

# Bose-Einstein condensation with Optimal Rate for Trapped Bosons in the Gross-Pitaevskii Regime

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joint work with C. Brennecke and B. Schlein.

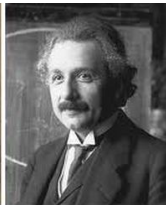


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# Bose-Einstein condensate



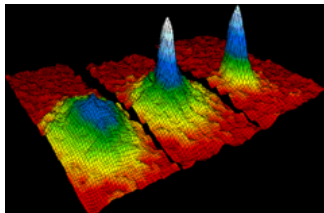
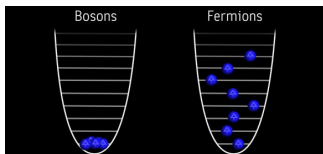
Satyendra Nath Bose



Albert Einstein

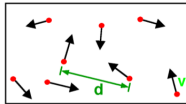


Symmetric wavefunction



Velocity-distribution rubidium

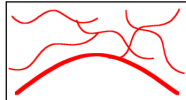
# Bose-Einstein condensation: Heuristics



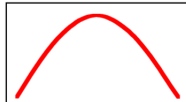
High  
Temperature  $T$ :  
thermal velocity  $v$   
density  $d^{-3}$   
"Billiard balls"



Low  
Temperature  $T$ :  
De Broglie wavelength  
 $\lambda_{dB} = h/mv \propto T^{-1/2}$   
"Wave packets"



$T = T_{crit}$ :  
Bose-Einstein  
Condensation  
 $\lambda_{dB} \approx d$   
"Matter wave overlap"



$T = 0$ :  
Pure Bose  
condensate  
"Giant matter wave"



# Reduced one-particle density matrix

Let  $\psi_N \in L_s^2(\mathbb{R}^{3N})$  be a normalized symmetric (i.e. bosonic) wavefunctions. We define the **reduced one-particle density** of  $\psi_N$  to be

$$\gamma_N^{(1)}(x, y) := \int_{\mathbb{R}^{3(N-1)}} \psi_N(x, X) \overline{\psi_N(y, X)} dx_2 \dots dx_N.$$

For an observable  $\mathcal{O}$  that only depends on 1-particle, i.e. is of the form

$$\mathcal{O} = \text{Sym}_N(\mathcal{O}^{(1)} \otimes Id^{\otimes(N-1)}),$$

then

$$\langle \psi_N, \mathcal{O} \psi_N \rangle = N \text{tr}(\mathcal{O}^{(1)} \gamma_N^{(1)}).$$

For example, the kinetic energy  $\mathcal{K} = \frac{\hbar^2}{2m} \sum_{j=1}^N -\Delta_{x_j}$  is a one-particle observable

However, the interaction energy  $\mathcal{V}_N = \sum_{1 \leq i < j \leq N} V_N(x_i - x_j)$  would be a 2-particle observable.

# Definition complete asymptotic B.E.C.

## Theorem 1

Let  $(\gamma_N^{(1)})_{N \geq 1}$  be a the sequence of one-particle reduced density matrices corresponding to the sequence  $(\gamma_N)_{N \geq 1}$  of  $N$ -body bosonic states on the Hilbert space  $L^2(\Omega)^{\otimes N}$ . Let  $\varphi \in L^2(\Omega)$  with  $\|\varphi\|_2 = 1$ . Then the following are equivalent:

- 1  $\gamma_N^{(1)} \xrightarrow{N \rightarrow \infty} |\varphi\rangle\langle\varphi|$  weakly-\* in trace class
- 2  $\gamma_N^{(1)} \xrightarrow{N \rightarrow \infty} |\varphi\rangle\langle\varphi|$  in the Hilbert-Schmidt norm
- 3  $\gamma_N^{(1)} \xrightarrow{N \rightarrow \infty} |\varphi\rangle\langle\varphi|$  in the trace norm
- 4  $\langle\varphi, \gamma_N^{(1)} \varphi\rangle \xrightarrow{N \rightarrow \infty} 1$ .

We say that  $(\gamma_N^{(1)})_{N \geq 1}$  exhibits **complete asymptotic B.E.C** in the condensate wave function  $\varphi \in L^2(\Omega)$  iff any of the equivalent statements in the above theorem hold.

# Gross-Pitaevskii regime

We consider a system of  $N$ -particles with Hamiltonian

$$H_N = \sum_{j=1}^N [-\Delta_{x_j} + V_{ext}(x_j)] + \sum_{1 \leq i < j \leq N} N^2 V(N(x_i - x_j))$$

- Here  $V_{ext}$  is a trapping potential, i.e.  
 $\lim_{R \rightarrow \infty} \sup_{|x| \leq R} V_{ext}(x) = \infty$  and spherically symmetric.
- Repulsive, radially symmetric and short-range interaction, i.e.  
 $V \geq 0$  is spherically symmetric with compact support.

More technical assumptions

- $V_{ext} \in C^3(\mathbb{R}^3; \mathbb{R})$
- $\exists C > 0 \forall x, y \in \mathbb{R}^3 : V_{ext}(x+y) \leq C(V_{ext}(x) + C)(V_{ext}(y) + C)$ ,  $\nabla V_{ext}, \Delta V_{ext}$  have at most exponential growth as  $|x| \rightarrow \infty$ .

We also have the **Gross-Pitaevskii functional**

$$\mathcal{E}_{GP}(\psi) = \int_{\mathbb{R}^3} (|\nabla \psi(x)|^2 + V_{ext}(x)|\psi(x)|^2 + 4\pi \mathbf{a}_0 |\psi(x)|^4) dx.$$

## Theorem 2 (Brennecke, Schlein, Schraven '21)

Let  $V$  and  $V_{ext}$  be "sufficiently nice" and let  $E_N$  denote the ground state energy of  $H_N$ . Then, there exists a constant  $C > 0$  such that

$$E_N \geq N\mathcal{E}_{GP}(\varphi) - C \quad (1)$$

and

$$H_N \geq N\mathcal{E}_{GP}(\varphi) + C^{-1} \sum_{i=1}^N (1 - |\varphi\rangle\langle\varphi|_i) - C.$$

In particular, if  $\psi_N \in L_s^2(\mathbb{R}^{3N})$  with  $\|\psi_N\| = 1$  is an sequence of approximate ground states such that

$$\langle\psi_N, H_N\psi_N\rangle \leq N\mathcal{E}_{GP}(\varphi) + \zeta,$$

for a  $\zeta > 0$ , then the reduced density  $\gamma_N^{(1)}$  associated with  $\psi_N$  satisfies

$$1 - \langle\varphi, \gamma_N^{(1)}\varphi\rangle \leq \frac{(C + \zeta)}{N}.$$

Our result extends the following previous results:

- Lieb, Seiringer '02: Show convergence, but without optimal rate (both in the box and in the trap).
- Boccato, Brennecke, Cenatiempo, Schlein '20: Optimal rate in the box.
- Nam, Napiórkowski, Ricaud, Triay '20: Optimal rate in both the box and the trap. Need to impose smallness condition for the scattering length for the lower bound.