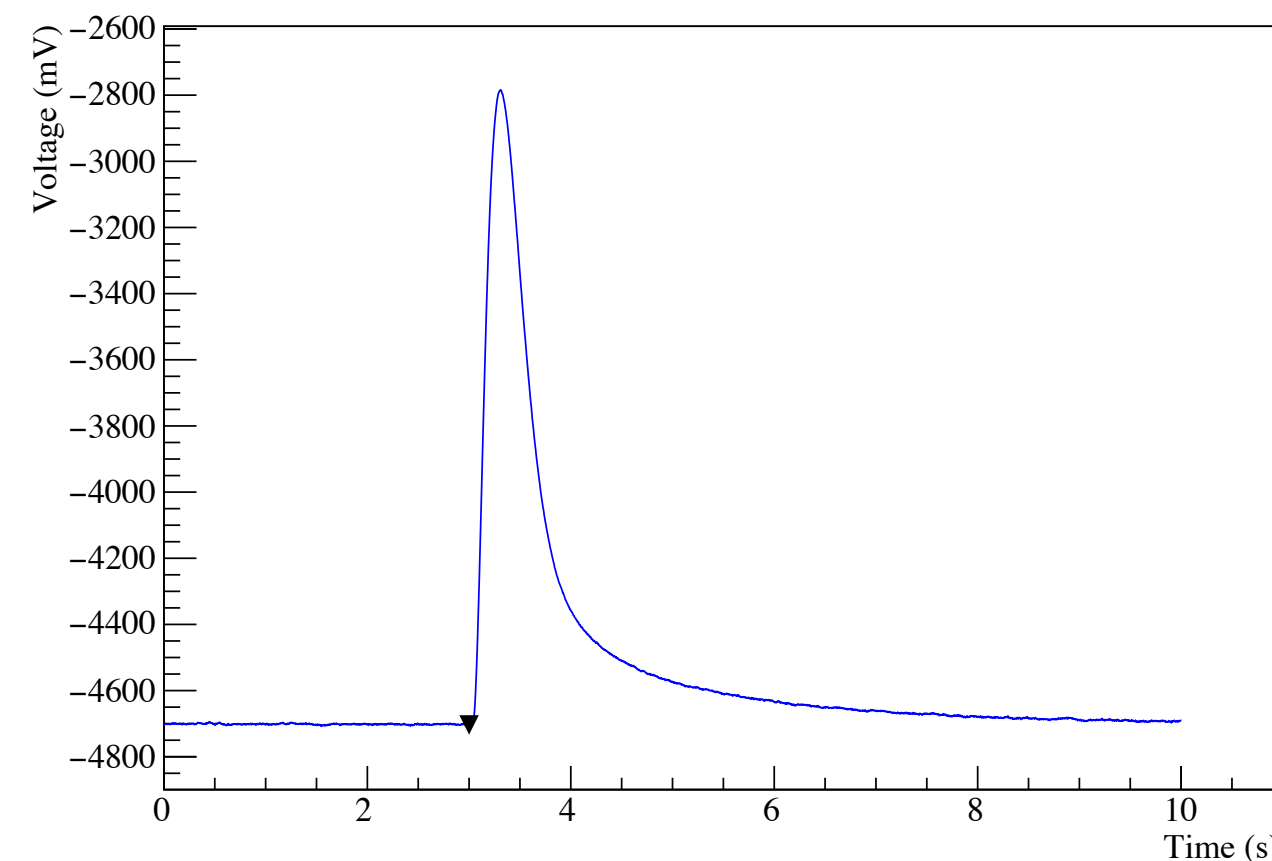
- A sample of clean noise waveforms is used to extract each channel's average noise
- A sample of clean particle pulses is used to compute the average response of each detector

PREPROCESS

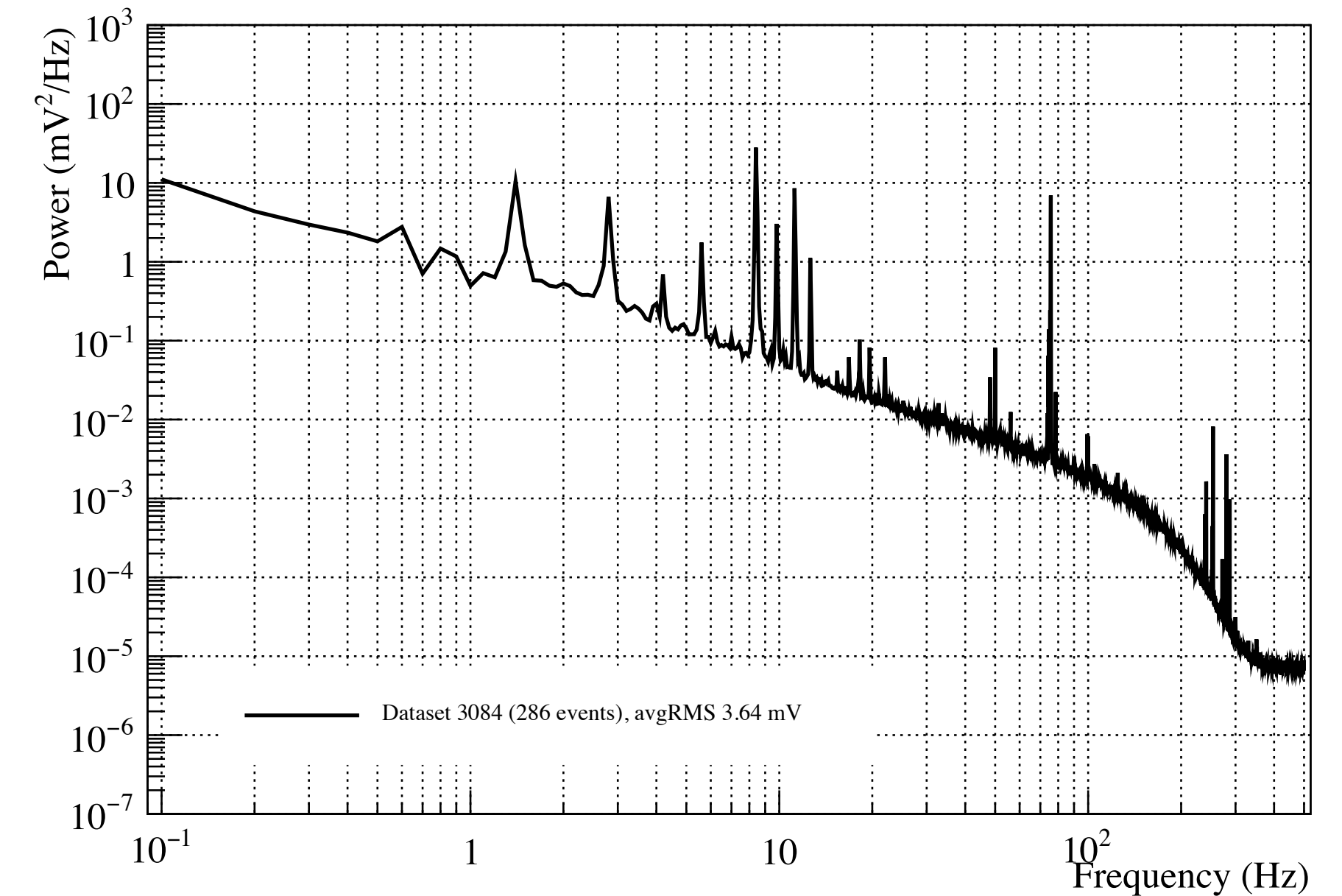
AVERAGE NOISE



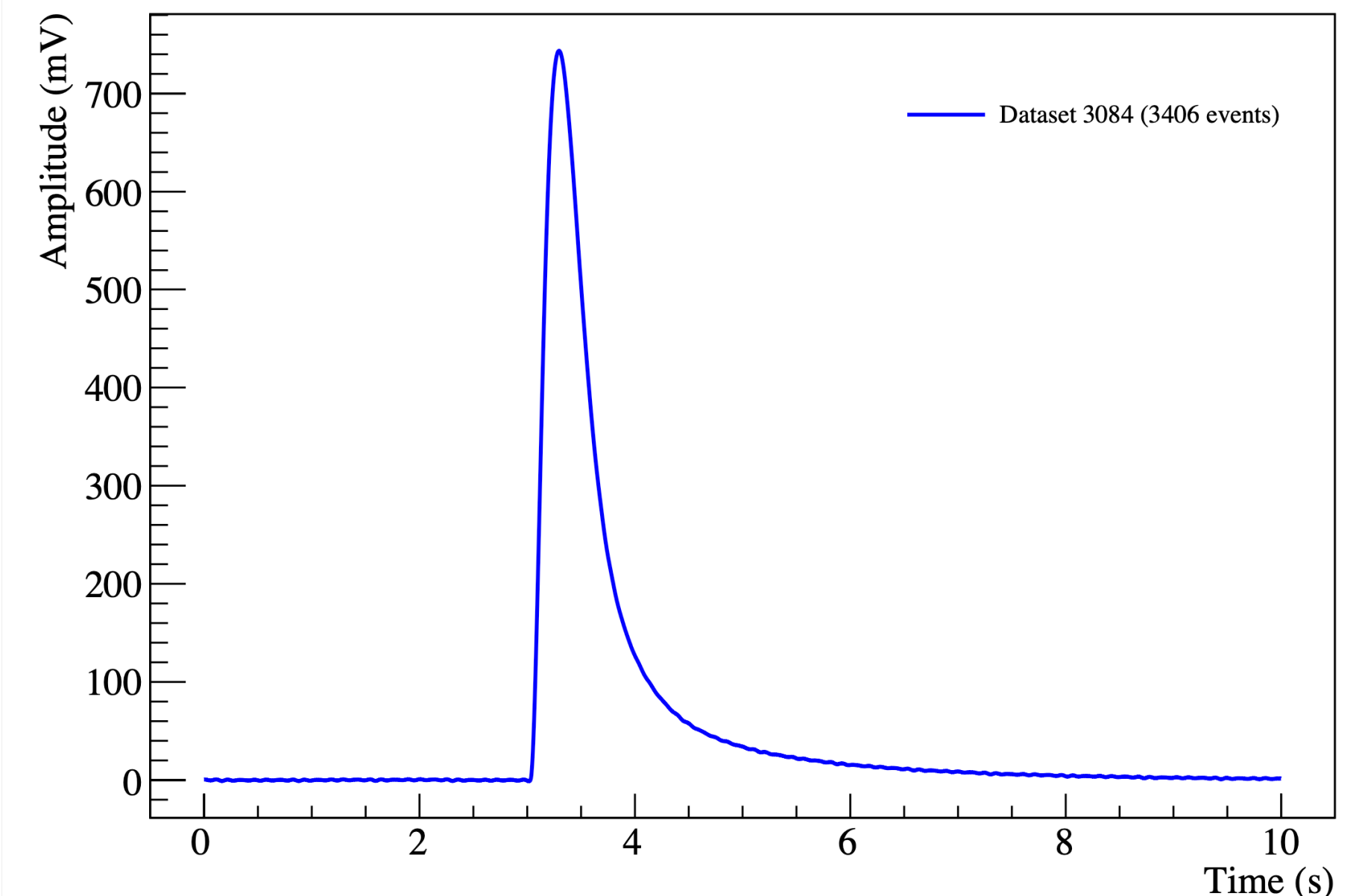
AVERAGE PULSE



Average Noise Power Spectrum: ch. 542 - dss 3084



Average Pulse: ch. 542 - ds 3084

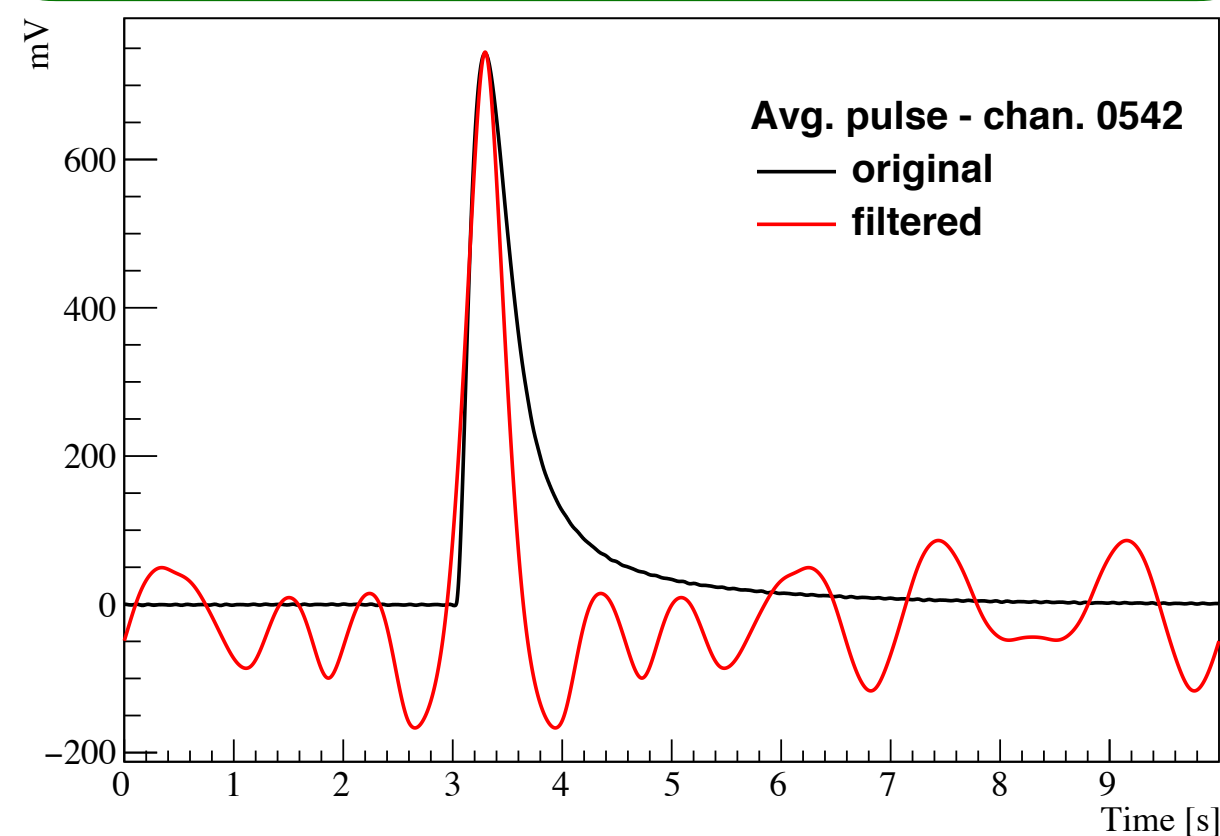




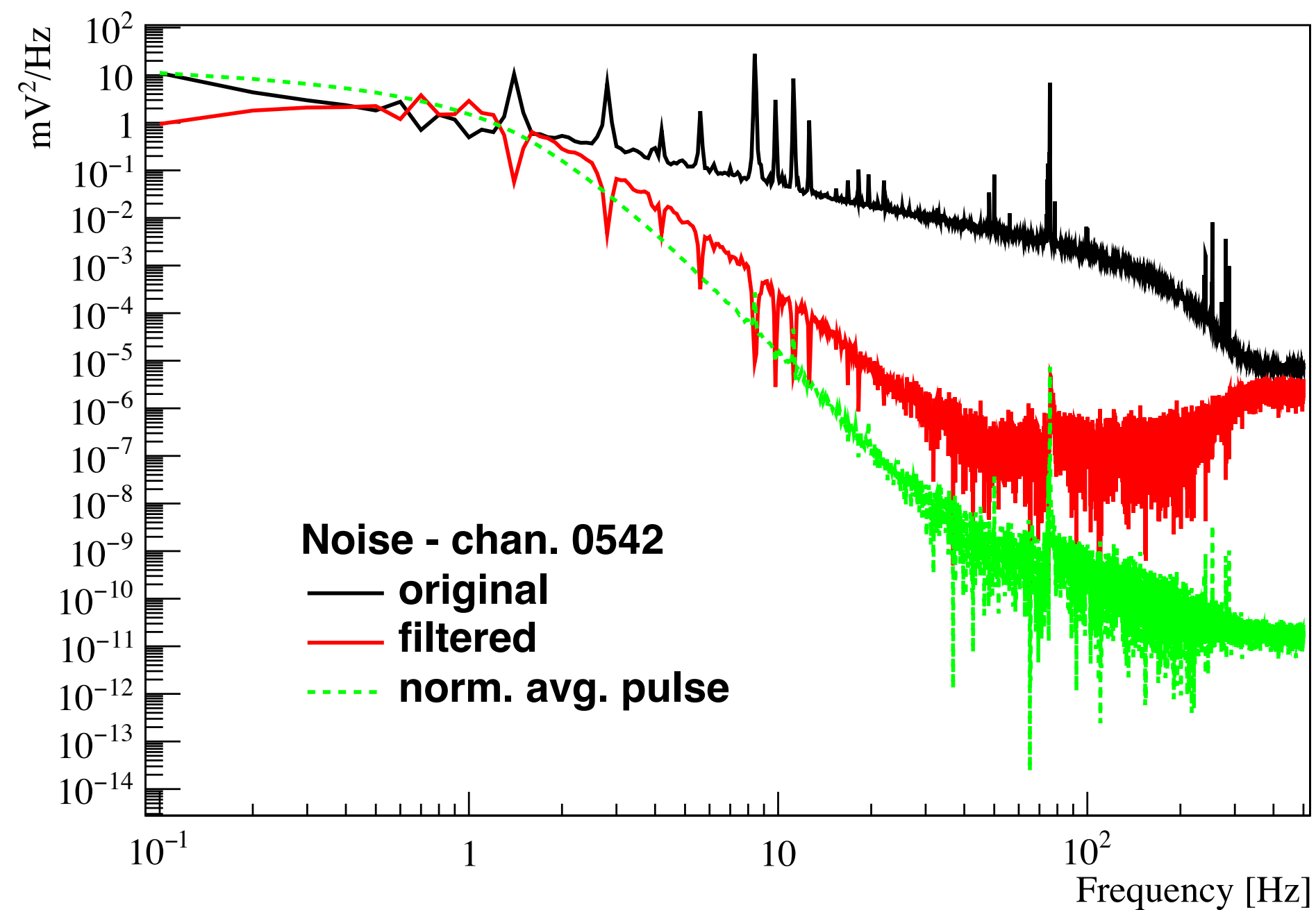
# The CUORE experiment

## Data processing

- PREPROCESS
- AVERAGE PULSE
- AVERAGE NOISE
- AMPLITUDE**



- STABILIZATION
- CALIBRATION
- BLINDING
- COINCIDENCES
- PULSE SHAPE ANALYSIS**

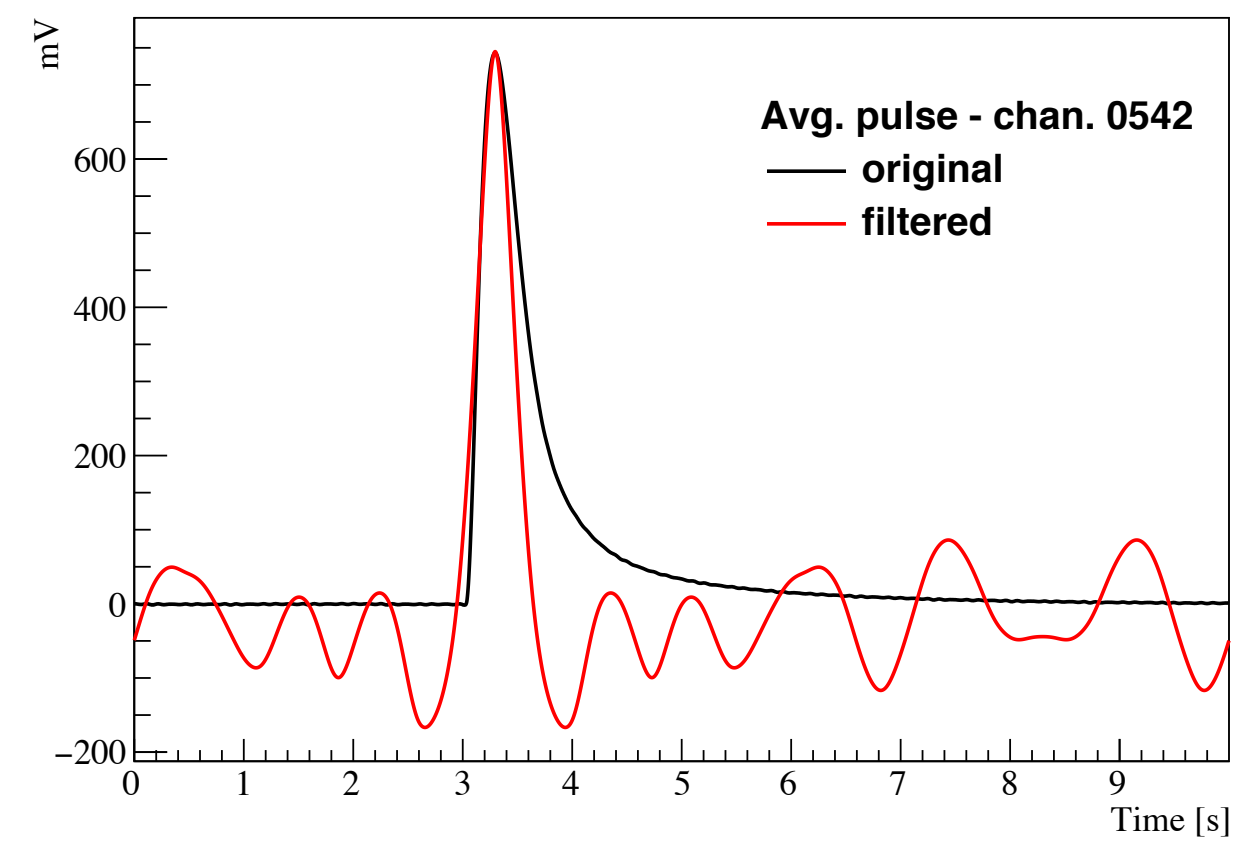


- Stationary noise assumption
- Energy independent pulse shape
- Amplitude extracted with matched filter

# The CUORE experiment

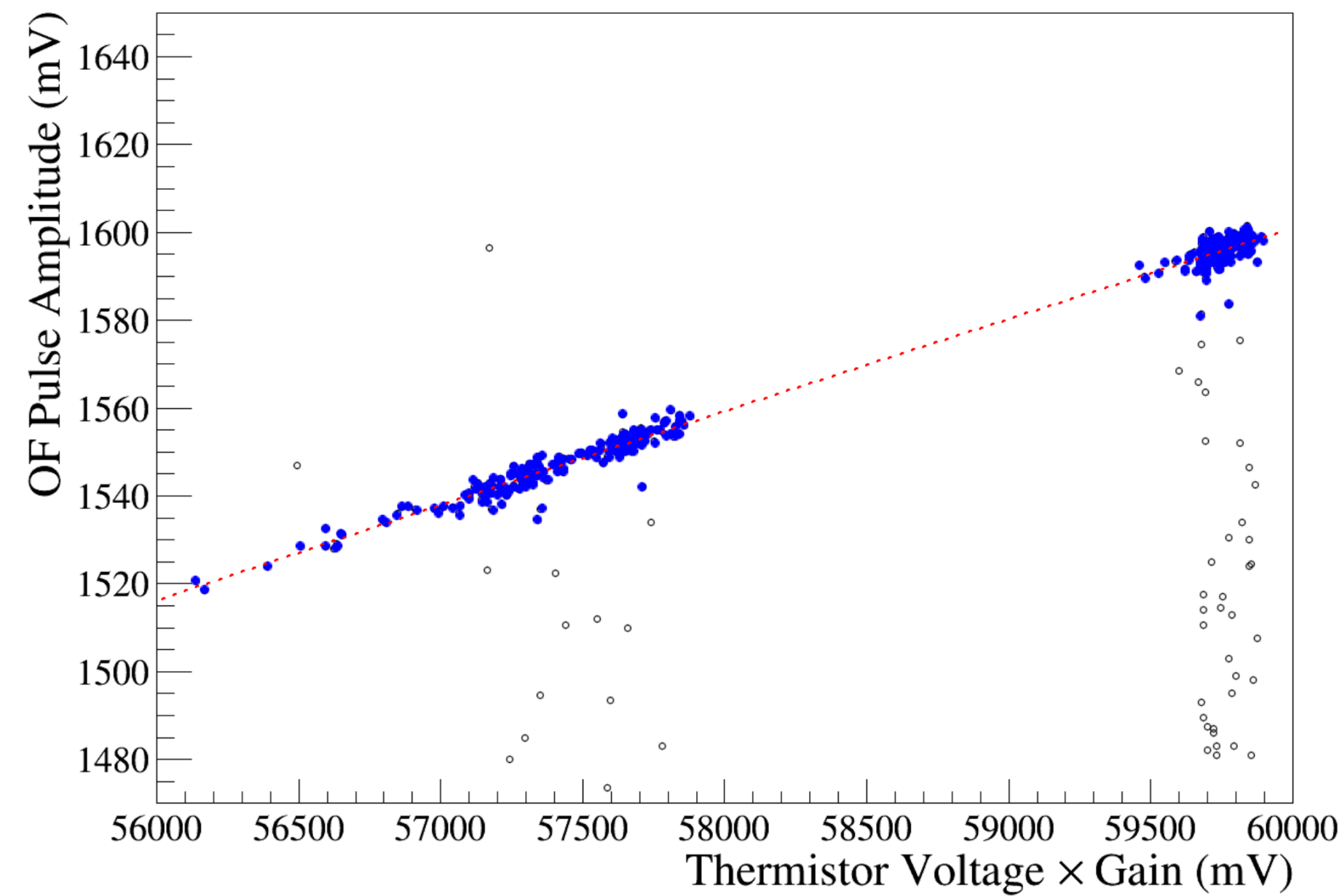
## Data processing

- PREPROCESS
- AVERAGE PULSE
- AVERAGE NOISE
- AMPLITUDE
- STABILIZATION**



- CALIBRATION
- BLINDING
- COINCIDENCES
- PULSE SHAPE ANALYSIS

[C. Alduino et al., Phys. Rev. C 93, 045503 (2016)]



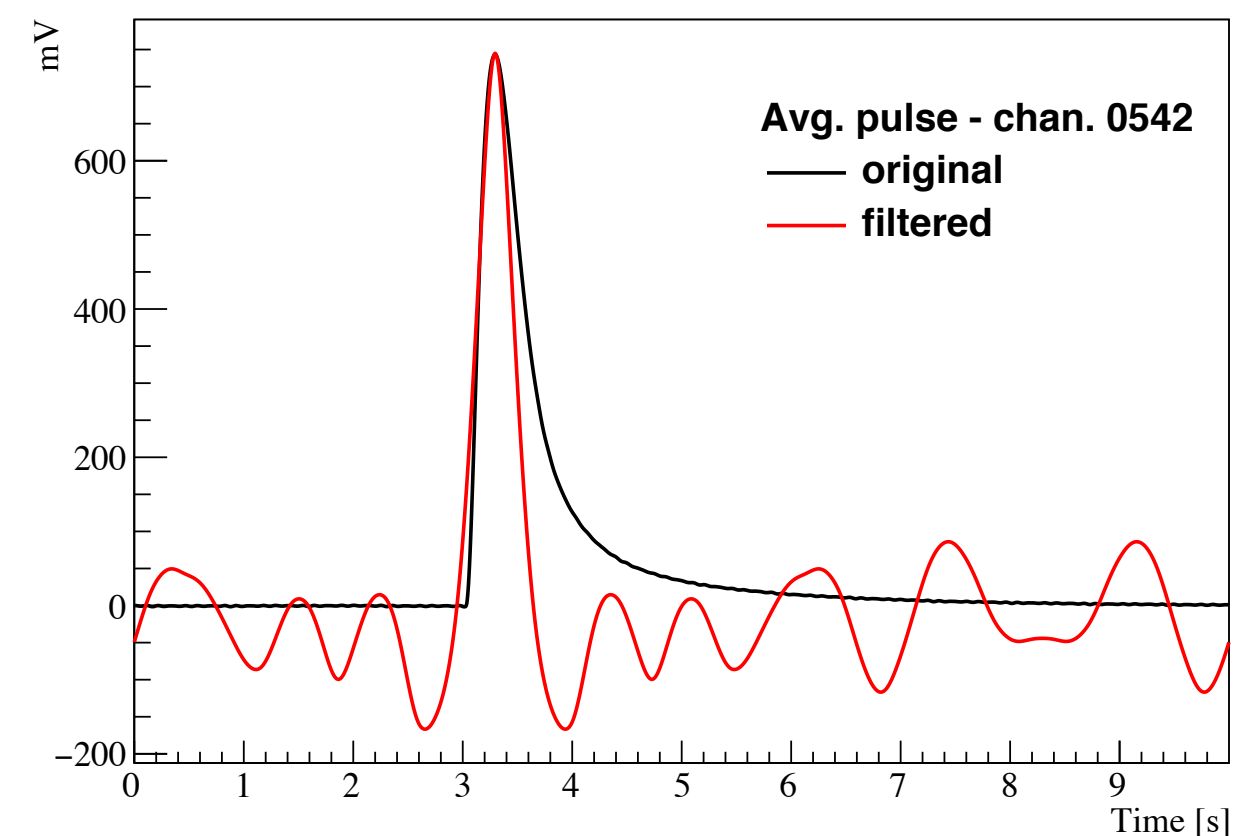
- Fixed amounts of energy periodically injected via heater pulses
- High-statistics 2615 keV line from calibration data
- Temperature fluctuations produce small variations in the heat capacity of the CUORE bolometers (thermal gain)
- A linear correction is applied



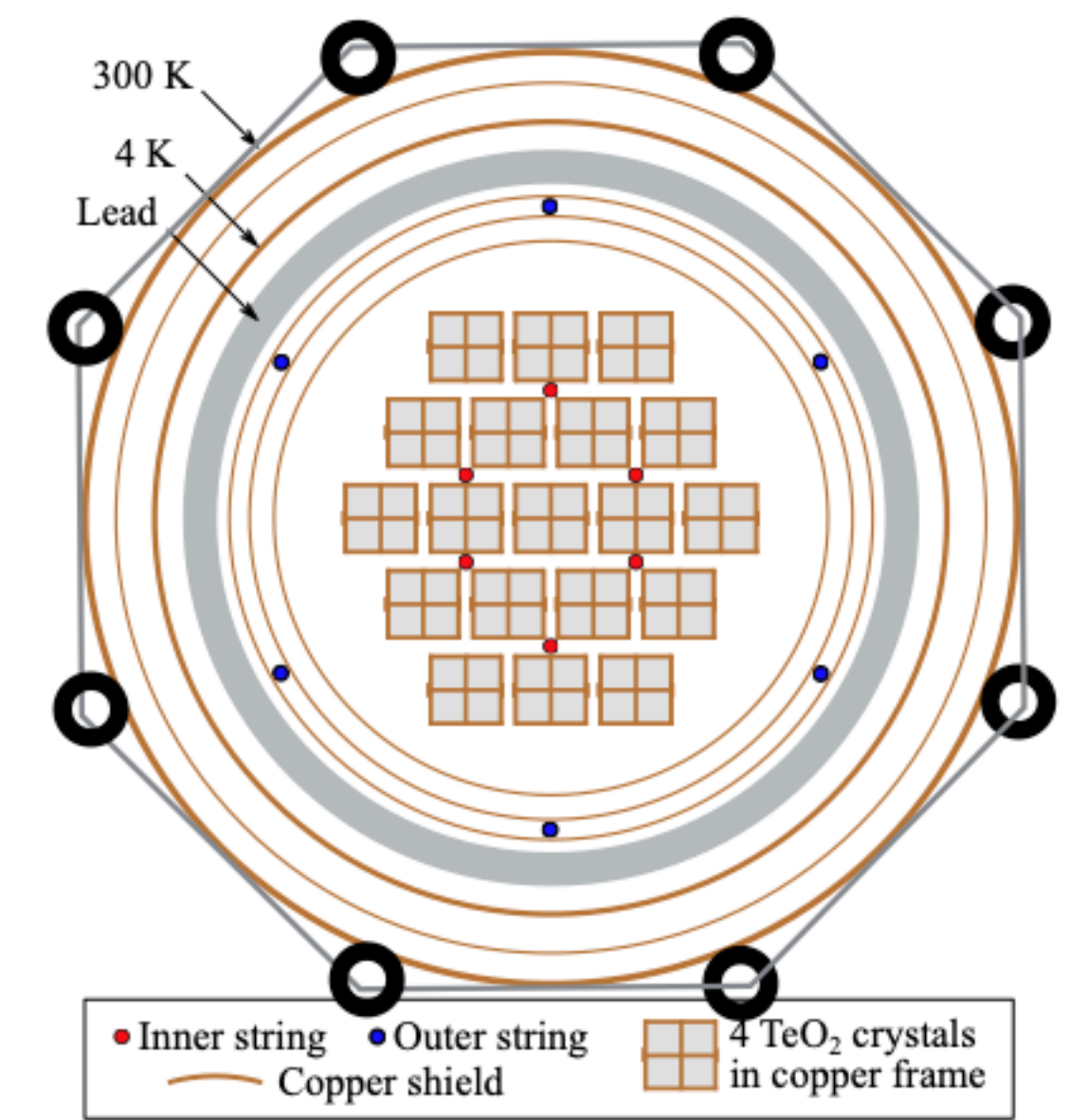
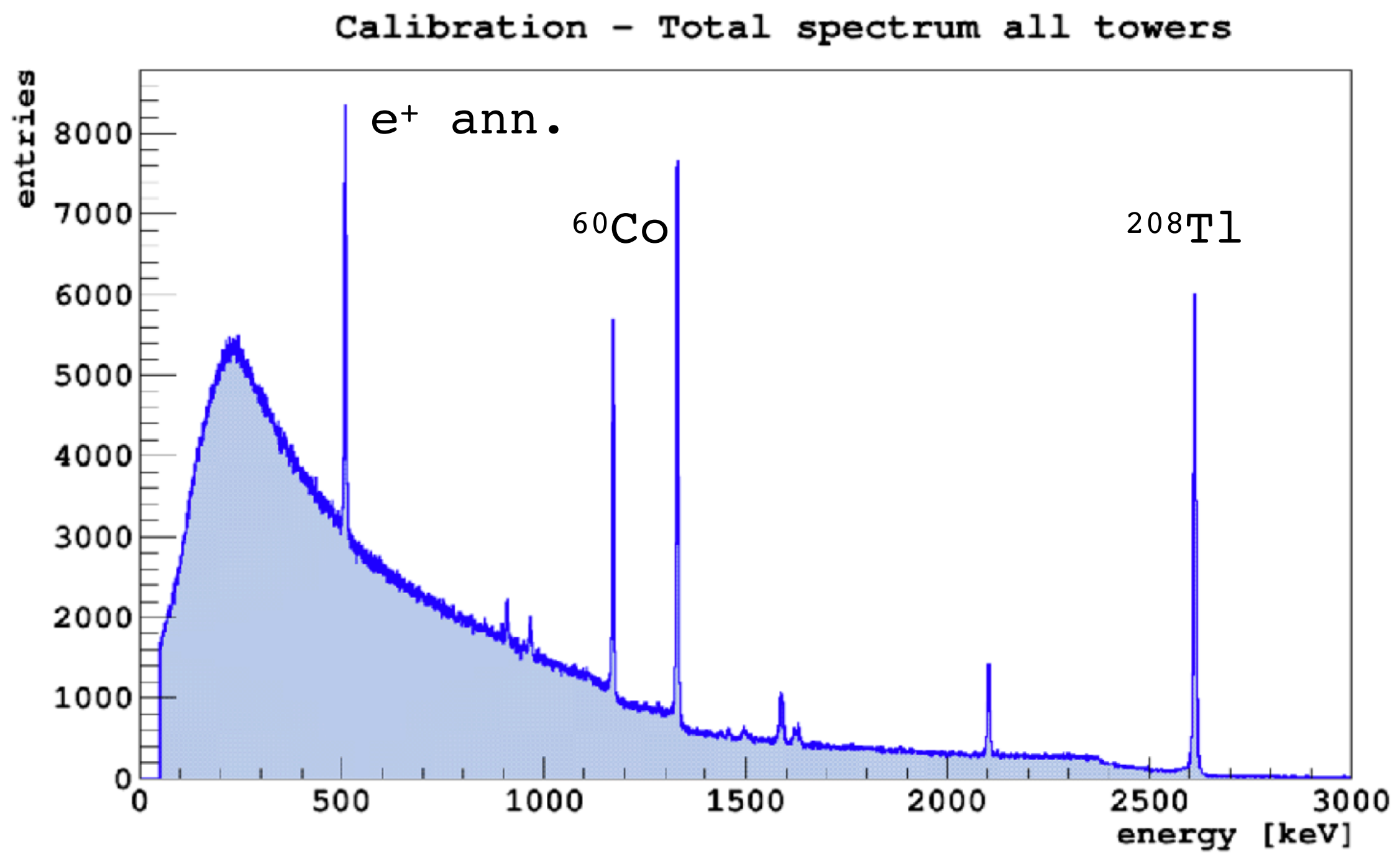
# The CUORE experiment

## Data processing

- PREPROCESS
- AVERAGE PULSE
- AVERAGE NOISE
- AMPLITUDE
- STABILIZATION
- CALIBRATION**



- BLINDING
- COINCIDENCES
- PULSE SHAPE ANALYSIS

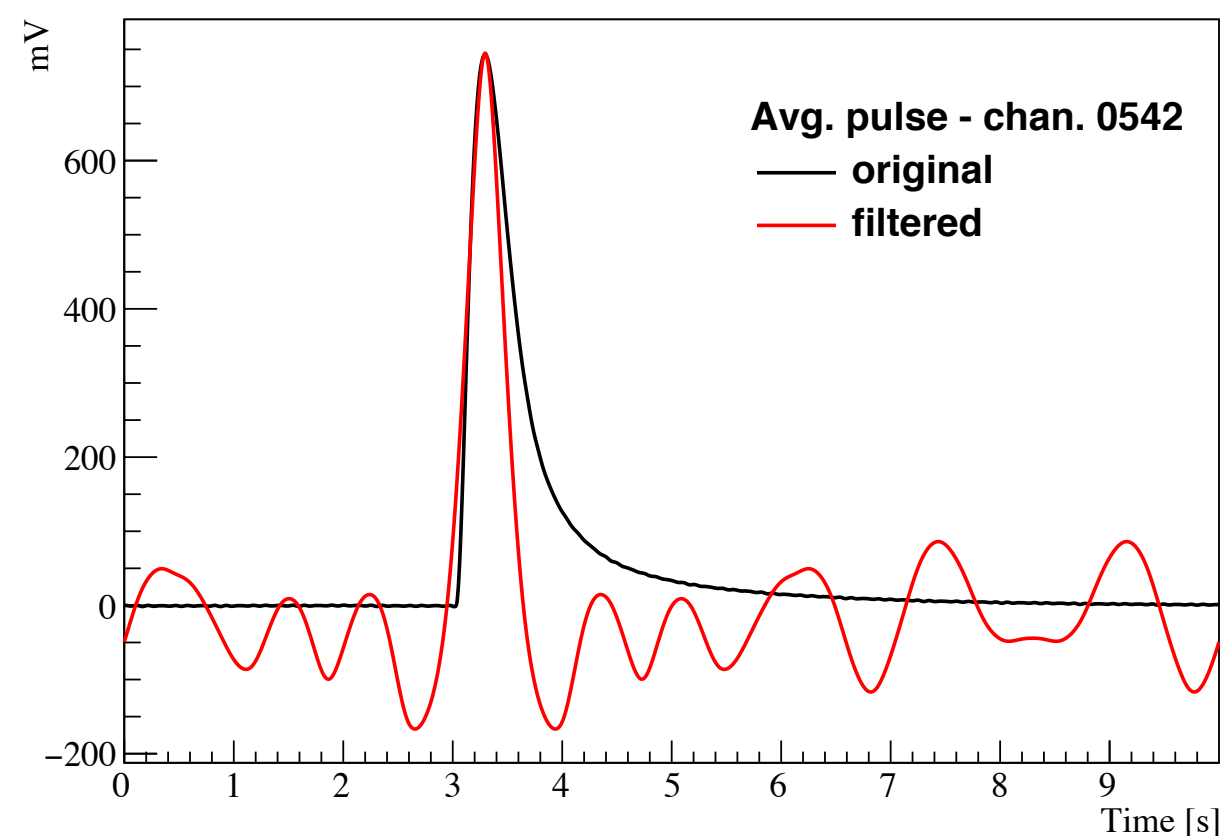


- 8 room temperature calibration strings
- $^{232}\text{Th} + ^{60}\text{Co}$  sources
- 2<sup>nd</sup> order polynomial calibration function

# The CUORE experiment

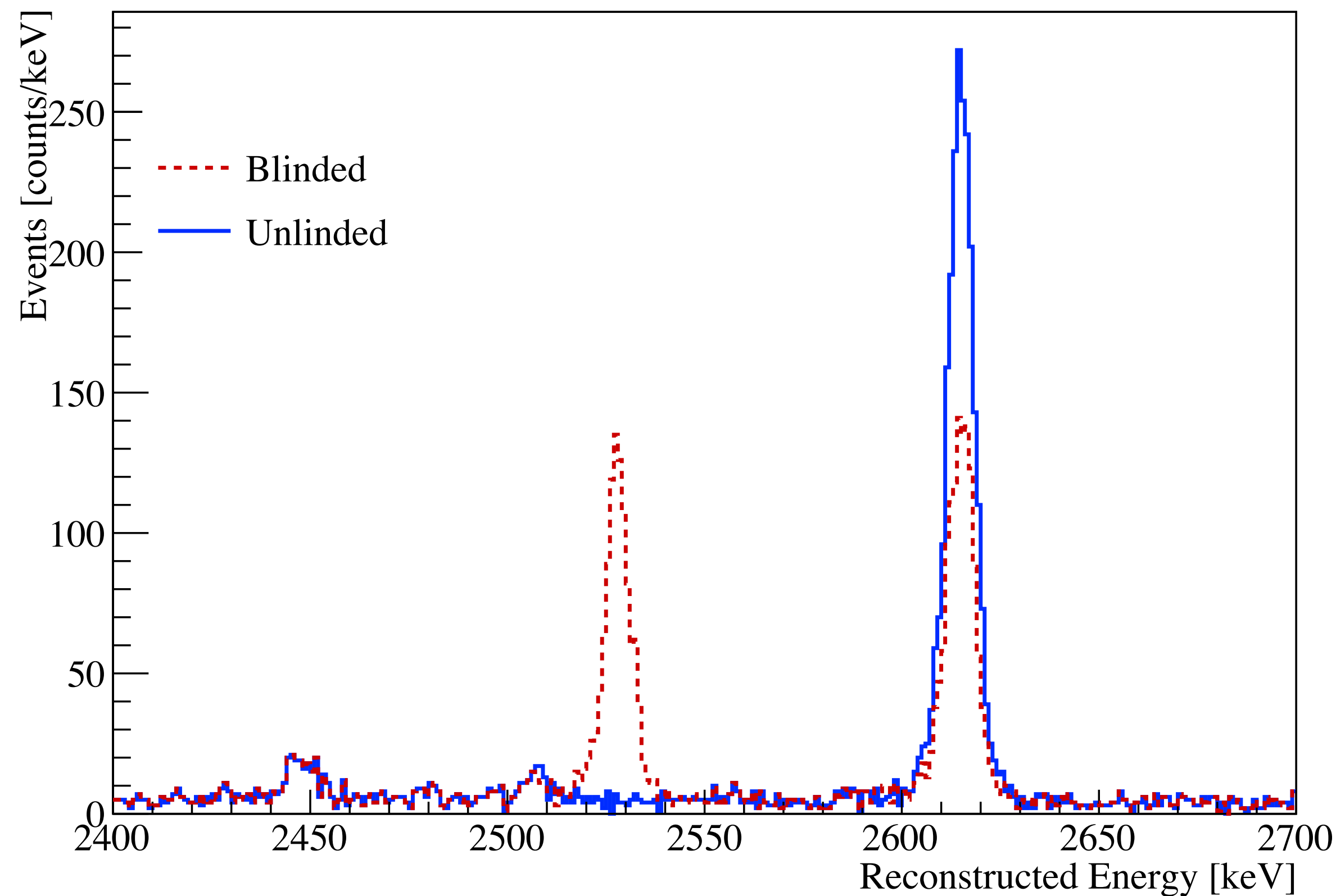
## Data processing

- PREPROCESS
- AVERAGE PULSE
- AVERAGE NOISE
- AMPLITUDE
- STABILIZATION
- CALIBRATION
- BLINDING**



- COINCIDENCES
- PULSE SHAPE ANALYSIS

Summed Spectrum



Given the seed, a random fraction of events is shifted from the 2615 keV  $^{208}\text{Tl}$  peak to  $Q_{\beta\beta}$

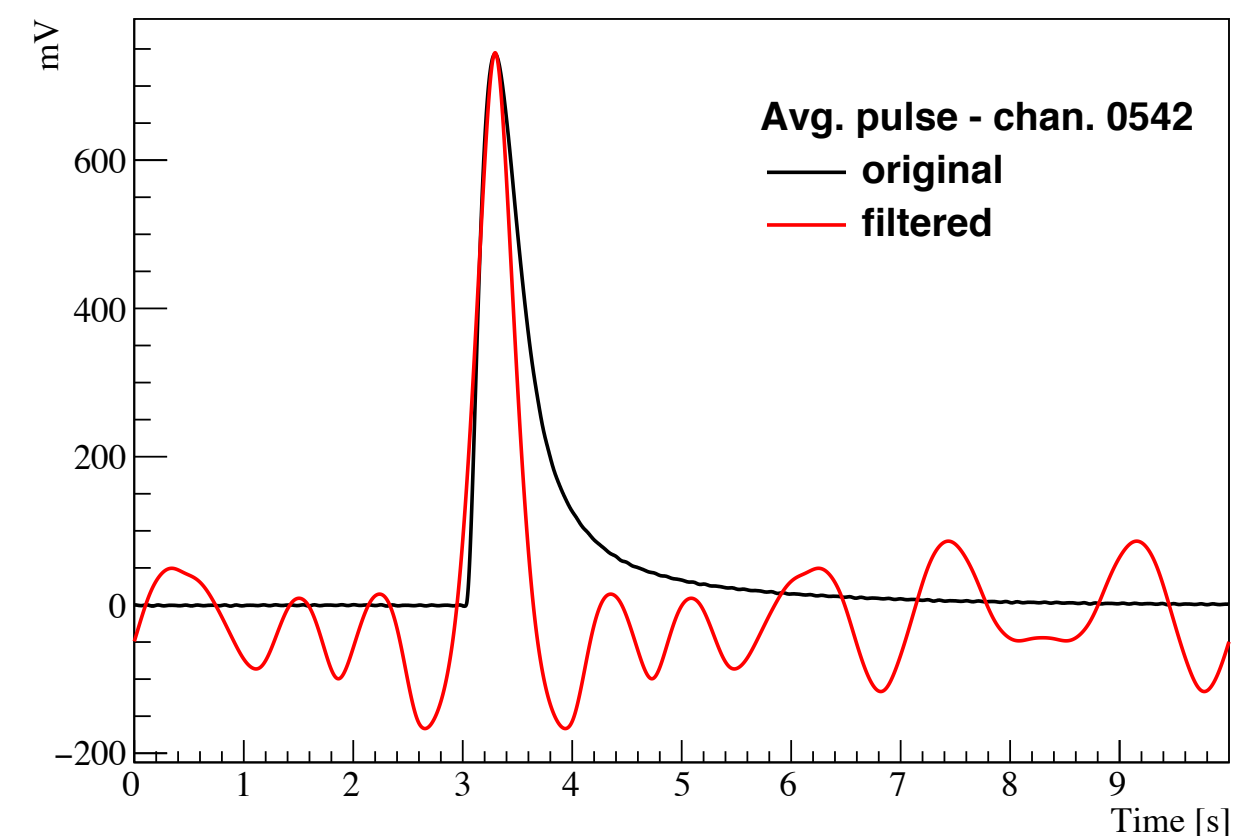
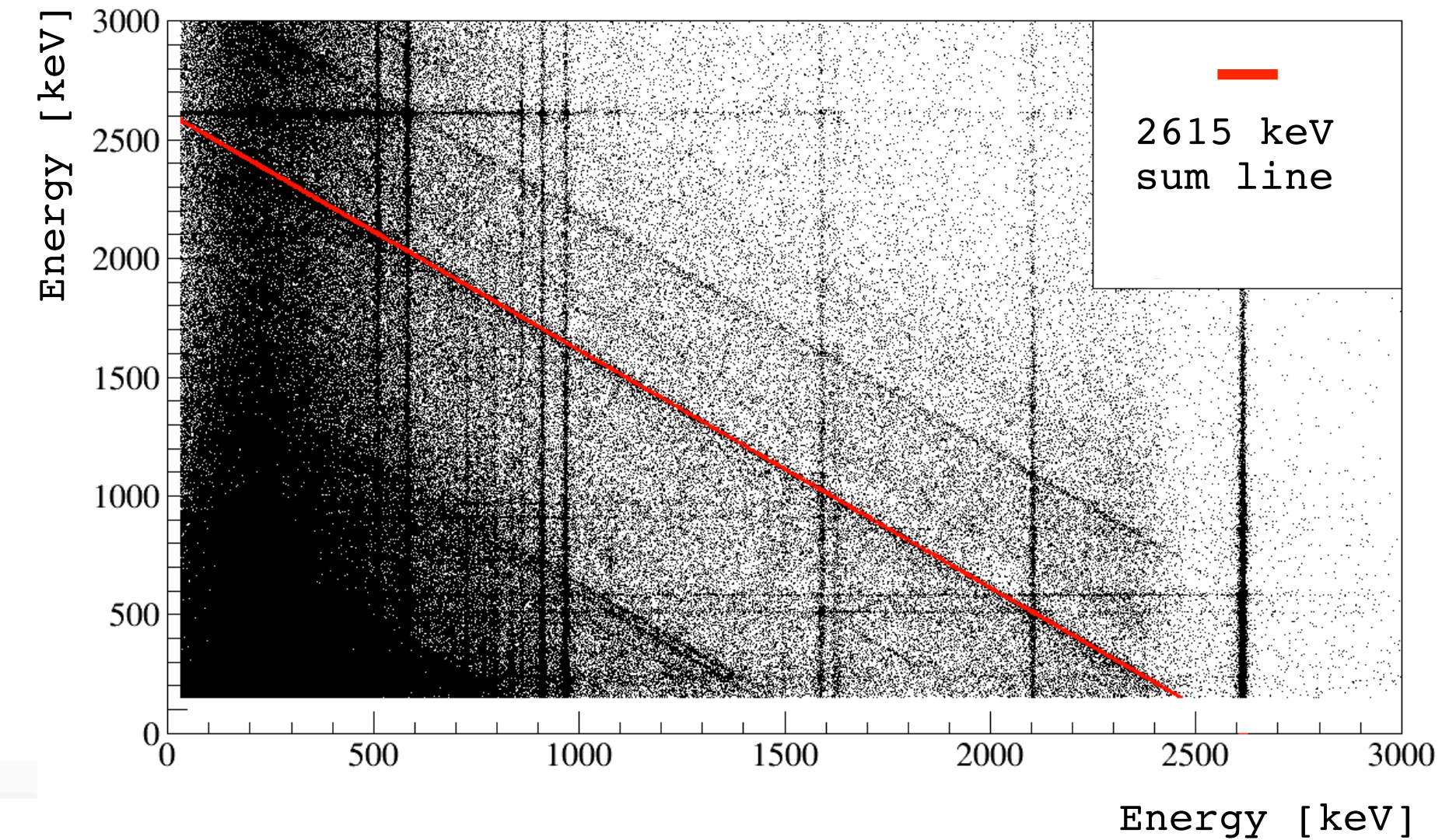
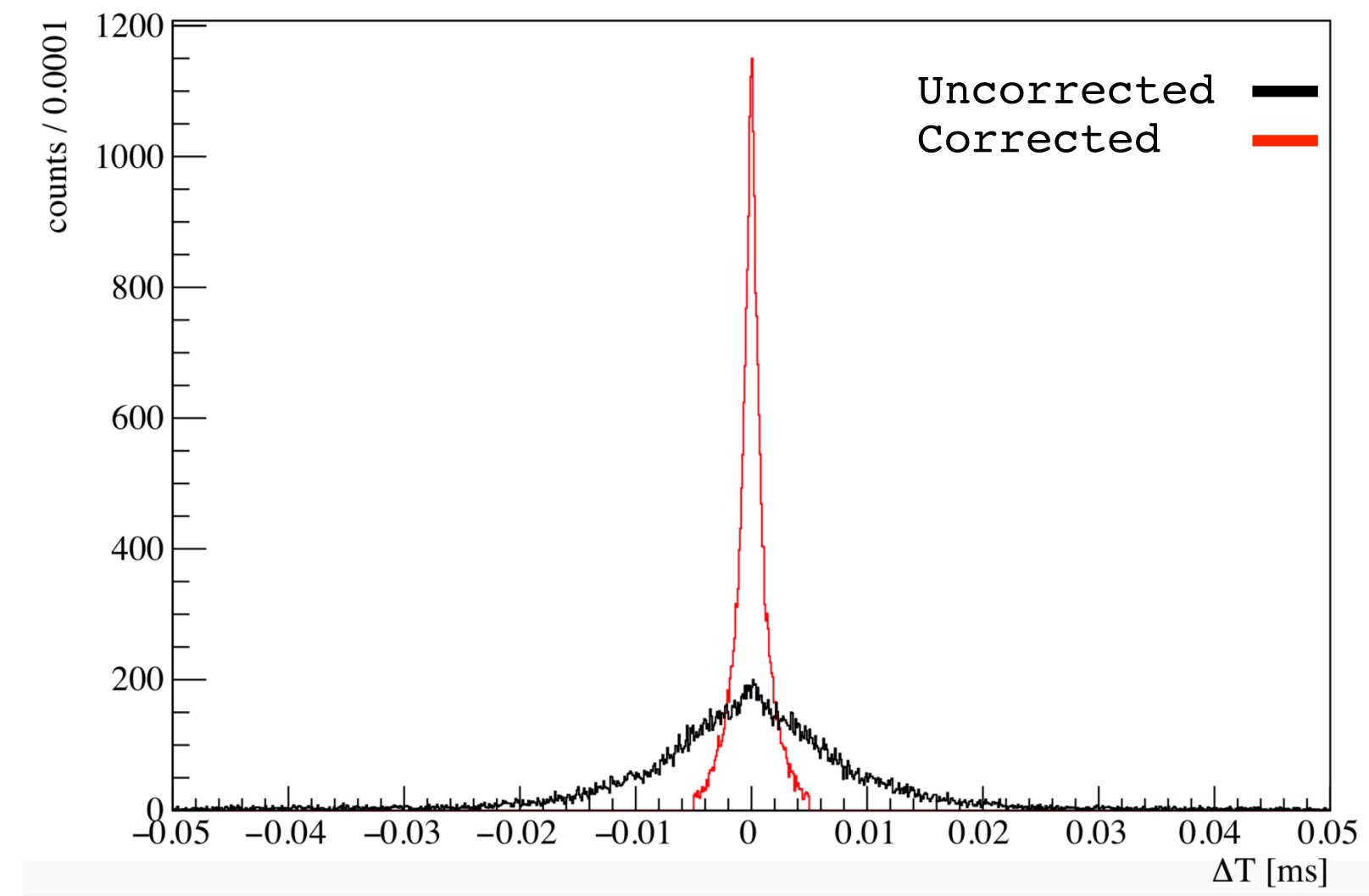


# The CUORE experiment

## Data processing

- PREPROCESS
- AVERAGE PULSE
- AVERAGE NOISE
- AMPLITUDE
- STABILIZATION
- CALIBRATION
- BLINDING
- COINCIDENCES**

Time difference of 2615 keV double-site events



## PULSE SHAPE ANALYSIS

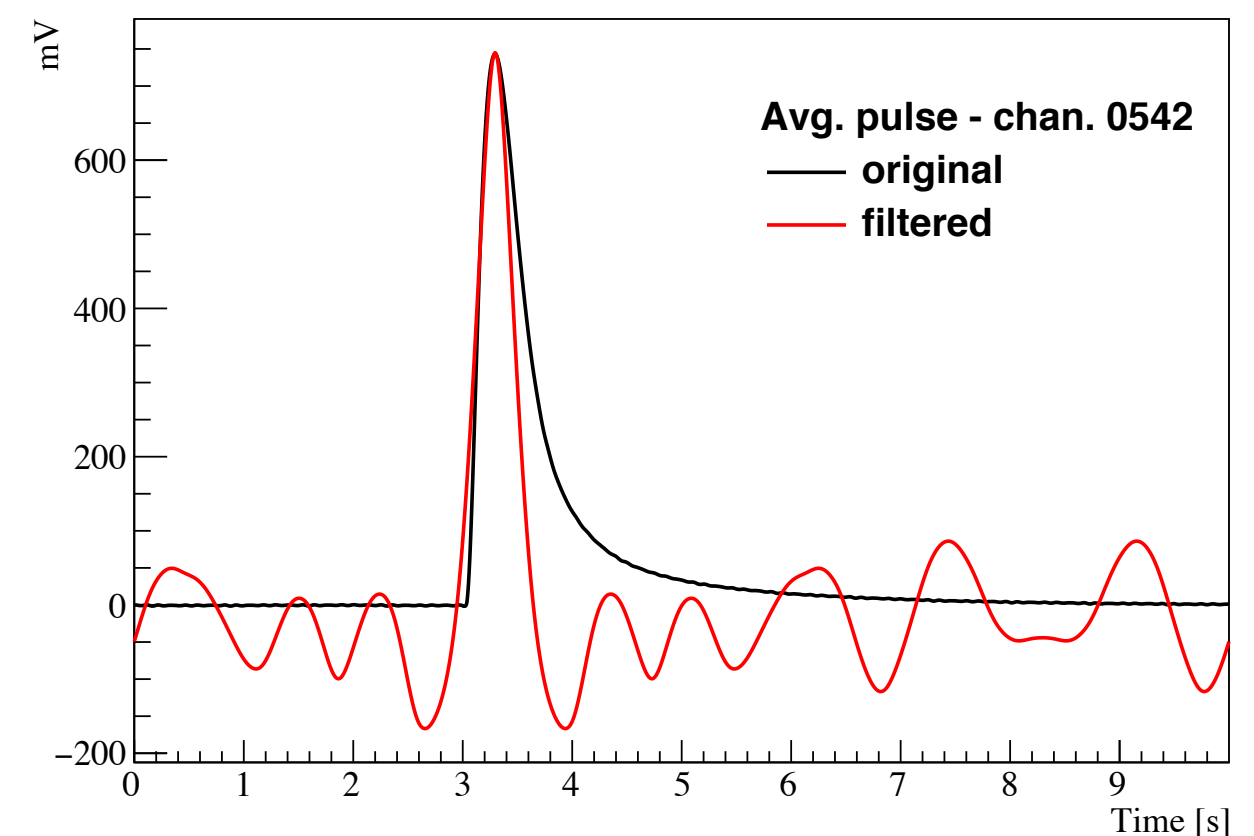
The event time reconstruction is affected by differences in the pulse rise time among different channels.

A by-channel correction allows to select coincident events in a time window as narrow as  $\pm 5$  ms

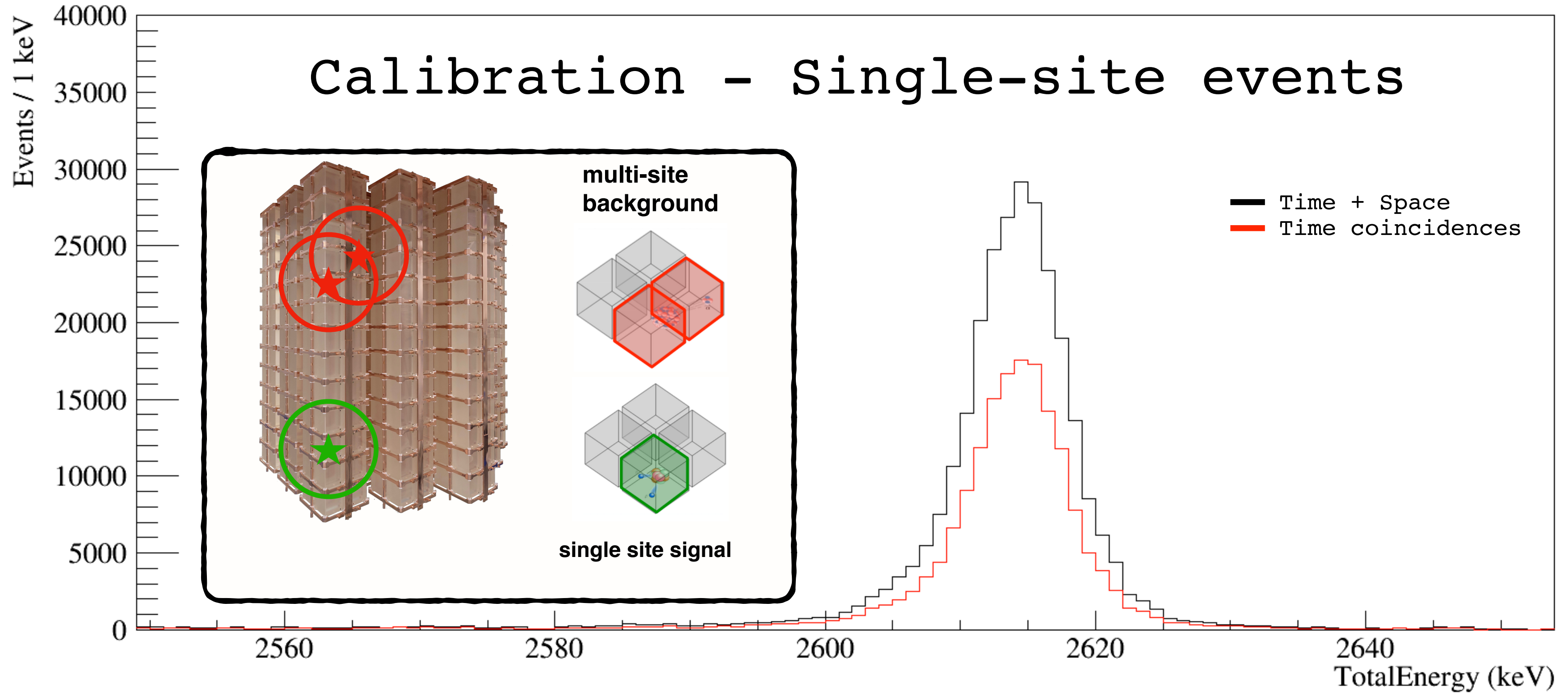
# The CUORE experiment

## Data processing

- PREPROCESS
- AVERAGE PULSE
- AVERAGE NOISE
- AMPLITUDE
- STABILIZATION
- CALIBRATION
- BLINDING
- COINCIDENCES**



## PULSE SHAPE ANALYSIS



A space-based definition of coincident events improves signal purity in coincident multiplets

Accidental coincidences strongly suppressed



# The CUORE experiment

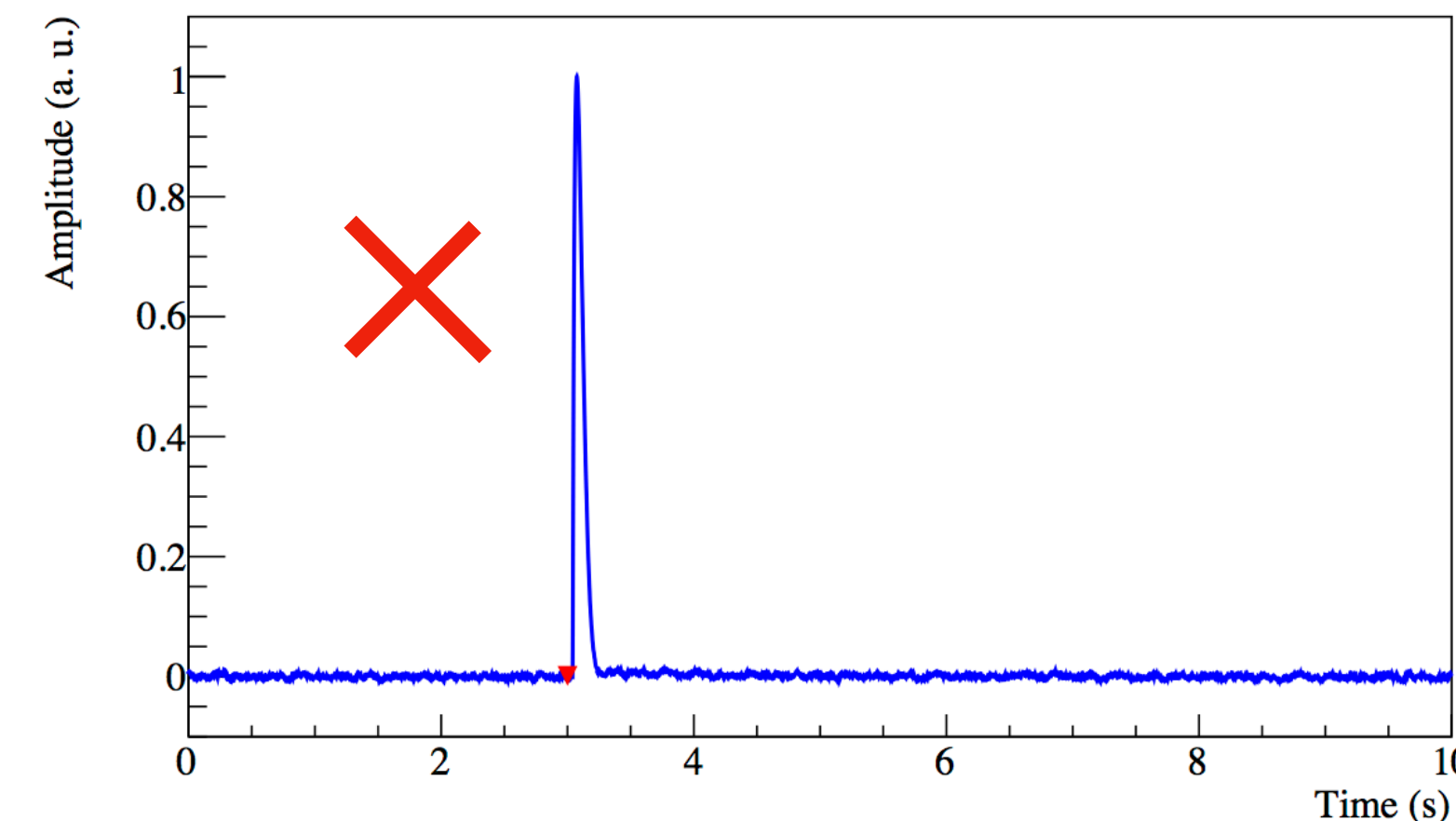
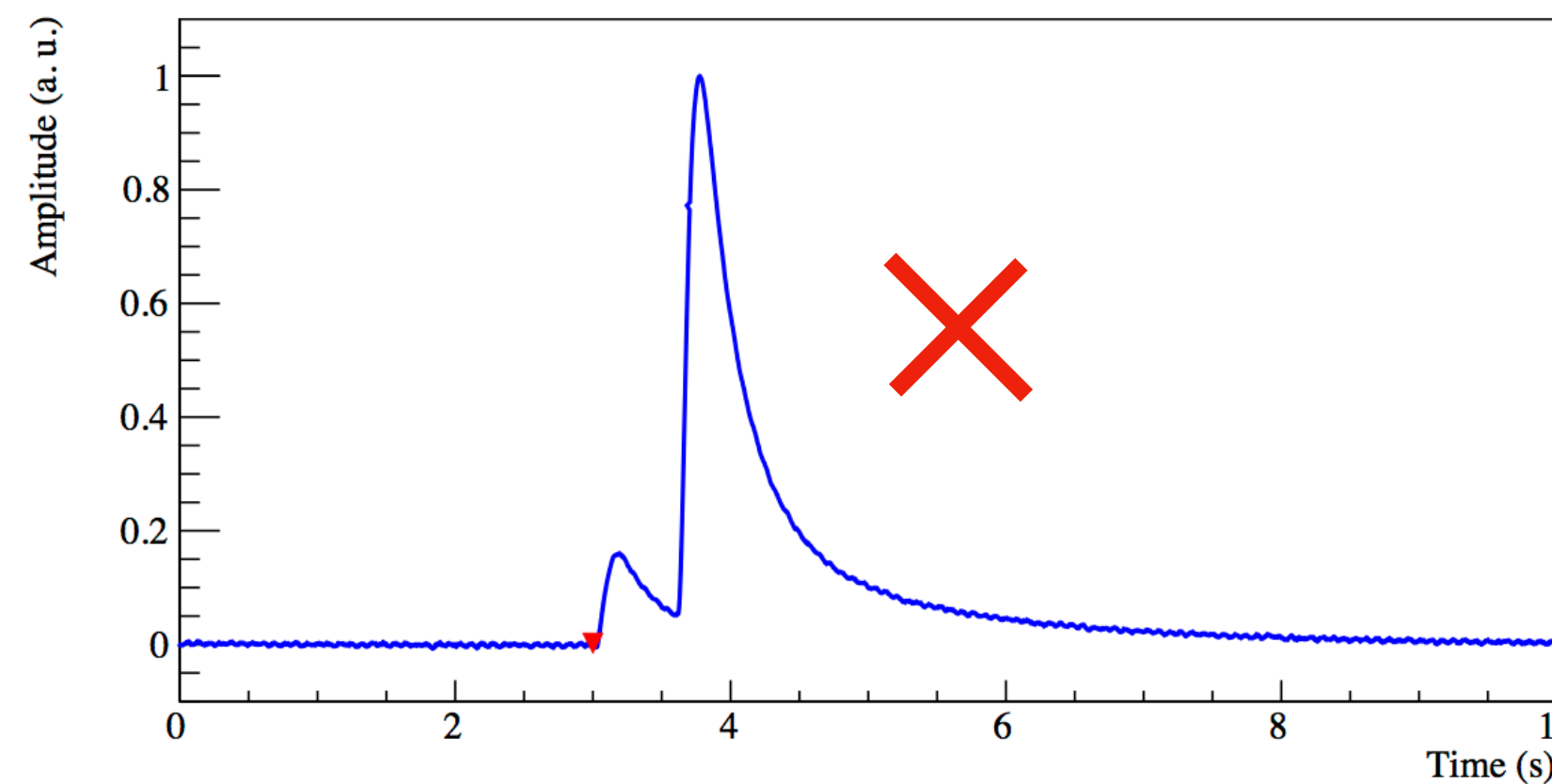
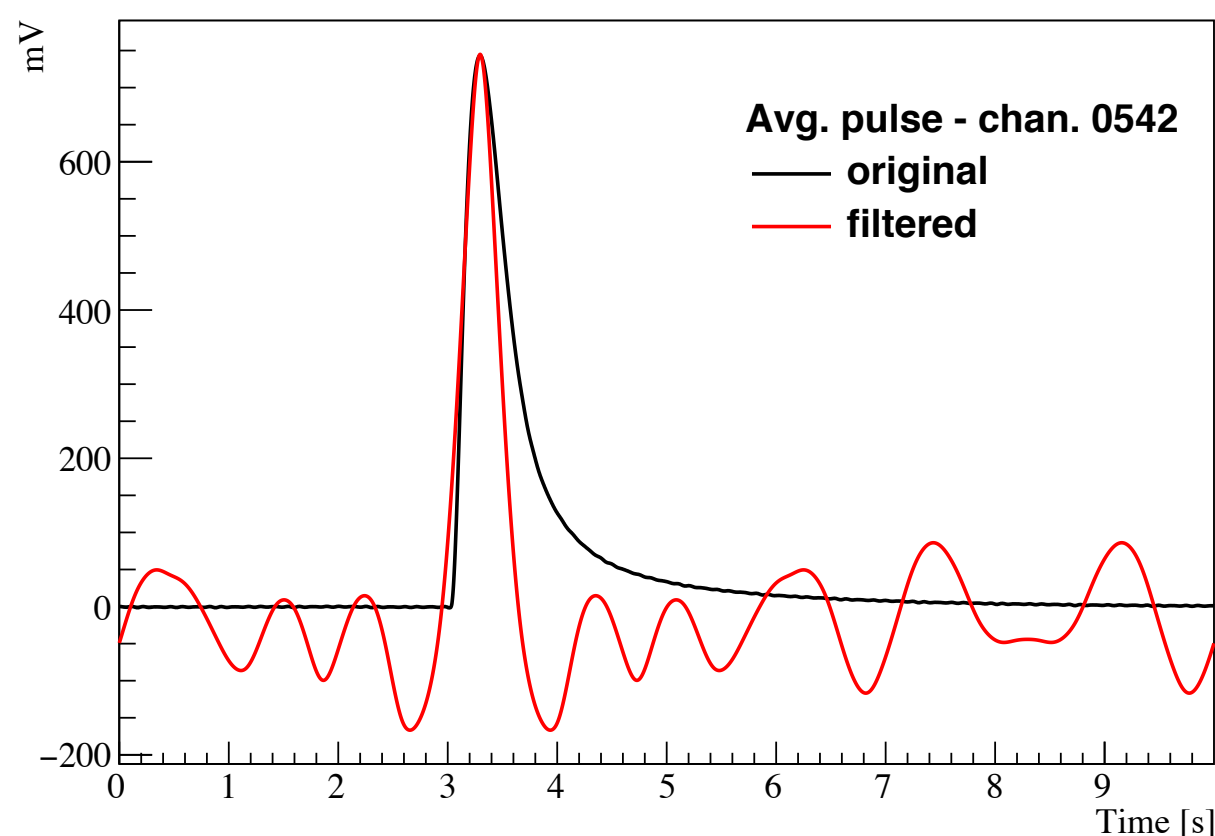
## Data processing

- PREPROCESS
- AVERAGE PULSE
- AVERAGE NOISE
- AMPLITUDE
- STABILIZATION
- CALIBRATION
- BLINDING
- COINCIDENCES
- PULSE SHAPE ANALYSIS**

6 pulse shape parameters are computed for each event (e.g. rise time) and a 6-dim distance from a control sample is evaluated (Mahalanobis distance)

Discards pile-up events

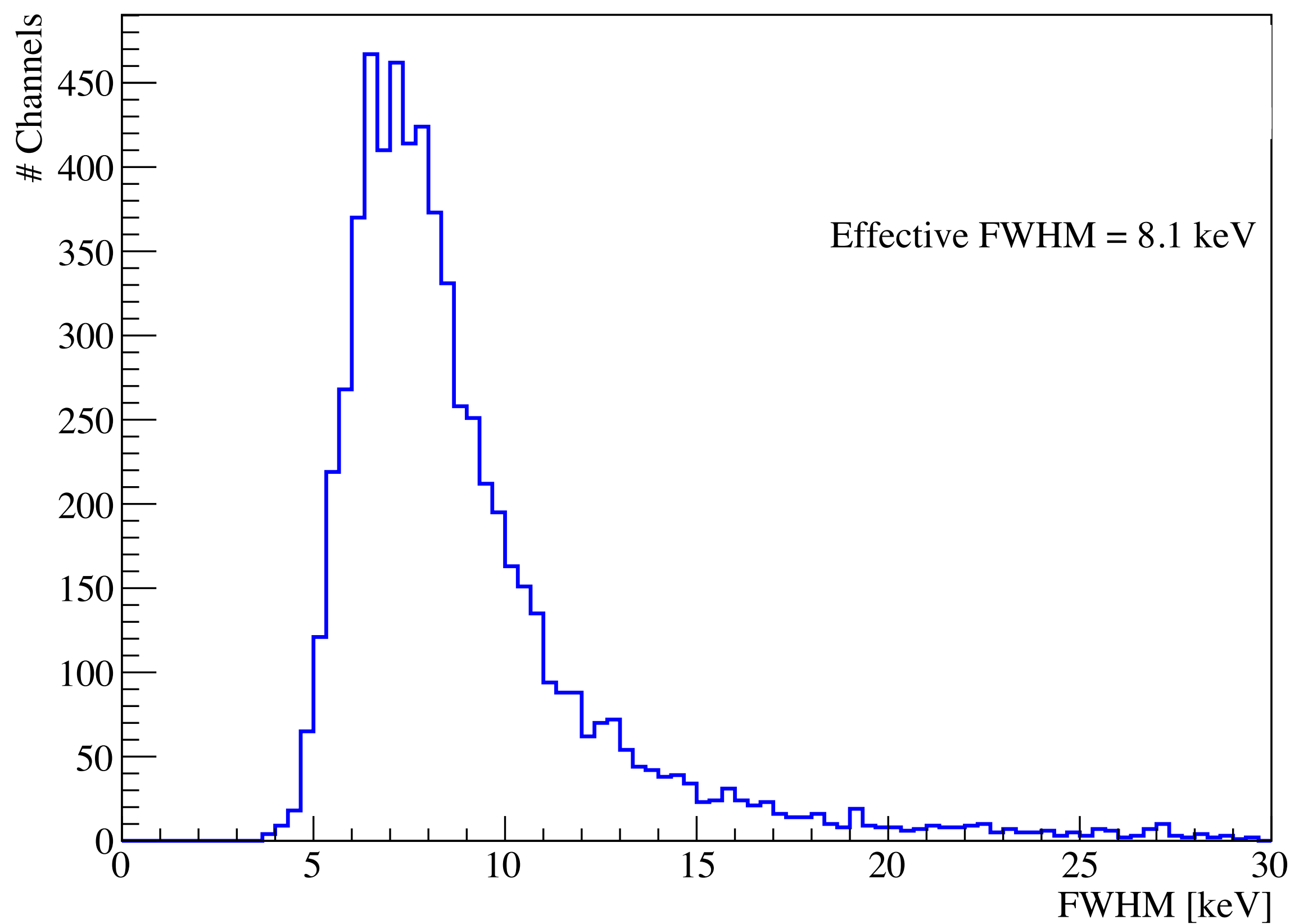
Discards NTD particle interactions



# The CUORE experiment

## Resolution

Calibration Resolution at 2615 keV



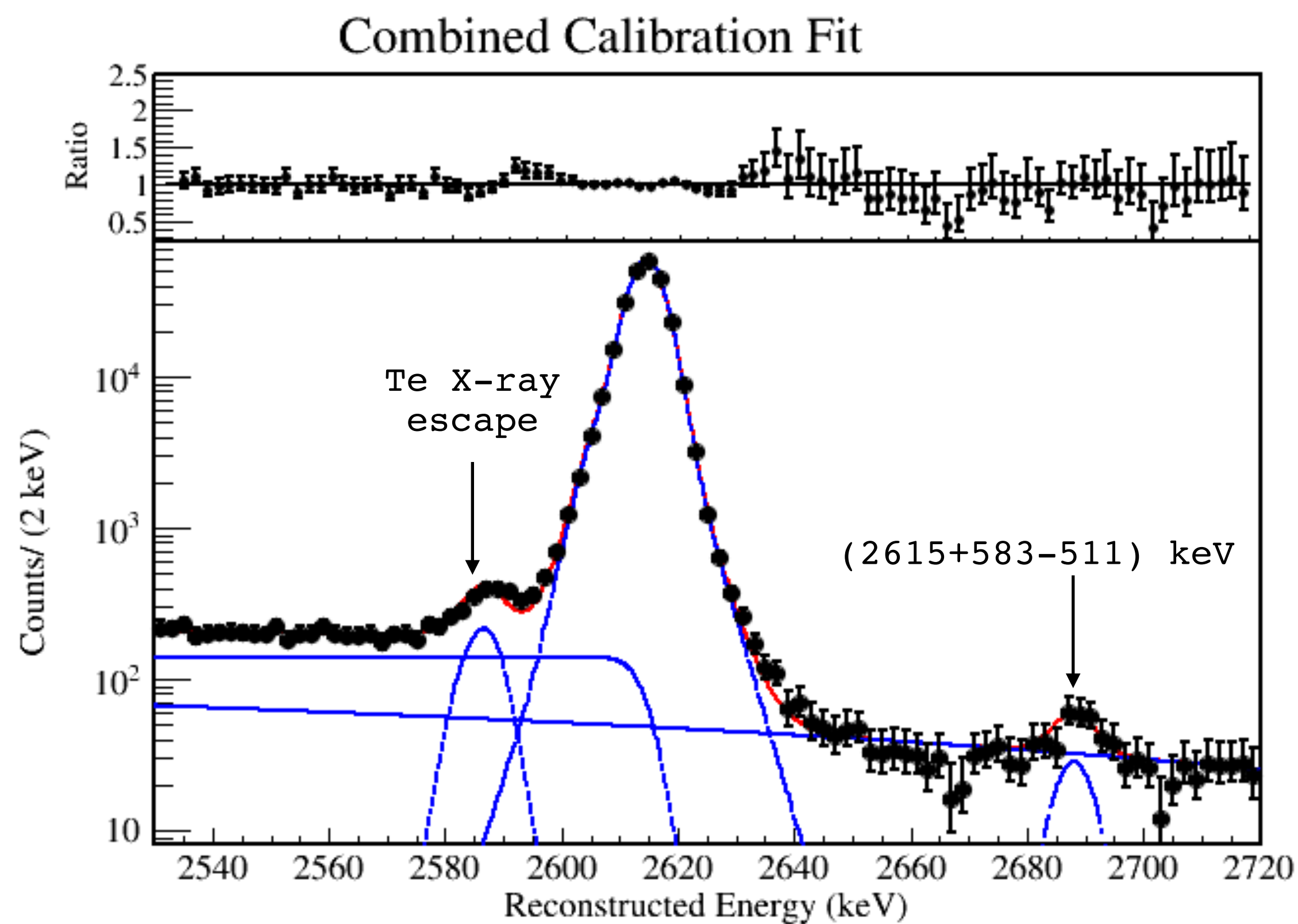
- Full Width Half Maximum (FWHM) resolution computed for each channel-dataset pair
- Exposure weighted harmonic mean effective resolution

$$\frac{1}{\Delta E} = \frac{\sum_{ch,ds} M \Delta t_{ch,ds} / \Delta E_{ch,ds}}{\sum_{ch,ds} M \Delta t_{ch,ds}}$$



# The CUORE experiment

## Detector response model

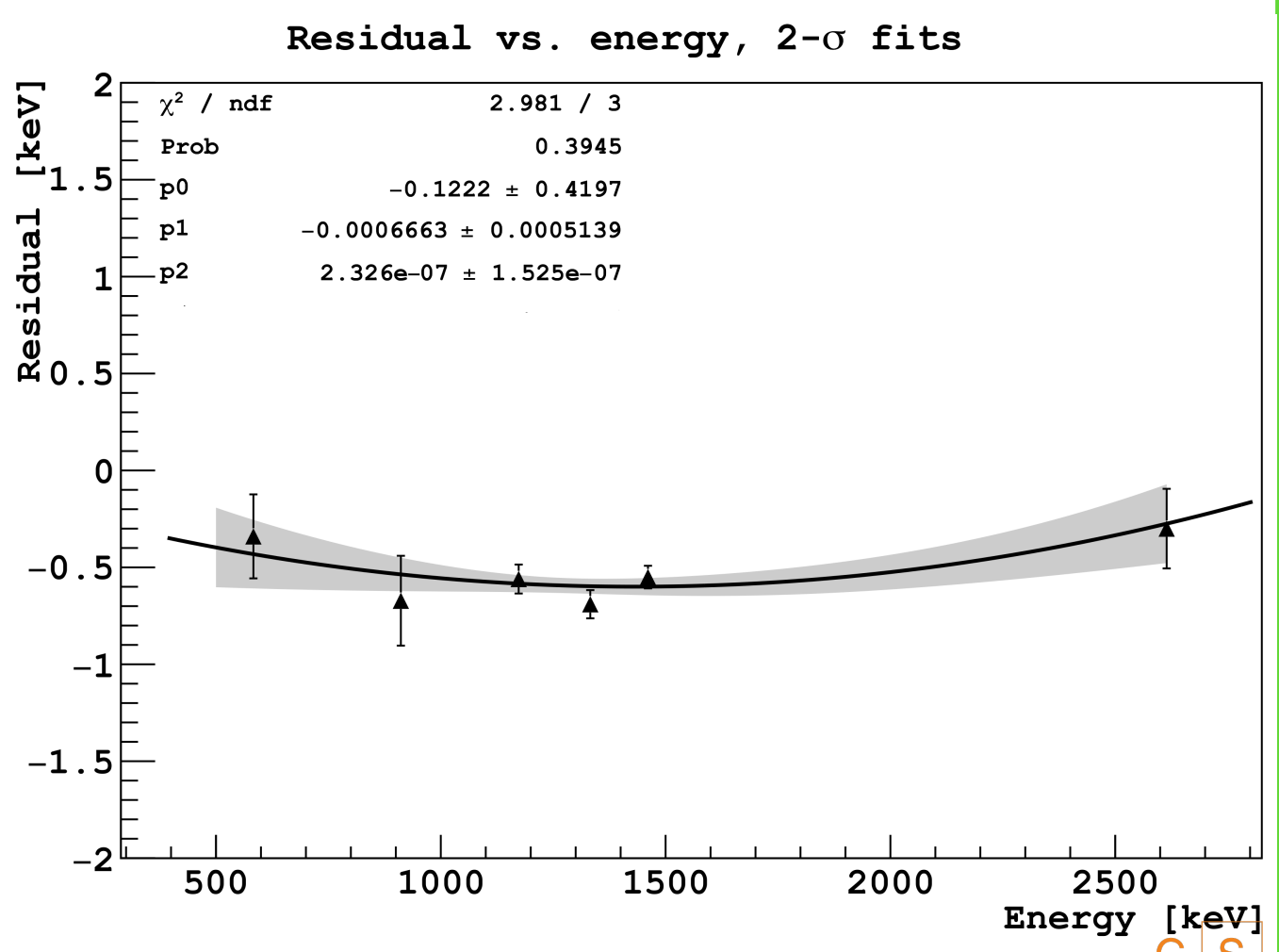
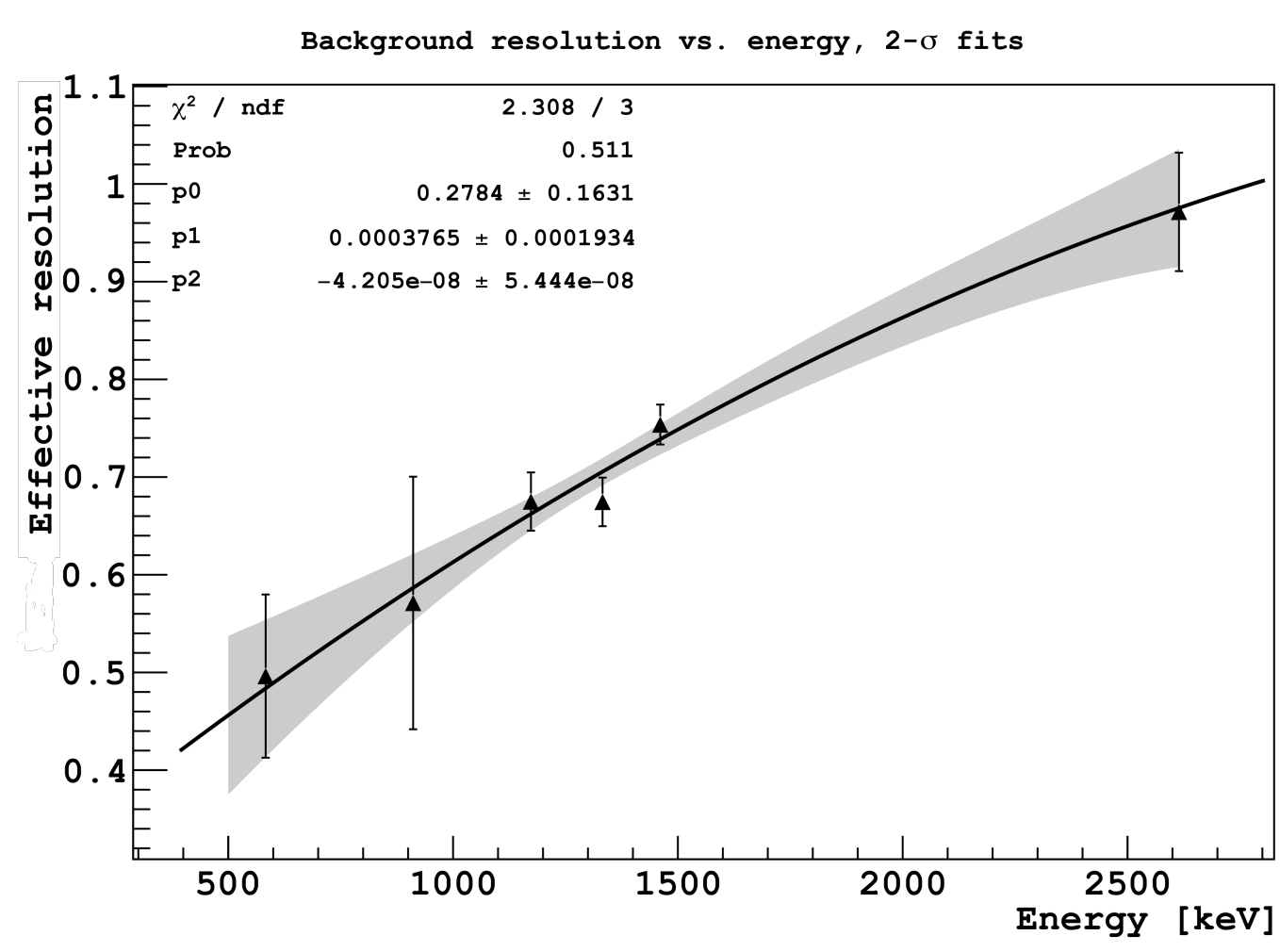
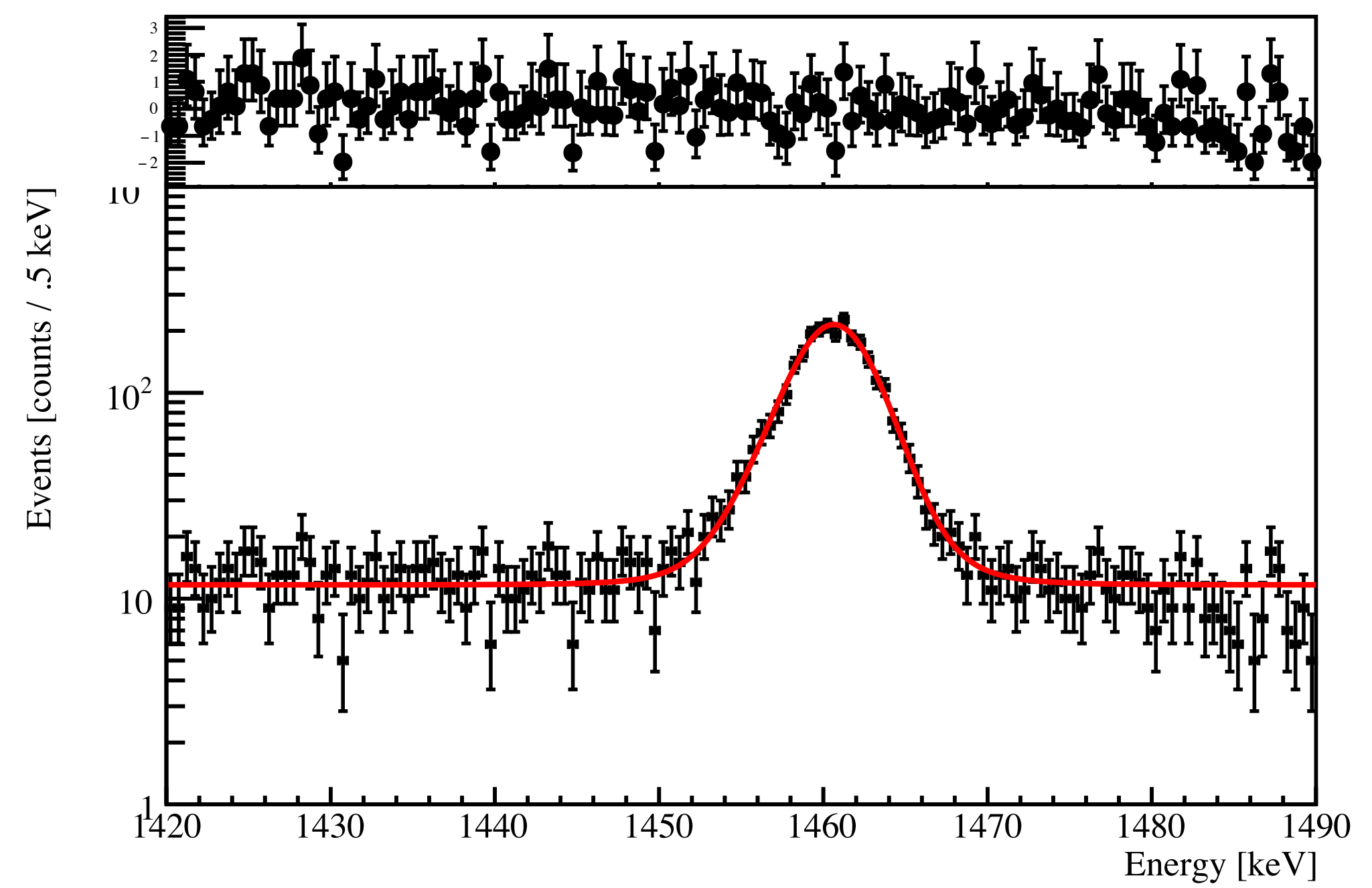


- Fit  $^{208}\text{Tl}$  line at 2614.5 keV in calibration spectra to precisely evaluate the lineshape on a channel-dataset basis
- Main peak parameterized with a 3-Gaussian superposition
- Fit run simultaneously on each tower to constrain backgrounds and side-structures

# The CUORE experiment

## Detector response model

- Extract FWHM from peak fit to physics data
- Fit with 2<sup>nd</sup> order polynomial and extrapolation to any energy (e.g.  $Q_{\beta\beta}$ )
- Compute peak position residual from literature values
- Fit with 2<sup>nd</sup> order polynomial as a function of energy
- Resolution and energy scale uncertainty (systematics)



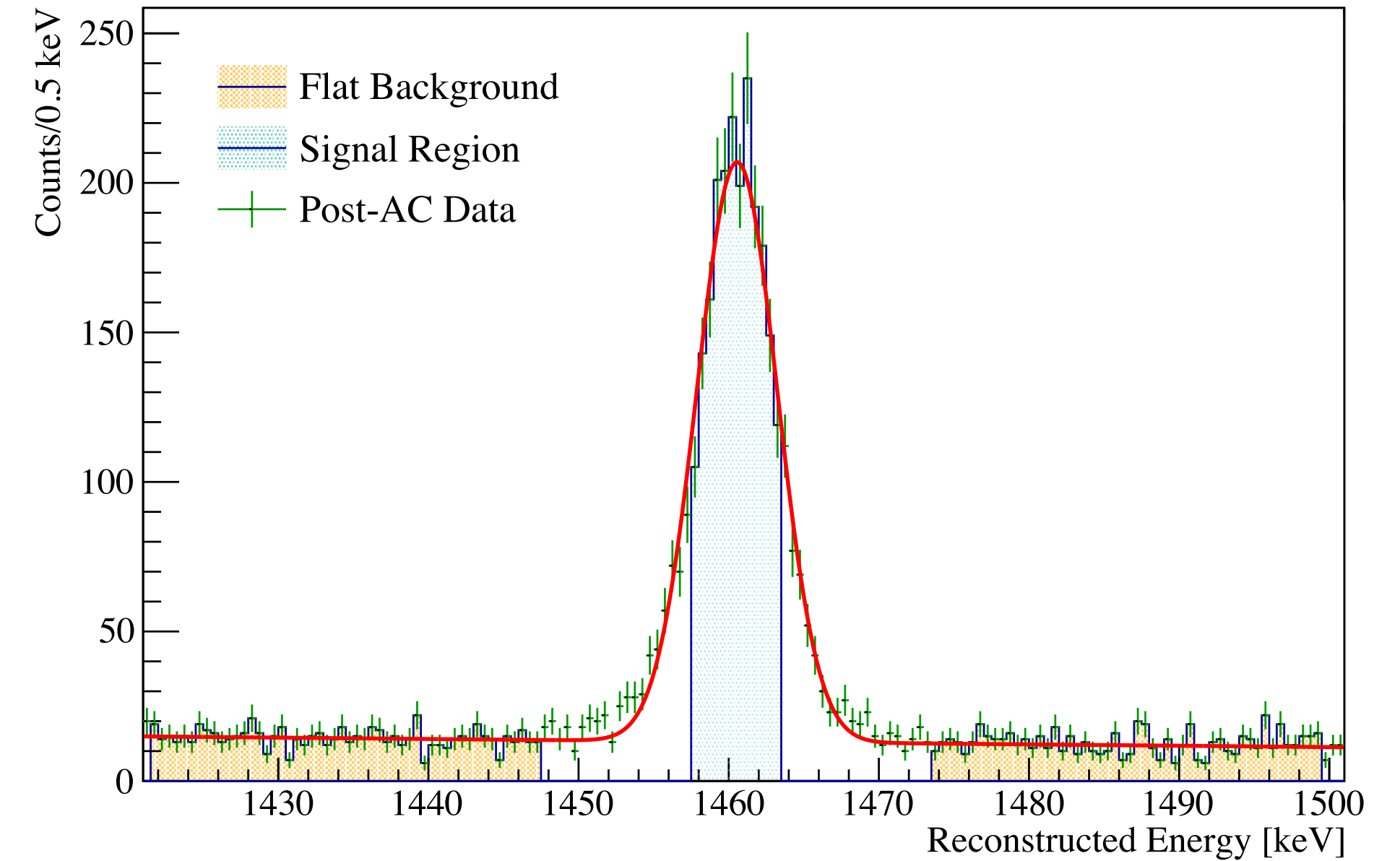


# The CUORE experiment

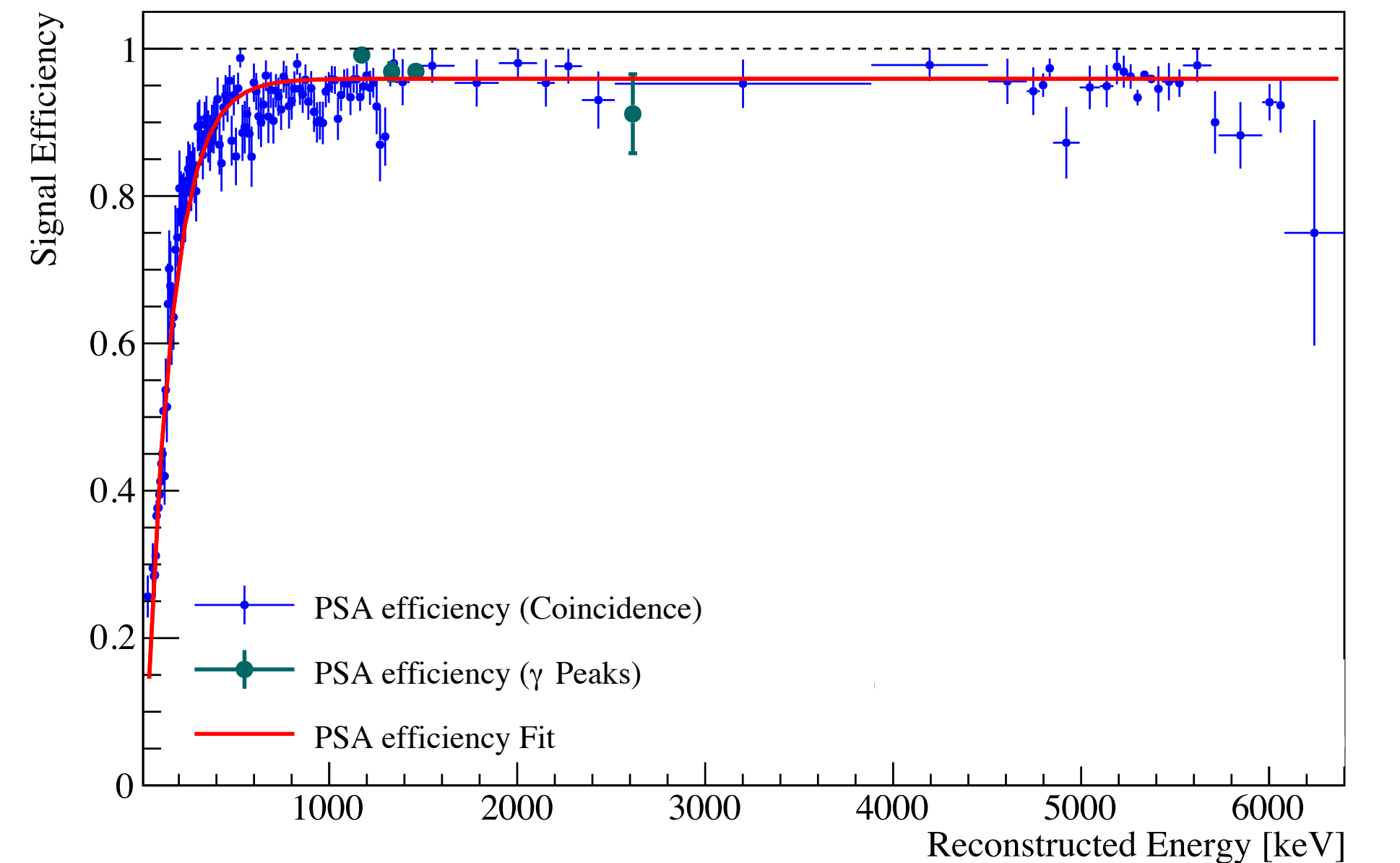
## Detector response model

- Reconstruction efficiency includes
  - Trigger
  - Event reconstruction
  - Pile-up rejection
- Anti-coincidence efficiency quantifies the probability of correctly identifying single-site events
- Pulse Shape Analysis (PSA) Efficiency: fraction of events selected in a 6-dimensional pulse-shape parameter space

Anti-Coincidence Efficiency



PSA Efficiency





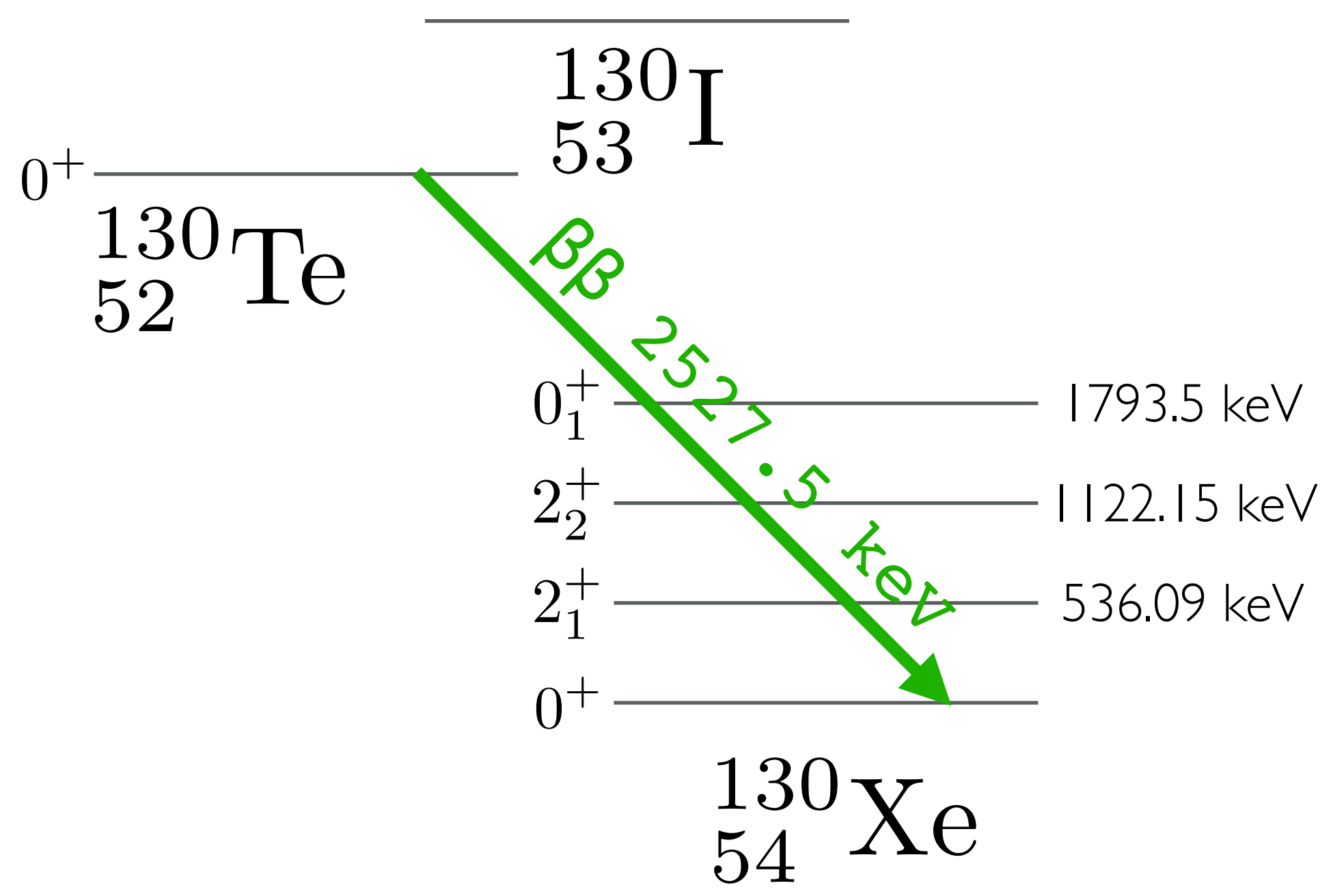
# $0\nu\beta\beta$ Search Ground State Analysis



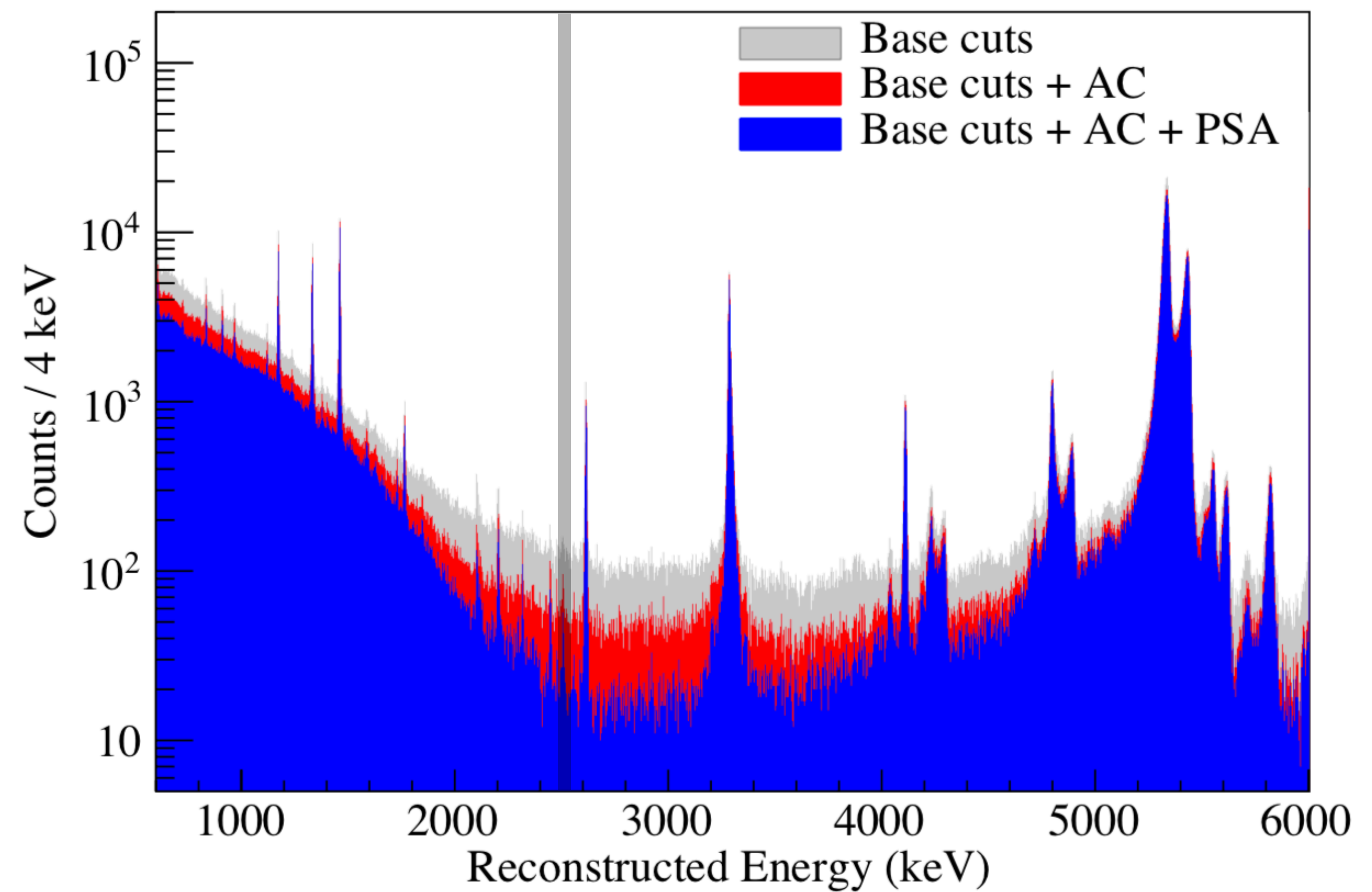


# Ground state analysis

[C. Alduino et al., Phys. Rev. Lett. 124, 12 122501 (2020)]



[B. Singh, Nuclear Data Sheets 93, 33 (2001)]

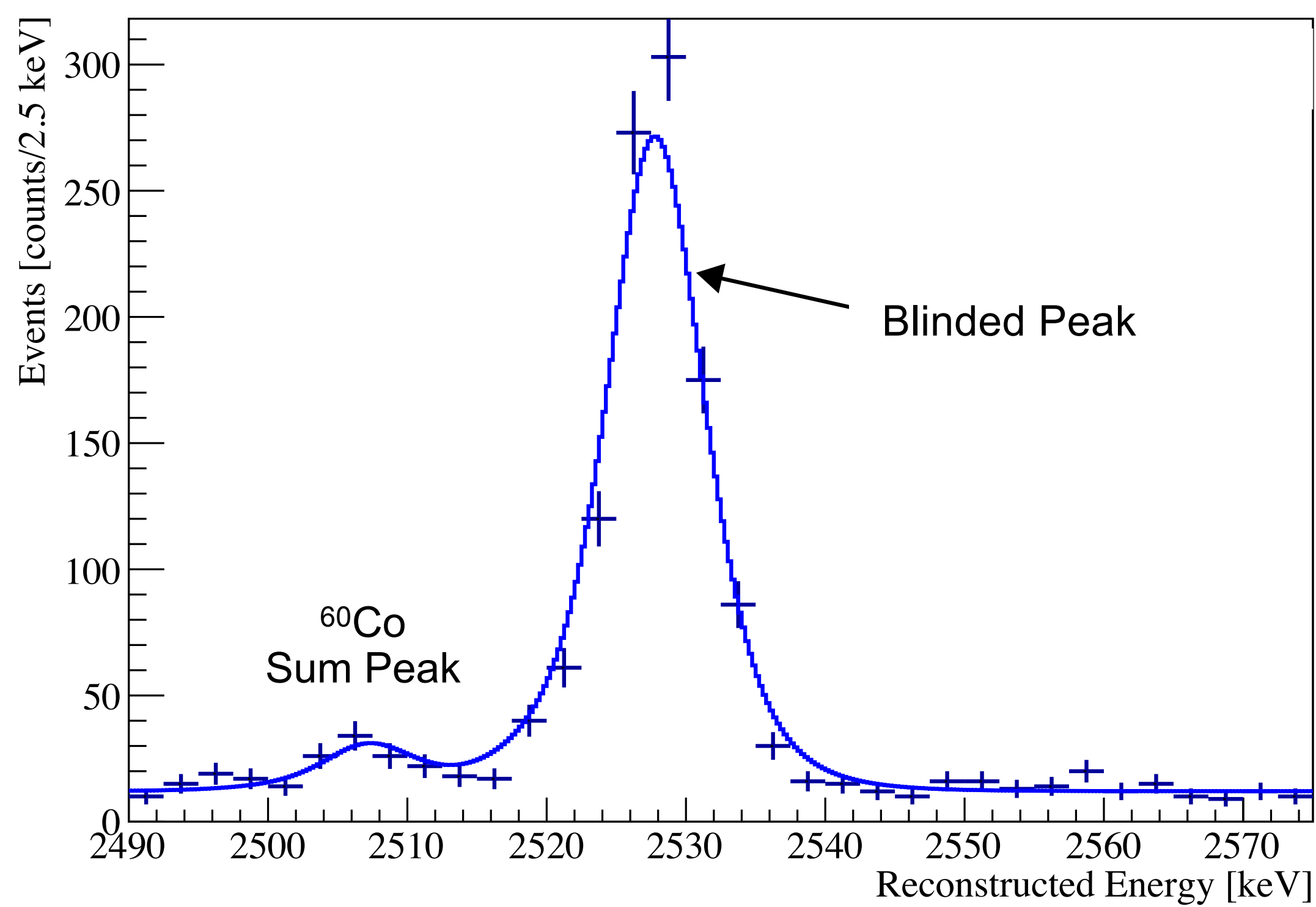


Signal events mostly single-site  $\epsilon = 88.3(5) \%$

# Ground state analysis

$$P(\vec{\theta} | \vec{E}, H_{S+B}) = \frac{\mathcal{L}(\vec{E} | \vec{\theta}, H_{S+B}) \cdot \pi(\vec{\theta} | H_{S+B})}{\int_{\Omega} \mathcal{L} \pi d\vec{\theta}}$$

$$\mathcal{L}(\vec{E} | \vec{\theta}, H_{S+B}) = \prod_{\text{dataset channel}} \prod \left[ \frac{e^{-\lambda} \lambda^n}{n!} \prod_{\text{event } i} \left( \frac{S}{\lambda} pdf_{0\nu\beta\beta}(E_i | \vec{\theta}) + \frac{C}{\lambda} pdf_{60\text{Co}}(E_i | \vec{\theta}) + \frac{b}{\lambda} \frac{1}{\Delta E} \right) \right]$$



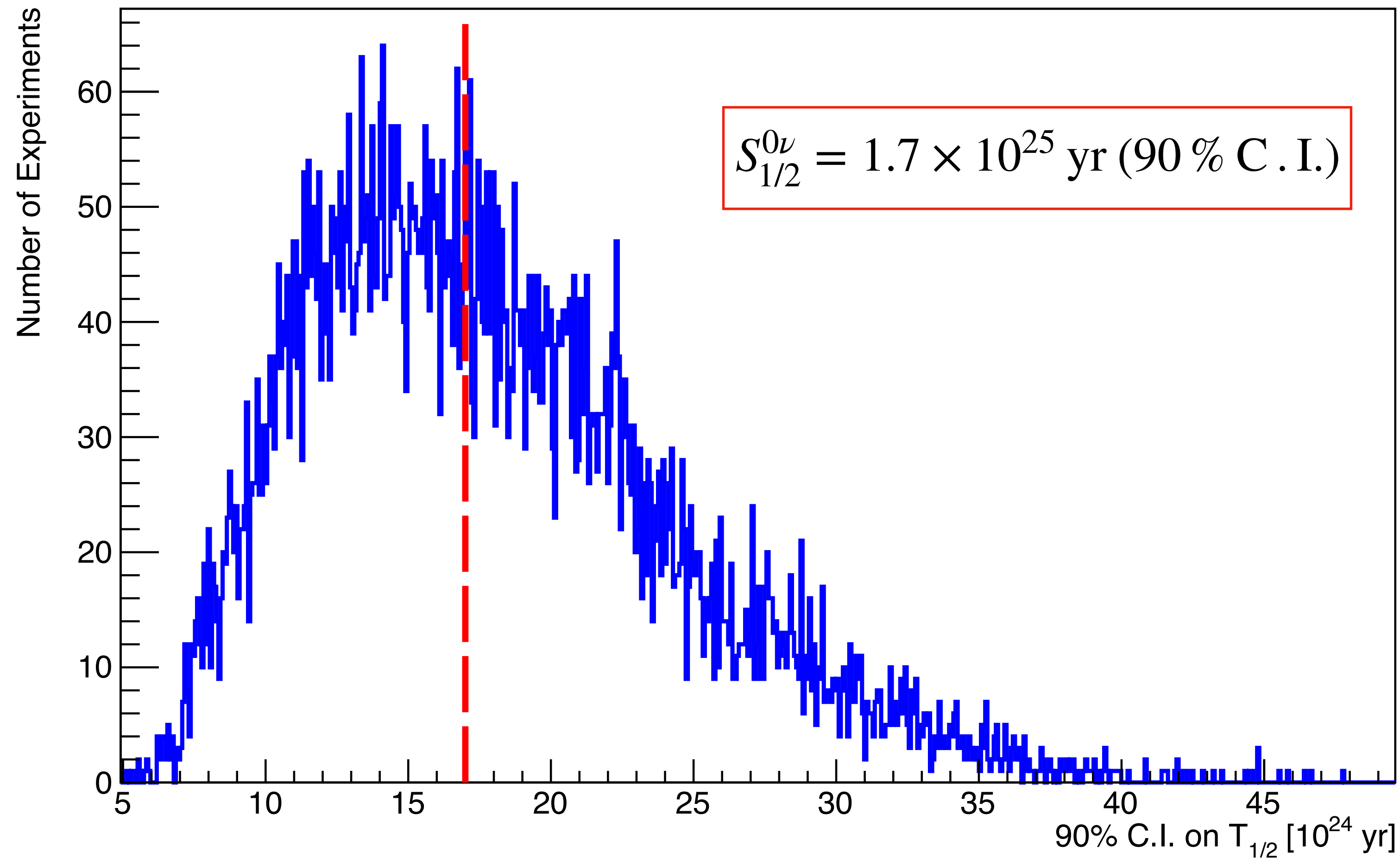
Region Of Interest (ROI)  
[2490, 2575] keV

[C. Alduino et al., Phys. Rev. Lett. 124, 12 122501 (2020)]



# Ground state analysis

## Projected Sensitivity

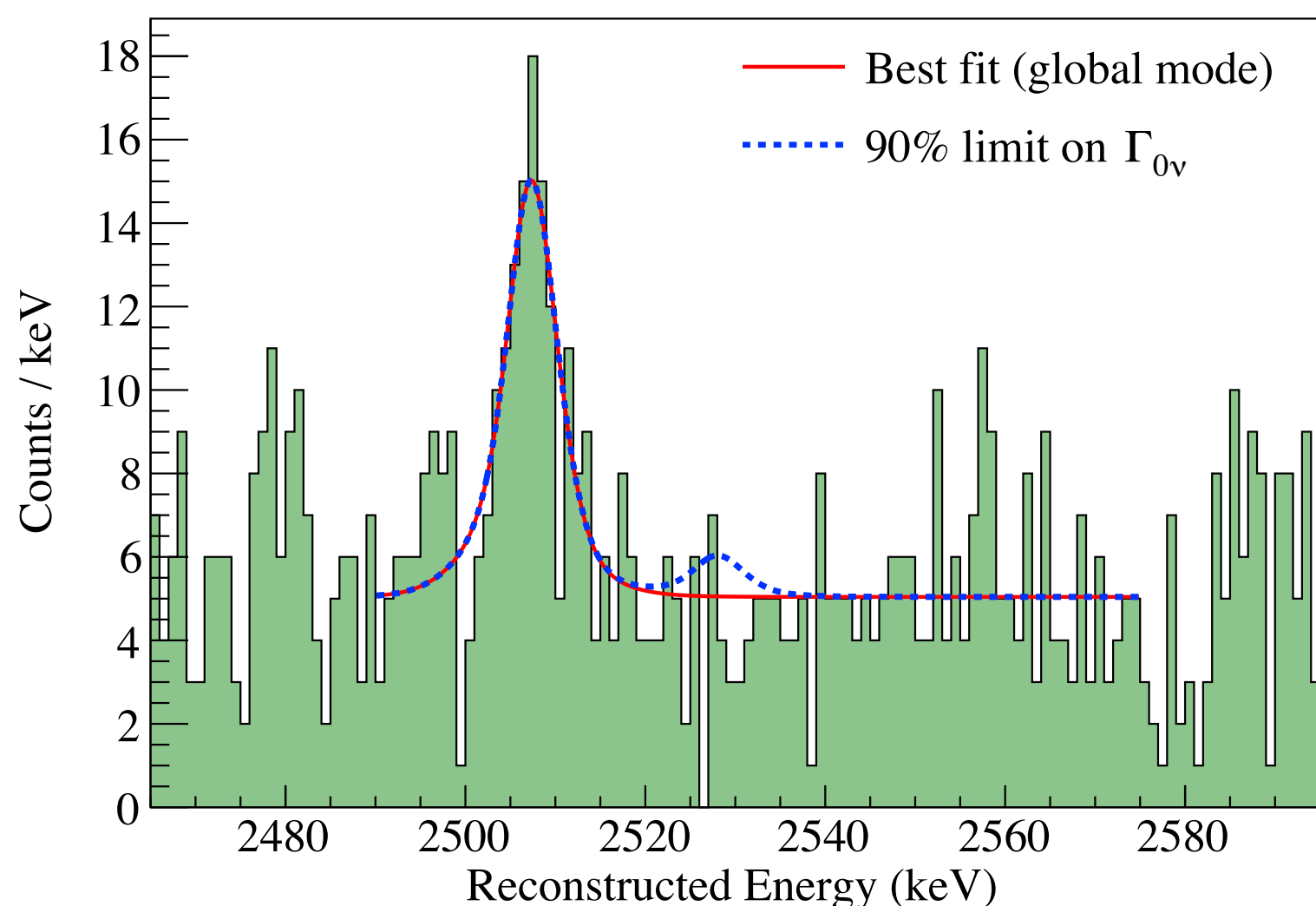


$10^4$  Toy Monte Carlo pseudo-experiments are generated and fit to extract the median sensitivity

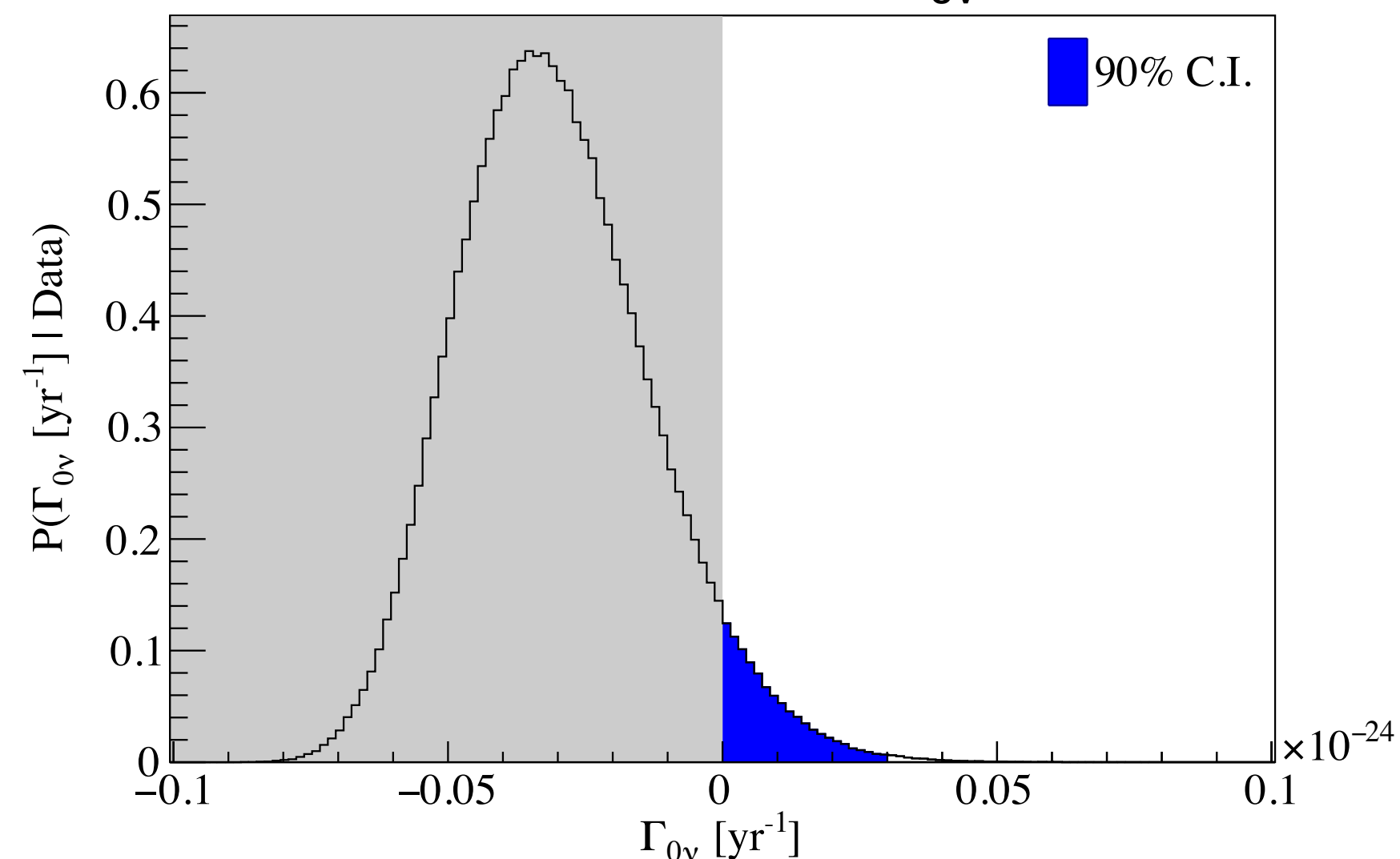
[C. Alduino et al., Phys. Rev. Lett. 124, 12 122501 (2020)]

# Ground state analysis

CUORE ROI Spectrum



Posterior for  $\Gamma_{0\nu}$



- No evidence for  $0\nu\beta\beta$  decay

$$T_{1/2}^{0\nu} > 3.2 \times 10^{25} \text{ yr (90\% C.I.)}$$

- Interpretation in context of light Majorana neutrino exchange

$$m_{\beta\beta} < 75 - 350 \text{ meV}$$

[C. Alduino et al., Phys. Rev. Lett. 124, 12 122501 (2020)]

- Total exposure  $\text{TeO}_2$ : 372.5 kg · yr
- Likelihood model: flat continuum (BI), posited peak for  $0\nu\beta\beta$  (rate), peak for  $^{60}\text{Co}$  (rate + position)
- Unbinned fit on physical range (rates non-negative), uniform prior on the rate parameter  $\Gamma_{0\nu}$
- Systematics: repeat fits with nuisance parameters, allow negative rates (<0.4% impact on limit)

## Detector Performance Parameters

Background Index

$$(1.38 \pm 0.07) \times 10^{-2} \text{ cnts}/(\text{keV} \cdot \text{kg} \cdot \text{yr})$$

Characteristic FWHM  $\Delta E$  at  $Q_{\beta\beta}$

$$7.0 \pm 0.3 \text{ keV}$$

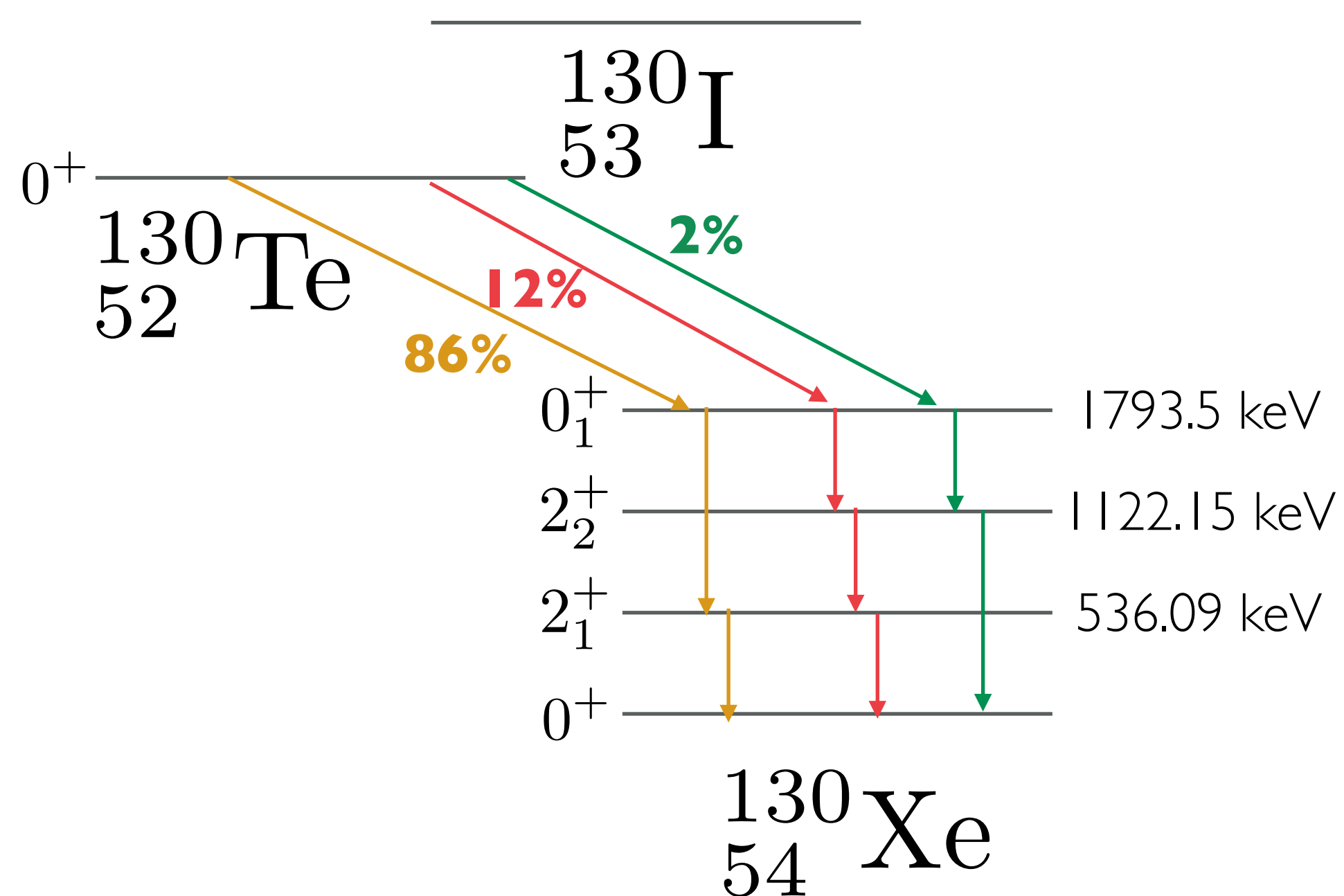




# $\beta\beta$ Search Excited State Analyses



# Excited states analysis



[B. Singh, Nuclear Data Sheets 93, 33 (2001)]

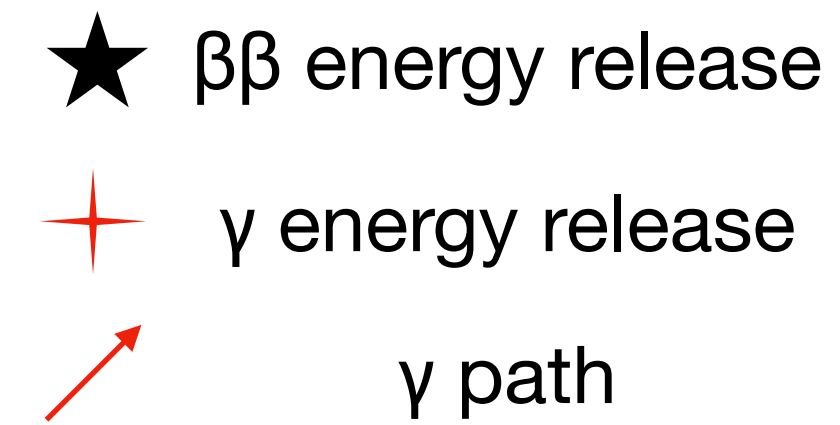
There are 3 possible de-excitation patterns:

- A. (86%)  $\beta\beta_{0-734} + \gamma_{1257} + \gamma_{536}$
- B. (12%)  $\beta\beta_{0-734} + \gamma_{671} + \gamma_{586} + \gamma_{536}$
- C. (2%)  $\beta\beta_{0-734} + \gamma_{1122} + \gamma_{671}$

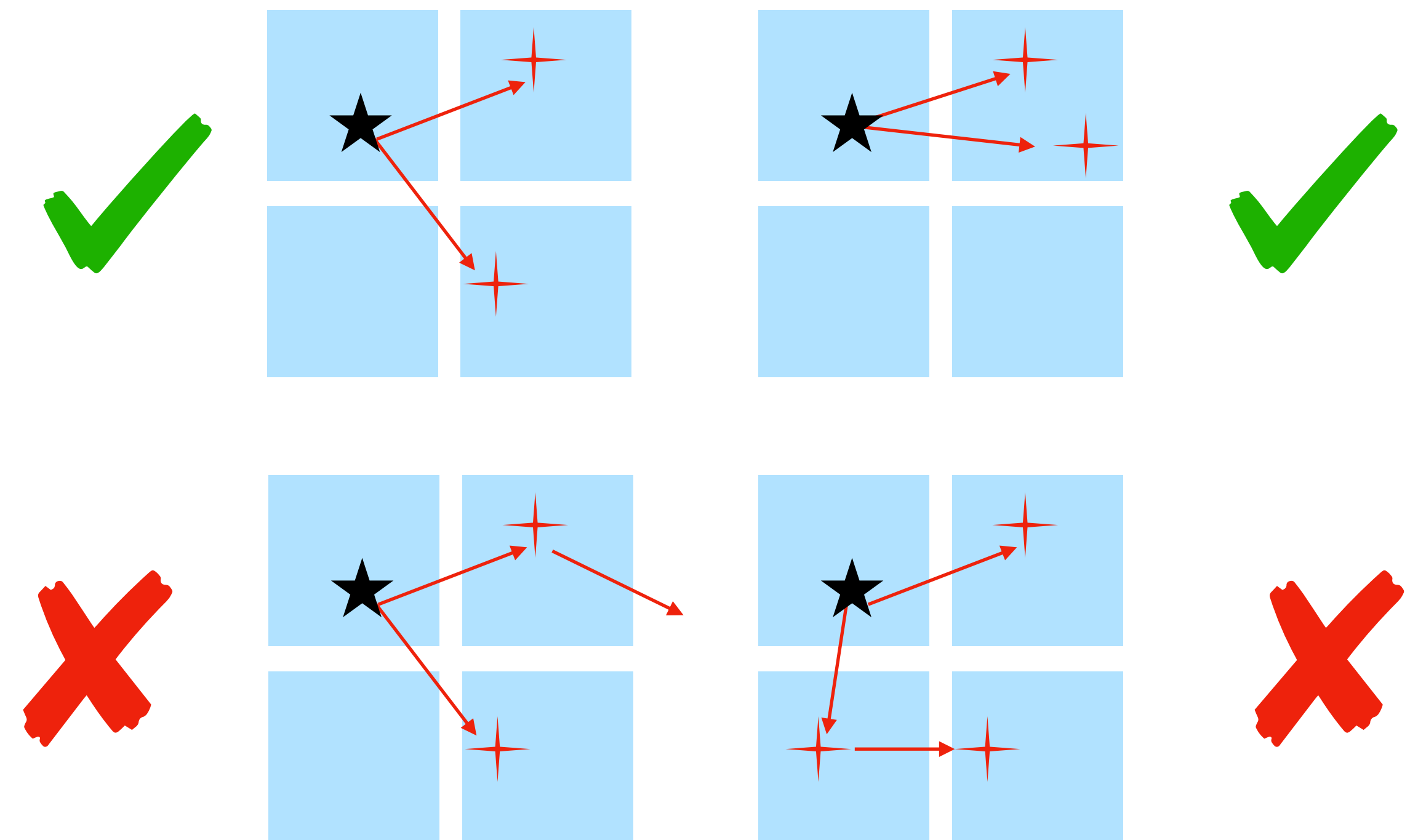


# Excited states analysis

## Full containment requirement



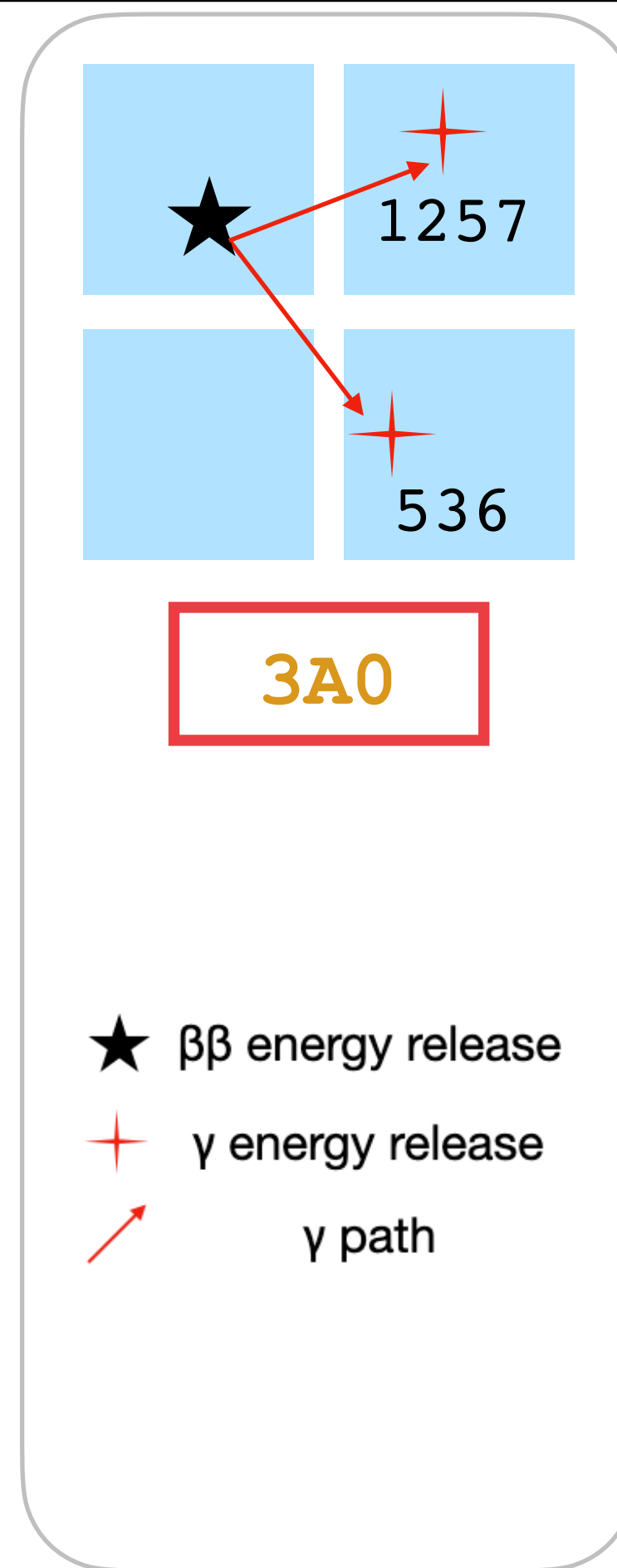
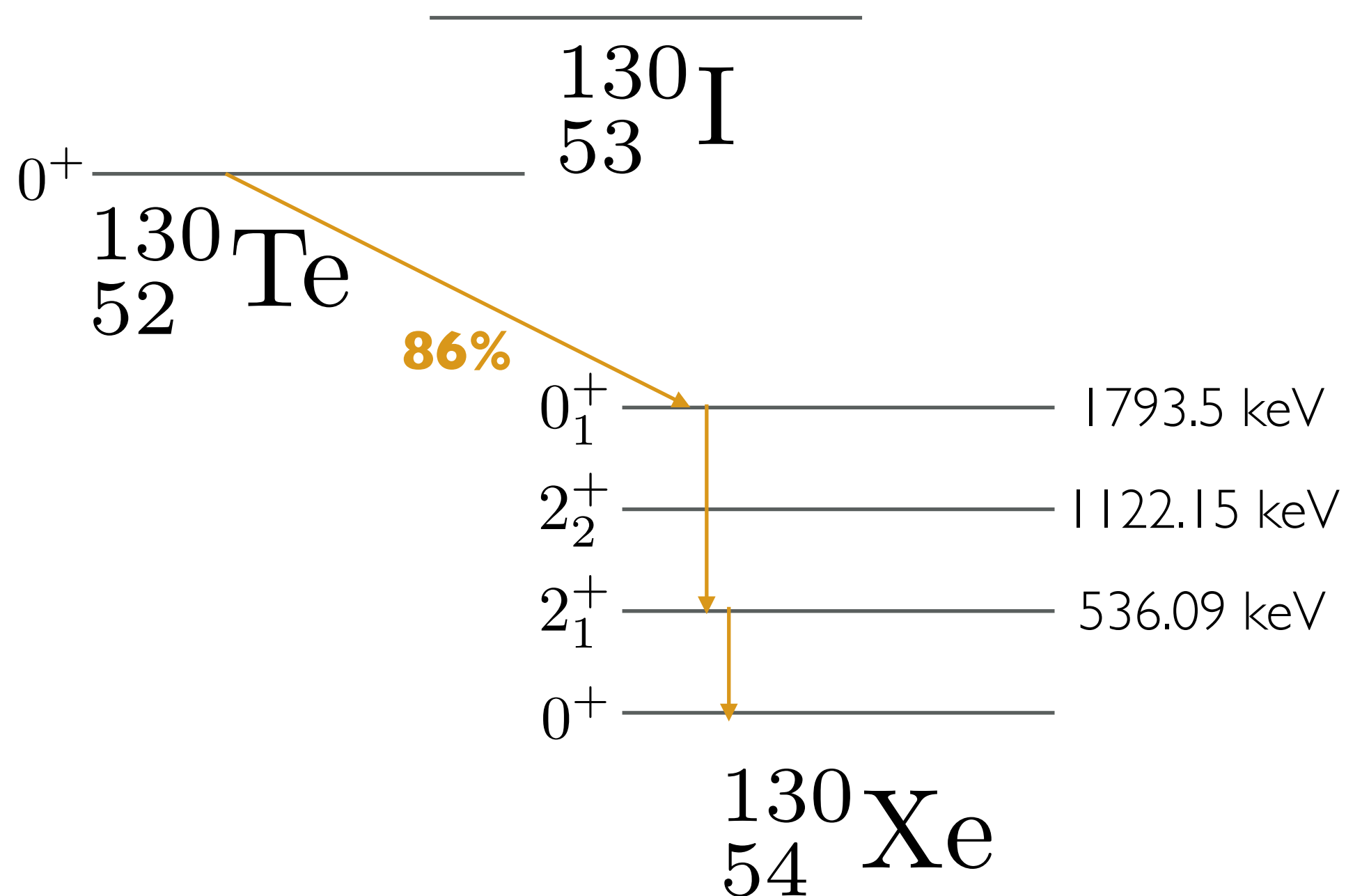
- Each final state particle (except neutrinos) must fully release its energy in no more than one crystal
- Not considering partially contained signatures
- Significantly simpler analysis
- Significant efficiency loss due to this choice



# Excited states analysis

## Labelling signatures

Let us build a **Multiplicity 3** (3 crystals involved) signature coming from de-excitation **pattern A**



Number of involved crystals

[Multiplicity] [Pattern] [integer index]

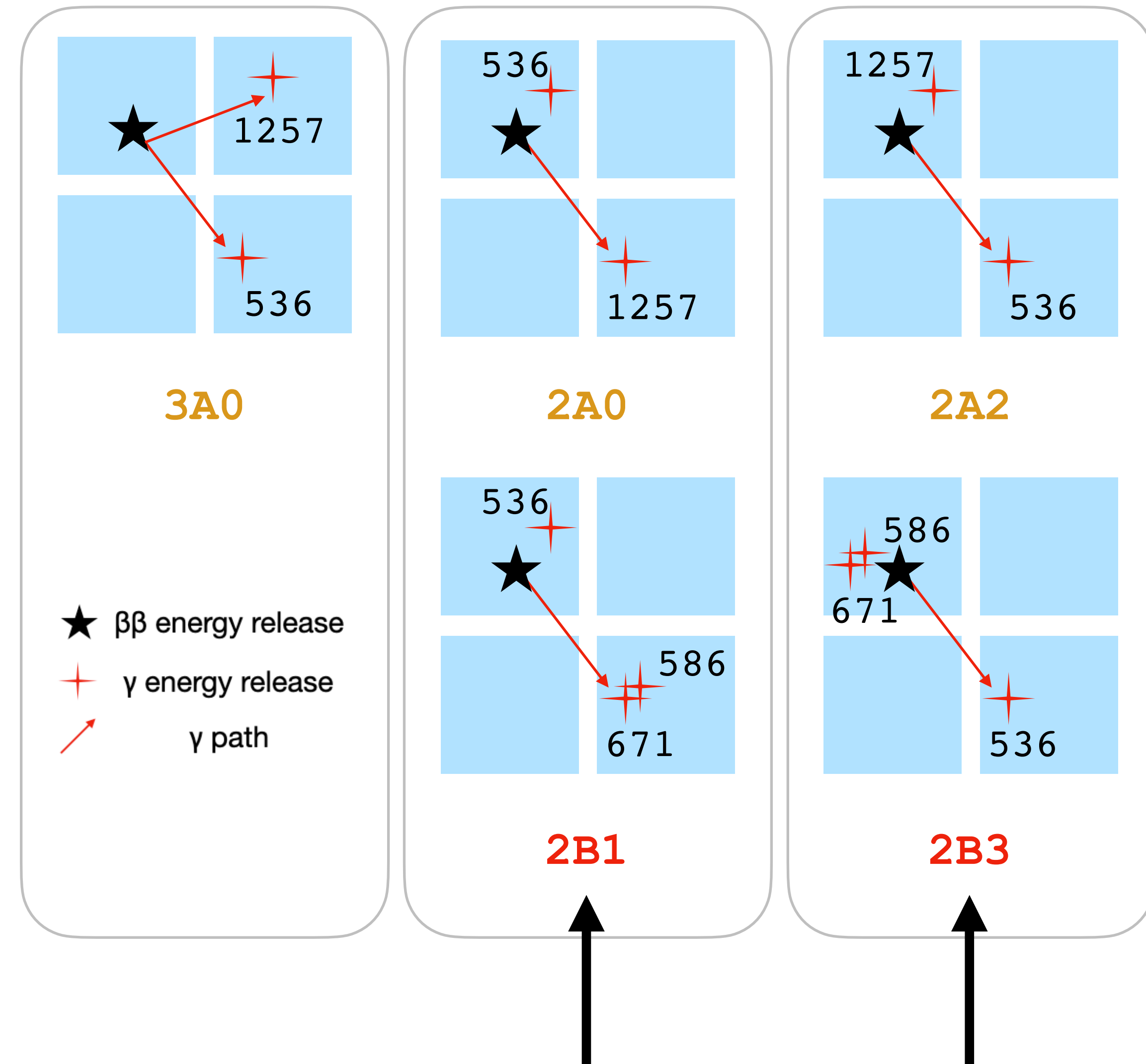
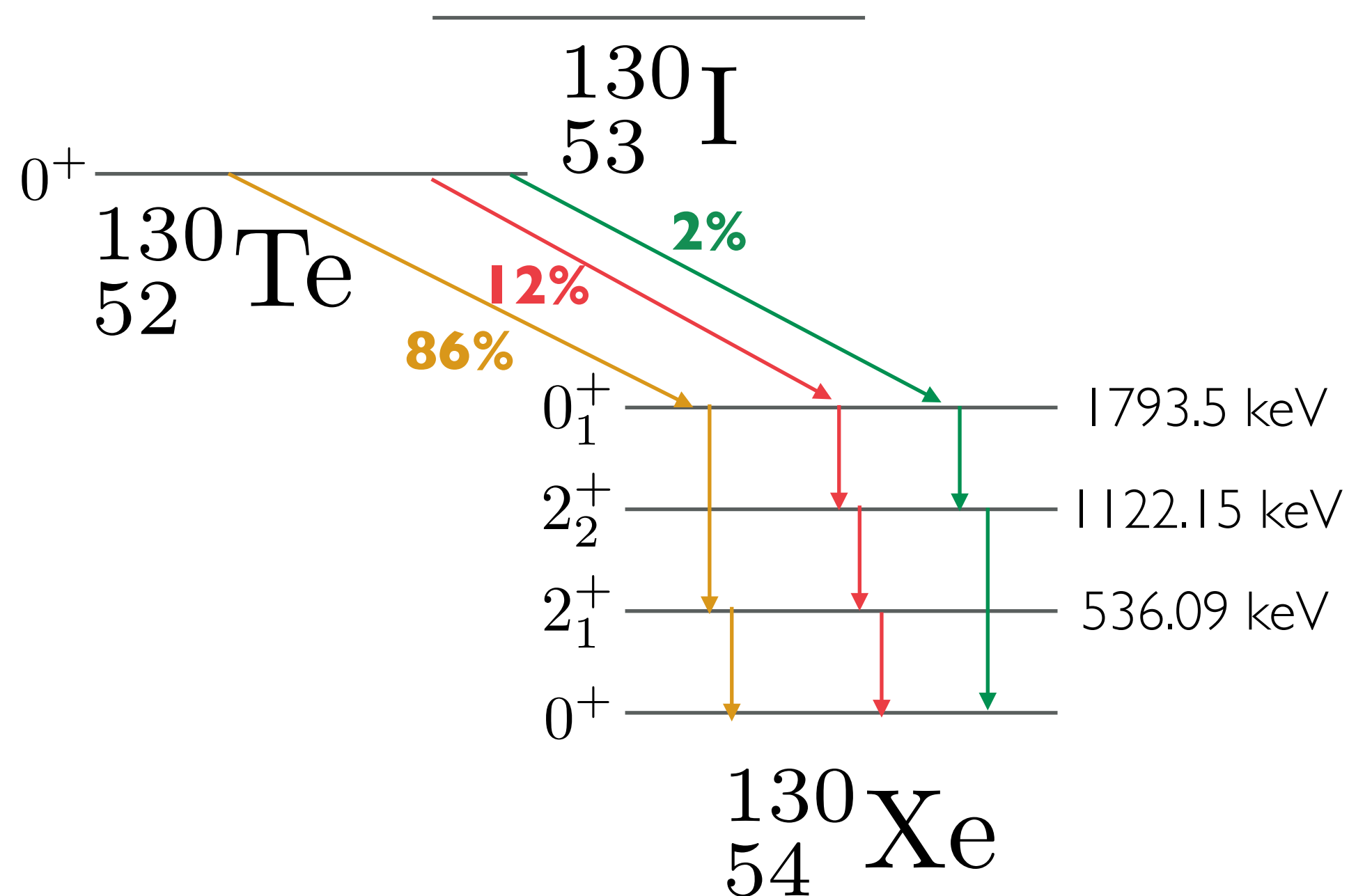
$^{130}\text{Xe}$  de-excitation gamma pattern

Indexing possible combinations



# Excited states analysis

## Experimental signatures



Same experimental signature, 2 different de-excitation patterns

# Excited states analysis

## Ranking experimental signatures

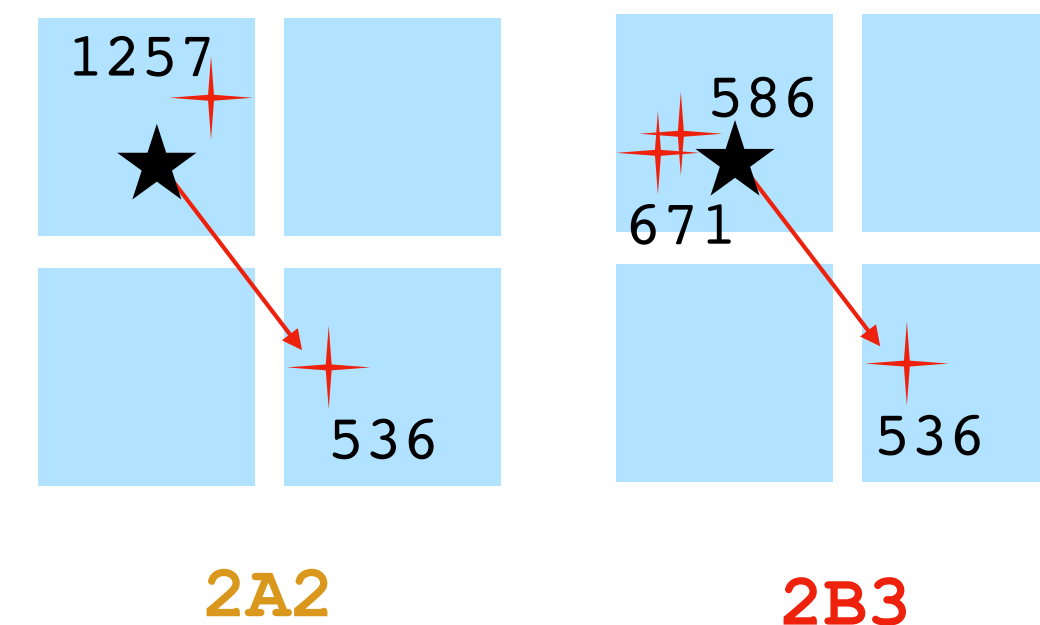
$$S(\epsilon_s, B_s) = \underbrace{\theta(B_s - 1) \frac{\epsilon_s}{\sqrt{B}}}_{\text{background-free}} + \underbrace{\theta(1 - B_s) \frac{5\epsilon_s}{-\ln(3 \cdot 10^{-7})}}_{\text{background-dominated}}$$

- Many possible signatures
- Most contribute negligibly to the sensitivity
- A threshold of  $S > 5\%$  was set
- Three signatures account for 80-90% of the overall sensitivity

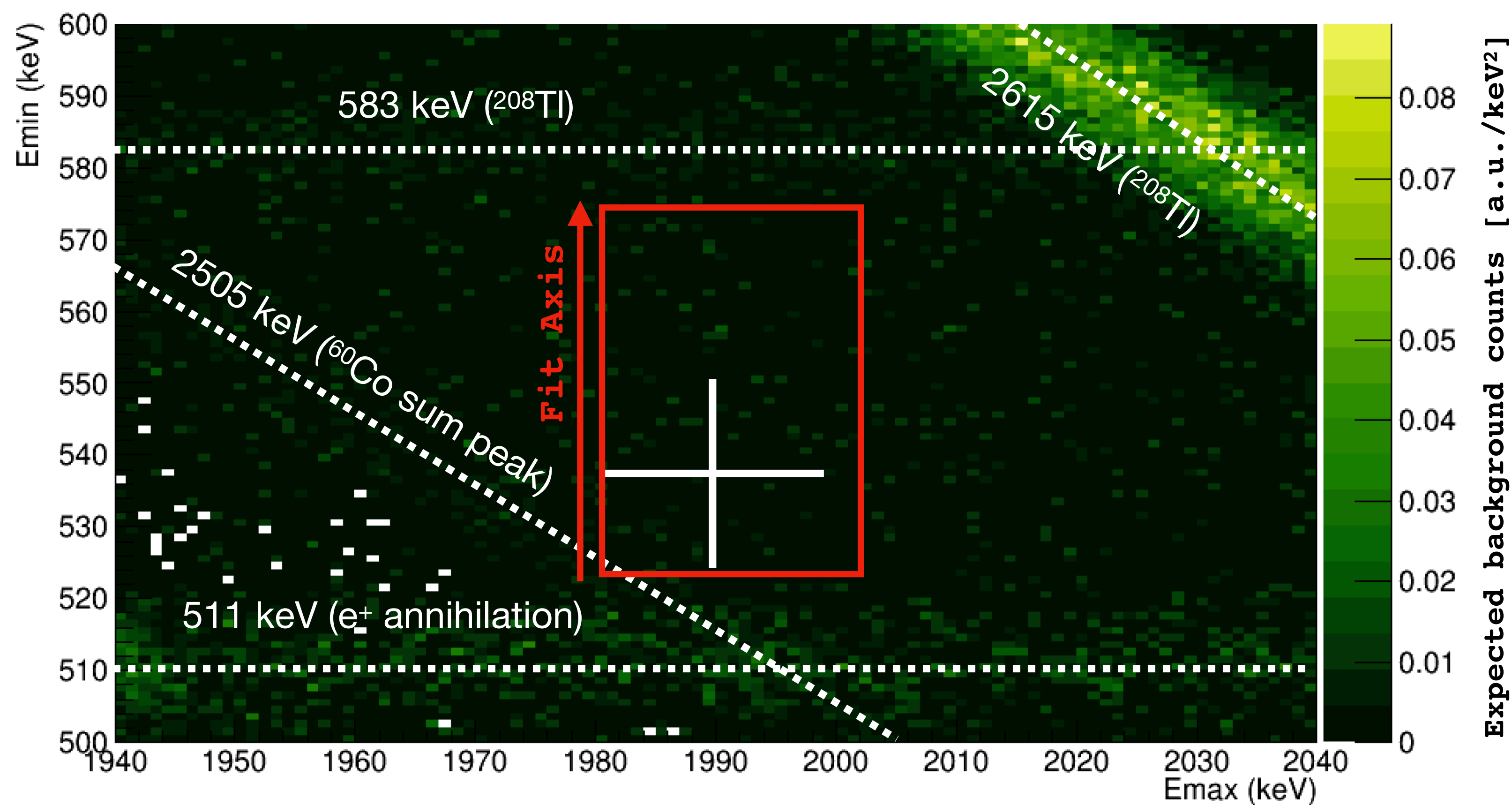
$S_{2\nu}(S_{0\nu})$ Signature	Crystal 1 [keV]	Crystal 2 [keV]	Crystal 3 [keV]	Crystal 4 [keV]
39.8% (38.5%) 2A0-2B1	$\beta\beta + \gamma_{(A2)}$ $\beta\beta + \gamma_{(B3)}$ 536 - 1270	$\gamma_{(A1)}$ $\gamma_{(B1)} + \gamma_{(B2)}$ 1257		
2.3% (2.3%) 2A1-2B2-2C1	$\gamma_{(A1)} + \gamma_{(A2)}$ $\gamma_{(B1)} + \gamma_{(B2)} + \gamma_{(B3)}$ $\gamma_{(C1)} + \gamma_{(C2)}$ 1793	$\beta\beta$ $\beta\beta$ $\beta\beta$ 0 - 734		
21.6% (24.7%) 2A2-2B3	$\beta\beta + \gamma_{(A1)}$ $\beta\beta + \gamma_{(B1)} + \gamma_{(B2)}$ 1257 - 1991	$\gamma_{(A2)}$ $\gamma_{(B3)}$ 536		
25.5% (20.3%) 3A0	$\gamma_{(A1)}$ 1257	$\beta\beta$ 0 - 734	$\gamma_{(A2)}$ 536	
2.4% (2.8%) 2B0-2C2	$\beta\beta + \gamma_{(B2)} + \gamma_{(B3)}$ $\beta\beta + \gamma_{(C1)}$ 1122 - 1856	$\gamma_{(B1)}$ $\gamma_{(C2)}$ 671		
0.2% (0.1%) 2B4	$\beta\beta + \gamma_{(B2)}$ 586 - 1320	$\gamma_{(B1)} + \gamma_{(B3)}$ 1207		
1.4% (2.0%) 2B5	$\beta\beta + \gamma_{(B1)} + \gamma_{(B3)}$ 1207 - 1941	$\gamma_{(B2)}$ 536		
1.0% (0.9%) 2B6-2C0	$\beta\beta + \gamma_{(B1)}$ $\beta\beta + \gamma_{(C2)}$ 671 - 1405	$\gamma_{(B2)} + \gamma_{(B3)}$ $\gamma_{(C1)}$ 1122		
1.1% (1.1%) 3B0	$\beta\beta + \gamma_{(B3)}$ 536 - 1270	$\gamma_{(B1)}$ 671	$\gamma_{(B2)}$ 586	
0.8% (0.5%) 3B1-3C0	$\gamma_{(B2)} + \gamma_{(B3)}$ $\gamma_{(C1)}$ 1122	$\beta\beta$ $\beta\beta$ 0 - 734	$\gamma_{(B1)}$ $\gamma_{(C2)}$ 671	
0.1% (0.1%) 3B2	$\gamma_{(B1)} + \gamma_{(B2)}$ 1257	$\beta\beta$ 0 - 734	$\gamma_{(B3)}$ 536	
1.3% (1.0%) 3B3	$\beta\beta + \gamma_{(B2)}$ 586 - 1320	$\gamma_{(B1)}$ 671	$\gamma_{(B3)}$ 536	
0.1% (0.1%) 3B4	$\gamma_{(B1)} + \gamma_{(B3)}$ 1207	$\beta\beta$ 0 - 734	$\gamma_{(B2)}$ 536	
1.5% (1.0%) 3B5	$\beta\beta + \gamma_{(B1)}$ 671 - 1405	$\gamma_{(B2)}$ 586	$\gamma_{(B3)}$ 536	
0.8% (0.6%) 4B0	$\beta\beta$ 0 - 734	$\gamma_{(B1)}$ 671	$\gamma_{(B2)}$ 586	$\gamma_{(B3)}$ 536

# Excited states analysis

## Selection cut definition



- The ROI for each signature is defined in an ordered energy space
- Selection cuts based on Monte Carlo (MC) simulations of the CUORE background
- Analytical description of the background based on simulations





# Excited states analysis

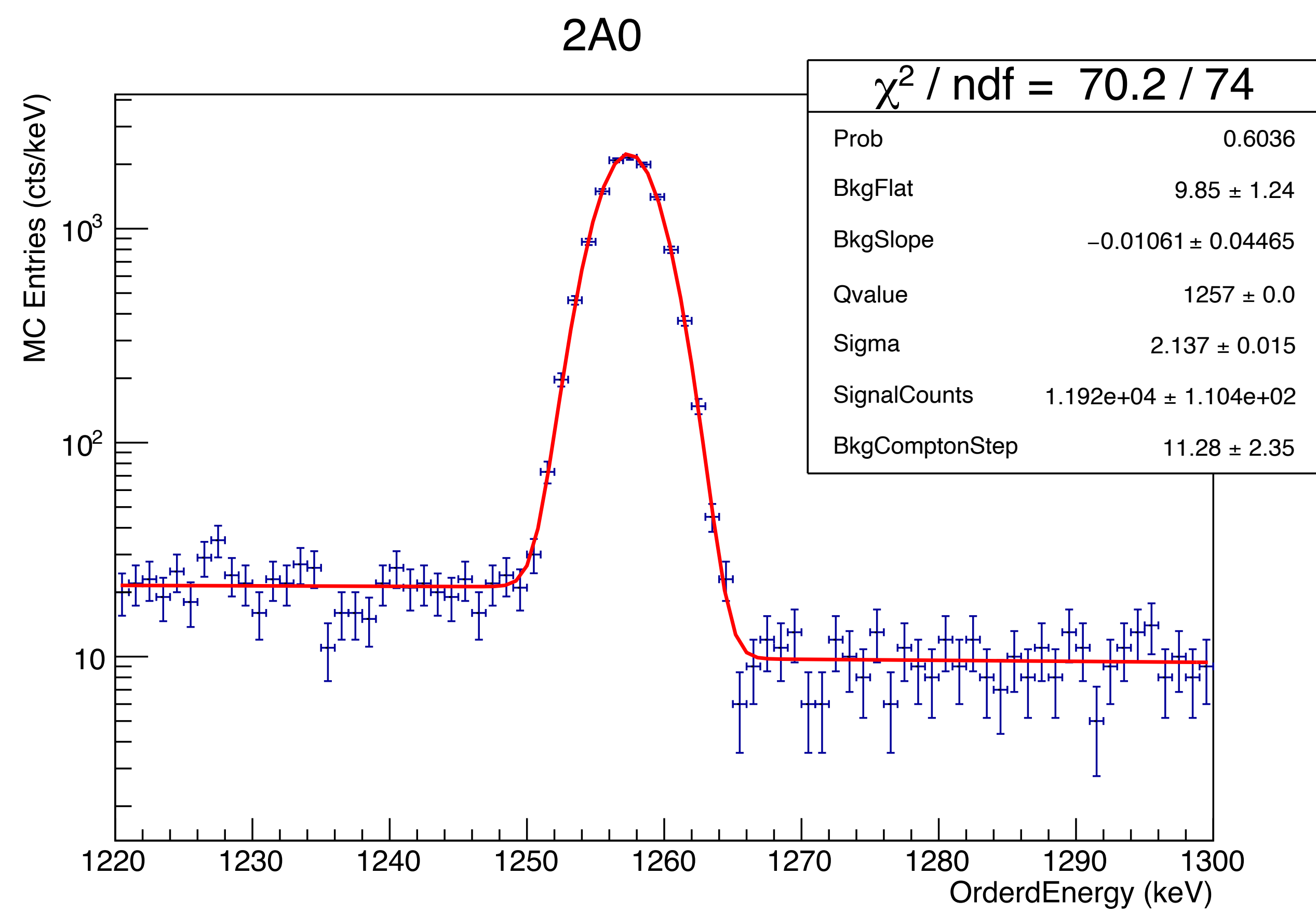
## Ranking experimental signatures

$0\nu\beta\beta$				
Scenario	Multiplicity	Score	Cut	Selection
2A0 - 2B1	2	39%	$1247 \text{ keV} < E_{\min} < 1280 \text{ keV}$	$E_{\min} = \gamma(1257)$
			$1247 \text{ keV} < E_{\max} < 1280 \text{ keV}$	$E_{\max} = \beta\beta(734) + \gamma(536)$
2A2 - 2B3	2	25%	$523 \text{ keV} < E_{\min} < 573 \text{ keV}$	$E_{\min} = \gamma(536)$
			$1981 \text{ keV} < E_{\max} < 2001 \text{ keV}$	$E_{\max} = \beta\beta(734) + \gamma(1257)$
3A0	3	20%	$526 \text{ keV} < E_{\min} < 546 \text{ keV}$	$E_{\min} = \gamma(536)$
			$700 \text{ keV} < E_{\text{med}} < 760 \text{ keV}$	$E_{\text{med}} = \beta\beta(734)$
			$1247 \text{ keV} < E_{\max} < 1267 \text{ keV}$	$E_{\max} = \gamma(1257)$
$2\nu\beta\beta$				
Scenario	Multiplicity	Score	Cut	Selection
2A0 - 2B1	2	40%	$620 \text{ keV} < E_{\min} < 1150 \text{ keV}$	$E_{\min} = \beta\beta(0-734) + \gamma(536)$
			$1220 \text{ keV} < E_{\max} < 1300 \text{ keV}$	$E_{\max} = \gamma(1257)$
2A2 - 2B3	2	22%	$523 \text{ keV} < E_{\min} < 573 \text{ keV}$	$E_{\min} = \gamma(536)$
			$1360 \text{ keV} < E_{\max} < 1990 \text{ keV}$	$E_{\max} = \beta\beta(0-734) + \gamma(1257)$
3A0	3	26%	$400 \text{ keV} < E_{\min} < 523 \text{ keV}$	$E_{\min} = \beta\beta(0-734)$
			$523 \text{ keV} < E_{\text{med}} < 573 \text{ keV}$	$E_{\text{med}} = \gamma(536)$
			$1779 \text{ keV} < E_{\text{med}} + E_{\max} < 1807 \text{ keV}$	$E_{\max} = \gamma(1257)$

- Blind selections  
(defined on MC simulations)  
in ordered energy space
- Expected background  
shape: flat or linear
- No background peaks  
expected

# Excited states analysis

## Containment efficiency



Once the selection cuts are fixed, the efficiency is obtained from a fit to Monte Carlo simulations of signal events

# Excited states analysis

## Model

$$\log \mathcal{L}_{s,ds}(\vec{E} | \vec{\theta}, H_{S+B}) = -(\lambda_{s,ds}^{(S)} + \lambda_{s,ds}^{(B)}) + \sum_{ev \in (s,ds)} \log \left[ \lambda_{s,ds}^{(S)} \frac{M\Delta t_{ch}}{M\Delta t_{ds}} pdf^{(S)}(\vec{E}) + \lambda_{s,ds}^{(B)} \frac{M\Delta t_{ch}}{M\Delta t_{ds}} pdf^{(B)}(\vec{E}) \right]$$

- Bayesian Analysis (BAT)
- Likelihood model: linear background, peak for  $\beta\beta/\gamma$
- Unbinned fit on physical range (non-negative rate), uniform prior on  $\Gamma$
- Systematics addressed repeating fit with additional nuisance parameters

$$\epsilon_s = \left[ \sum_p BR_p \cdot \frac{[N_{MC}^{(sel)}]_p^{(s)}}{[N_{MC}^{(tot)}]_p} \right] \epsilon_{cut}^M \epsilon_{acc}$$

leading efficiency term

Process	Containment	Cut	Accidentals	Total
$0\nu\beta\beta$	10.0%	88.7%	98.7%	8.7%
$2\nu\beta\beta$	6.8%			5.9%

**Effective efficiencies**

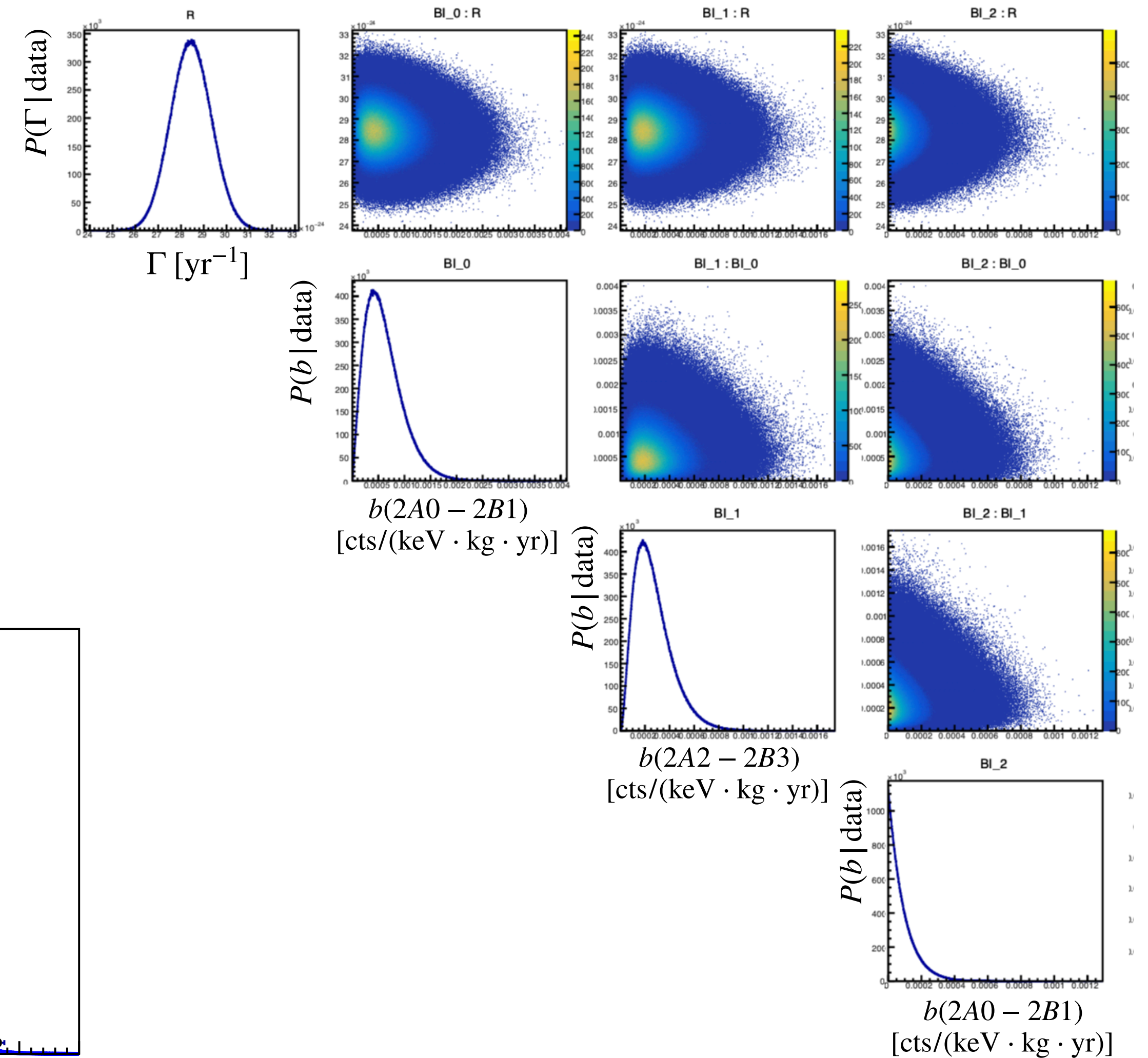
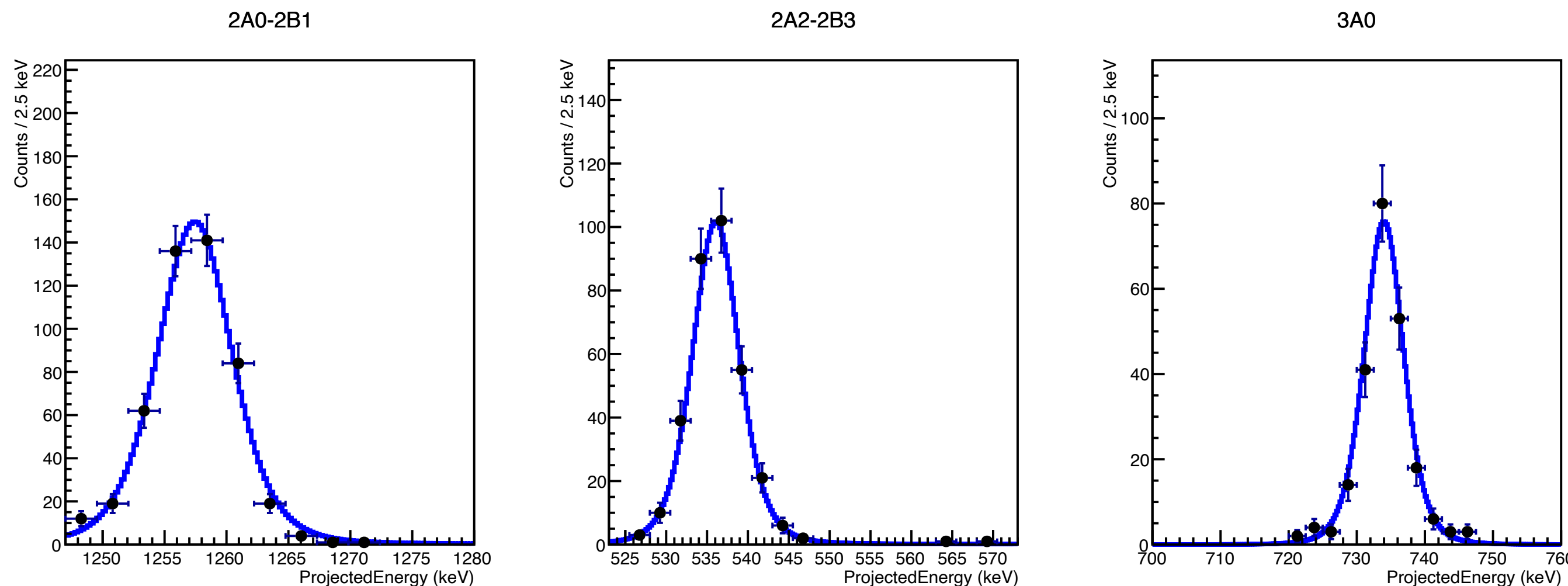
containment summed over signatures  
other components exposure averaged



# Excited states analysis

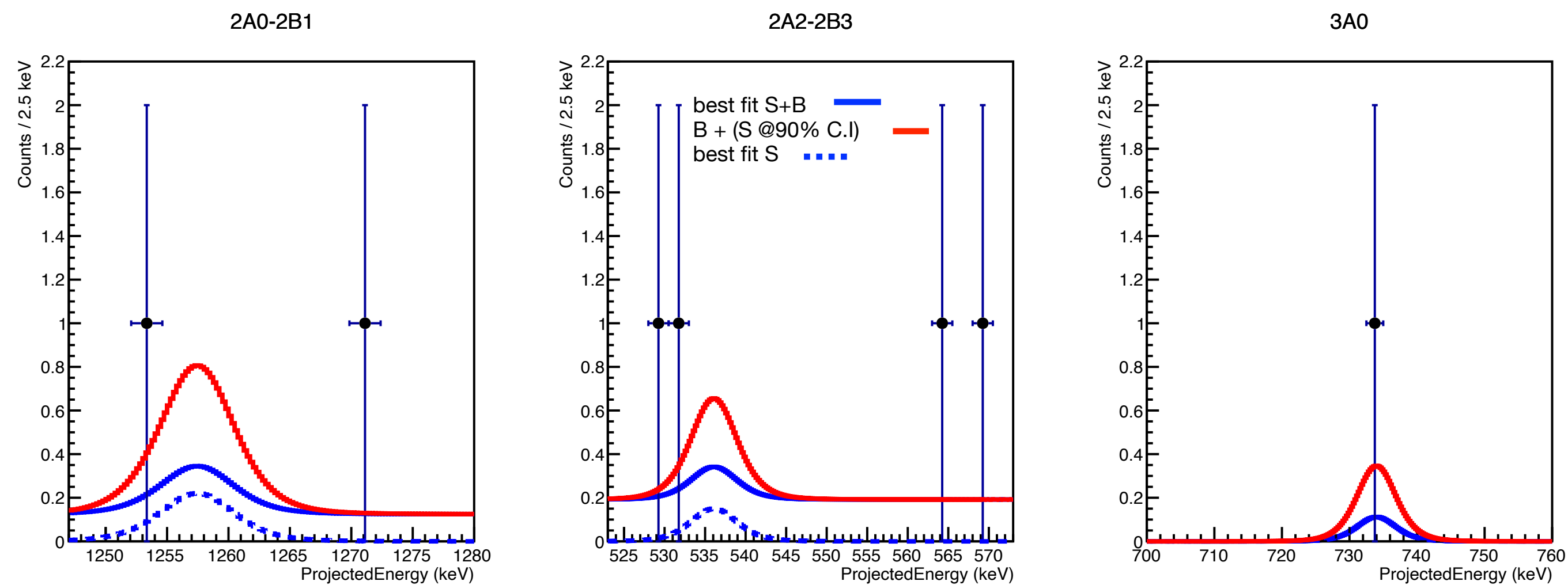
## Blinded fit ( $0\nu\beta\beta$ )

- Monte Carlo signal events injected into physics data
- Background shape validation
- Unbiased background estimation
- Allows a preliminary sensitivity computation

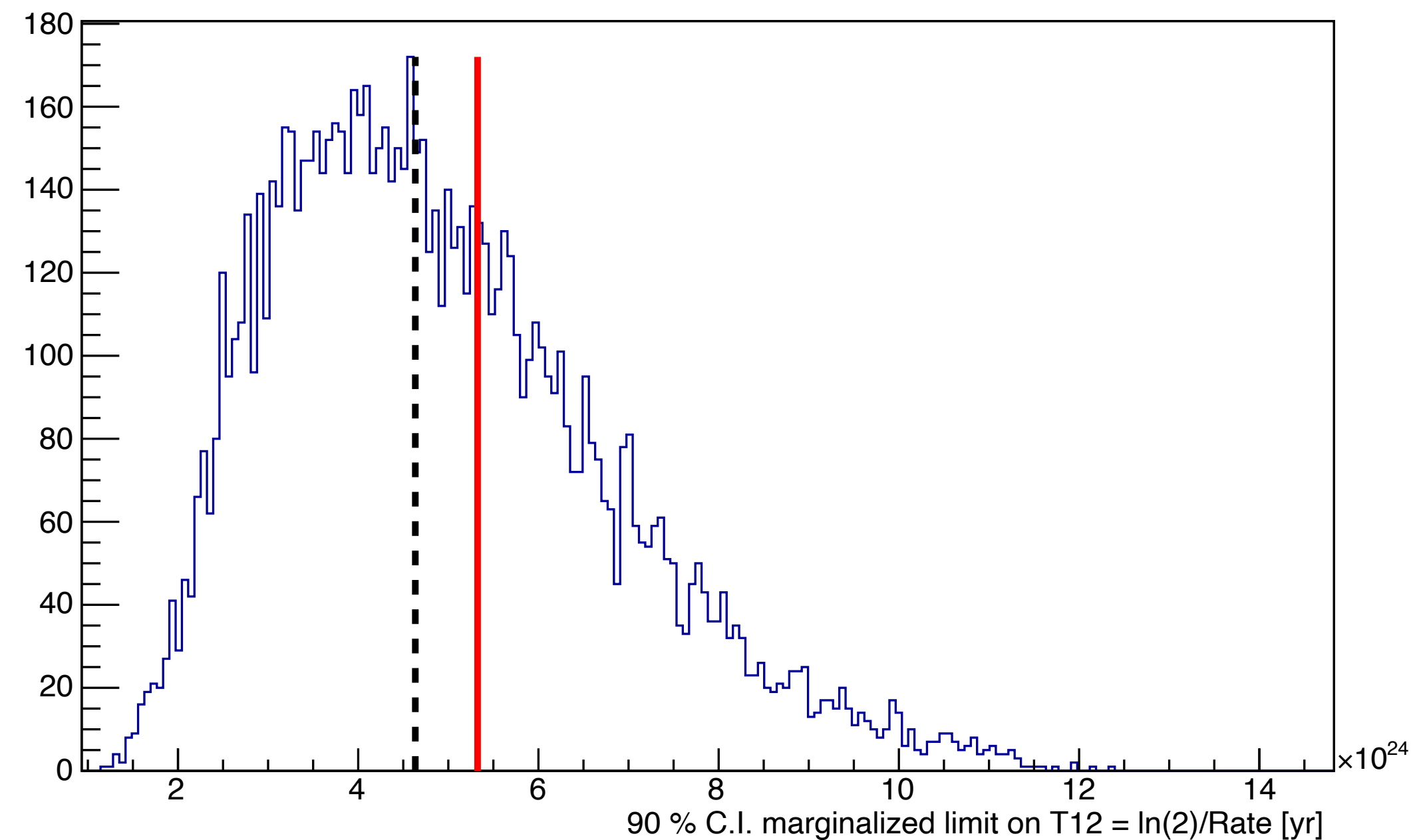


# Excited states analysis

## $0\nu\beta\beta$ unblinded fit and sensitivity



Expected limit setting sensitivity - Median:  $4.63 \times 10^{24}$  [yr] MAD:  $1.25 \times 10^{24}$  [yr]



Marg. Post.	Mean	St. Dev.	Units
Rate	6.37	4.85	$10^{-26}$ 1/yr
b (2A0-2B1)	2.15	1.38	$10^{-4}$ cts/(keV kg yr)
b (2A2-2B3)	2.56	1.20	$10^{-4}$ cts/(keV kg yr)
b (3A0)	5.67	5.42	$10^{-5}$ cts/(keV kg yr)

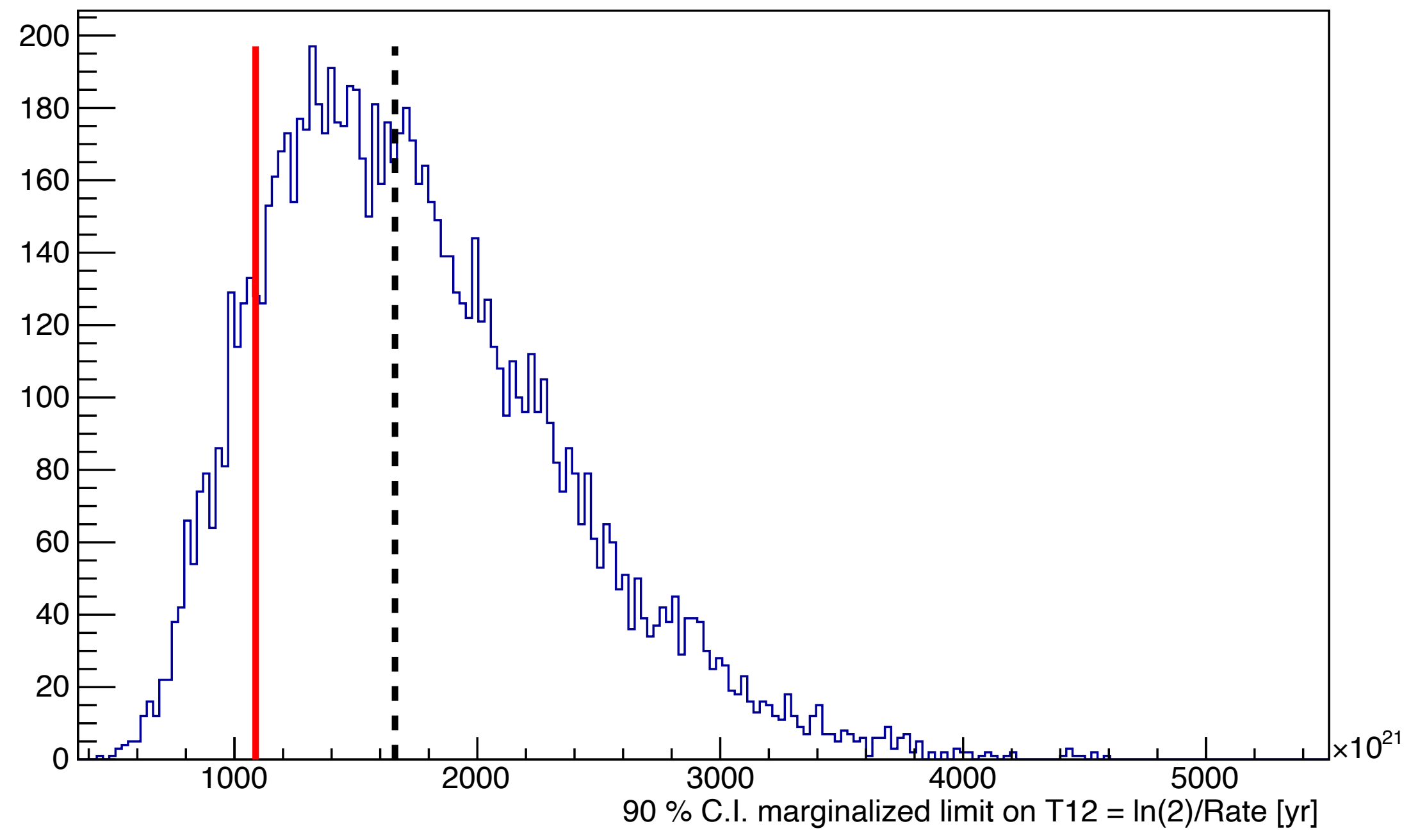
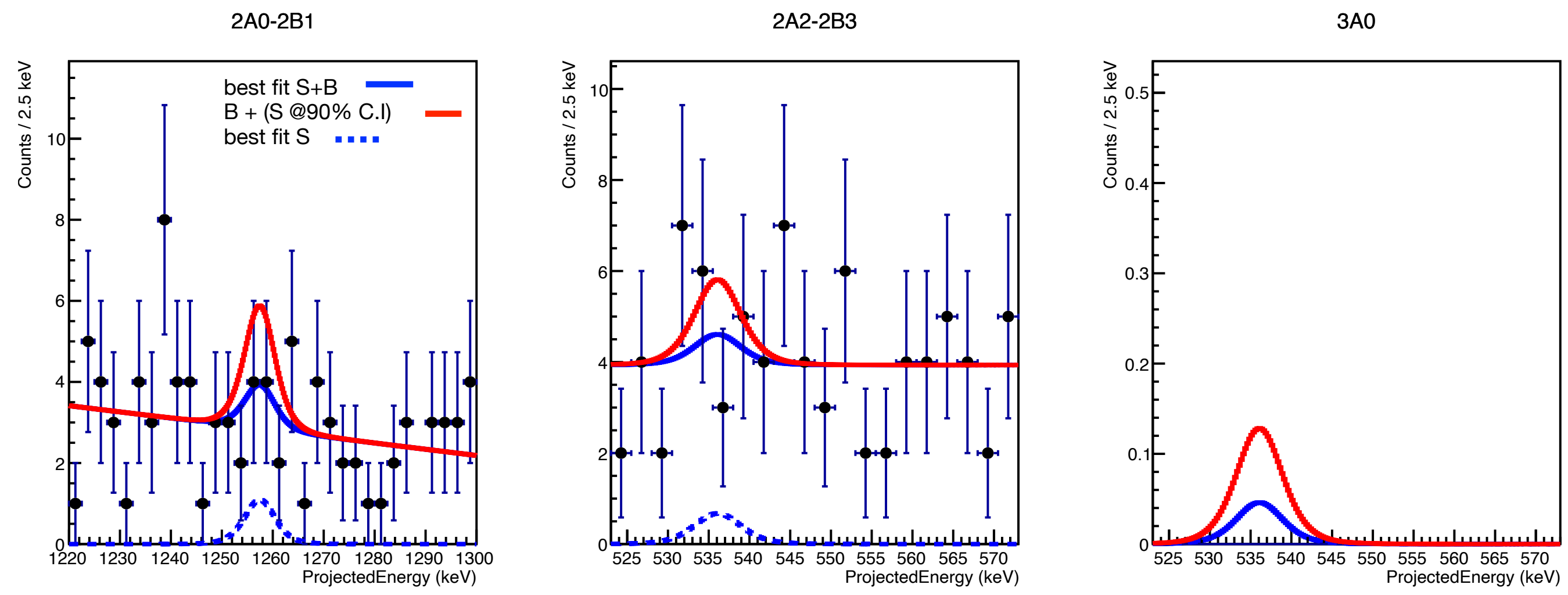
$$T_{1/2}^{0\nu} > 5.4 \times 10^{24} \text{ yr (90 \% C.I.)}$$

systematics not included

# Excited states analysis

## $2\nu\beta\beta$ unblinded fit and sensitivity

Expected limit setting sensitivity - Median:  $1.66 \times 10^{24}$  [yr] MAD:  $3.99 \times 10^{23}$  [yr]



Marg. Post.	Mean	St. Dev.	Units
Rate	3.33	2.18	$10^{-25}$ 1/yr
b (2A0-2B1)	3.01	0.33	$10^{-3}$ cts/(keV kg yr)
b (2A2-2B3)	4.23	0.49	$10^{-3}$ cts/(keV kg yr)
b (3A0)	5.37	5.36	$10^{-5}$ cts/(keV kg yr)
slope (3A0)	-5.17	4.23	$10^{-3}$ 1/keV

$$T_{1/2}^{2\nu} > 1.1 \times 10^{24} \text{ yr (90 \% C.I.)}$$

systematics not included

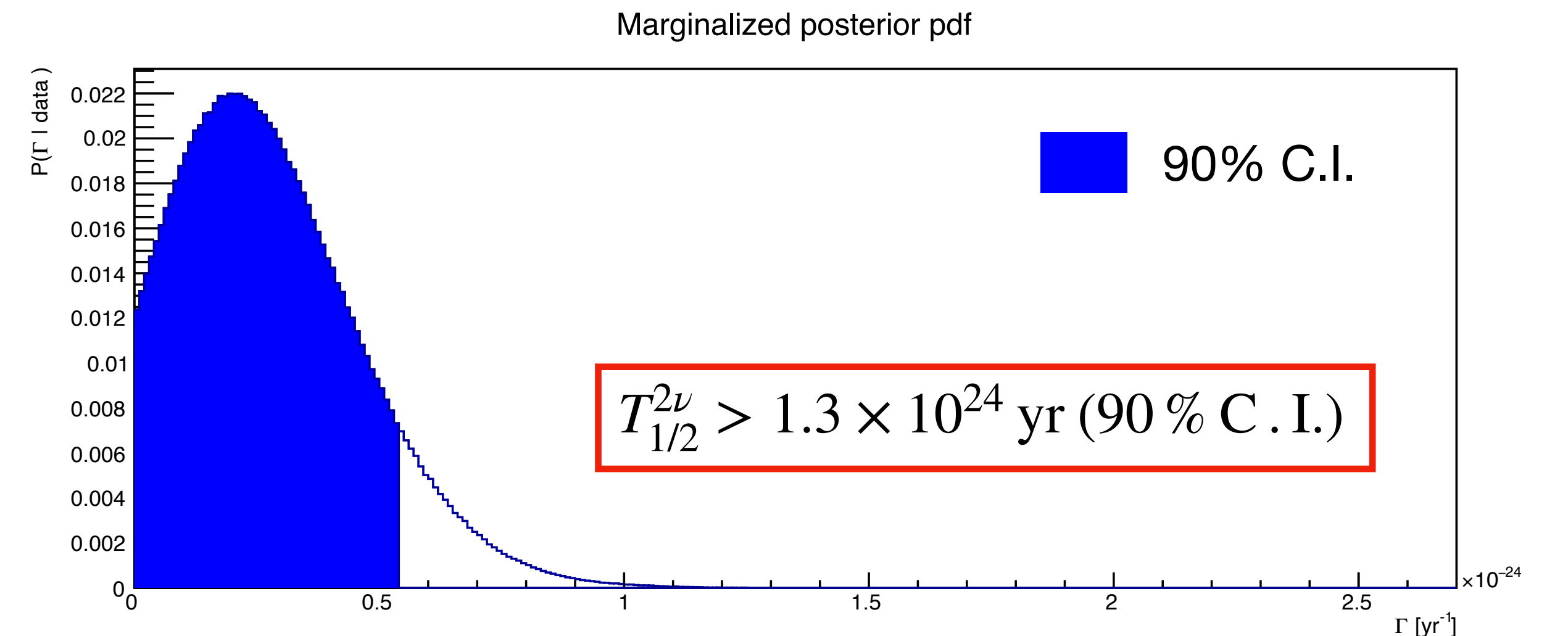
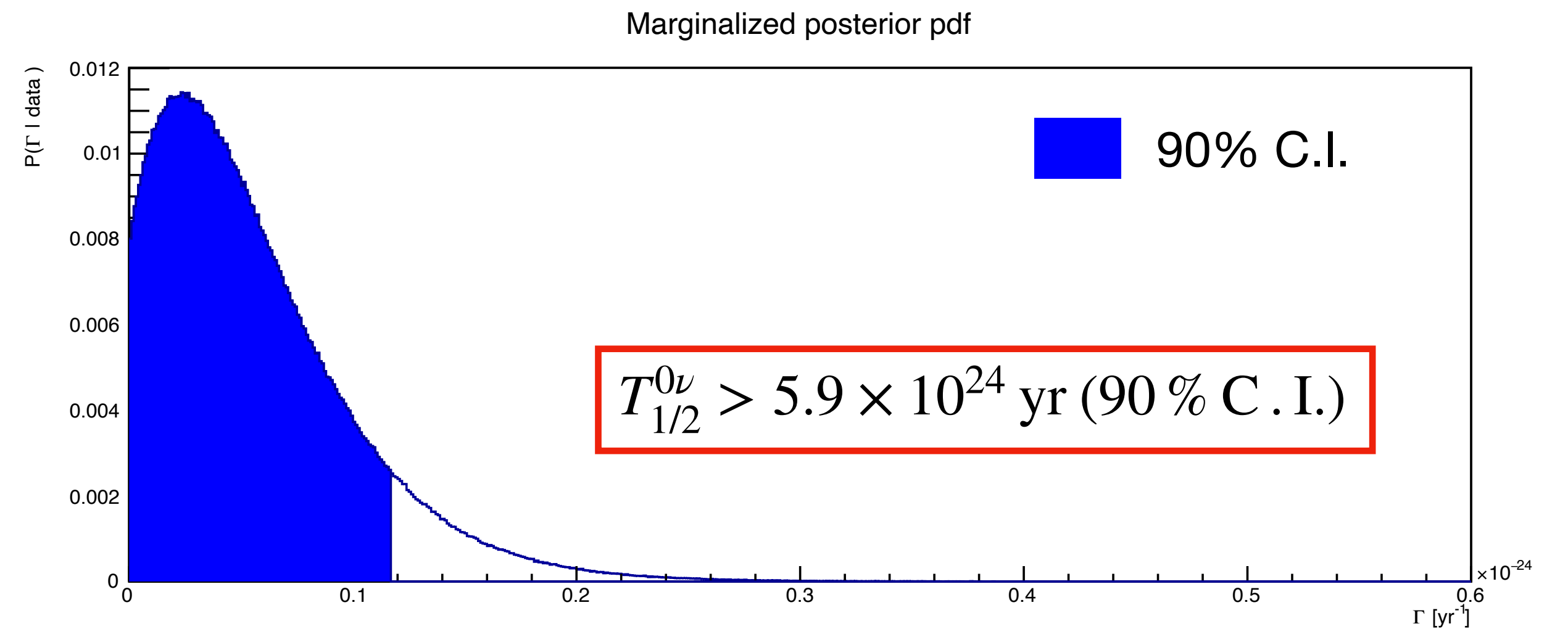


# Excited states analysis

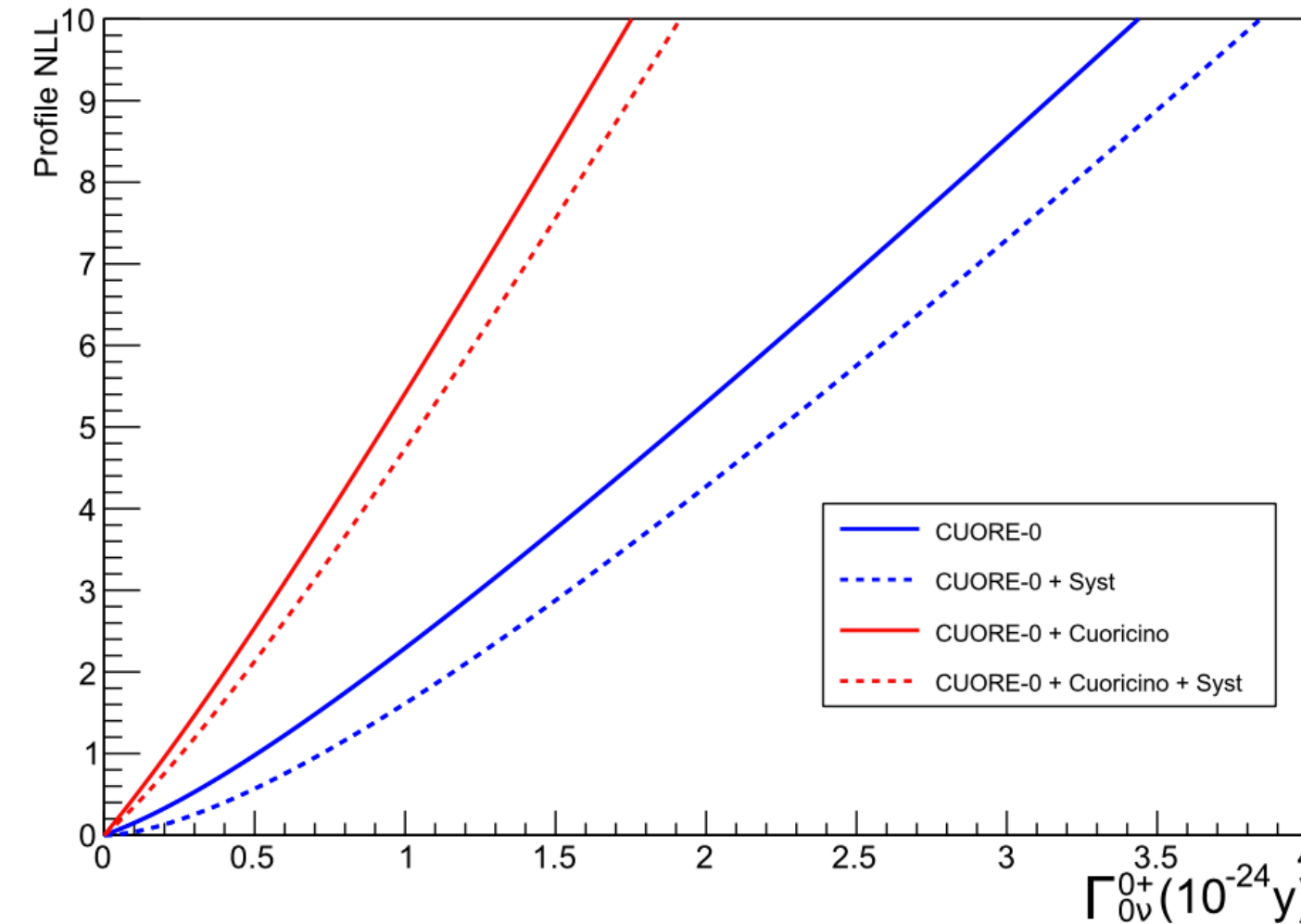
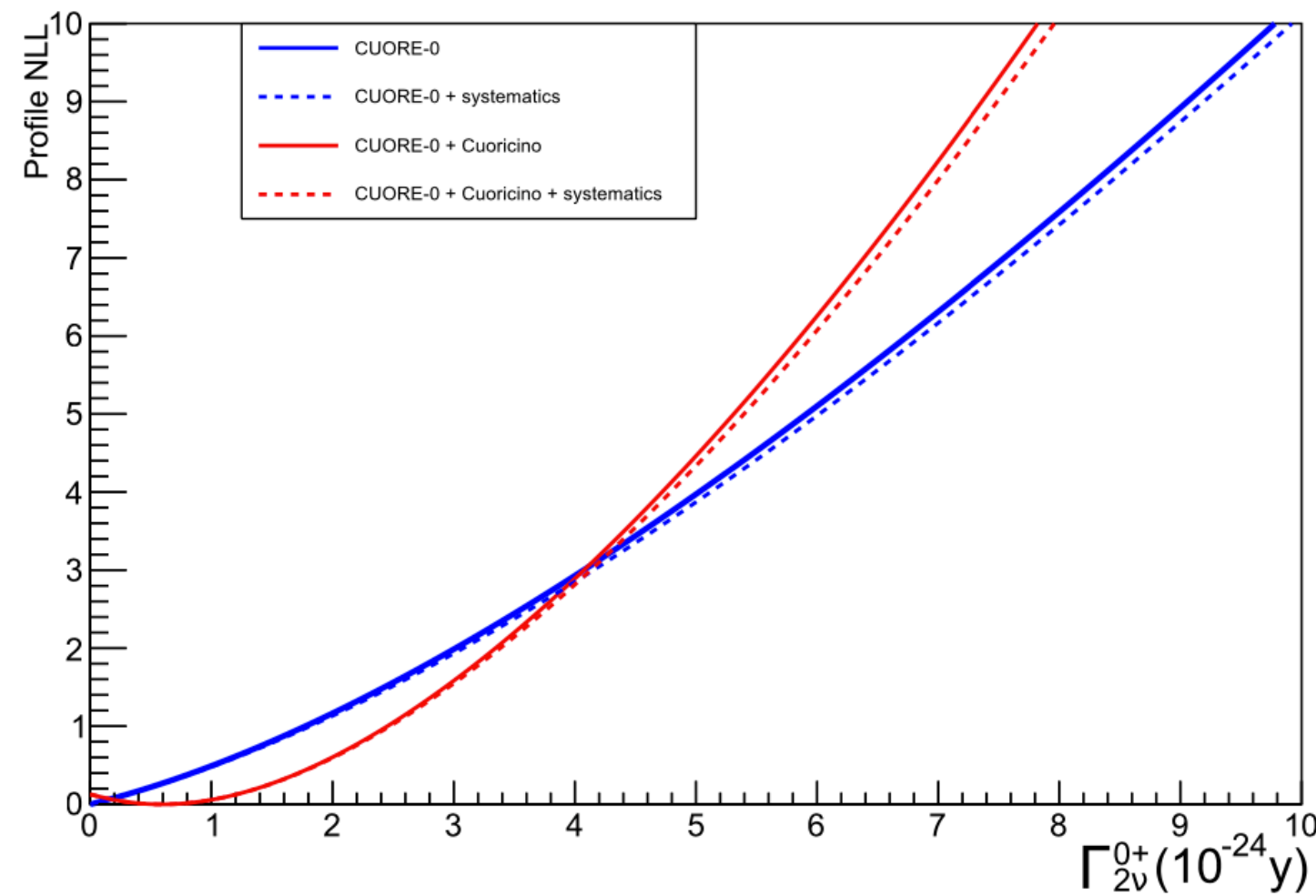
## Systematics

Nuisance Parameter	Prior	$0\nu\beta\beta$ $\Delta T_{1/2}^{90}$	$2\nu\beta\beta$ $\Delta T_{1/2}^{90}$
Detector response	multivariate	10.1%	17.6%
Cut efficiency	gaussian	-0.1%	< 0.1%
PSA efficiency	uniform	-0.7%	-0.2%
Accidental coincidences	gaussian	-0.4%	0.1%
Containment efficiency	gaussian	-0.6%	-0.4%
Isotopic abundance	gaussian	-0.1%	-0.2%
Combined	multivariate	10.1%	17.2%

[G. Fantini et al., paper in preparation]



# Excited states analysis



$$T_{1/2}^{2\nu} > 2.5 \times 10^{23} \text{ yr, 90 \% C.L.}$$

$$T_{1/2}^{0\nu} > 1.4 \times 10^{24} \text{ yr, 90 \% C.L.}$$

[C. Alduino *et al.*, *Eur. Phys. J. C* 79, 795 (2019)]

$$T_{1/2}^{2\nu} > 1.3 \times 10^{24} \text{ yr (90 \% C.I.)}$$

$$T_{1/2}^{0\nu} > 5.9 \times 10^{24} \text{ yr (90 \% C.I.)}$$

This work: [G. Fantini *et al.*, paper in preparation]

$$T_{1/2}^{2\nu} = (7.2 - 16) \times 10^{24} \text{ yr} \quad \text{QRPA}$$

[P. Pirinen, J. Suhonen, *Phys. Rev. C* 91, 054309 (2015)]

$$T_{1/2}^{2\nu} = 2.2 \times 10^{25} \text{ yr}$$

IBM-II + experimental input

[B. Lehnert, 10.1051/epjconf/20159301025]

# Summary

- Original contributions to the CUORE analysis
- New result on the double beta decay to excited states both in the neutrino-less and Standard Model mode
- No evidence for signal in either mode
- A factor 5 more stringent Bayesian limit with respect to the current most sensitive search (CUORE-0)



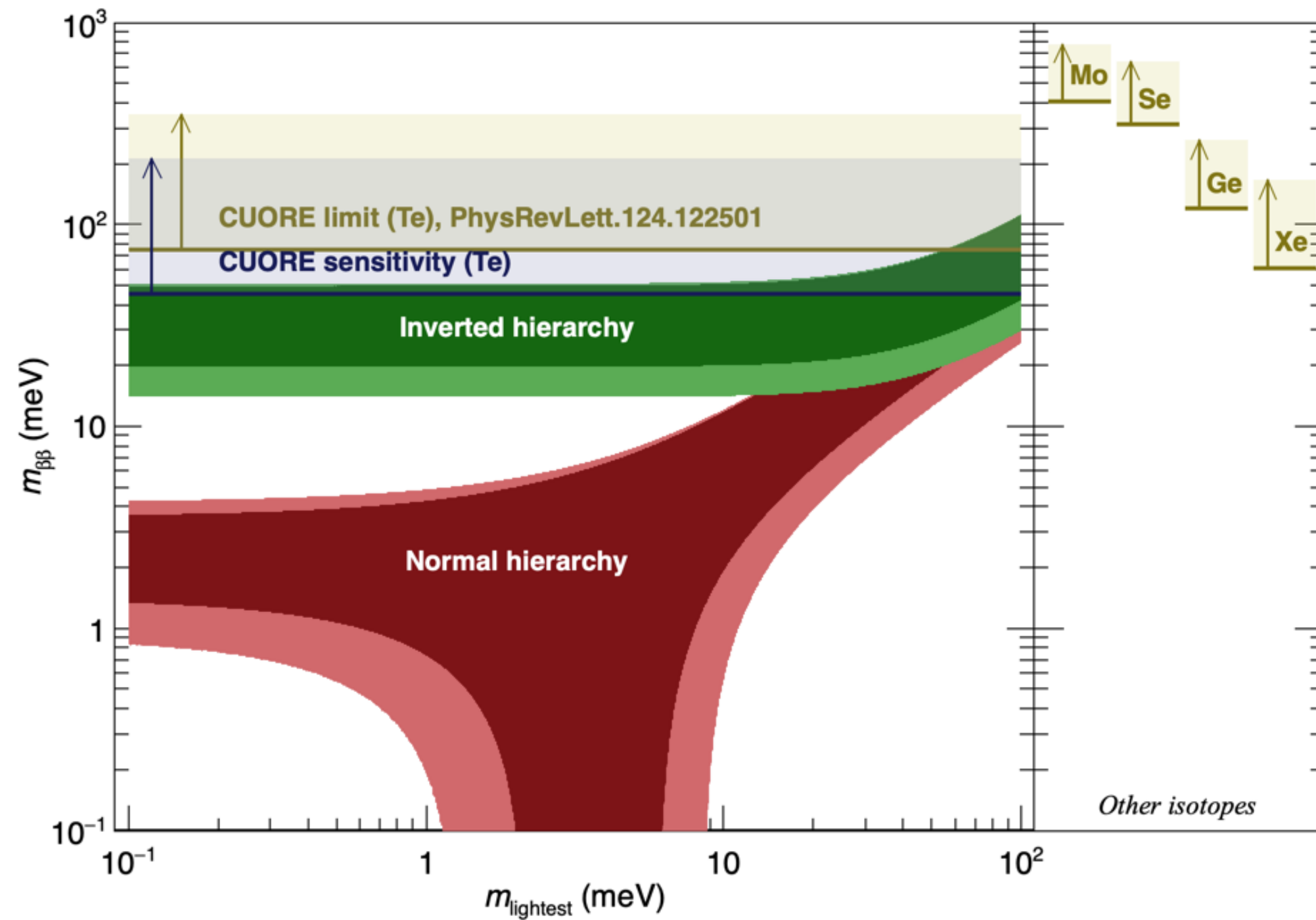




# Backup

# Effective Majorana Mass Interpretation

## NMEs Used



$$m_{\beta\beta} < 75 - 350 \text{ meV}$$

- [57] J. Engel and J. Menéndez, *Rept. Prog. Phys.* **80**, 046301 (2017).
- [58] J. Barea, J. Kotila, and F. Iachello, *Phys. Rev.* **C91**, 034304 (2015).
- [59] F. Šimkovic *et al.*, *Phys. Rev.* **C87**, 045501 (2013).
- [60] J. Hyvärinen and J. Suhonen, *Phys. Rev.* **C91**, 024613 (2015).
- [61] J. Menéndez *et al.*, *Nucl. Phys.* **A818**, 139 (2009).
- [62] T. R. Rodriguez and G. Martinez-Pinedo, *Phys. Rev. Lett.* **105**, 252503 (2010).
- [63] N. López Vaquero, T. R. Rodríguez, and J. L. Egido, *Phys. Rev. Lett.* **111**, 142501 (2013).
- [64] J. M. Yao *et al.*, *Phys. Rev.* **C91**, 024316 (2015).
- [65] M. T. Mustonen and J. Engel, *Phys. Rev.* **C87**, 064302 (2013).
- [66] A. Neacsu and M. Horoi, *Phys. Rev.* **C91**, 024309 (2015).
- [67] A. Meroni, S. T. Petcov, and F. Simkovic, *JHEP* **02**, 025 (2013).



# See-saw mechanism

$$\mathcal{L}_{D+M} = -\bar{N}_L^c M N \quad N_L = \begin{bmatrix} \nu_L \\ \nu_R^c \end{bmatrix} \quad M = \begin{bmatrix} m_L & m_D \\ m_D & m_R \end{bmatrix}$$

Assuming  $SU(3)_C \times SU(2)_L \times U(1)_Y$  gauge invariance and renormalizability  $m_L = 0$

$m_D$  of the order of the Higgs v.e.v.,  $m_R$  arbitrarily large

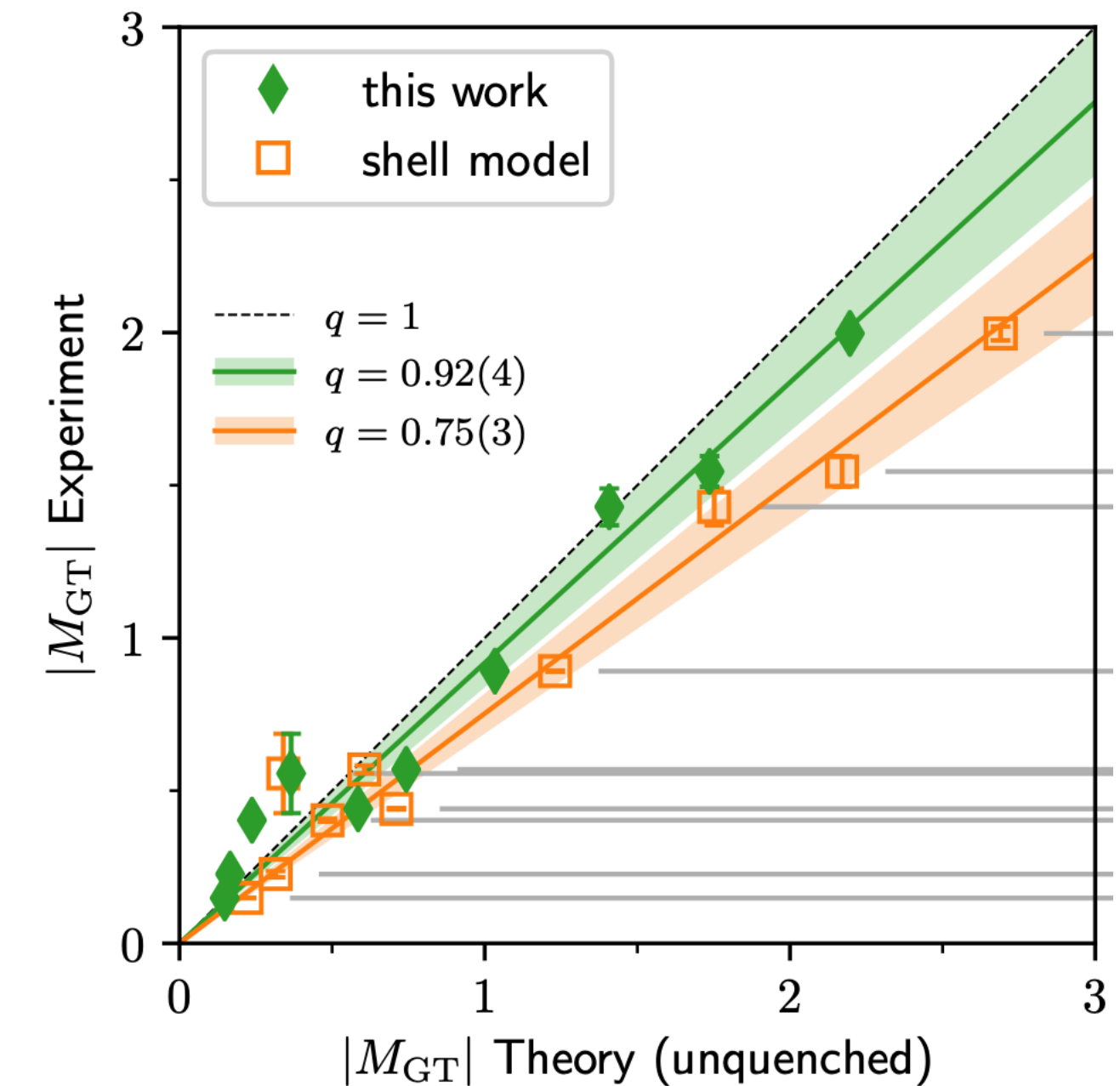
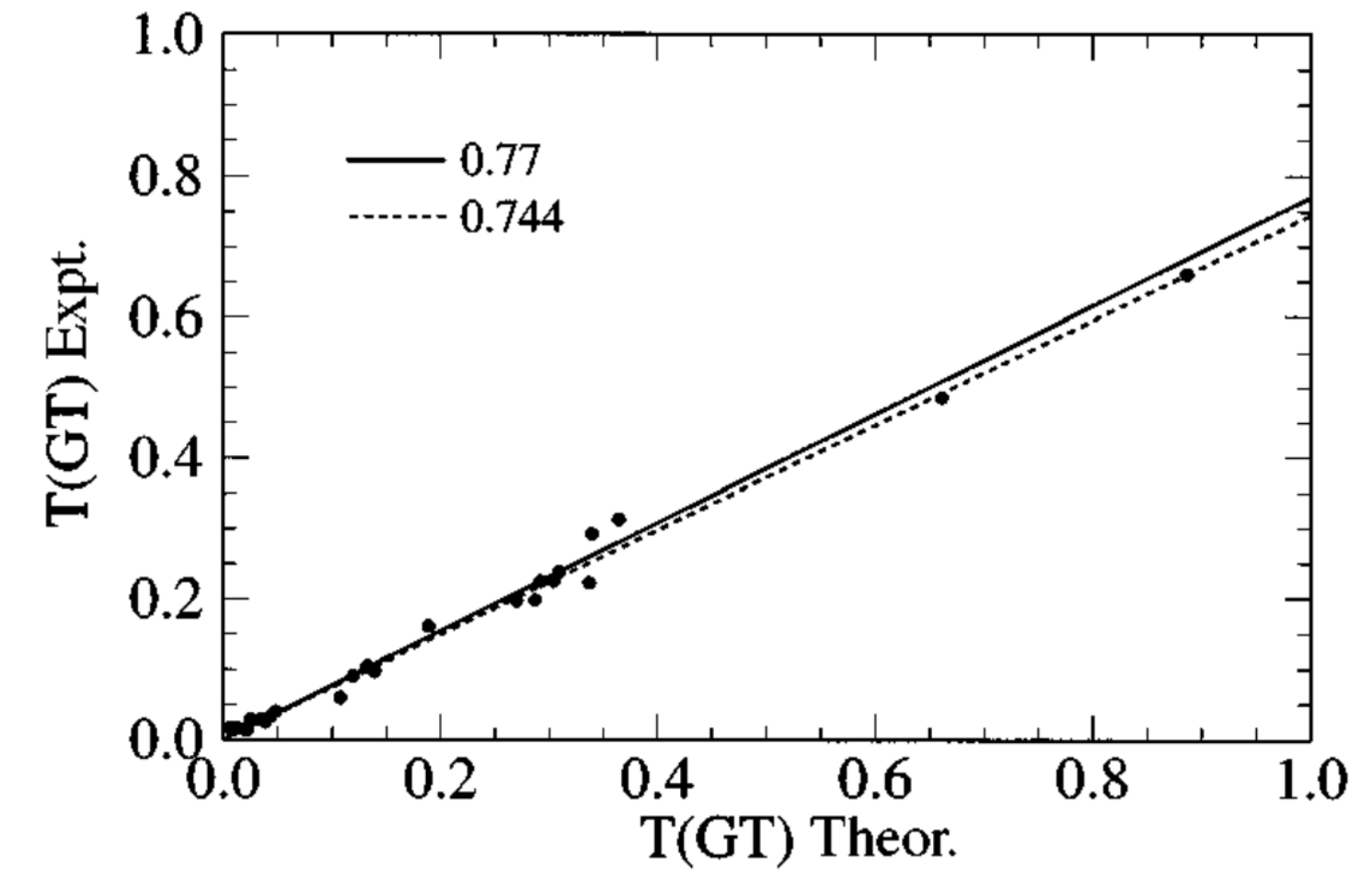
$$m_{\pm} = \frac{m_R}{2} \pm \frac{m_R}{2} \sqrt{1 + \frac{4m_D^2}{m_R^2}} \quad \xrightarrow{m_R \gg m_D} \quad \begin{aligned} m_+ &\sim m_D^2/m_R \\ m_- &\sim m_R \end{aligned}$$

# $g_A$ quenching

- Parameterization of theory vs experiment mismatch
- Difference between the Gamow-Teller calculation and experimental results in beta-decay motivates quenching

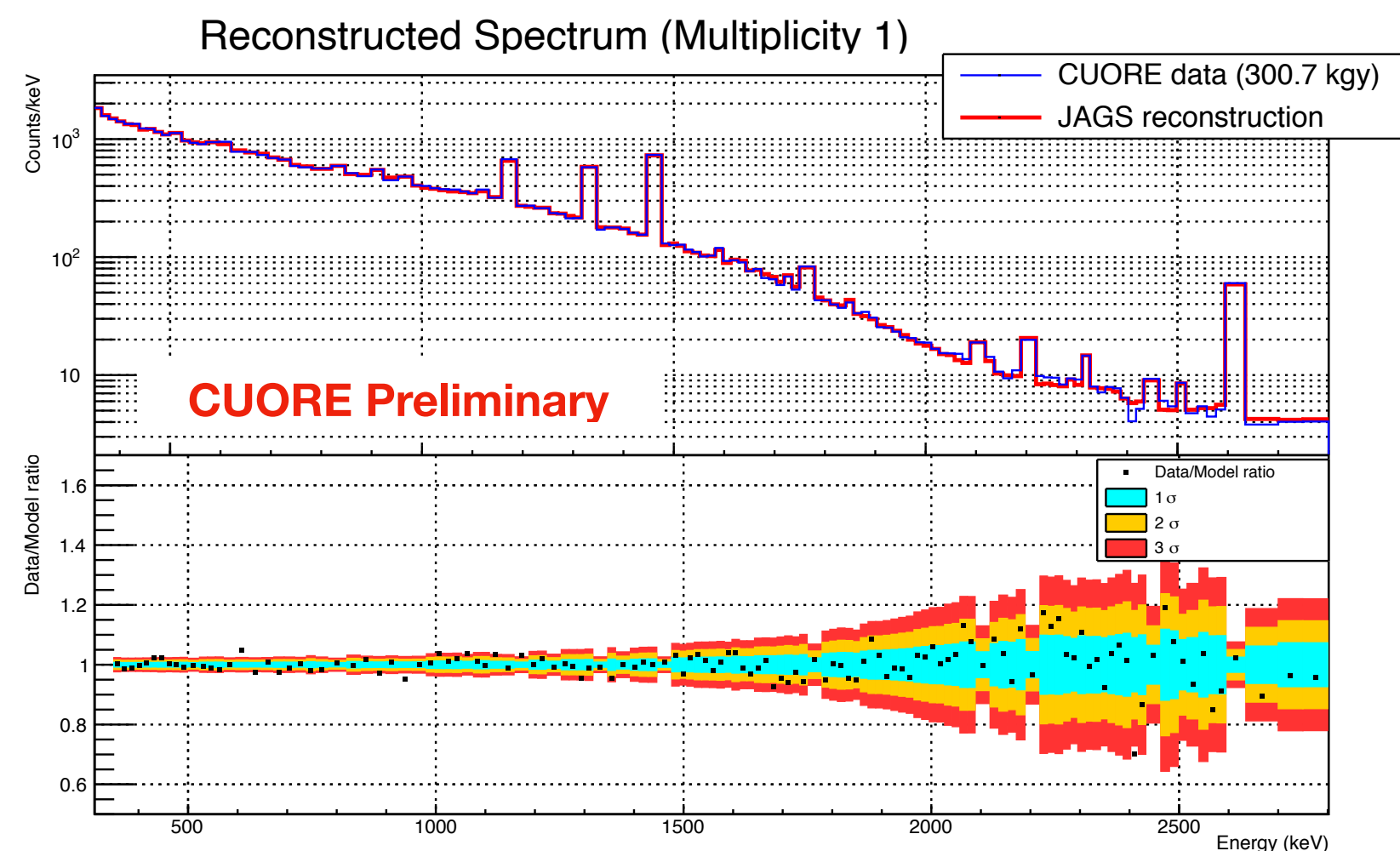
$$g_A^{eff} = q g_A$$

- Important uncertainty in neutrino-less double beta decay NME
- Ab-initio calculations including two-body / meson-exchange currents and additional nuclear correlations do not need any “quenching”

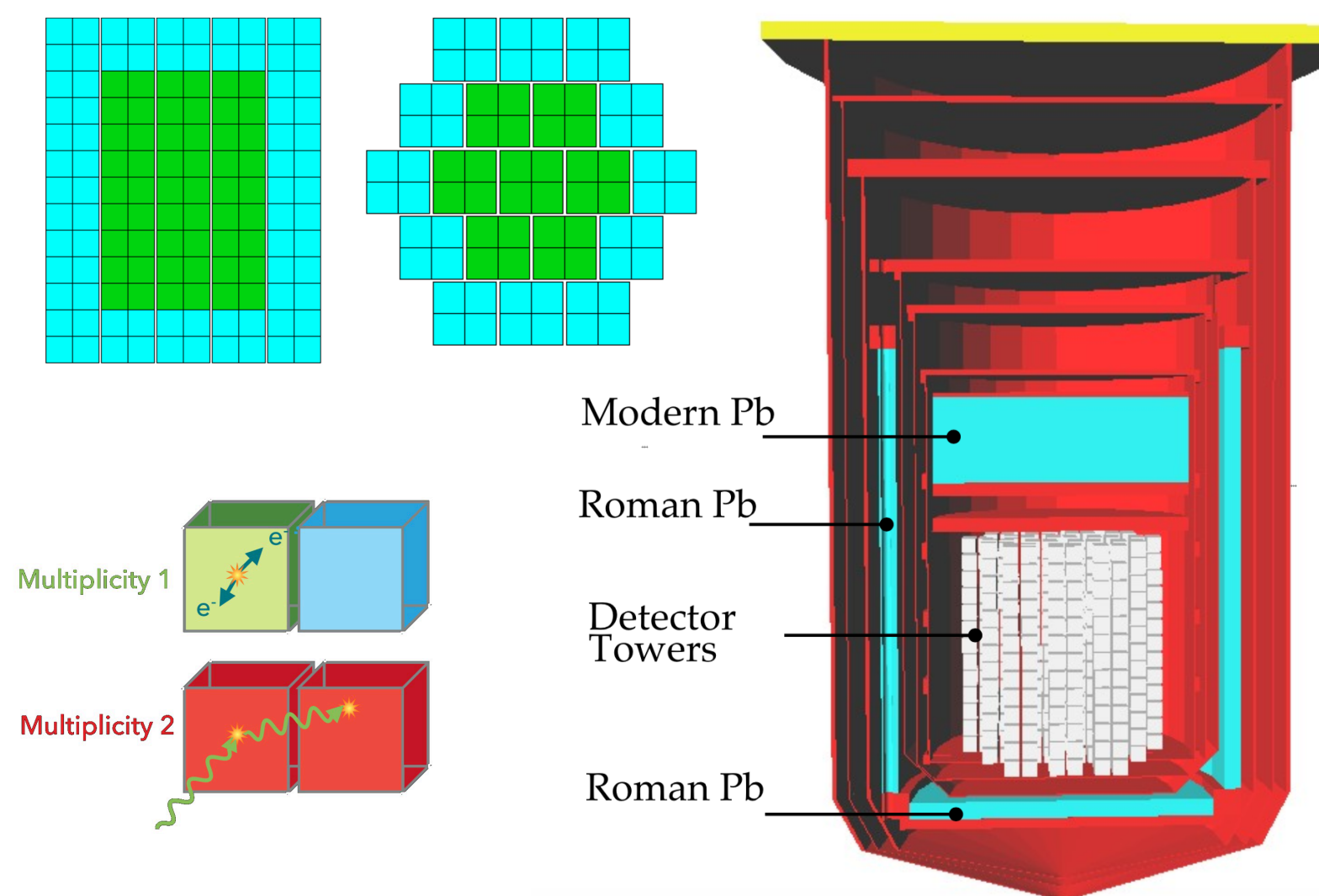
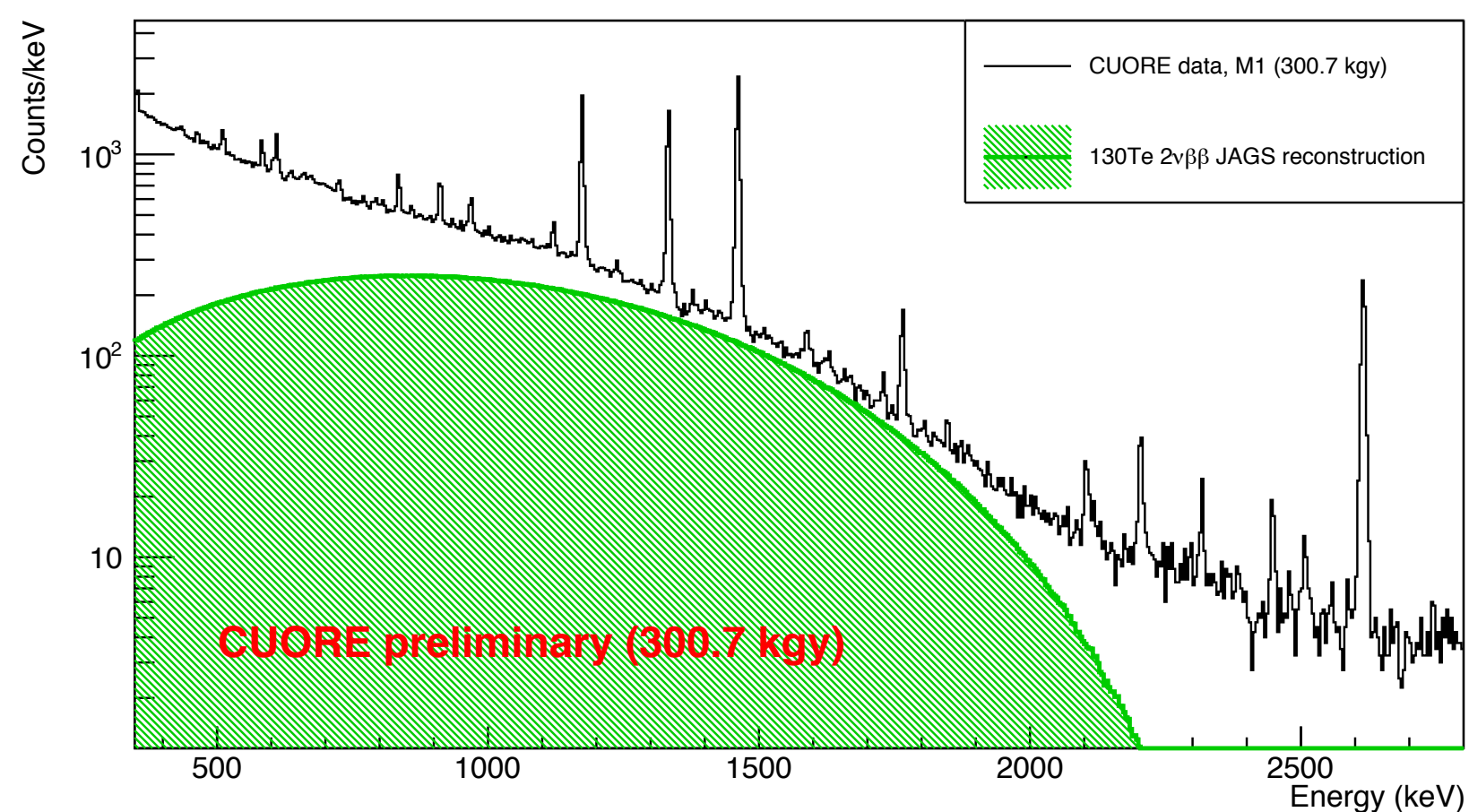


[Gysbers *et al.*, Nature Physics 15, 428-431 (2019)]

# Background model



$^{130}\text{Te } 2\nu\beta\beta - \text{M1}$



- GEANT4 simulation + detector response function to produce expected spectra
- 62 sources considered, Bayesian fit flat priors (except muons)
- Coincidences and self-shielding exploited to constrain source position

$$T_{1/2}^{2\nu} = [7.71^{+0.08}_{-0.06}(\text{stat.})^{+0.17}_{-0.15}(\text{syst.})] \times 10^{20} \text{ yr}$$



# Baryon Asymmetry in the Universe

## Sakharov Conditions

- Baryon number violation
- C-symmetry and CP symmetry violation
- Interactions out of thermal equilibrium

$$\eta = \frac{n_B - \bar{n}_B}{n_\gamma} \sim 10^{-10}$$

CP violation in the quark sector is not enough to explain the observed asymmetry

$$\frac{D_{CKM}}{T_{sph}^{12}} \sim 10^{-20} \ll \eta$$

Baryogenesis via leptogenesis is a possibility